# EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

# Proposal to the ISOLDE and Neutron Time-of-Flight Committee

Measurement of the <sup>238</sup>U(n,γ) cross section at n\_TOF

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#### **Abstract**

The <sup>238</sup>U(n,γ) cross section is one of the key reaction channels for nuclear applications, and small changes in it have a major impact on the results of many neutronic calculations. In recent years, an effort has been made to improve knowledge of this cross section, and four new measurements (two of them at n\_TOF) and new evaluations have been performed. Despite this, there is still inconsistent data and sizeable differences between evaluated cross sections, and new evaluations are currently being worked on. Here we propose a new measurement at n\_TOF with a new experimental setup focused on resolving some of the existing inconsistencies, thereby being able to improve the evaluations.

**Requested protons**: 3.2·10<sup>18</sup> protons on target **Experimental Area**: EAR1

### **1 INTRODUCTION**

Neutron capture on  $^{238}$ U is one of the key reaction channels for nuclear applications [\[1\]](#page-5-0)[-\[4\].](#page-5-1) In particular, <sup>238</sup>U is the major component of the light water reactor fuels, so the <sup>238</sup>U(n,y) cross section is one of the most relevant to perform calculations of many kinds: inventory, reactivity, criticality, etc. for both present and advanced concept reactors.

Due to the importance of  $^{238}$ U in nuclear applications, recent efforts have been made to improve existing evaluations. In particular, new neutron capture measurements have been performed at Los Alamos Neutron Science Center [\[5\],](#page-5-2) GELINA [\[6\]](#page-5-3) and n\_TOF [\[7\]\[8\];](#page-5-4) and two new evaluations of  $^{238}$ U have been released: JEFF-3.3 [\[9\]](#page-5-5) and IAEA-CIELO [\[10\],](#page-5-6) both using the mentioned new datasets. The IAEA CIELO evaluation was then adopted by ENDF/B-VIII.0.

Despite all these efforts, the work is currently continuing to improve the evaluations of <sup>238</sup>U. Proof of this is the new JENDL-5.0 evaluation [\[11\],](#page-5-7) which is based on the IAEA-CIELO but with some corrections; and the existence of the IAEA INDEN project [\[12\],](#page-5-8) a continuation of IAEA-CIELO which continues studying some of the most important isotopes for nuclear applications, including <sup>238</sup>U. Moreover, the JEFF-3.3 and the IAEA-CIELO<sup>238</sup>U evaluations present sizeable differences [\[13\].](#page-5-9)

Regarding the <sup>238</sup>U(n,γ) measurements carried out at n\_TOF, both were made with the same sample and in the same experimental area (EAR1), but different detectors were used. One of the measurements was performed with two  $C_6D_6$  detectors [\[7\],](#page-5-4) and the other with the n\_TOF Total Absorption Calorimeter (TAC) [\[8\].](#page-5-10) Although the <sup>238</sup>U(n,y) has already been measured twice at n\_TOF, we propose to perform an additional measurement mainly due to the following two reasons.

The first one is related with the  $^{238}U(n,y)$  cross section in the lower energy part of the Resolved Resonance Region (RRR). Both the JEFF-3.3 and the IAEA-CIELO evaluations take the resonance parameters up to 1200 eV from Kim et. al [\[6\],](#page-5-3) which suggests a new average radiation width of 22.5 meV, which is a bit lower than the previous evaluated value of 23.0 meV. This apparently causes problems in the interpretation of some integral benchmarks. A collaborative effort was made between JRC-Geel and PSI to verify this problem, and it was found that by adopting a radiation width of about 22.7 meV for all resonances a consistent description of the benchmark is obtained. Additional <sup>238</sup>U(n,γ) experimental data will help to verify the average radiation width of 22.7 meV. The sample used for the two <sup>238</sup>U(n,γ) measurements at n\_TOF was rather thick (0.375(2) g/cm<sup>2</sup>, i.e. 9.56(5) ·10<sup>-4</sup> atoms/barn) so the resonances at low energies were saturated; and saturated resonances are not the most appropriate for measuring resonance widths, since the fitted values become very dependent on the resolution function and multiple interaction corrections. A new measurement with a thin sample would solve these problems and allow the parameters of the largest low-energy resonances to be measured much more precisely.

The second reason is related with the cross section above 100 keV. Only two of the four new neutron capture measurements mentioned at the beginning of this document extends above 100 keV: the ones from Ullmann et al. [\[5\]](#page-5-2) (Los Alamos) and Mingrone et al. [\[7\]](#page-5-4) (n\_TOF). Both of them have reported cross sections larger than the evaluated ones above 100 keV. This is discussed in the IAEA-CIELO evaluation publication (see Fig. 15 of [\[10\]](#page-5-6) and associated text and Fig. 2 of [\[14\]](#page-5-11) and associated text), where they conclude that both the data from Los Alamos and n\_TOF are probably wrong due to missing corrections. We show Fig. 2 o[f \[14\]](#page-5-11) in [Figure 1,](#page-2-0) which illustrates this situation.



<span id="page-2-0"></span>**Figure 1 Average capture cross section for <sup>238</sup>U(n,γ) as a function of neutron energy. The recent measurement at Los Alamos**  (Ullmann et al[. \[5\]\)](#page-5-2), GELINA (Kim et al. [\[6\]\)](#page-5-3) and n\_TOF (Mingrone et al. [\[7\]](#page-5-4) with  $C_6D_6$  and Wright et al. [\[8\]](#page-5-10) with the TAC) are **compared to three different evaluations: GMA, Calrson et al, and JEFF-3.2. The figure has been obtained fro[m \[14\].](#page-5-11)**

In addition to the thin sample, we propose to use a thick sample to perform a new measurement of the <sup>238</sup>U(n,γ) cross section in order to solve these discrepancies. One of the most important corrections in the keV region at n\_TOF is the one related to the background produced by in-beam γ-rays (~40% of the background, according to Fig. 2 of [\[7\]\)](#page-5-4). The correction is performed by measuring a <sup>nat</sup>Pb sample and using black-resonance filters in the neutron beam. As mentioned in [\[7\],](#page-5-4) it was possible to use the filters to determine the background only below 100 keV, and the determination of the background above 100 keV relies in the γ-ray energy spectra obtained from Monte Carlo calculations. The new measurement we propose will improve this situation for three reasons:

- 1. The new n TOF spallation target produces less γ-rays (~50% less according to Monte Carlo calculations performed during the design phase) [\[15\].](#page-5-12)
- 2. We will use a different experimental setup, consisting also in  $C_6D_6$  detectors, but located at different forward and backward angles with respect to the neutron beam. The previous measurement used two detectors located at  $125^\circ$  with respect to the beam. For the same threshold in the detectors, such a setup more sensitive to scattered γ-rays than one made of detectors placed backwards. With the new setup proposed, the contribution of the in-beam γrays to the background will be different for each detector, thus allowing the determination of this background component better estimated and corrected for.
- 3. We will use additional filters in the neutron beam.

#### **2 EXPERIMENTAL SETUP**

The measurement will be performed at EAR1, since it has a better resolution function than EAR2.

For the detection system, we are considering two different possibilities. The first one is to use the same setup as in the  $^{181}$ Ta(n,γ) [\[16\]](#page-5-13) and  $^{nat}$ Er(n,γ) [\[17\]](#page-5-14) measurements performed in May-June 2023, which is shown in the left panel of [Figure 2.](#page-3-0) This setup consists in three  $C_6D_6$  BICRON detectors, one of them located at 90<sup>°</sup> with respect to the neutron beam and the other two at 125<sup>°</sup>; and three sTED small  $C_6D_6$ detectors located at 90 $^{\circ}$ , 110 $^{\circ}$  and 130 $^{\circ}$ . According to the analysis work carried out to date, the results seem excellent. The detection efficiency of this setup for <sup>238</sup>∪(n,γ) cascades is expected to be ~17%.



**Figure 2 On the left, experimental setup used in the <sup>181</sup>Ta(n,γ) and natEr(n,γ) measurements in May-June 2023. On the right, geometry of a cluster of 24 C6D<sup>6</sup>** *sliceTED* **modules implemented in the Geant4 Monte Carlo transport code, an experimental setup which is at the moment in the design phase.**

<span id="page-3-0"></span>The second possibility for the experimental setup would be to use a new detector, for the moment called *sliceTED*, which is still in the design phase. It consists in a high efficiency  $C_6D_6$  setup made of several  $C_6D_6$  modules. An illustration of a possible sliceTED detector is presented in the right panel of [Figure 2.](#page-3-0) The operational principle of this detector, described in detail in [\[18\],](#page-5-15) is similar to the standard  $C_6D_6$  total energy detectors, but with a larger (n,y) detection efficiency. The design presented i[n Figure 2,](#page-3-0) made of 24  $C_6D_6$  modules with 3"x3"x3" dimensions each, is expected to have a detection efficiency of ~60% for <sup>238</sup>U(n,γ) cascades. In addition to the high efficiency, it has the advantage of detecting y-rays emitted at different angles, making the results little sensitive to anisotropies in the  $\gamma$ -ray emission. The main drawback will be probably an increase in the *neutron sensitivity*.

Concerning the samples, we will use a thin and a thick sample. For the thick sample, we are planning to use the same as in the previous measurements performed at n\_TOF [\[7\]\[8\],](#page-5-4) which is a 99.999% pure <sup>238</sup>U metallic sample provided by EC-JRC-Geel with 6.125(2) grams and 53.90  $\times$  30.30 mm<sup>2</sup> (9.56(5) $\cdot$ 10<sup>-4</sup> atoms/barn). For the thin sample, we estimated that a  $n$ <sup>at</sup>U sample with an areal density of 1.8·10<sup>-5</sup> atoms/barn (7.07 mg/cm<sup>2</sup>) will be appropriate for our purposes.

The number of protons requested for this measurement is  $3 \cdot 10^{18}$ , which is broken down in the different measurement configurations in [Table 1.](#page-5-16) The number of protons for the thin sample measurement has been estimated to achieve sufficient statistics at the largest resonances so that the uncertainties in the resonance parameters are dominated by systematic effects. An example of the expected results for two resonances is shown in [Figure 3.](#page-4-0) For the thick sample measurement, the number of requested protons leads to ~3000 counts per bin in the keV region when using 100 bins per decade. This will lead to uncertainties due to counting statistics of ~3% in each bin after subtracting the background (~2% before the subtraction). The expected number of counts in the measurement when measuring the thick target is shown in [Figure 4.](#page-4-1) Additional measurements with <sup>197</sup>Au, natpb and nat<sub>C</sub> samples and with filters, needed as a reference measurement and for subtracting the background, have been also considered.



<span id="page-4-0"></span>**Figure 3 Expected <sup>238</sup>U(n,γ) yield when measuring the thin sample, for the first (left) and fourth (right) strongest resonances. The statistical fluctuations have been modelled assuming 0.5·10<sup>18</sup> protons on target when using the same experimental setup as in the <sup>181</sup>Ta(n,γ) and natEr(n,γ) measurements. The background also comes from the same experiment.**



<span id="page-4-1"></span>**Figure 4 Expected total number of counts in the <sup>238</sup>U(n,γ) measurements (100 bpd) when delivering 0.9·10<sup>18</sup> protons on target. The background comes from the <sup>181</sup>Ta(n,γ) and natEr(n,γ) measurements.**

# **Summary of requested protons:**





**Table 1 – Summary of the number protons on target requested for each of the configurations to be measured.**

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- <span id="page-5-7"></span>[11] O. Iwamoto et al., Japanese evaluated nuclear data library version 5: JENDL-5, Journal of Nuclear Science and Technology, 60:1, 1-60 (2023)[. https://doi.org/10.1080/00223131.2022.2141903](https://doi.org/10.1080/00223131.2022.2141903)
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# **Appendix**

## **DESCRIPTION OF THE PROPOSED EXPERIMENT**

Please describe here below the main parts of your experimental set-up:



## **HAZARDS GENERATED BY THE EXPERIMENT**

Additional hazard from flexible or transported equipment to the CERN site:



