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Measurements of the $pp \rightarrow W \gamma \gamma$ and $pp \rightarrow Z \gamma \gamma$ cross sections at $\sqrt{s} = 13$ TeV and limits on anomalous quartic gauge couplings

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Abstract

The cross section for W or Z boson production in association with two photons is measured in proton-proton collisions at a centre-of-mass energy of 13 TeV. The data set corresponds to an integrated luminosity of 137 fb^{-1} collected by the CMS experiment at the LHC. The $W \rightarrow e \nu_e$ and $Z \rightarrow e^+ e^-$ decay modes (where e, ν_e) are used to extract the W and Z cross sections in a phase space defined by electron (muon) with transverse momentum larger than 35 (30) GeV and photon transverse momentum larger than 20 GeV. The measured cross sections in this phase space are $\sigma_{W \gamma \gamma} = 13.6^{+1.9}_{-1.9} \text{ (stat)}^{+4.0}_{-4.0} \text{ (syst)} \pm 0.08 \text{ (PDF + scale) fb}$ and $\sigma_{Z \gamma \gamma} = 5.41^{+0.58}_{-0.55} \text{ (stat)}^{+0.64}_{-0.70} \text{ (syst)} \pm 0.06 \text{ (PDF + scale) fb}$. Limits on anomalous quartic gauge couplings are set in the framework of an effective field theory with dimension-8 operators.

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1 Introduction

The measurement of the associated production of a vector boson $V = W, Z$ and two photons in proton-proton (pp) collisions, denoted as $V\gamma\gamma$, is a powerful test of the standard model (SM). The nonabelian nature of the electroweak interaction predicts the presence of self-interacting vector boson vertices. The strength of the interaction is set by the values of triple and quartic gauge couplings predicted by the SM. The measurement of possible deviations from the theoretical predictions could provide indirect evidence of new particles or new interactions. Discrepancies at high photon momentum, where new physics might give a measurable deviation from the SM cross section, would produce evidence for the possible existence of anomalous quartic gauge couplings (aQGCs). A parametrisation of predictions involving anomalous couplings, independent of any specific new physics model, can be calculated in an effective field theory (EFT) framework [1]. Triboson production is also an important background for several SM and beyond the SM processes, such as the Higgs boson production in association with vector bosons (with HVV). Thus, studies of the $V\gamma\gamma$ interactions provide an excellent opportunity for a deeper understanding of electroweak interactions.

Some of the elementary processes resulting in the production of a massive vector boson in association with two photons at the CERN LHC are presented in the leading-order (LO) Feynman diagrams of Fig. 1.

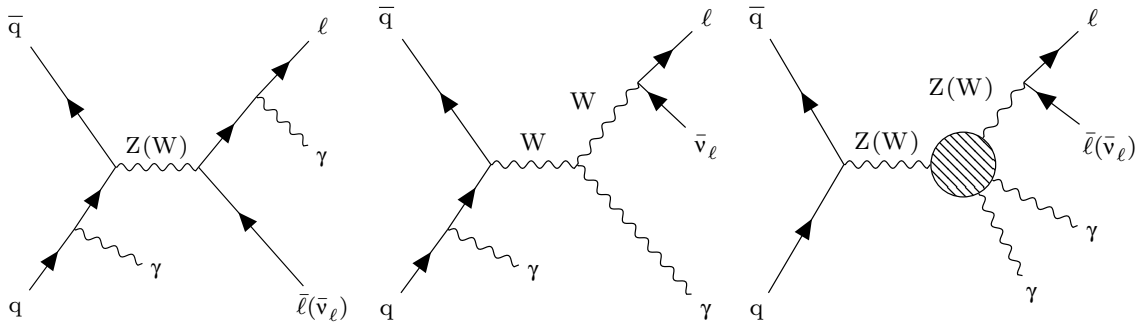


Figure 1: Representative Feynman diagrams for the $V\gamma\gamma$ production in the SM (left and centre) and beyond the SM (right).

Previous measurements of the $V\gamma\gamma$ production cross sections have been performed by the ATLAS and CMS Collaborations at the CERN LHC in pp collisions at a centre-of-mass energy of $\sqrt{s} = 8$ TeV [2–4]. Limits on the presence of aQGCs were also reported in these papers.

In this paper, the first measurements of the $pp \rightarrow W\gamma\gamma$ and $pp \rightarrow Z\gamma\gamma$ cross sections at $\sqrt{s} = 13$ TeV are presented using data collected between 2016 and 2018 by the CMS experiment, corresponding to an integrated luminosity of 137 fb^{-1} . For these measurements, only direct decays into electrons or muons are considered. Measurements are compared with the latest available calculations at next-to-LO (NLO) in perturbative quantum chromodynamics (QCD) [5–7]. Limits on the aQGCs are presented in the framework of an electroweak EFT with dimension-8 operators.

2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward

calorimeters extend the pseudorapidity coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionisation chambers embedded in the steel flux-return yoke outside the solenoid.

Events of interest are selected using a two-tiered trigger system [8]. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz with a latency of about 4 μ s. The second level, known as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimised for fast processing that reduces the event rate to around 1 kHz before data storage.

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, is reported in Ref. [9].

3 Event simulation

The associated production of a W Z boson and at least two photons is searched for in events with one lepton (two opposite-sign, same-flavour leptons) and two photons. Only electron and muon decay channels are used, while the τ decays are treated as a background. The V signal samples are generated at NLO with MADGRAPH5_aMC@NLO [10] (2.2.2 for 2016 samples, 2.2.6 for 2017 and 2018 samples). The NLO NNPDF 3.0 set [11] (for 2016 samples) and the next-to-NLO NNPDF 3.1 set [12] (for 2017 and 2018 samples) are used as parton distribution function (PDF) sets.

The main background contribution comes from the misidentification of jets as photons, which is estimated in single-photon control regions. Thus, various background samples involving a single photon are needed, including V samples. Other single-photon and diphoton processes (such as $t\bar{t}$ produced in association with one or two photons or a photon and a jet) contribute as backgrounds and are estimated using Monte Carlo (MC) simulations.

The V single-photon samples, the t and the $t\bar{t}$ samples are generated at NLO with MADGRAPH5_aMC@NLO. The τ plus jets samples are generated at LO with MADGRAPH5_aMC@NLO. The same PDF sets as for the signal samples are used. The V samples, which are generated with SHERPA v2.2.6 [13, 14] at NLO precision with up to two additional jets and at LO precision for the three-jet computation using the NNPDF 3.1 set, are used for consistency checks and systematic uncertainties evaluations.

Other processes that do not have a photon in the matrix element calculation are exploited for consistency tests in control regions. The W and Z inclusive samples and the triboson samples VVV (WWZ , WZZ and ZZZ) are generated at NLO with MADGRAPH5_aMC@NLO. Single top quark events are generated at NLO with POWHEG v2 [15–18] and decays are simulated with the MADSPIN package [19]. Diboson samples VV (WW , ZZ and WZ) are generated with PYTHIA v8.219 for the 2016, v8.226 for the 2017, and v8.230 for the 2018 samples [20].

The PYTHIA v8.226 (v8.230) package version is used for hadronisation with the CUETP8M1 tune [21] (CP5 tune [22]) for the 2016 (2017 and 2018) samples.

Photons can be present also in other processes because of the hadronisation phase of the generation performed with PYTHIA (v8.2) even if not explicitly produced at matrix element level. To avoid possible double counting effects in the event selection, a procedure for the removal of the overlapping phase space region between inclusive and exclusive samples is implemented. Photons are selected at the generator level following a selection as close as possible to the

one performed at reconstruction level. The total number of selected photons at the generator level is then used to remove overlapping phase space between different samples. For inclusive samples (such as W +jets or Drell–Yan+jets), the event is discarded if one or more photons are selected. For single-photon processes (such as W or Z), the event is discarded if the total number of selected photons at the generator level is different from one. The event is discarded from the diphoton processes if it has less than two photons selected at the generator level.

Predictions for aQGC signals are obtained by including a set of weights, corresponding to the presence of the anomalous couplings, to the V reference samples simulated with MADGRAPH5_aMC@NLO. For this purpose, an aQGC model [23] is used.

Additional pp interactions in the same or adjacent bunch crossings (known as pileup) is included by adding simulated minimum bias events to the hard scattering. The events in the MC simulations are weighted so the distribution of the number of pileup interactions matches the one measured in data. The interaction of the particles with the CMS detector is simulated with GEANT4 [24].

4 Event selection

Events for the V analysis are selected using isolated single-lepton trigger requirements [8]. Single-electron trigger algorithms have a transverse-momentum p_T threshold of 27 GeV (for the 2016 data-taking period) and 32 GeV (for the 2017 and 2018 data-taking periods); single-muon trigger algorithms require p_T above 24 GeV for all three years.

All measured particles are reconstructed using the particle-flow (PF) algorithm [25]; this algorithm reconstructs and identifies each individual particle in an event with an optimised combination of information from the various elements of the CMS detector. The reconstructed vertex with the largest value of the sum of the p_T^2 of the physics objects is the primary pp interaction vertex. The photon energies are obtained from the ECAL measurement. The electron energies are determined from a combination of the electron momentum at the primary interaction vertex as determined by the tracker, the energy of the corresponding ECAL cluster, and the energy sum of all bremsstrahlung photons spatially compatible with the ones originating from the electron track. The muon energies are obtained from the curvature of the muon tracks.

Electrons candidates are required to have $p_T \geq 15$ GeV in the pseudorapidity ranges that exclude the barrel-endcap transition region, $1.44 \leq \eta \leq 1.57$ and $2.50 \leq \eta \leq 2.50$. A variety of criteria is used to separate genuine electrons from misidentified ones. A tight identification is used to select prompt electrons (produced at the primary vertex) and isolated electrons in the final state [26]. Background contributions from misidentified jets or electrons inside a jet are rejected applying electron isolation criteria, which exploit the PF-based event reconstruction. The electron isolation variables are obtained by summing the p_T of charged hadrons compatible with the primary vertex I_{chg} , of neutral hadrons I_{neu} , and of photons I_{pho} inside a cone of radius $R = \sqrt{\frac{\Delta x^2 + \Delta y^2}{2}} = 0.3$ around the electron direction, where ϕ is the azimuthal angle in radians. Additional photons and neutral hadronic contributions to the isolation variable, coming from pileup, are subtracted using the jet area approach [27].

Muons candidates are required to have $p_T \geq 15$ GeV in the pseudorapidity range $2.4 \leq \eta \leq 2.4$. Muon identification criteria are based on the fit quality for tracks measured in the tracker and muon detectors. A tight muon identification is used to reconstruct muons in the final state [28]. To distinguish between prompt muons and those from hadron decays within jets, muons are required to be isolated with respect to all nearby PF reconstructed particles. For the computa-

tion of the PF isolation, the I_{chg} , I_{neu} and I_{pho} components are summed in a cone of $R = 0.4$ around the muon direction. The corrected energy sum is obtained by subtracting the pileup contribution to I_{neu} and I_{pho} , which is estimated as half of the corresponding charged hadronic component.

Photons are selected with $p_T > 20 \text{ GeV}$ in the pseudorapidity range $|\eta| < 1.44$ and $1.57 < |\eta| < 2.38$.

2.5. Photon identification is based on the sequential application of several selections. A medium photon identification is used to reconstruct prompt photons (i.e. not from hadron decays) in the final state [26]. The average efficiency for this selection is 80%. Photons selected in this analysis are required to have a narrow transverse shape of the electromagnetic shower, a minimal energy deposit in the HCAL, and to be isolated with respect to other particles. The same isolation variable previously described for the electron selection is used.

The reconstruction, identification, and isolation efficiencies of leptons and photons and the trigger efficiencies of leptons are measured with the “tag-and-probe” technique [29], as a function of particle p_T and p_T in both data and simulation. A sample of events containing a Z boson decaying into e^+e^- or $\mu^+\mu^-$ is used for these measurements. The photon efficiencies are derived using a sample of electrons from Z decays with no requirement on the track and charge of the candidate. These efficiencies are used to correct for the differences between data and simulation.

An event is categorised as a W decaying to leptons if exactly one electron (muon) with $p_T > 35 \text{ GeV}$ is selected. The selected lepton must match the one that triggered the event and must be associated with the primary vertex of the collision. If the event contains any additional different-flavour leptons or opposite-sign same-flavour leptons, it is excluded from the W boson candidate sample, but is further checked for the presence of a Z boson candidate.

An event is categorised as a Z boson candidate decaying to leptons if two opposite-sign leptons of the same flavour are selected. The leading p_T electron (muon) is required to have $p_T > 35 \text{ GeV}$, and the subleading electron or muon is required to have $p_T > 15 \text{ GeV}$. Only the leading lepton is required to match the one that triggered the event, although both are required to be associated with the primary vertex of the collision. The invariant mass of the dilepton system is required to be $m_{\ell\ell} > 55 \text{ GeV}$.

Events are selected if they have a single W or Z boson candidate and at least two photons. All reconstructed photons must be separated from each other and from each reconstructed lepton by $R = 0.4$. Photons are discarded if $m_{e,\text{lead}} < m_Z - 5 \text{ GeV}$ (where $m_{e,\text{lead}}$ is the invariant mass of an electron and the leading photon and m_Z is the Z boson invariant mass) or if $m_e < m_Z - 5 \text{ GeV}$ (where m_e is the invariant mass of an electron and the two photons). In this way, photons likely coming from final-state bremsstrahlung radiation are removed and, therefore, the contribution from electrons misidentified as photons is reduced as well.

5 Background estimation

The backgrounds in both the W and Z signal regions are categorised as events with a genuine photon or with another object misidentified as a photon. The largest contribution in both channels comes from the misidentification of jets as photons. Another important source of background originates from electrons that are reconstructed as photons because the deposit in the calorimeter is not associated with a track in the tracker. This contribution is particularly relevant in the W electron channel. Both of these background processes are estimated by exploiting a control sample in data. The remaining minor contributions from processes that have

genuine photons (t , $t\bar{t}$, $t\bar{t}$, and VV) are estimated using MC simulations and referred to as “others”.

The background from events containing nonprompt photons is estimated following the method described in Refs. [3, 4]. A W or a Z boson is selected together with one photon that passes the standard selection except for the isolation requirement both in data and simulation. Events are categorised as “tight” or “loose” if the photon candidate passes or fails the isolation requirement. The probabilities for photon and for a jet f to be isolated are computed as

$$\frac{N_{,MC}^T}{N_{,MC}^T + N_{,MC}^L} \text{ and } f = \frac{N_{,data}^T}{N_{,data}^T + N_{,data}^L},$$

where $N_{,MC}^T$ ($N_{,MC}^L$) is the number of simulated events with a tight (loose) photon while $N_{,data}^T$ ($N_{,data}^L$) is the number of events with a tight (loose) photon candidate in data after the subtraction of the prompt photon contribution from simulation. These probabilities are calculated separately for photons in the ECAL barrel and endcap regions as a function of the photon p_T . The jet-photon misidentification background in the diphoton phase space is then estimated by solving the system

$$\begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix} = \begin{pmatrix} 1 & 2 & 1f_2 & f_1 2 & f_1 f_2 \\ 1 & 1 & f_2 & f_1 1 & f_2 \\ 1 & 1 & f_2 & 1 & f_1 2 \\ 1 & 1 & 1 & f_2 & 1 & f_1 1 & f_2 \end{pmatrix} \begin{pmatrix} j \\ j \\ j \\ jj \end{pmatrix},$$

where the indices of the and f coefficients refer to the leading and the subleading photon. The N_{XY} vector contains the number of events where two (TT), one (TL and LT) or zero (LL) photon candidates pass the isolation requirement. The A_{AB} vector contains the number of signal (j) and background (j, j and jj) events. This method is validated in a control region enriched in the jet-photon misidentification background where both photons fail the isolation selection.

In the W channel, a contamination is present from $Z \rightarrow ee$ events where an electron from the Z boson decay is misclassified as a photon. To estimate this contribution, a correction factor is computed for Z events and is then applied to the simulation. The invariant mass of an electron and a photon is reconstructed in data and MC simulation while removing the requirement $m_{e,\text{lead}} > m_Z + 5 \text{ GeV}$. This mass distribution is fitted with the sum of a signal template, derived from MC simulation, and a background function, which has an exponential decay distribution at high mass (above the Z peak) and a turn-on (linked to the electron and photon p_T thresholds) described by an error function at low mass. A correction factor, obtained in intervals of the photon p_T and η , is then computed as

$$\mathcal{F}(p_T, \eta) = \frac{N_{\text{inv}}^{\text{data}} / N_Z^{\text{data}}}{N_{\text{inv}}^{\text{MC}} / N_Z^{\text{MC}}}, \quad (1)$$

where $N_{\text{inv}}^{\text{data}}$ is the number of events in the electron-photon invariant mass peak obtained by integration of the fitted signal shape and N_Z^{data} is the number of events for the $Z \rightarrow ee$ invariant mass distribution obtained by integration of a fitted double-sided Crystal-Ball function [30] in the data. The same procedure is used to calculate the number of events in MC simulation. By fitting all the distributions in the different p_T , η bins, a set of correction factors is computed. The MC simulation is then corrected for these factors on an event-by-event basis if a reconstructed photon matches a generator-level electron. The event-by-event correction factors are on average about 20%.

The pre-fit (i.e. before the fitting procedure described in Section 7) diphoton p_T distributions for the $W\gamma\gamma$ and $Z\gamma\gamma$ analyses, separated in the electron and muon channels, are shown in Fig. 2. The same distributions obtained in the control region enriched in misidentified photons are shown in Fig. 3 for the $W\gamma\gamma$ and $Z\gamma\gamma$ electron channels. The data and prediction agree thus validating the jet–photon misidentification background estimation procedure. A similar level of agreement is observed in the muon channel.

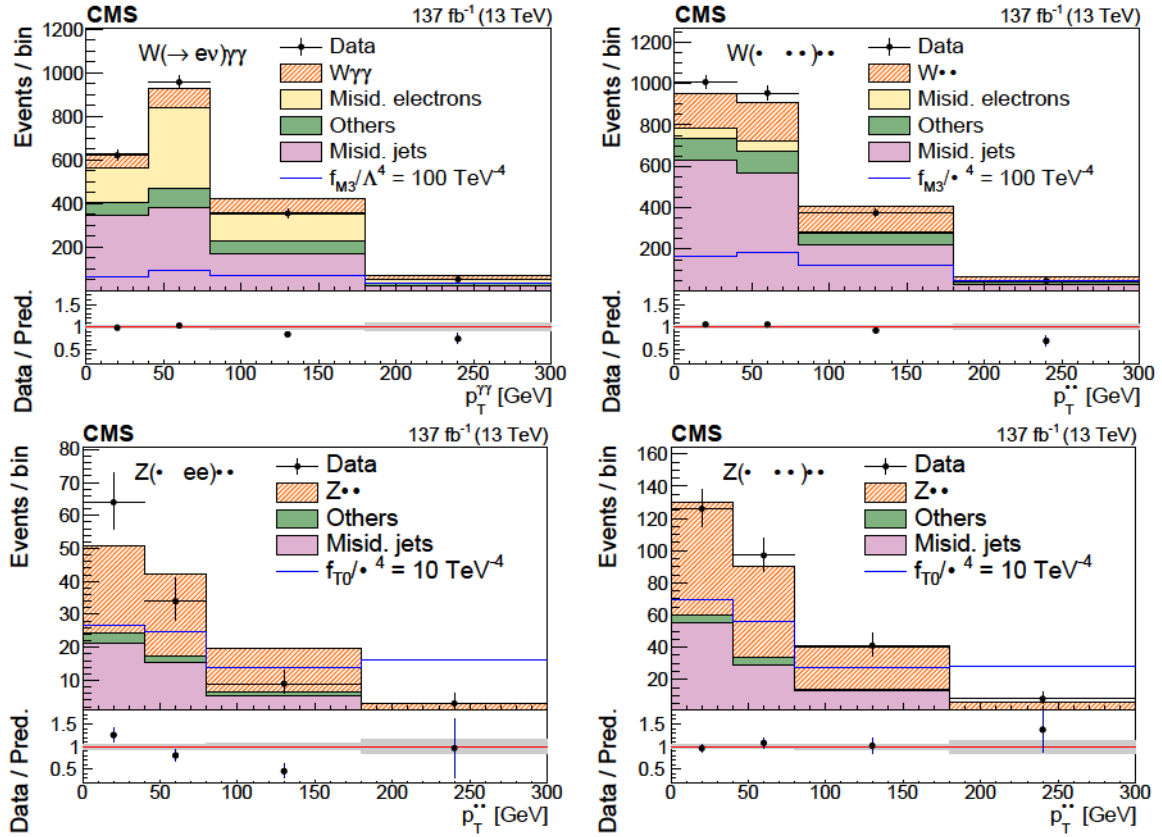


Figure 2: Distribution of the transverse momentum of the diphoton system for the $W\gamma\gamma$ electron (upper left) and muon (upper right) channels and for the $Z\gamma\gamma$ electron (lower left) and muon (lower right) channels. The predicted yields are shown with their pre-fit normalisations. The black points represent the data with error bars showing the statistical uncertainties. The hatched histogram shows the expected signal contribution. The background estimate for electron (jet) misidentified as photons, obtained from control samples in data, is shown in light brown (purple). The remaining background, derived from MC simulation, is shown in green. In the ratio plots, the grey hashed area is the statistical uncertainty on the sum of signal and backgrounds, while the uncertainty in the black dots is the statistical uncertainty of the data. In blue, the expected distribution for an example value of the anomalous coupling parameters f_{M3}/Λ^4 and f_{T0}/Λ^4 is also shown (see Section 8 for the details).

6 Systematic uncertainties

Systematic effects can affect the rates and distributions of both data and simulation. To estimate these uncertainties, the full analysis procedure is repeated by varying each quantity by plus or minus its standard deviation uncertainty. In this procedure, correlations between the systematic uncertainties are included where appropriate.

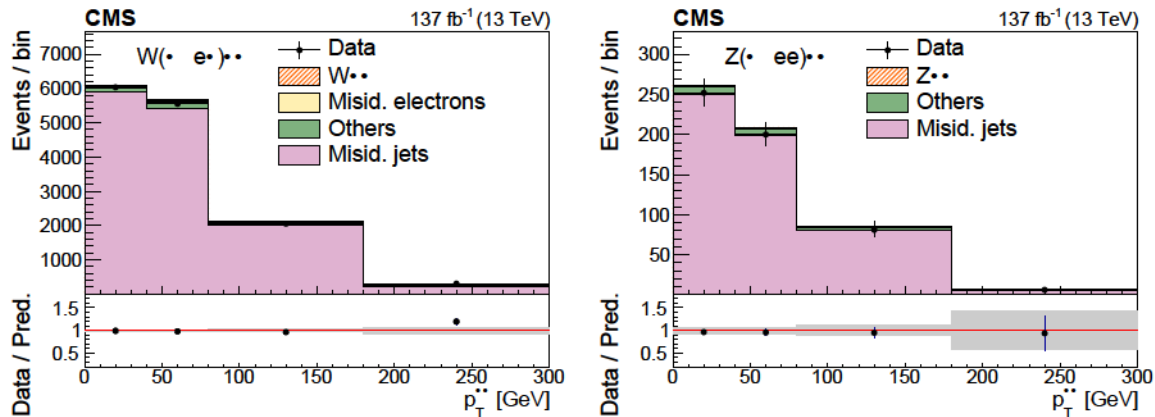


Figure 3: Distribution of the transverse momentum of the diphoton system, obtained in the control region enriched in misidentified photons, for the $W\gamma\gamma$ and for the $Z\gamma\gamma$ electron channels. The black points represent the data with error bars showing the statistical uncertainties. The hatched histogram shows the expected negligible signal contribution. The background estimate for electron (jet) misidentified as photons, obtained from control samples in data, is shown in light brown (purple). The remaining backgrounds, derived from MC simulation, are shown in green. In the ratio plots, the grey hashed area is the statistical uncertainty on the sum of signal and backgrounds, while the uncertainty in the black dots is the statistical uncertainty of the data.

The dominant systematic uncertainties come from the estimation of the backgrounds. The uncertainty in the correction factor related to the background of electrons misidentified as photons is determined by propagating the estimated uncertainty in the correction factor \mathcal{F} of Eq. 1. The latter has two components: a statistical one, that comes from the uncertainty in the fitting procedure; and a systematic one that is computed by taking half the difference between the \mathcal{F} factors obtained by performing the fit with a double-sided Crystal-Ball function and with the nominal method.

To determine the systematic uncertainty coming from the jet-photon misidentification background, the same strategy is applied to a QCD control sample that is obtained using the $W\gamma$ selection but inverting the isolation requirement on the leptons while keeping the photon selection identical to the one for the signal region. This sample is used to obtain an alternative estimate of the jet-photon misidentification background contribution in the W channel. For the Z channel, the QCD control sample resulting from the $Z\gamma$ selection with the inversion of the lepton isolation has insufficient events. Hence, a transfer factor from the $W\gamma$ selection is computed and applied for the determination of the alternative estimate of the jet-photon misidentification background contribution in the Z channel. The systematic uncertainty is computed as half the difference in the distributions between the standard method or the one just described.

Another source of uncertainty in the jet-photon background is related to the modelling of the initial- and final-state radiation and of the energy spectra of the final state particles. An alternative MC simulation, obtained with SHERPA, is used to evaluate this uncertainty.

The uncertainties in the lepton and photon reconstruction and selection efficiencies are included by computing the cross section with these efficiencies varied up and down by one standard deviation. The uncertainty related to these data-to-simulation corrections is estimated by including the uncertainty in the tag-and-probe method. Uncertainties in the trigger efficiencies are negligible.

Table 1: Summary of the systematic uncertainties (in percent) for the W and Z cross section measurements. The numbers indicate the impact of each systematic uncertainty in the value of the measured cross section in the corresponding channel.

Systematic source	$e e$	%	%	%	ee	%	%	%
Integrated luminosity	1	2	2	2	3	1	3	3
Pile-up	2	1	1	1	2	1	1	1
Electron efficiencies	4	—	1	1	3	—	1	1
Muon efficiencies	1	1	1	1	2	1	1	1
Photon efficiencies	18	13	12	12	6	5	5	5
Jet-photon misid.	25	22	21	21	6	5	6	6
Electron-photon misid.	4	1	1	1	—	—	—	—
W theoretical cross section	3	3	3	3	1	1	1	1
Z theoretical cross section	4	1	1	1	7	5	6	6
Other bkg theoretical cross section	5	2	2	2	1	1	1	1
Simulated sample event count	18	7	8	8	7	3	4	4

The uncertainty in the value of the theoretically computed cross section is accounted for during the subtraction from data of the background processes estimated from MC. Furthermore, the value of the expected cross section has an impact on the estimation of the jet-photon background because the contribution from prompt photons is subtracted from the distribution in data using the W and Z simulations. To estimate these contributions, the cross sections of the W , Z , and of the other minor backgrounds are varied independently. The uncertainty in the Z cross section is estimated as half the difference between the next-to-NLO and the NLO values computed with MATRIX [31], and amounts to 2.5%. The same uncertainty is assumed for the W cross section, and a value of 7.5% is used for the other simulated backgrounds.

The total inelastic cross section is varied by 4.6% [32] to estimate the impact on the final result of the pileup reweighting procedure. The uncertainty because of the integrated luminosity measurement is equal to 2.5, 2.3, and 2.5% for the 2016, 2017 and 2018 data taking periods, respectively [33–36]. Because of the uncorrelated time evolution of some systematic uncertainties, the total integrated luminosity has an uncertainty of 1.8% and is applied to all the processes estimated with an MC simulation. The effect of the uncertainty in the integrated luminosity affects the estimation of the jet-photon misidentification background as well as the MC contributions in the diphoton distributions.

For the extraction of the results, each systematic uncertainty is represented by a nuisance parameter, which affects the shape and the normalisation of the distribution of the various background contributions. The variation of the nuisance parameter results in a continuous perturbation of the spectrum, following a Gaussian probability density function. The impact of each systematic uncertainty is obtained by freezing the set of associated nuisance parameters to their best-fit values and comparing the total uncertainty in the measured cross section with the result from the nominal fit [37]. The contributions of the different systematic uncertainties for both the W and Z processes are presented in Table 1.

7 Cross section measurements

The cross sections for the W and Z processes are measured separately in the electron and muon channels using a sample of events corresponding to an integrated luminosity of 137 fb^{-1} (LHC Run 2 data). The observed and predicted numbers of events are presented in Table 2.

The measured yields in the electron and muon channels are extrapolated to a common fiducial phase space determined from simulated signal events at the generated particle level. Generated

Table 2: Summary of the pre-fit predicted and observed numbers of events for 137 fb^{-1} for the W (upper Table) and Z (lower Table) selections in the electron and muon channels. The systematic uncertainties of the individual backgrounds and the total background are obtained by summing the contributions of different systematic uncertainties in quadrature. The statistical uncertainties are those related to the MC event samples and control region statistical uncertainties.

Process	$e e$					
Misid. jets	918	23 (stat)	180 (syst)	1441	27 (stat)	280 (syst)
Misid. electrons	669	28 (stat)	34 (syst)	107	9 (stat)	7 (syst)
Others	217	11 (stat)	20 (syst)	286	11 (stat)	25 (syst)
Total backgrounds	1804	38 (stat)	180 (syst)	1834	30 (stat)	280 (syst)
Expected signal	248	6 (stat)	17 (syst)	500	8 (stat)	33 (syst)
Total prediction	2052	38 (stat)	180 (syst)	2334	31 (stat)	280 (syst)
Data	1987			2384		
Process	ee					
Misid. jets	42	4 (stat)	9 (syst)	98	5 (stat)	27 (syst)
Others	6	1 (stat)	1 (syst)	11	2 (stat)	1 (syst)
Total backgrounds	48	4 (stat)	9 (syst)	109	6 (stat)	27 (syst)
Expected signal	68	2 (stat)	5 (syst)	157	3 (stat)	11 (syst)
Total prediction	116	4 (stat)	8 (syst)	266	6 (stat)	23 (syst)
Data	110			272		

particles are considered stable if their mean decay length is larger than 1 cm. Generated leptons are required to have a $p_T \geq 15 \text{ GeV}$ and $\eta \leq 2.5$. The momenta of photons in a cone of $R = 0.1$, the same cone size as the one applied to reconstructed data, are added to the charged lepton momentum to correct for final-state radiation. Generated photons are required to have $p_T \geq 15 \text{ GeV}$ and $\eta \leq 2.5$. Additionally, the candidate photons are required to have no selected leptons or photons in a cone of radius $R = 0.4$ and no other stable particles, apart from photons and neutrinos, in a cone of radius $R = 0.1$. Events are then selected in the W channel by requiring exactly one electron (muon) with $p_T \geq 30 \text{ GeV}$ and at least two photons with $p_T \geq 20 \text{ GeV}$. Events are selected in the Z channel by requiring two electrons (muons), at least one of them with $p_T \geq 30 \text{ GeV}$, and not less than two photons, each of them with $p_T \geq 20 \text{ GeV}$. Additionally, the invariant mass of the dilepton system is required to be $m_{ll} \geq 55 \text{ GeV}$.

The expected theoretical cross sections are predicted at NLO and their uncertainties come from the finite MC sample event count used to compute them, the PDF set, the factorisation and renormalisation scales. Statistical uncertainties are estimated to be of the order of 0.2% in both the W and Z channels. Uncertainties related to the PDF set are estimated using a set of 100 replicas of the NNPDF 3.1 PDF set, following the Ref. [12] prescription, and are estimated to be of the order of 0.3% in the $e e$ and $e \mu$ and of 0.8% in the ee and $e\mu$ channels. Uncertainties related to the renormalisation and factorisation scale choice are estimated by independently varying μ_R and μ_F by a factor of 0.5 and 2, with the condition that $1/2 \leq \mu_R/\mu_F \leq 2$. The uncertainties are defined as the maximal differences from the nominal values and are estimated to be of the order of 0.6 (0.5)% in the $e e$ ($e \mu$) and of 0.6 (0.7)% in the ee ($e\mu$) channels. Uncertainties related to the value of the strong coupling are estimated to be of the order of 0.03 (0.02)% in the $e e$ ($e \mu$) and of 0.4% in the ee and $e\mu$ channels.

Binned maximum likelihood fits to the diphoton p_T distributions in Fig. 2 are performed to

extract the signal strength μ and the significance of the results [38, 39]. The results are obtained separately in the electron, muon and lepton channels. The systematic uncertainties and the statistical uncertainty in the MC predictions are treated as nuisance parameters in the fits and profiled. The high p_T bins in the distributions are the more relevant ones for the determination of the limits.

The measured cross sections are obtained by multiplying the observed signal strength μ by the expected theoretical cross section of the signal MC simulated sample. The theoretical cross section for the $W\gamma\gamma$ and $Z\gamma\gamma$ signals obtained from MADGRAPH5_aMC@NLO at NLO accuracy are 18.70 ± 0.03 (MC stat) ± 0.12 (PDF + scale) fb and 5.96 ± 0.01 (MC stat) ± 0.06 (PDF + scale) fb, respectively.

In the electron channel, the best fit value for the $W\gamma\gamma$ signal strength is $0.23_{-0.22}^{+0.22}$ (stat) $_{-0.30}^{+0.32}$ (syst) and for the $Z\gamma\gamma$ signal strength is $0.73_{-0.17}^{+0.18}$ (stat) $_{-0.13}^{+0.12}$ (syst). The measured cross sections are:

$$\begin{aligned}\sigma(W\gamma\gamma)_{ev}^{\text{meas}} &= 4.4_{-4.1}^{+4.1} \text{ (stat)}_{-5.5}^{+6.0} \text{ (syst)} \pm 0.03 \text{ (PDF + scale) fb,} \\ \sigma(Z\gamma\gamma)_{ee}^{\text{meas}} &= 4.35_{-0.99}^{+1.05} \text{ (stat)}_{-0.77}^{+0.71} \text{ (syst)} \pm 0.05 \text{ (PDF + scale) fb.}\end{aligned}$$

In the muon channel, the best fit value for the $W\gamma\gamma$ signal strength is $0.74_{-0.11}^{+0.11}$ (stat) $_{-0.22}^{+0.23}$ (syst) and for the $Z\gamma\gamma$ signal strength is $1.06_{-0.11}^{+0.11}$ (stat) $_{-0.10}^{+0.10}$ (syst). The measured cross sections are:

$$\begin{aligned}\sigma(W\gamma\gamma)_{\mu\nu}^{\text{meas}} &= 13.8_{-2.1}^{+2.1} \text{ (stat)}_{-4.2}^{+4.3} \text{ (syst)} \pm 0.08 \text{ (PDF + scale) fb,} \\ \sigma(Z\gamma\gamma)_{\mu\mu}^{\text{meas}} &= 6.29_{-0.64}^{+0.67} \text{ (stat)}_{-0.58}^{+0.57} \text{ (syst)} \pm 0.07 \text{ (PDF + scale) fb.}\end{aligned}$$

The results of the fit for the electron and muon channels separately are compatible within two sigmas. In the combined electron and muon channel, the best fit value for the $W\gamma\gamma$ signal strength is $0.73_{-0.10}^{+0.10}$ (stat) $_{-0.22}^{+0.22}$ (syst) and for the $Z\gamma\gamma$ signal strength is $0.91_{-0.09}^{+0.10}$ (stat) $_{-0.12}^{+0.11}$ (syst). The measured cross sections are:

$$\begin{aligned}\sigma(W\gamma\gamma)^{\text{meas}} &= 13.6_{-1.9}^{+1.9} \text{ (stat)}_{-4.0}^{+4.0} \text{ (syst)} \pm 0.08 \text{ (PDF + scale) fb,} \\ \sigma(Z\gamma\gamma)^{\text{meas}} &= 5.41_{-0.55}^{+0.58} \text{ (stat)}_{-0.70}^{+0.64} \text{ (syst)} \pm 0.06 \text{ (PDF + scale) fb.}\end{aligned}$$

The sensitivity for the $W\gamma\gamma$ cross section measurement is dominated by the muon channel. The measured signal strengths are summarised in Fig. 4.

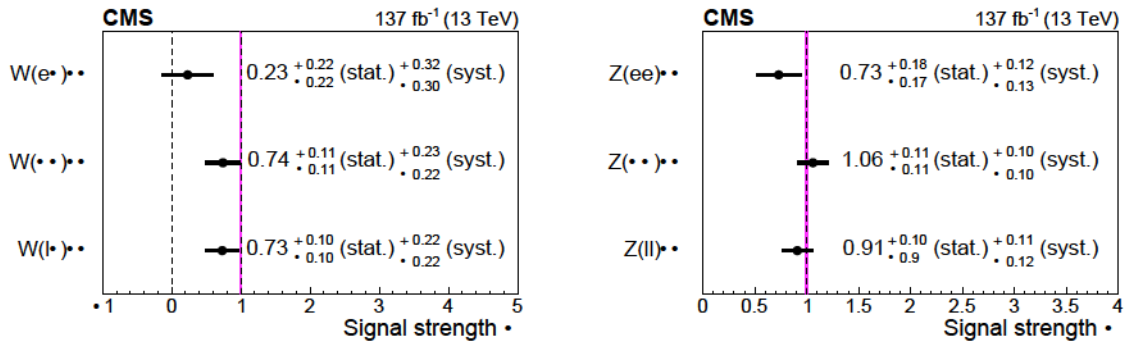


Figure 4: Best fit values of the signal strengths for the $W\gamma\gamma$ (left) and $Z\gamma\gamma$ (right) channels. The error bars represent the total uncertainty while the magenta bands represent the theoretical uncertainty in the MADGRAPH5_aMC@NLO cross section.

The significance of the cross section measurement for both the W and Z channels is quantified using the background-only hypothesis under the asymptotic approximation [40]. The observed (expected) significance for the W signal is 0.6–2.7 in the electron channel and 3.0–4.3 in the muon channel; for the Z is 3.4–5.0 in the electron channel and 5.4–5.1 in the muon channel; the combined significance for the W is 3.1–4.5 and for the Z is 4.8–5.8.

8 Limits on anomalous quartic gauge couplings

Studies of the anomalous gauge couplings can be performed in the EFT framework [1] by expanding the SM Lagrangian to include terms with dimension higher than four. In particular, both the W and Z processes are sensitive to the presence of dimension-6 and dimension-8 operators [41]. Because of the available statistics in the V channel, the sensitivity to dimension-6 operators is expected to be lower than the one in the diboson production. The contribution of each operator is proportional to a coupling constant f and to the inverse of the energy scale at which the new phenomena appear.

In the generation of the anomalous couplings samples, a calculation using 10 (8) different dimension-8 operators was performed for the W (Z) process. The operators can be divided into two subsets: the \mathcal{L}_{M0} – \mathcal{L}_{M7} ones, that contain both the $SU(2)_L$ and $U(1)_Y$ field strengths and the covariant derivative of the Higgs doublet, and the \mathcal{L}_{T0} – \mathcal{L}_{T9} ones, that contain only the two field strengths. In particular, the W channel is especially sensitive to the $M2$, $M3$, $T0$, $T1$, $T2$, $T5$, $T6$, and $T7$ operators, whereas the Z channel is especially sensitive to the $T0$, $T1$, $T2$, $T5$, $T6$, $T7$, $T8$, and $T9$ operators.

The distribution of the p_T of the diphoton system (shown in Fig. 2) is used to constrain the aQGC parameters under the hypothesis of absence of anomalies in triple gauge couplings. The contribution of aQGCs is enhanced at high values of the p_T of the diphoton system. The distribution of the aQGCs as a function of the couplings themselves has a quadratic behaviour, and hence a parabolic fit is implemented to interpolate between the different values obtained via the parameter scan. The fitting procedure is performed bin-by-bin to exploit the shape of the distribution to set the limits and include the different systematic uncertainties. To further increase the sensitivity, electron and muon channels are combined. Each operator coefficient is scanned independently with all other operators set to zero. The extraction of the 95% confidence level upper and lower limits on the aQGCs is performed by exploiting the procedure described in Ref. [39]. The expected and measured limits for both the W and Z processes are presented in Table 3.

In particular, the intervals computed for the f_{T5} and f_{T6} parameters are the most constraining ones in the W channel and are comparable to the most stringent results obtained by the W [42] analysis of the CMS Collaboration at 13 TeV. The intervals computed for the f_{T0} and f_{T5} parameters in the W channel are more stringent than the ones obtained by the Z [3] and Z [43] analyses of ATLAS at 8 TeV. For the Z channel, the most stringent interval is the one computed for the f_{T9} parameter, which is competitive with the results obtained by the Z and ZZ [44] analyses of CMS at 13 TeV. The intervals computed for the f_{T8} and f_{T9} parameters in the Z channel are more stringent than the ones obtained by the Z and Z analyses of ATLAS at 8 TeV.

Table 3: Expected and observed 95% confidence level intervals for the different anomalous couplings in both the W and Z channels.

Parameter	W TeV ⁴		Z TeV ⁴	
	Expected	Observed	Expected	Observed
$f_{M2}/^4$	57.3, 57.1	39.9, 39.5	—	—
$f_{M3}/^4$	91.8, 92.6	63.8, 65.0	—	—
$f_{T0}/^4$	1.86, 1.86	1.30, 1.30	4.86, 4.66	5.70, 5.46
$f_{T1}/^4$	2.38, 2.38	1.70, 1.66	4.86, 4.66	5.70, 5.46
$f_{T2}/^4$	5.16, 5.16	3.64, 3.64	9.72, 9.32	11.4, 10.9
$f_{T5}/^4$	0.76, 0.84	0.52, 0.60	2.44, 2.52	2.92, 2.92
$f_{T6}/^4$	0.92, 1.00	0.60, 0.68	3.24, 3.24	3.80, 3.88
$f_{T7}/^4$	1.64, 1.72	1.16, 1.16	6.68, 6.60	7.88, 7.72
$f_{T8}/^4$	—	—	0.90, 0.94	1.06, 1.10
$f_{T9}/^4$	—	—	1.54, 1.54	1.82, 1.82

9 Summary

The cross sections for both the W and Z processes are measured in proton-proton collisions by the CMS experiment at a centre-of-mass energy of 13 TeV corresponding to an integrated luminosity of 137 fb⁻¹.

The cross sections are measured in a fiducial region where simulated signal events are selected at generator level in the W channel by requiring exactly one electron or muon with transverse momentum $p_T \geq 30$ GeV and at least two photons, each with $p_T \geq 20$ GeV. Events are selected in the Z channel by requiring two oppositely charged electrons or muons, at least one of them with $p_T \geq 30$ GeV, and at least two photons, each with $p_T \geq 20$ GeV. All leptons and photons are required to have pseudorapidity $|\eta| \leq 2.5$. Additionally, the invariant mass of the dilepton system is required to exceed $m_{\ell\ell} \geq 55$ GeV.

The measured cross sections are $13.6^{+1.9}_{-1.9}$ (stat) $^{+4.0}_{-4.0}$ (syst) ± 0.08 (PDF + scale) fb for the W channel and $5.41^{+0.58}_{-0.55}$ (stat) $^{+0.64}_{-0.70}$ (syst) ± 0.06 (PDF + scale) fb for the Z channel. These results are in agreement with the theoretical cross sections computed at next-to-leading order. The corresponding signal significances are 3.1 and 4.8 standard deviations. Limits on anomalous quartic gauge couplings are set using both channels.

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