

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

CERN LIBRARIES, GENEVA



CM-P00053199

PH I/COM-73/39

3 July 1973

PHYSICS I

ELECTRONICS EXPERIMENTS COMMITTEE

P R O P O S A L

A STUDY OF THE EXCITED STATES OF  ${}^4_{\Lambda}\text{H}$  AND  ${}^4_{\Lambda}\text{He}$  HYPERNUCLEI

M. Bedjidian <sup>\*</sup>), A. Filipkowski <sup>\*\*</sup>), J. Y. Grossiord <sup>\*)</sup>  
A. Guichard <sup>\*)</sup>, M. Gusakow <sup>\*</sup>), S. Majewski <sup>\*\*</sup>), H. Piekarz <sup>\*\*</sup>),  
J. Piekarz <sup>\*\*</sup>), J. R. Pizzi <sup>\*</sup>), J. Pniewski <sup>\*\*</sup>) and J. Zaknewski <sup>\*\*</sup>)

\* ) Institut de Physique Nucléaire, Université Claude Bernard de Lyon et Institut National de Physique Nucléaire et de Physique des Particules (IN2 P3) - (France)

\*\* ) Institute of Experimental Physics, University of Warsaw, Institute of Nuclear Research, Warsaw, (Poland)

As it was pointed out in ref. 1 we are interested in a continuation of the hypernuclear  $\gamma$ - spectroscopy studies. The experiment recently performed at CERN by the CERN-Heidelberg-Warsaw Collaboration (2, 3) has shown that features of the  $\Lambda$ N S-State interaction can be examined by this method. The present proposal outlines the physical basis and the experimental procedure of a new experiment.

## I. PHYSICAL BASIS OF THE PROPOSED EXPERIMENT

In the experiment described in ref. 2 and 3, the first hypernuclear  $\gamma$ - transitions ascribed to the mass number 4 hypernuclei ( ${}^4_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{He}$ ) have been distinguished. In the  $\gamma$ -spectra obtained from  ${}^6\text{Li}$  and  ${}^7\text{Li}$  targets irradiated with stopping  $\text{K}^-$  mesons, two  $\gamma$ -lines at 1.09 MeV and 1.42 MeV have been observed. The first one was firmly established as a hypernuclear line whereas the identification of the second one might still be questioned. Nevertheless, it was proved that both lines considered as hypernuclear lines cannot be ascribed to any other hypernuclei but  ${}^4_{\Lambda}\text{H}$  and  ${}^4_{\Lambda}\text{He}$ .

The observation of the 1.09 MeV line made it possible to reconsider the properties of the  $\Lambda$ N interaction in terms of various phenomenological hard core potentials with attractive wells of exponential shape outside the core. As the input data the binding energies of  ${}^3_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{H}$  and  ${}^4_{\Lambda}\text{He}$  were used, in addition to the 1.09 MeV  $\gamma$ -transition energy which was alternatively ascribed to the excitation of either the  ${}^4_{\Lambda}\text{H}$  or the  ${}^4_{\Lambda}\text{He}$  hypernucleus. Knowing the binding energy values of the ground and excited states of the s-shell hypernuclei, one can determine the appropriate  $\Lambda$ - nucleus well depth parameters for these states. This can be achieved by means of the interpolation formulas obtained by Herndon and Tang<sup>4)</sup> with the variational method. Each of these formulas connects the binding energy of a hypernucleus with the  $\Lambda$ -core nucleus

potential well depth parameter. The  $\Lambda N$  potential depths are the linear combinations of the corresponding  $\Lambda$ -core nucleus potential depths. Having the  $\Lambda p$  potential parameters, one can estimate the  $\Lambda p$  elastic scattering cross-sections and compare them with the experimental data obtained at low energies<sup>3)</sup>. The results of these calculations are shown in fig. 1. It is seen that only the  $\Lambda N$  potentials, marked C and D, with core radii  $d = 0,45$  fm and  $0,60$  fm and an intrinsic range  $b = 1,5$  fm can be fitted to the experimental data. This is inconsistent with what was obtained in ref. 4 ( $d = 0,45$  fm,  $b = 2,0$  fm) when the excitation energy of mass number 4 hypernuclei was not known. The new experimental data made it possible to use a more general form of the charge symmetry breaking term in  $\Lambda N$  interaction containing both spin-dependent and spin-independent parts. The spin-independent part turned out to be significant.

The estimated values of the singlet and triplet state scattering parameters  $a_s$  and  $a_t$  differ markedly for both  $\Lambda n$  and  $\Lambda p$  potentials. This indicates a rather strong spin-dependence of the  $\Lambda N$  interaction. In addition to this - at present still rather simple approach - for both assignments of the 1.09 MeV line, the excitation energy of a second mass number 4 hypernucleus has been predicted to be higher than 1 MeV for all fitted potentials

The confirmation of the existence of a hypernuclear line at 1.42 MeV and the proper assignment of both these lines to  ${}^4_{\Lambda}\text{H}$  and  ${}^4_{\Lambda}\text{He}$  would make it possible to gain better information about the  $\Lambda N$  potential parameters.

## II. EXPERIMENTAL PROGRAMME

The aim of the proposed experiment is the identification of the 1.42 MeV line and the proper assignment of the 1.09 MeV line to one of the mass number 4 hypernuclei :  ${}^4_{\Lambda}\text{H}$  or  ${}^4_{\Lambda}\text{He}$ .

In the vicinity of the 1.42 MeV line, a background line at 1.32 MeV was also present. It has been found that this line was due to the interaction of fast secondary neutrons with  ${}^{23}\text{Na}$ , a constituent of the  $\gamma$ -detector<sup>3)</sup>. The 1.32 MeV line was observed in a spectrum slightly delayed with respect to the lines attributed to  $\gamma$ -rays emitted directly from the target. The delay time amounted to about 2ns in the previous experiment, where the distance between the  $\gamma$ -detector and the target was equal to 15 cm only. In the proposed experiment this distance will be enlarged to 45 cm to increase the delay time up to 5-6 ns. This enlargement seems to be sufficient for a good sepa-

ration of the 1.42 MeV and 1.32 MeV lines, tentatively ascribed to a  $\gamma$ -transition in the target or to the fast neutron interaction with the  $\gamma$ -detector, respectively.

A correct assignment of the 1.09 MeV line to the  $\gamma$ -transition in either the  ${}^4_{\Lambda}\text{H}$  or  ${}^4_{\Lambda}\text{He}$  hypernucleus can be done if the  $\gamma$ -spectra are taken in coincidence with decay pions. The decay of the  ${}^4_{\Lambda}\text{H}$  hypernucleus is accompanied in 50% of all the cases by the emission of a  $\pi^-$ -meson with the kinetic energy of 53 MeV. The relative intensity measurements of the 1.09 MeV line observed both in coincidence with the 53 MeV pions and without it should lead to the proper assignment of this line.

### III. EXPERIMENTAL SET-UP

#### 3.1 - $K^-$ - beam

In the proposed experiment the solid angle subtended by the  $\gamma$ -detector will be diminished by a factor of about 10 as compared with the previous set-up<sup>2, 3)</sup>. The same statistics of the observed hypernuclear events can still be obtained with the 10 times more intense  $K^-$ -beam e. g. with the short version of the K12 beam where 7000  $K^-$ /burst could be stopped in the target<sup>5)</sup>. (The background conditions are discussed separately).

#### 3.2 - Counters and electronics

A schematic diagram of the experimental set-up is presented in fig. 2. A telescope consisting of the veto-Čerenkov counter  $\bar{C}_1$  and three plastic scintillation counters  $S_{1-4}$  will be used to obtain the K-meson trigger. A threshold Čerenkov veto-counter made of plexiglass has a high rejection power ( $> 99,5\%$ ) for the 500 MeV/c pions. A further improvement of the kaon - to - pion ratio can be achieved with the help of the scintillation counters  $S_{1-4}$  which are only sensitive to a specific energy loss of the passing kaons. The Cu and C moderators for the kaon beam are placed between the telescope counters. The beam will be stopped both in  ${}^6\text{Li}$  and  ${}^7\text{Li}$  targets.

The NaI(Tl) crystal, of the size 4" x 3", will be used as a  $\gamma$ -detector. The aluminium cover of the crystal has been made especially thin to reduce the background induced by neutrons. Seven XP1010 phototubes will collect the scintillation photons produced in the crystal. It is more advantageous to use a great number of small tubes instead of a one large tube since high counting rates are expected.

A range telescope will be used to detect the 53 MeV decay pions to be expected in the case of the  ${}^4_{\Lambda}\text{He}$  hypernucleus. The pions emitted from the target will pass through a properly shaped Cu-degrader (the differences in range for various directions are reduced in this way) and will be stopped in six layers of the scintillation counters  $S_{7-12}$ . The hodoscope counters  $H_1$  (arranged in a matrix  $5 \times 5$ ), the MWPC  $P_5$  and  $P_6$  ( $10 \times 10 \text{ cm}^2$ , 42 wires in each chamber) and the hodoscope counters  $H_2$  (arranged in a matrix  $10 \times 10$ ) will make it possible to determine the direction of the pion. Four MWPC  $P_1$ ,  $P_2$ ,  $P_3$  and  $P_4$  ( $10 \times 10 \text{ cm}^2$ , 48 wires in each chamber) will determine the K-meson direction. The information from  $P_1$ ,  $P_2$ ,  $P_3$ ,  $P_4$ ,  $P_5$ ,  $P_6$ ,  $H_1$  and  $H_2$  counters will allow the determination of the interaction point in the target. The accuracy of the determination of this point is sufficiently high so that after corrections for energy loss of pions in the target, their energy will be determined within limits of straggling dispersion. It is expected that energy resolution for pions of 50 MeV will be not worse than 3 MeV.

The simplified diagram of electronics is shown in fig. 3. At first, pulses from all scintillation counters are fed to the fast decision-making logic. Then the data containing information about energy of  $\gamma$ -rays, time-of-flight of  $\gamma$ -rays and energy of pions are stored in the OMM of the K-202 computer with the help of CAMAC electronics. After preliminary selection made on-line, the data will be stored on MTU. The off-line analysis will make it possible to correct for energy loss of pions in the target as well as to choose a time delay corresponding to  $\gamma$ -rays emitted directly from the target.

In parallel to computer data-taking system, a 2-dimensional pulse-height analyser TRIDAC-C will be used to measure the energy of  $\gamma$ -rays and their time-of-flight. In this way a fast control of the incoming data will be possible.

### 3.3 - Counting rates and background

The increased intensity of the kaon beam compensates for the reduction of the counting rates obtained with a smaller angle subtended by the  $\gamma$ -detector in the proposed set-up. Therefore the rates observed in a given time unit and ascribed to the hypernuclear lines will not be changed with respect to those obtained previously<sup>2, 3)</sup>. In a week of the machine time (ca 100 h) one

should register about 5000  $\gamma$ -quanta corresponding to the 1.09 MeV line in the case of the  ${}^7\text{Li}$  target and some 1500 counts for each of the two lines, the 1.09 MeV and the 1.42 MeV, in the case of the  ${}^6\text{Li}$  target.

In a short version of the K 12 beam, the total background counting rate observed in the  $\gamma$ -detector is expected to be about  $2 \times 10^5$  counts per second. It will be twice as high as the rate measured previously with the extended version of the K 12 beam and the  $\gamma$ -detector placed at a distance of 15 cm from the beam axis. In addition, with the increased intensity of the beam, one may expect an increase of the incidental counts which, however should be compensated when the time-of-flight base is much larger (The rate of the incidental events in the previous experiment did not exceed 1%). Thus one should expect that the detection of the hypernuclear  $\gamma$ -spectra will be feasible under the new experimental conditions.

In a 100 hour run, the total number of the 1.09 MeV  $\gamma$ -quanta observed for the  ${}^7\text{Li}$  target in coincidence with the 53 MeV decay pions should amount to 500 if they are emitted from the  ${}^4_{\Lambda}\text{H}^x$  and to 50 only if it were the  ${}^4_{\Lambda}\text{He}^x$  hypernucleus. These predictions can be deduced from an evaluation of the solid angle ( $\sim 20\%$ ) subtended by the range telescope for the 53 MeV pions. Here account has also been taken of the above mentioned reduction of the effect by the factor of about 10 and of the 53 MeV pion decay rate of the  ${}^4_{\Lambda}\text{H}$  hypernucleus.

#### IV. MACHINE TIME AND MANPOWER

Nine weeks of the machine time during 1974 would be desired.

Eleven physicists and two technicians from Lyon and Warsaw could participate in the experiment during the running time.

REFERENCES

1. A. FILIPKOWSKI et al.  
Letter of Intent., PH I/COM/73-1
2. A. BAMBERGER et al.  
Phys. Lett., 36 B, (1971), 412
3. A. BAMBERGER et al.  
Excited states of light hypernuclei  
to be published
4. R.C. HERNDON, Y.C. TANG  
Phys. Rev., 159, (1967), 853
5. A. BAMBERGER et al.  
CERN Report, No 72-2, (1972)

FIGURE CAPTIONS

- Figure 1       $\Lambda$ -p total cross-section as a function of the c.m energy.  
The lines are calculated by Bamberger et al.<sup>3)</sup> for various  
 $\Lambda$ N potentials fitted to the elastic scattering and hypernuclear  
data and to the excitation energy  $E_{\gamma} = 1.09$  MeV ascribed to  
(a)  ${}^4_{\Lambda}\text{H}^*$ , (b)  ${}^4_{\Lambda}\text{He}^*$
- Figure 2      The experimental set-up
- Figure 3      A simplified scheme of electronics

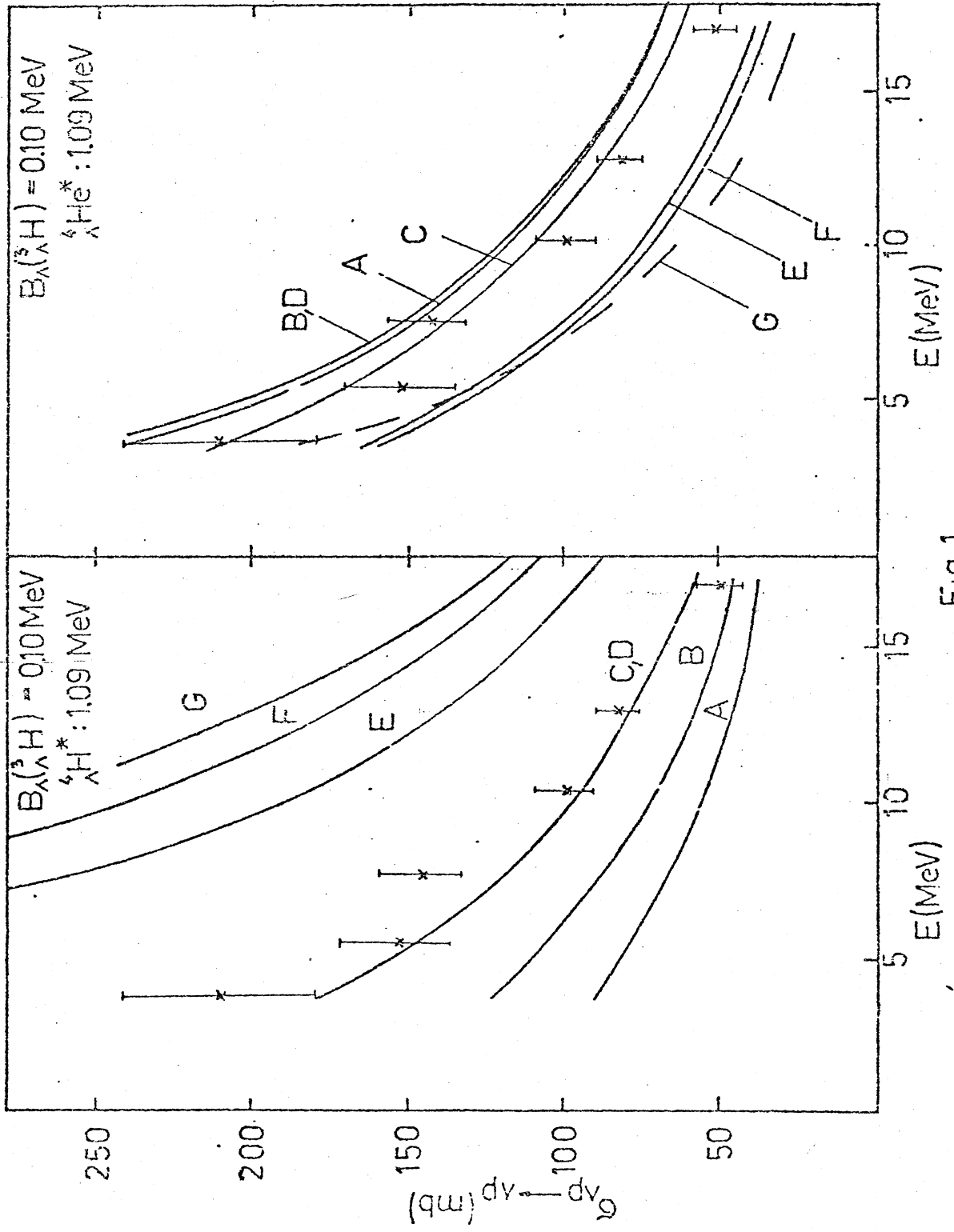


Fig 1



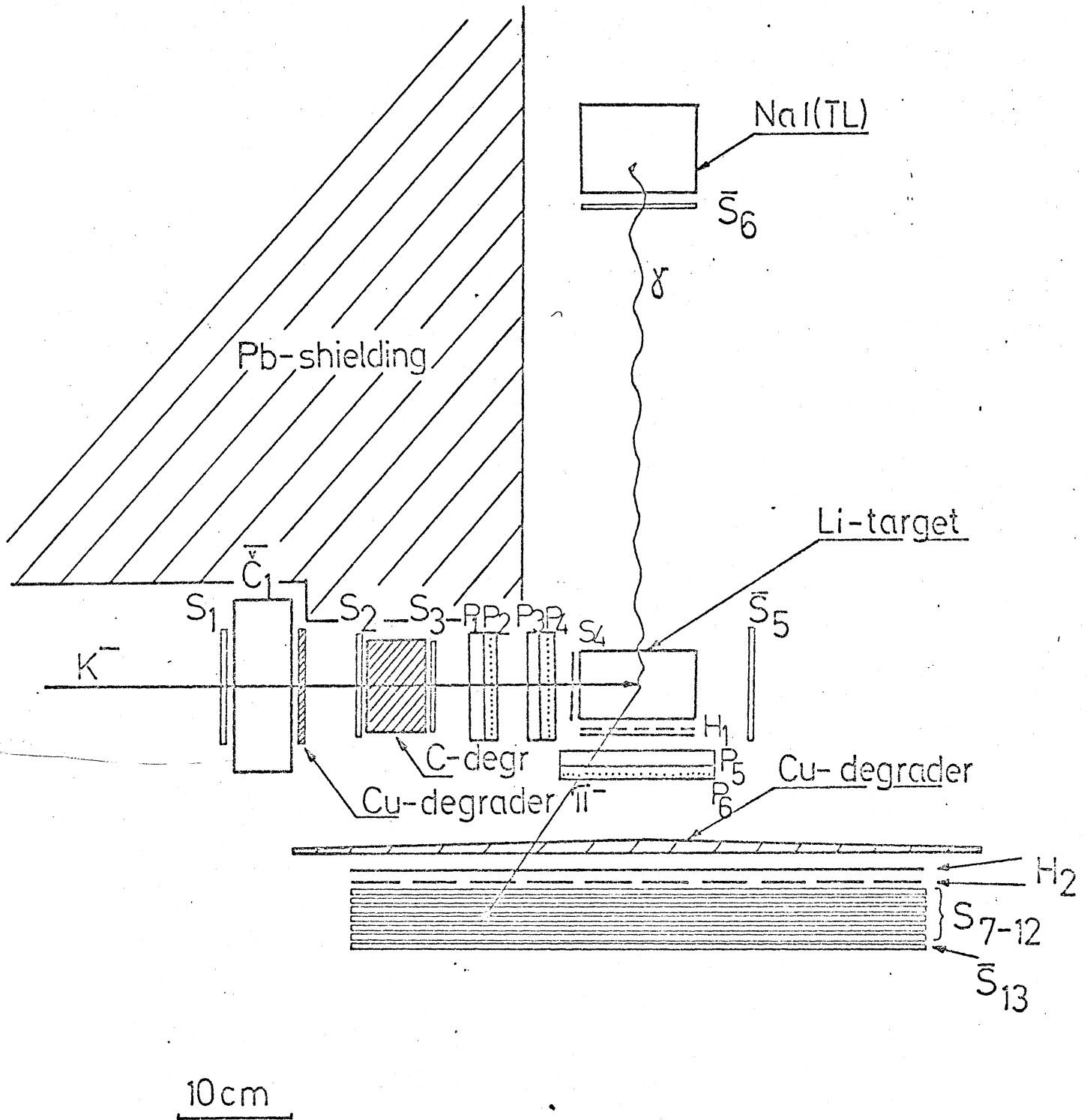


Fig 2

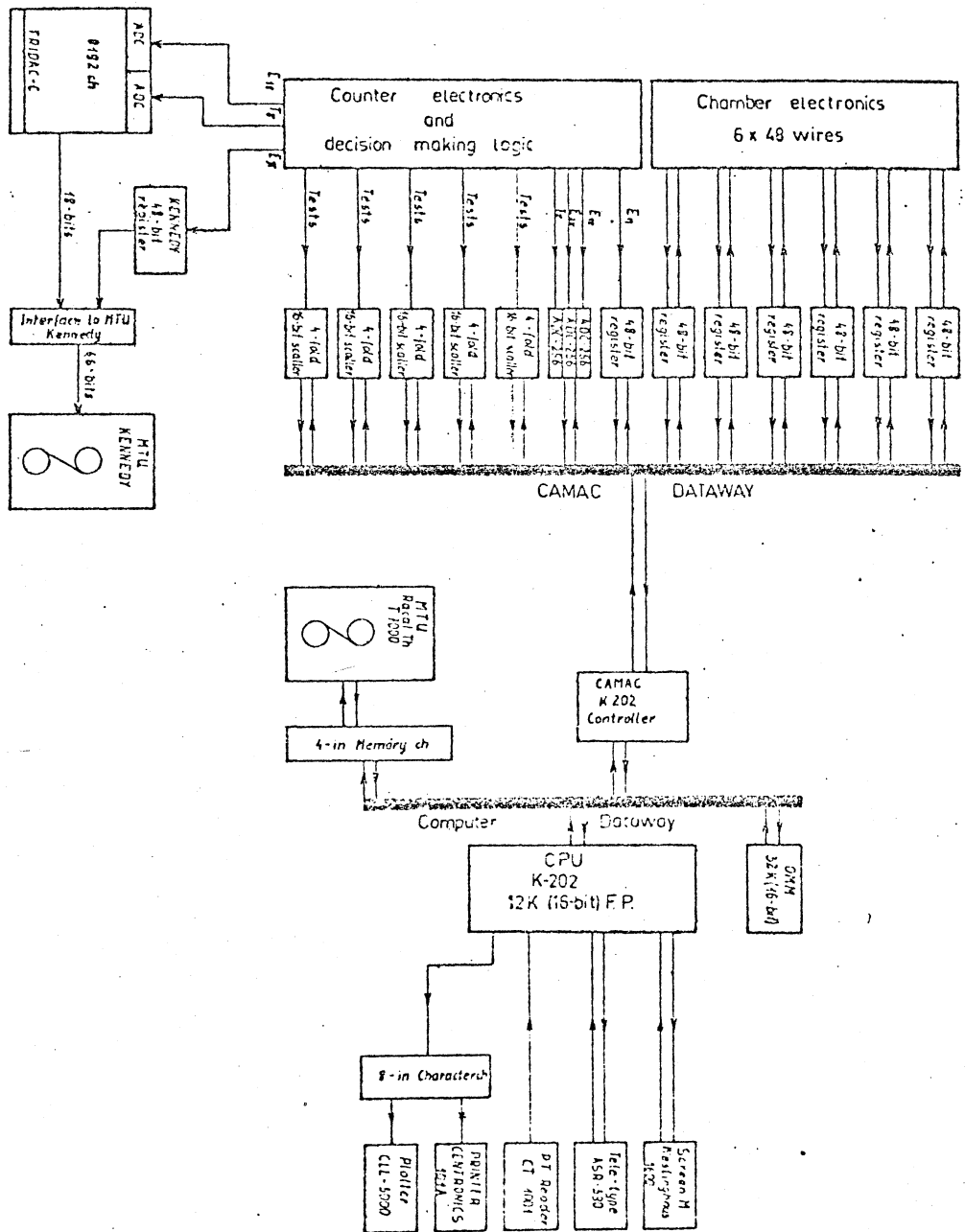


FIGURE 3