

The impact of axon wiring costs on small neuronal networks

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Abstract. Recent papers by D. Chklovskii and E.M. Izhikevich suggest that wiring costs may play a significant role in the physical layout and function of neuronal structures. About eighty years ago, in his paper on the relationship between diameter and branching angles in trees, C.D. Murray proposed the volume as the cost function which dictates growth. O. Shefi et al. grafted this idea into neuroscience as a possible optimisation mechanism. Our paper presents computational experiments on the impact of wiring cost functions proposed by D. Chklovskii and O. Shefi et al. when applied to interneuronal connections in small ML neuronal networks.

1 Introduction

The simulation of networks of spiking neurons is a one of the major topics in neuroscience. A recent overview and rigorous evaluation of different models of spiking neurons can be found in [7]. E.M. Izhikevich discusses twenty of the most prominent features of biological spiking neurons (tonic spiking, phasic spiking, threshold variability, etc.) against the expressiveness of eleven models of spiking neurons, including an assessment of the “biological plausibility” of neuron models. The author highlights three models as particularly suited for simulations of networks of spiking neurons, where the high degree of approximation of biophysical properties plays an important role. The three models are the Hodgkin-Huxley model, the Morris-Lecar model, and a new model developed by the author himself. Due to the high accuracy of the Hodgkin-Huxley model (four defining equations with tens of parameters), the model is computationally expensive and allows the simulation of small networks only with currently available hardware. This is why, at the moment, the Morris-Lecar model (two defining equations and a moderate number of parameters) is very popular in the computational neuroscience community.

As pointed out in [8], in most of the research on networks of spiking neurons synchrony of firing is emphasized, i.e. if two or more neurons have a common post-synaptic target and fire synchronously, then their spikes arrive at the target at the same time, thereby evoking potent post-synaptic responses. This implies an implicit assumption that the axonal conduction delays are negligible or equal. E.M. Izhikevich [8] (see also [9]) investigates the effect of signal delays and pre-synaptic (asynchronous) firing sequences on post-synaptic responses,

i.e. different (initial) pre-synaptic firing sequences result in the activation of different subsets of neurons. The underlying approximation of neuronal activity is the model introduced in [7], which comprises of three variables and four parameters. To distinguish the effect from random asynchronous firing, the notion of “polychronization” is introduced, and subsets of activated neurons are called “polychronous groups.” For constant, randomly chosen conduction delays, results from simulations of a 10^3 neuronal network are presented and discussed in [8]. Polychronous groups are identified from their topological relation and firing activity in simulations after convergence of the network. The author identifies 5,269 different polychronous groups in the network of 10^3 neurons, which suggests an unprecedented memory capacity of such systems. A natural extension of this work is to investigate specific cost functions associated with conduction delays, instead of constant, randomly chosen values. The present paper presents preliminary research into this direction. A future goal is to optimize the 2D and 3D placement of neurons in such a way that the number of polychronous groups is maximized.

2 Modelling conduction delays

Various models have been proposed recently for numeric simulations of axonal conduction delays; see [1, 2, 3, 4, 13] and the literature therein. We intend to employ the connection cost function as proposed in [2, 3, 4] and in [13], since the authors provide a justification of their model that is based on experimental evidence in the context of the Optimal Neuronal Layout Problem.

2.1 The Chklovskii delay function

A common assumption is that wiring costs are related to the wiring volume. However, in his paper [4], D. Chklovskii attempts to tackle the problem of wiring costs in neuronal networks by proposing a cost function that is determined by the square of the wire length. A volume cost function would result in axons tending to be extremely thin. Although this may decrease energy and other requirements, it would certainly hinder the propagation of signals. These conflicting requirements tend to suggest some optimum combination that must be factored in any plausible solution.

The Chklovskii model is motivated by a comparison of results from neuronal layout calculations to the arrangements in the macaque pre-frontal cortex and the *Caenorhabditis elegans*, in the latter case for 279 neurons [2, 3]. The experimental analysis suggests that wiring costs are indeed proportional to the “wire volume” for a fixed diameter, and the costs grow linearly with the wire length. Furthermore, the wiring cost is a function of the propagation delay $T = L/k \cdot d$, where L is the length of the connection and $k \cdot d$ is the signal speed for the diameter d ; $k = \text{const}$. This implies the approximation

$$\text{Cost} = \alpha \cdot \frac{\pi}{4} \cdot d^2 \cdot L + \beta \cdot \left(\frac{L}{k \cdot d} \right)^\gamma, \quad (1)$$

where $\alpha, \beta, \gamma = \text{const.}$ If the cost function is minimized with respect to the diameter, i.e. $\partial \text{Cost} / \partial d = 0$, one obtains $\text{Cost} \asymp L^{3 \cdot \gamma / (\gamma + 2)}$, which means $\text{Cost} \asymp L$ for $\gamma = 1$ and $\text{Cost} \asymp L^3$ as $\gamma \rightarrow \infty$. As a compromise, $\text{Cost} \asymp L^2$ is assumed in [4]. One of the main objectives of the present paper is to analyse whether this assumption is justified, where the basic parameter is the intensity of spiking activities of neurons in the Morris-Lecar model.

2.2 The Shefi et al. delay function

In [13], Shefi et al. propose an alternative approach to the optimal layout problem and the laws that dictate the formation of neural structures, drawing on the earlier works of C.D. Murray [11, 12]. Murray's papers originally dealt with volumes and the relationship between the diameters of a given pair of sub-branches of a blood vessel or tree. This relationship basically stated that the optimal diameters of the bifurcating branches depended on the diameter of the parent branch or blood vessel and the angle between them, given by the following equations:

$$\cos(\alpha_1) = \frac{d_0^4 + d_1^4 - d_2^4}{2d_1^2 \cdot d_0^2}, \quad \cos(\alpha_2) = \frac{d_0^4 + d_2^4 - d_1^4}{2d_2^2 \cdot d_0^2}, \quad \cos(\alpha_1 + \alpha_2) = \frac{d_0^4 + d_1^4 - d_2^4}{2d_1^2 \cdot d_2^2}, \quad (2)$$

where d_0 is the parent neurite diameter, d_1 and d_2 being the child neurite diameters, and α_1 and α_2 are the angles between the central axis of the parent neurite and their respective child neurites.

Complimentary to the diameter and angles relationship, Shefi et al. [13] considered the balance of tensions between the main neurite and the sub-neurites. This ensures the optimal layout and diameters and is mathematically equivalent to Murray's equations. However, for the scope of this paper we shall limit ourselves to (2). Both Murray's and tension approaches were applied to data from cultured networks of *locust ganglia* [13]. Although neither was conclusive, results were overall satisfactory for both models [13].

For the axon delay experiments described in this paper, the total wiring cost φ for each neurite shall be taken as the volume of the parent neurite up to the bifurcation point added to that of the child neurite as suggested in [13]:

$$\varphi = \pi \cdot \frac{(d_0^2 \cdot l_0 + d_1^2 \cdot l_1)}{4}, \quad (3)$$

where d_0 and d_1 respectively represent the parent and child neurite diameters and l_0 and l_1 are the parent and child "wire" lengths.

3 The neuronal network

The networks devised for the experiments are based on the assumption that neurons in nature tend to communicate largely through spike trains. The network simulation time is divided into iterations, where during an iteration every neuron can simulate the impact of the efferent spikes on its membrane and generate

its own action potentials as a result. These new potentials are propagated at the end of the current iteration but will only have effect in the next one. All spikes are stored and processed in order of arrival at the neuron. This strategy, although artificial, ensures that all spikes generated in the network will be processed in a synchronised manner, which does not affect the study of the axon delays.

The physical design for the Chklovskii experiment network is shown in Figure 1. The 24 ML-type neurons are linked in an arbitrary random manner. Consequently, a spike generated by one particular neuron is transmitted to a few (1 to 2) other neurons. Although the amplitude for a give spike will be the same for all, the arrival time will vary depending on the axon length.

A different approach is taken for the Shefi experiment where the neural connections, as seen in Figure 2, follow a pattern dictated largely by the bifurcations of the axons. In this case the axon delay is calculated by the total axon-neurite volume.

Depending on the cost function applied, axon length or volume directly affects the arrival of a pulse transmitted from one neuron to another.

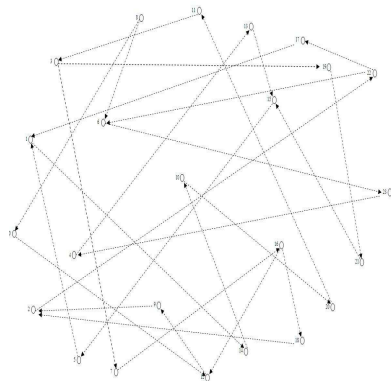


Fig. 1: Layout and connectivity for the Chklovskii experiments.

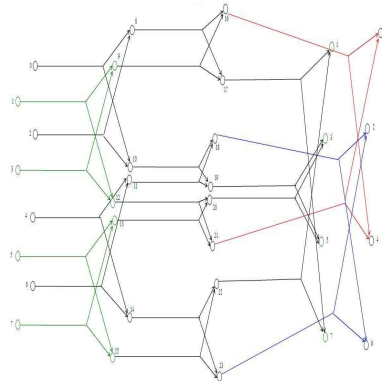


Fig. 2: A simplified network diagram for Shefi et al. experiments.

4 Computational experiments

The simulation and integration of the ML equations is handled by the XPPAUT program package. This freeware, written by G.B. Ermentrout [5], is a graphical tool that we utilized for solving the ML's differential equations, thus allowing each neuron to change its state in accordance with the input spikes it receives, see also [6, 10]. The following parameters are common to all ML neurons:

$$\begin{aligned}
 v_k &= -84\text{mV}; l = -60\text{mV}; v_{ca} = 120\text{mV}; g_k = 8; g_l = 2; g_{ca} = 4; \\
 c &= 20\mu\text{F}; v_1 = -1.2\text{mV} (M_\infty = 0.5); v_2 = 18\text{mV}; \\
 v_3 &= 12\text{mV} (N_\infty = 0.5 \text{ mV}); v_4 = 17.4\text{mV}; \phi = .23; i = 30\mu\text{A}.
 \end{aligned}$$

In each experiment, we executed 14 iterations with each iteration simulating 400ms, i.e. the total simulation time of a single experiment was 5,600ms. The settings were the results of preliminary experiments with the XPAUT program package.

4.1 Chklovskii wiring cost results

As mentioned in Section 2.1, the main concern was to analyse the impact of the power γ in (1) on the spiking activity. Six separate experiments were conducted on the same network of 24 Morris-Lecar neurons for the Chklovskii wiring cost according to $\text{Cost} \propto L^{3\gamma/(\gamma+2)} = L^p$ and variable p (actually, γ). The settings for p were $p = 0.0$ (no wiring cost), 1.0, 1.5, 2.0, 2.5 and 3.0.

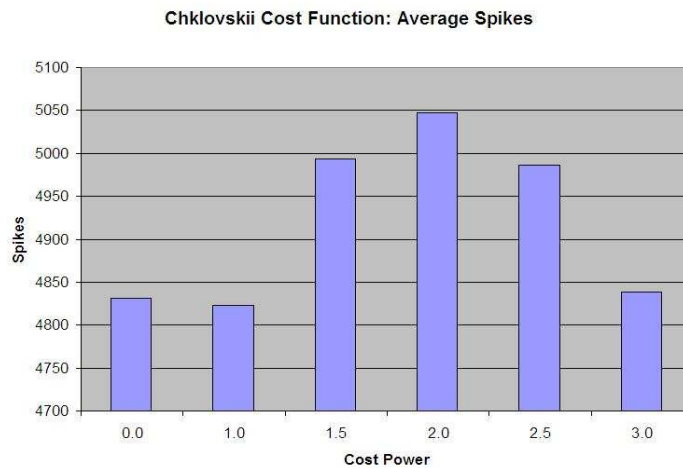


Figure 3: Average number of spikes produced by the same network.

Each set of 6 experiments was conducted with 4 different initial spike sets. The sets consist of mixed patterns of repeated 3, 2, or single spikes.

Fig. 3 shows the average over 4 different input spike sets for each of the six values for the power p in $\text{Cost} \propto L^p$. We note that by far the largest values are obtained for $1.5 \leq p \leq 2.5$.

4.2 Shefi et al Wiring Cost Results

Two simulations were conducted to demonstrate the impact of the Shefi et al. cost function: one simulation with cost function φ from (3), and one with conductance delay equal to zero. The same 4 sets of initial spikes were used as in the Chklovskii simulations. Fig. 4 shows the significant difference between the two cases.

Furthermore, we note that for different topologies, but the same number of ML neurons and the same simulation time, the average number of spikes produced within the 14 iterations of equal length is approximately the same.

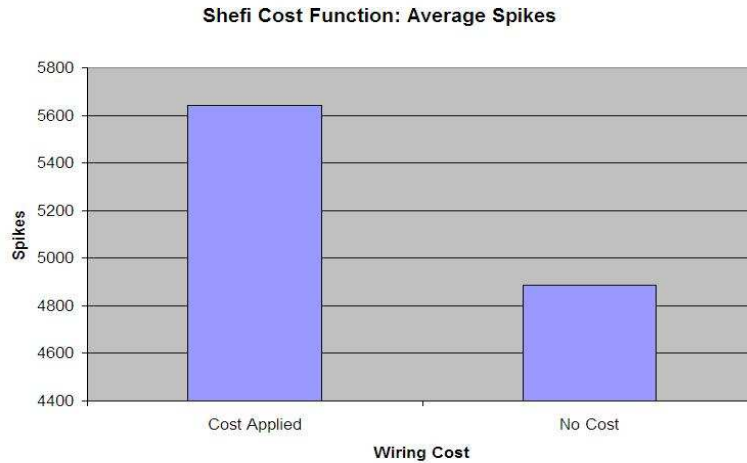


Figure 4: Average number of spikes produced by the network from Fig. 2.

Unfortunately, due to the limited space, the presentation of additional data and a more detailed analysis of our experimental data must be omitted.

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