

Measurement of the CKM angle γ in $B^\pm \rightarrow DK^{*\pm}$ decays

Motivation

- Cabibbo-Kobayashi-Maskawa (CKM) matrix describes the quark mixing

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{V_{CKM}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- Unitarity of V_{CKM} represented by a triangle in the complex plane

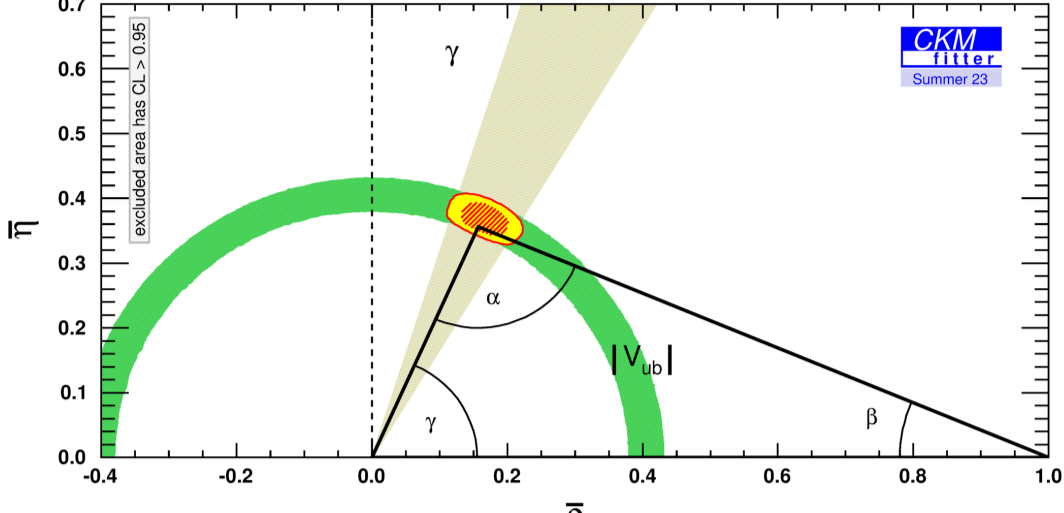
$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

- Weak phase γ is the only angle easily accessible at tree level

$$\gamma = \arg\left(-\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*}\right)$$

- Hadronic parameters can be determined from data => theoretical uncertainty on γ is negligible

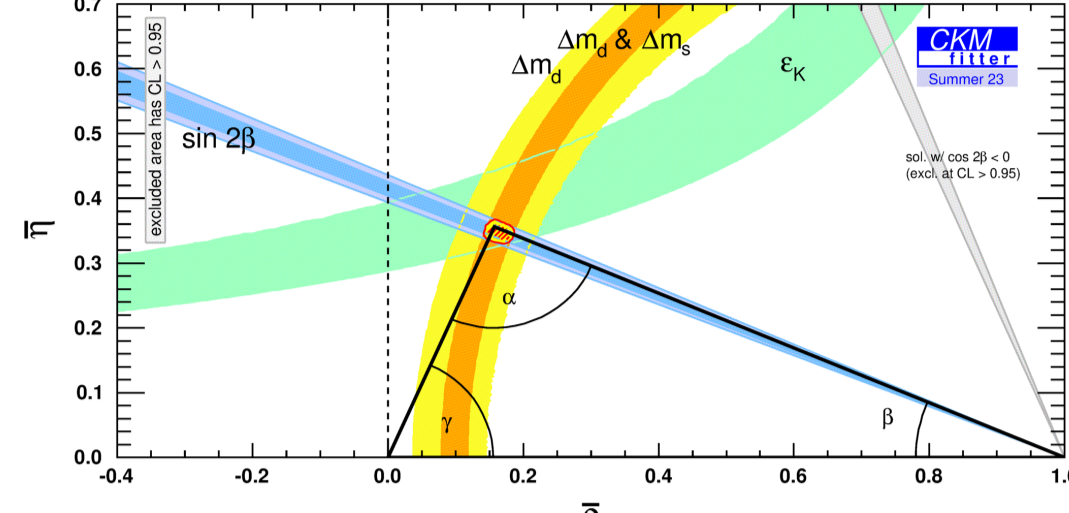
Tree-level (direct measurement)¹



$$\gamma = (66.4_{-3.0}^{+2.8})^\circ$$

World average (HFLAV²)

Loop-level (indirect measurement)¹



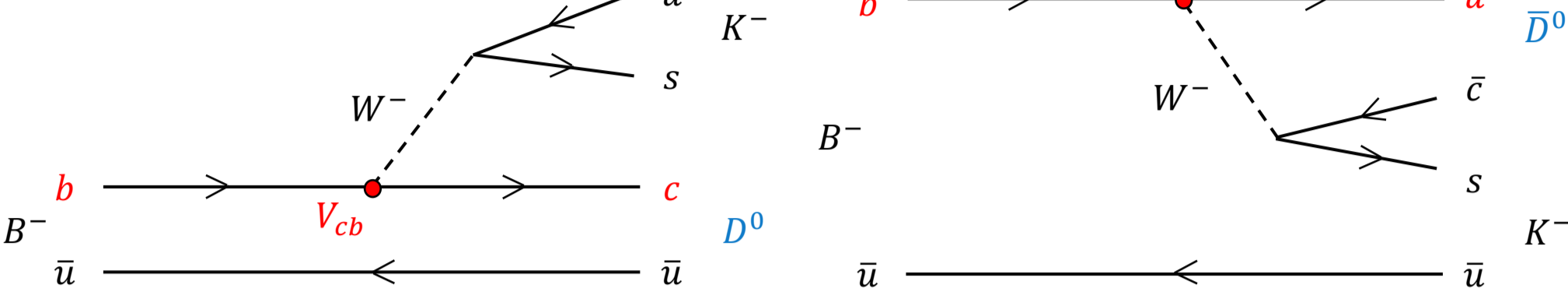
$$\gamma = (66.29_{-1.86}^{+0.72})^\circ$$

- Direct measurements of γ at tree level are expected to be benchmarks of the Standard Model
- Indirect measurements: global fits to the unitary triangle. Inputs include loop processes, where New Physics effects are expected to contribute

A discrepancy between direct and indirect measurements would be a clear sign of New Physics

- Direct measurements of γ in $B \rightarrow DK$ like decays. The D meson is a superposition of D^0 and \bar{D}^0 states, which are reconstructed in common final states. Both D^0 and \bar{D}^0 need to be able to decay to the same final state

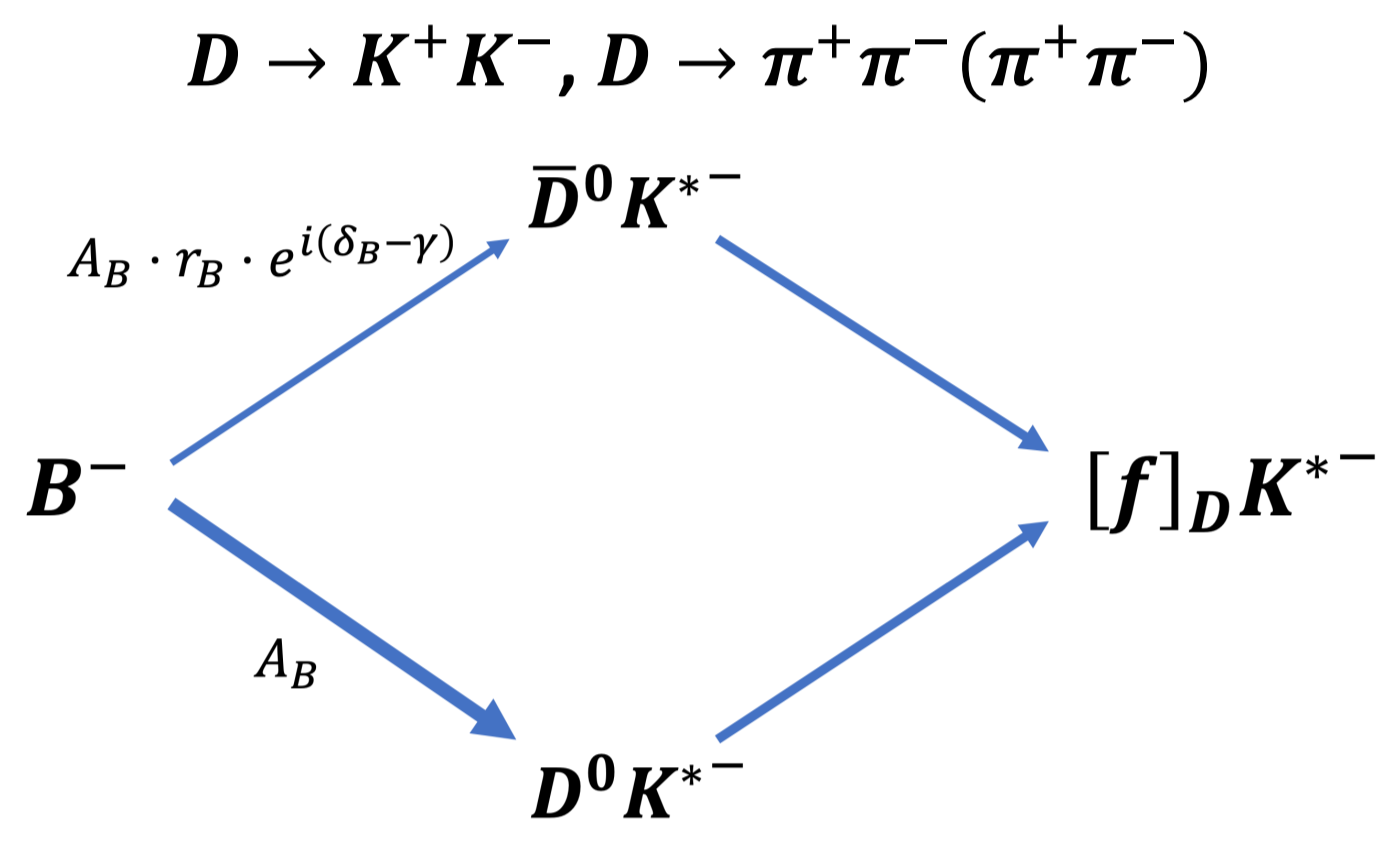
- Interference between $b \rightarrow cW$ and $b \rightarrow uW$ transitions gives sensitivity to γ



- Strategy to cover all B and D decay combinations to improve overall sensitivity to γ

- Measurement of γ in $B^\pm \rightarrow DK^*(892)^\pm$ decays. Comprehensive study of a range of D decay modes profiting from information between the different modes

GLW modes^{3,4}: (quasi)-CP eigenstates



- Measure rate ratios to the favoured mode
- Measure rate asymmetries between B^- and B^+
- Relatively smaller observable CP violation due to amplitudes of different sizes

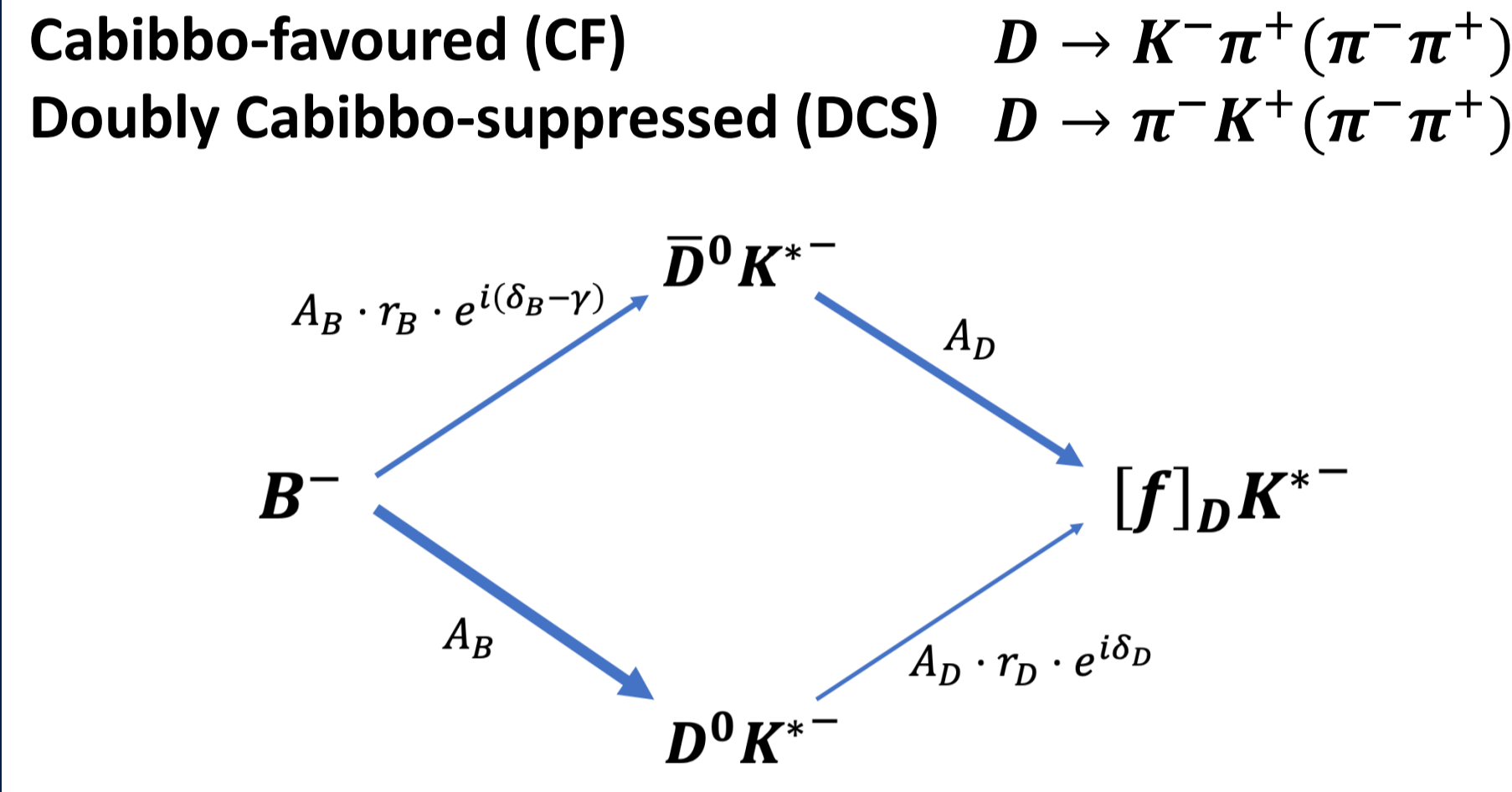
$$R_{CP} = \frac{\Gamma(B^- \rightarrow [h^+ h^-]_D K^{*\mp}) + \Gamma(B^+ \rightarrow [h^+ h^-]_D K^{*\mp})}{\Gamma(B^- \rightarrow [K^- \pi^+]_D K^{*\mp}) + \Gamma(B^+ \rightarrow [K^+ \pi^-]_D K^{*\mp})} \frac{B(D^0 \rightarrow K^- \pi^+)}{B(D^0 \rightarrow h^+ h^-)}$$

$$= 1 + r_B^2 + 2 \kappa r_B \cos(\delta_B) \cos(\gamma)$$

$$A_{CP} = \frac{\Gamma(B^- \rightarrow [h^+ h^-]_D K^{*\mp}) - \Gamma(B^+ \rightarrow [h^+ h^-]_D K^{*\mp})}{\Gamma(B^- \rightarrow [h^+ h^-]_D K^{*\mp}) + \Gamma(B^+ \rightarrow [h^+ h^-]_D K^{*\mp})}$$

$$= \frac{2 \kappa r_B \sin(\delta_B) \sin(\gamma)}{R_{CP}}$$

ADS modes⁵: non-CP eigenstates



- External inputs: D decay parameters r_D, δ_D
- Maximal interference due to similar sized amplitudes

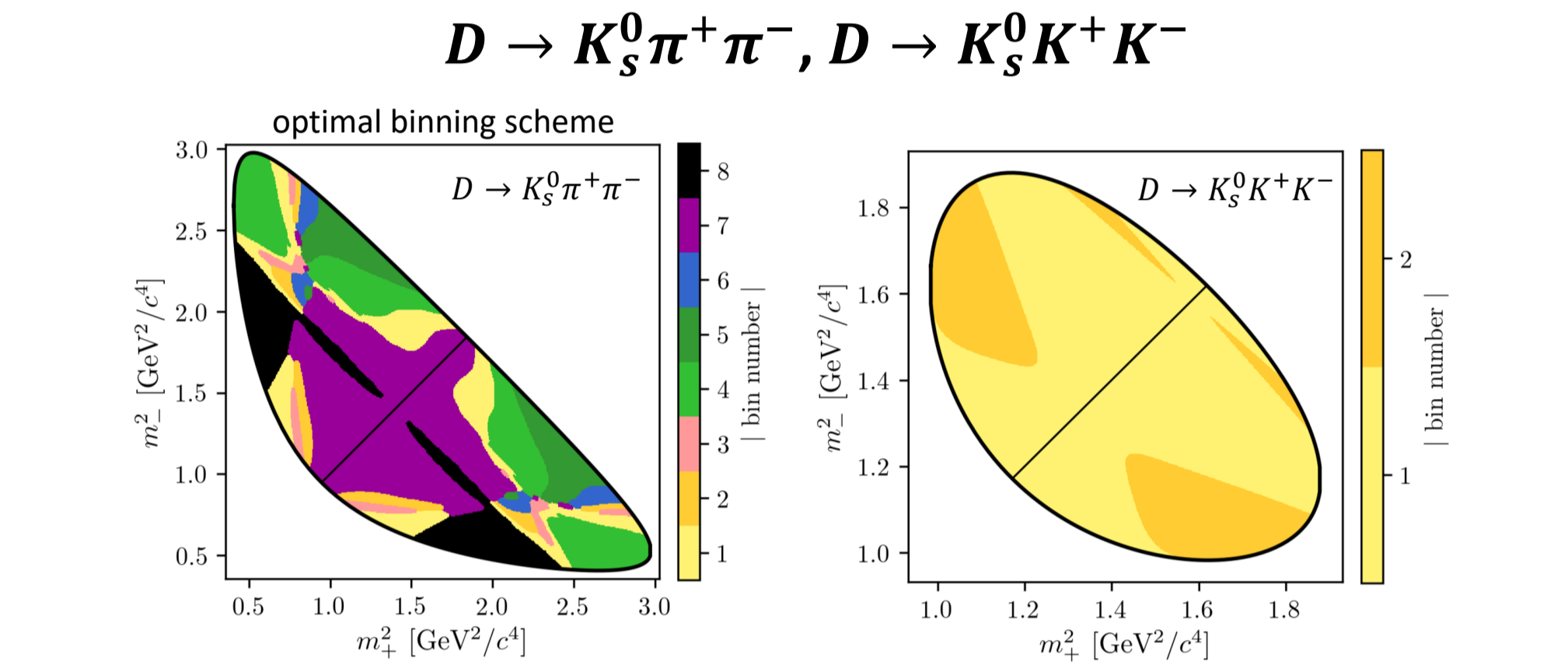
$$R_{ADS} = \frac{\Gamma(B^- \rightarrow [\pi^- K^+]_D K^{*\mp}) + \Gamma(B^+ \rightarrow [\pi^+ K^-]_D K^{*\mp})}{\Gamma(B^- \rightarrow [K^- \pi^+]_D K^{*\mp}) + \Gamma(B^+ \rightarrow [K^+ \pi^-]_D K^{*\mp})}$$

$$= r_B^2 + r_D^2 + 2 \kappa r_B r_D \cos(\delta_B + \delta_D) \cos(\gamma)$$

$$A_{ADS} = \frac{\Gamma(B^- \rightarrow [\pi^- K^+]_D K^{*\mp}) - \Gamma(B^+ \rightarrow [\pi^+ K^-]_D K^{*\mp})}{\Gamma(B^- \rightarrow [\pi^- K^+]_D K^{*\mp}) + \Gamma(B^+ \rightarrow [\pi^+ K^-]_D K^{*\mp})}$$

$$= \frac{2 \kappa r_B r_D \sin(\delta_B + \delta_D) \sin(\gamma)}{R_{ADS}}$$

BPGGSZ modes^{6,7}: three-body self-conjugate states



- Complex system of resonances
- Interference appears as different distributions of the D meson Dalitz plot for B^- and B^+
- First measurement at LHCb in this channel

$$x_\pm = r_B \cdot \cos(\delta_B \pm \gamma)$$

$$y_\pm = r_B \cdot \sin(\delta_B \pm \gamma)$$

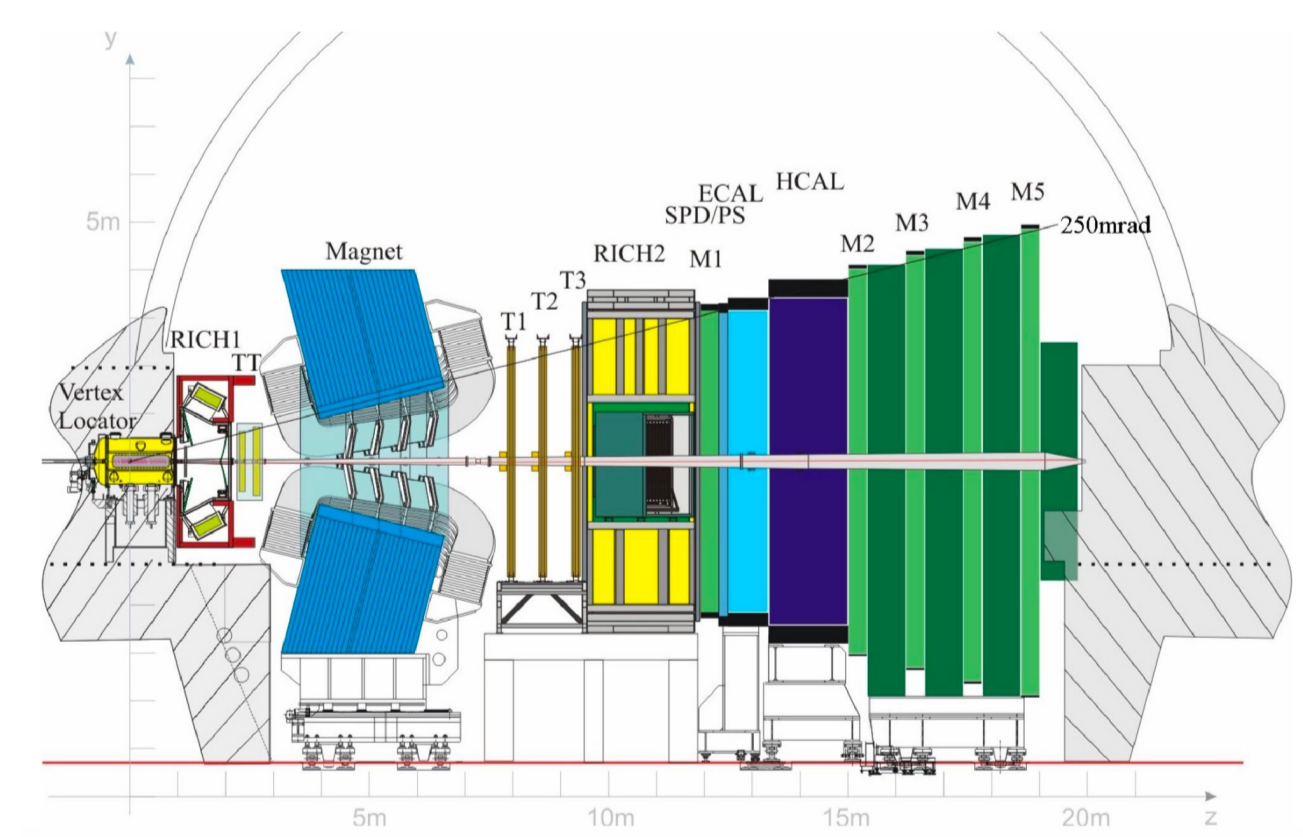
Yields in each Dalitz bin

$$N_{\pm i} \propto F_{\pm i} + (x_\pm^2 + y_\pm^2) F_{\mp i} + 2 \kappa \sqrt{F_i F_{-i}} (x_\pm c_{\pm i} + y_\pm s_{\pm i})$$

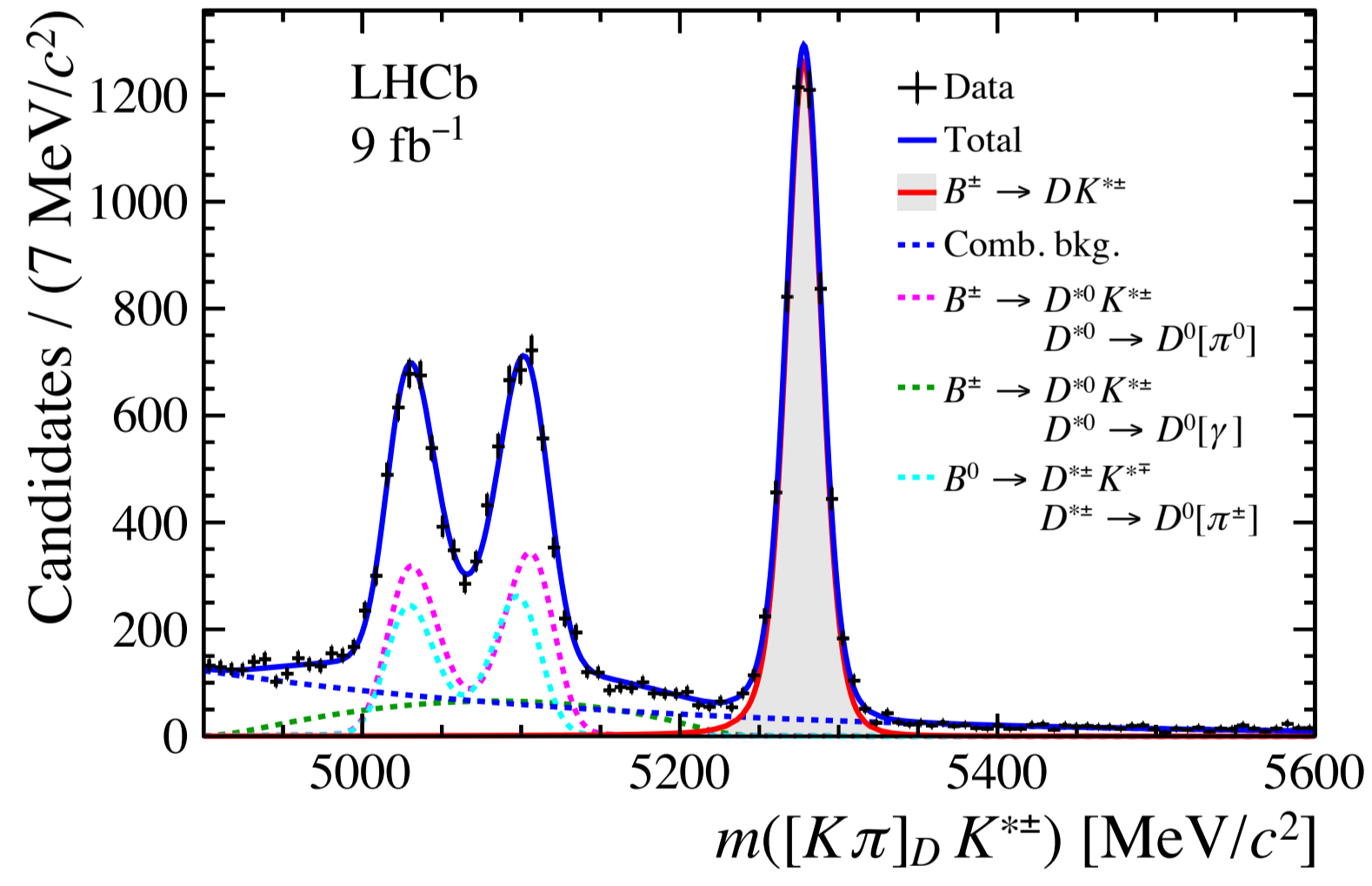
Fractional yield of flavour-tagged D^0 (for non-resonant contribution) Coherence factor (input from CLEO⁸+BESIII^{9,10,11}) Strong-phase differences

Candidate selection and background determination

- LHCb detector¹²:
 - Vertex Locator (VELO): secondary vertices measured with good resolution
 - Tracking system: good momentum resolution
 - Particle Identification (PID): distinguish pions from kaons



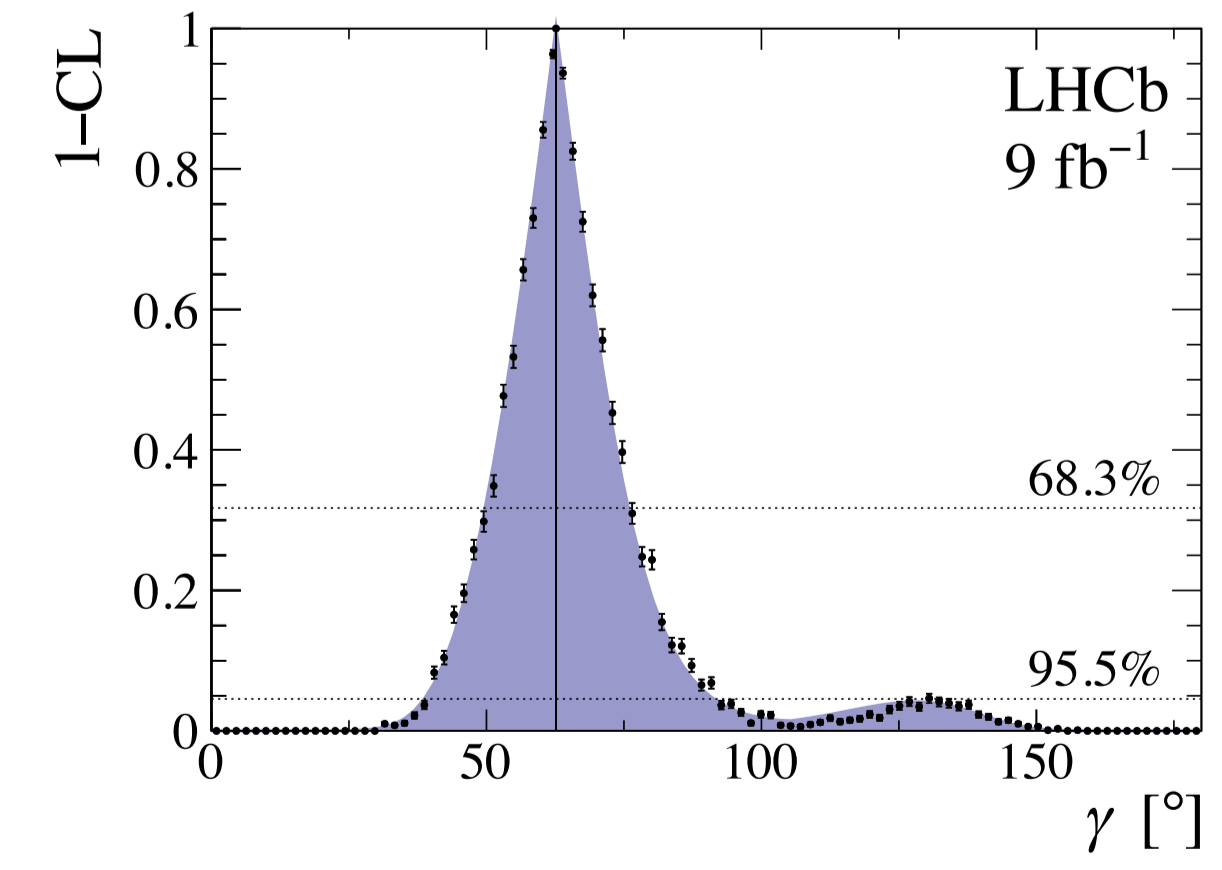
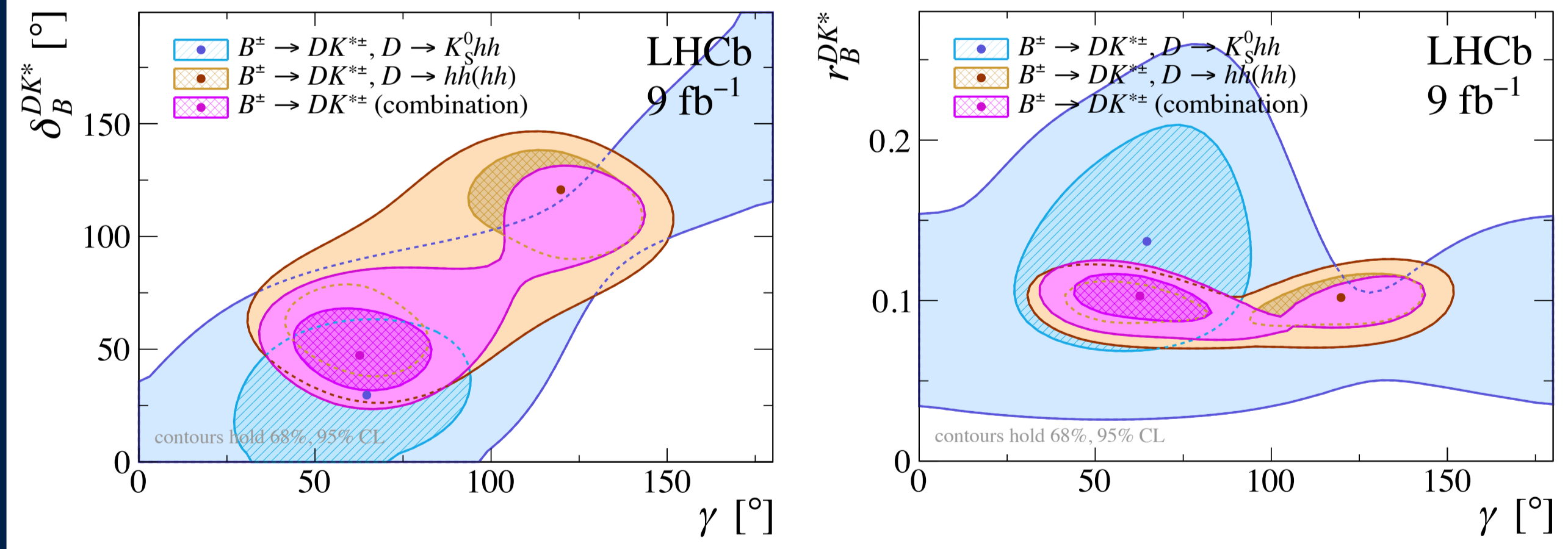
- Boosted Decision Tree (BDT) efficiently reduces combinatorial background
- Peaking backgrounds removed by flight distance cuts



- Mass parametrisation benefits from the high statistics favoured mode
- Advantages of this channel:
 - Clean signal peak, well separated from partially reconstructed backgrounds
 - Large purity
 - No counterpart of misidentified background from $B^\pm \rightarrow D\pi^\pm$ in $B^\pm \rightarrow DK^\pm$

Results and interpretation

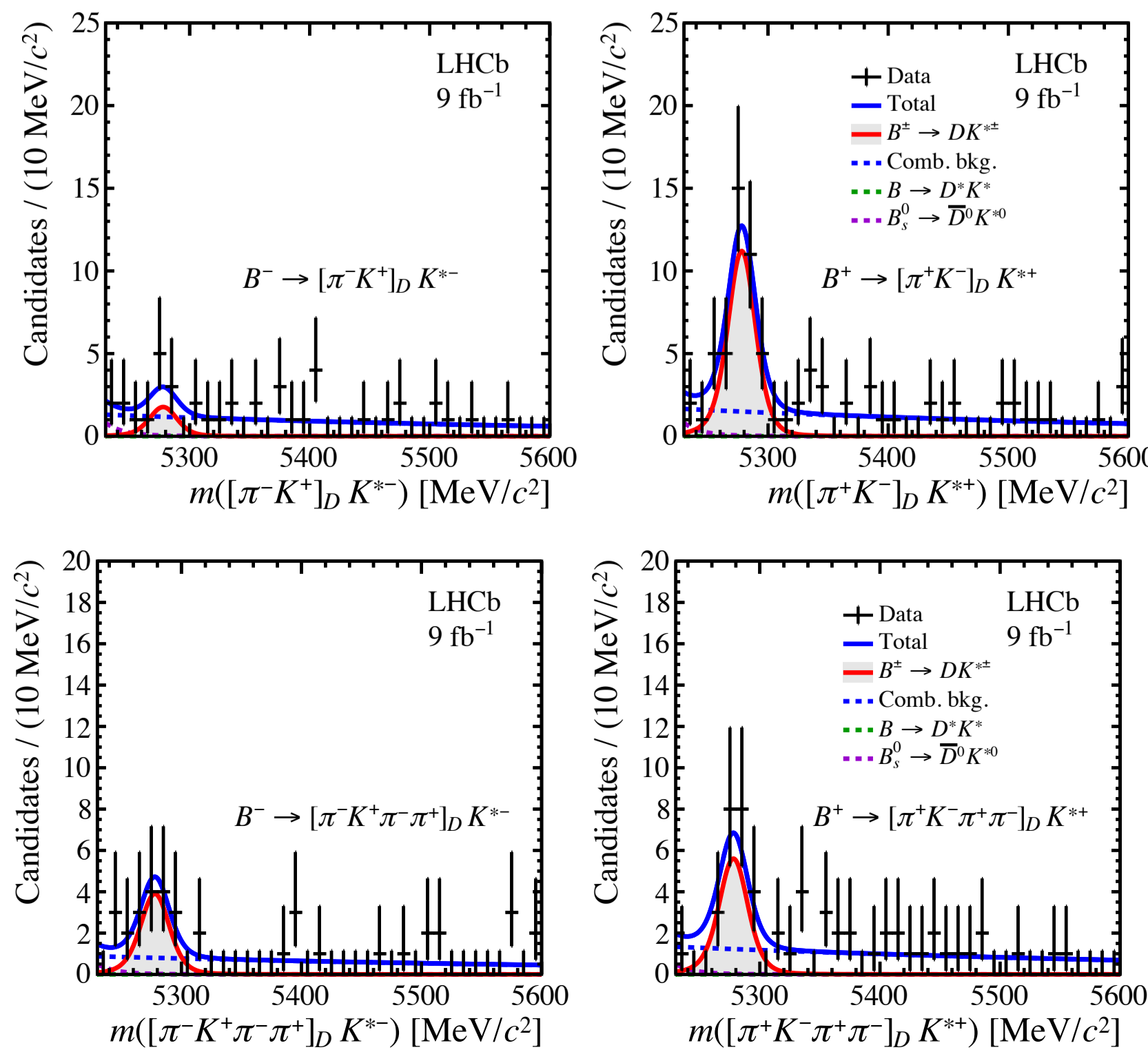
- Determination of the physics parameters of interest from the measured CP observables
- Multiple solutions obtained from two- and four-body decay modes, due to trigonometric equations relating CP observables to physics parameters
- Three-body decay modes lead to a single solution



- Combination yields $\gamma = (63 \pm 13)^\circ$
- Result for γ consistent with world average. Model-independent results, using external strong-phase inputs
- Valuable addition to γ measurements, included in the latest LHCb combination¹³

Determination of CP-violating observables

- Simultaneous fit for the different categories defined by B charge and D decay mode to measure the CP observables
- Small asymmetries within the favoured modes
- Larger asymmetries observed for the suppressed modes and CP-eigenstates modes
- CP observables measured for all three types of D decay modes considered in this analysis
- First observation of the suppressed $B^\pm \rightarrow [\pi^\pm K^\mp]_D K^{*\pm}$ and $B^\pm \rightarrow [\pi^\pm K^\mp \pi^\pm \pi^\mp]_D K^{*\pm}$ decays



References

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