EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

Proposal to the ISOLDE and Neutron Time-of-Flight Committee

Measurement of 28,29,30 Si (n, γ) capture cross sections to explain isotopic abundances in presolar grains

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Abstract: Neutron capture reactions on the stable silicon isotopes ^{28,29,30}Si are of key importance to understand the production of silicon in massive stars and isotopic abundances in presolar stardust grains. However, there are only few experimental data

at astrophysical neutron energies, and results from previous measurements are discrepant. We propose to measure neutron capture cross sections on ^{28,29,30}Si at EAR-1 and EAR-2, with the aim to obtain high accuracy data on neutron resonances, and at thermal neutron energies (25.3 meV), respectively.

Requested protons: 7×10^{18} (EAR-1) and 1.1×10^{18} (EAR-2) protons on target **Experimental Area:** EAR1 and EAR2

1 Introduction

The aim of this proposal is to determine accurate neutron capture cross sections on the stable silicon isotopes 28,29,30 Si, to increase accuracy and resolve discrepancies between previous measurements. The most recent measurement of 28,29,30 Si neutron capture cross sections was performed via the time-of-flight technique by Guber et al [1]. Maxwellian averaged cross sections were deduced from the measured resonance parameters combined with calculations of the direct capture component. For all isotopes, they found smaller stellar cross sections compared to previous data. Concerning the 28 Si(n, γ) and 29 Si(n, γ) cross sections, the only other data are more than 40 years old [2, 3] and while stellar cross sections for 29 Si are agreeing with [1] within uncertainties, there is a factor of about 2 difference for 28 Si. For the case of 30 Si(n, γ), the MACS in [1] was found to be almost 2 times smaller than results from an activation measurement by Beer *et al.* [4]. At present, the source of such a big difference for the 30 Si(n, γ) cross section between the two experiments is unknown.

A new measurement of these reactions is of high importance for astrophysical studies. First of all, the bulk of the ²⁹Si and ³⁰Si present today in the Milky Way and in the Sun are mostly made in the convective carbon-shell in massive stars at about 1 billion kelvin, and ejected by the supernova explosion [5, 6, 7, 8]. In these stellar conditions, the ^{28,29,30}Si neutron capture cross sections are crucial to shape the final yields and the relative abundances of ²⁹Si and ³⁰Si (see Fig. 1).

However, the relevance of these cross sections is not only limited to galactic chemical evolution and the origin of the Si isotopes in stars. Accurate cross section data are crucial for understanding the isotopic abundances of presolar grains. Presolar stardust grains are tiny inclusions found in meteorites which condensed from outflowing gasses of stellar winds and stellar explosions, shortly before the formation of the solar system [9, 10]. The unique isotopic abundance signatures of these grains can offer clues of their astrophysical origin, and they are essential to study stellar processes and improve the modelling of stellar evolution [11, 12, 13]. Presolar SiC grains of type Mainstream, type Y and type Z are made in low-mass Asymptotic Giant Branch (AGB) stars [9, 14, 15, 16]. In these stars, carbon freshly synthesized in the interior is dredged up to the surface and ejected through stellar winds. Then, SiC grains can form within the circumstellar envelope of these stars and their Si isotopic abundances reflect the signature of the internal neutron-capture nucleosynthesis superimposed on the composition of the pre-stellar gas, as modeled by the galactic chemical evolution. For these grains, isotopic silicon ratios have been accurately measured in the laboratory, with uncertainties lower than 5% (see, e.g., [17]). A precise knowledge of neutron-capture cross sections on the stable silicon isotopes is of paramount importance to disentangle the intrinsic stellar contribution from that of galactic chemical evolution. According to their analysis, comparing the Si abundances of presolar SiC grains of Type Y and Z to AGB models, Zinner et al. [17] favoured the value proposed by Guber et al. [1]. A new measurement of the Si isotope cross sections and in particular of the mysterious ${}^{30}\text{Si}(n,\gamma)$ is crucial to confirm those results.

SiC grains of type X and C likely originate from core-collapse supernovae [18, 19, 9]. Those belonging to the rare type C sub-group show peculiar abundances of Si and S isotopes. In particular, Pignatari et al. [19] showed that the strongly enhanced abundance of ${}^{32}S$

S 29	S 30	S 31	S 32	S 33	S 34	S 35	S 36
187 ms	1.18 s	2.58 s	94.99	0.75	4.25	87.5 d	0.01
β ⁺ γ 1384 βp 5.44; 2.13	β ⁺ 4.4; 5.1 γ 678	β ⁺ 4.4 γ 1266	σ 0.55 σ _{n, α} <0.0005	σ 0.46 σ _{n, α} 0.12 σ _{n, p} 0.002	σ 0.25	β ⁻ 0.2 no γ	or 0.24
P 28 268 ms β ⁺ 11.5 γ 1779; 4497	P 29 4.1 s	P 30 2.50 m	P 31 100	P 32 11.26 d	P 33 25.34 d	P 34 12.4 s	P 35 47.4 s
βp 0.680; 0.956	β ⁺ 3.9	β ⁺ 3.2	σ 0.17	β ⁻ 1.7	β ⁻ 0.2	β 5.4	β ⁻ 2.3
βα 2.105; 1.434	γ 1273	γ (2235)		no γ	no γ	γ 2127	γ 1572
Si 27	Si 28	Si 29	Si 30	Si 31	Si 32	Si 33	Si 34
4.16 s	92.223	4.685	3.092	3.62 h	172 a	6.18 s	2.77 s
β ⁺ 3.8 γ (2210)	σ 0.17	σ 0.12	or 0.107	β 1.5 γ (1266) σ 0.073	β 0.2 no γ σ<0.5	β 3.9; 5.8 γ 1848	β 3.1 γ 1179; 429; 1608
AI 26	Al 27	AI 28	AI 29	Al 30	Al 31	AI 32	Al 33
6.35 s 7.16 · 10 ⁵ a	100	2.346 m	6.6 m	3.60 s	644 ms	33 ms	41.7 ms
β ⁺ 3.2 γ 1809; 1130 σ _{n, a} 0.34 σ _{n, p} 1.97	σ 0.230	β 2.9 γ 1779	β 2.5 γ 1273; 2426; 2028	β 5.1; 6.3 γ 2235; 1263; 3498	β 5.6; 7.9 γ2317; 1695	β γ 1941; 3042; 4230	β βn γ 1941*; 4341; 1010

Figure 1: s-process flow in the mass region of interest for the measurement.

can only be explained by the presence of the long-lived ³²Si ($t_{1/2} = 132$ yr) in the ejecta, produced by neutron capture processes starting from the stable Si isotopes. ³²S is the stable radiogenic product of ³²Si and is not directly implanted in the grains. However, the neutron density has to be high enough to overcome the rather short-lived ³¹Si ($t_{1/2} = 2.6$ hr). The abundance of ³²S in the grains can therefore provide constraints on the neutron density reached during the SN explosion in the C-rich He shell material. All the neutron capture cross sections along the chain of Si isotopes (see Fig. 1) are an essential input in nucleosynthesis calculations to explain presolar grain abundances [19].

Finally, a detailed knowledge of silicon neutron capture rates is fundamental to properly compare theoretical silicon isotopic ratios with SiO maser radio observations of OH/IR stars, commonly linked to intermediate mass AGB stars (M> 5 M_{\odot}).

2 Experimental Setup

We propose to accurately measure 28,29,30 Si (n, γ) cross sections at n_TOF EAR-1 and EAR-2. At EAR-1, we will take advantage of the excellent neutron energy resolution and high instantaneous neutron flux, allowing a measurement up to a few hundred keV neutron energy. At EAR-2, we plan to measure the cross section at thermal neutron energies (25.3 meV), and the excitation function of the cross section at low energy. Such a study is not possible at EAR-1, due to uncertainties in the neutron beam profile at very low energy, where gravity effects play a role. In addition, the thermal neutron flux in EAR-1 is strongly suppressed by the borated water moderator and, therefore, a measurement of the thermal cross section is much less efficient compared to EAR-2. In both areas, we plan to detect the prompt γ radiation following a capture event using the well-established C₆D₆ detection system, consisting of 4 Legnaro type C₆D₆ detectors which have been specifically optimised for low neutron sensitivity [20]. In addition, we will complement the C₆D₆ detectors at EAR-2 with sTED detectors to measure the 4.9

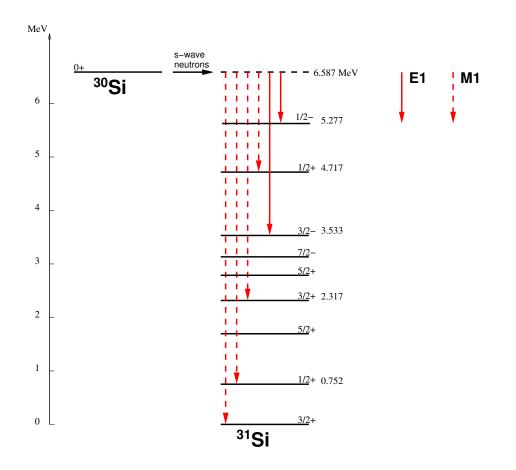


Figure 2: Primary γ -rays emitted following neutron capture events by ³⁰Si.

eV resonance in Au, since for this type of detector dead-time-corrections will be smaller, and as a consequence, a reliable normalisation to obtain absolute cross section values can be obtained. An example of the expected primary γ -rays emitted by s-wave neutrons capture by ³¹Si is shown in Figure 2. Additional γ 's are expected, obviously, from the decay cascade.

Each of the Si samples will consist of 1g of enriched (all three over 99%) metal, in form of a cylindrical disc of 2 cm diameter. Similar to other capture experiments, the data will be normalised using the saturated resonance technique on Au, measuring a Au sample of the same diameter as the Si samples. Backgrounds will be determined in runs with an empty sample holder, and the neutron scattering background will be studied by measuring with a natural Carbon sample.

For all silicon isotopes, the stellar cross section is determined by only very few resonances. This implies that these resonances need to be studied with high accuracy, due to their high impact on the stellar cross section. In addition, the direct capture component could play a significant role (details provided in [21]).

For our count rate estimates, we have taken into consideration the most recent evaluation of the neutron flux at EAR-1, the resolution function, and the cross sections available from the ENDF/B-VIII evaluation. For ³⁰Si, the first resonance at 2.235 keV has been removed as the presence of this resonance has been questioned and its non-existence confirmed in preliminary data taken with a natural-Si sample measured at n_TOF in 2021, during the commissioning. For the background estimation, we have taken data from a previous measurement with a similar setup. The panels to the left in Figure 3 show the count rate estimation with the assumptions described above for all three stable Si isotopes, assuming 2×10^{18} protons on target each. Only the resonant energy region is shown in the Figure. All resonances of interest are well resolved and have good counting statistics, as indicated by the star symbol which shows the total number of counts in each resonance. For background and normalisation measurements, we request 1×10^{18} protons. This will allow us to determine the background accurately between resonances, which may indicate if there is a strong, direct p-wave contribution to the cross section and high neutron energy.

Due to the 30 times higher neutron flux, we only require a modest amount of beam time for the EAR-2 measurement, 2×10^{17} protons per isotope. The right column of Figure 3 shows the expected count rate at low energy for each isotope. We have also included the expected background for this measurement due to reactions of neutrons in the experimental area (data from a recent capture measurement on Mo isotopes). The cross section at EAR-2 can be measured up to a few eV, which will provide important information on the contribution of sub-threshold resonances or the direct capture component. For normalisation and background studies, we request a further 5×10^{17} protons on target.

Summary of requested protons: 7×10^{18} protons EAR-1, 1.1×10^{18} at EAR-2 (can be split over 2 years)

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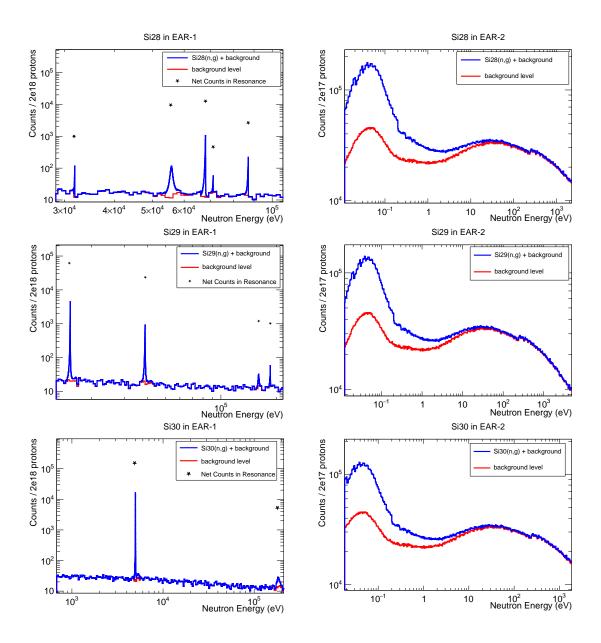


Figure 3: (Panels to the left) Count rate estimate for each of the stable Si isotopes for 2×10^{18} protons on target in EAR-1. The number of counts expected in each resonance is shown by a star symbol. (Panels to the right) Count rate estimate for each of the stable Si isotopes for 2×10^{17} protons on target in EAR-2. For both experimental areas, the background level has been estimated from previous measurements.

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Appendix

DESCRIPTION OF THE PROPOSED EXPERIMENT

Please describe here below the main parts of your experimental set-up:

Part of the experiment	Design and manufacturing						
If relevant, write here the name of	\boxtimes To be used without any modification						
the $\underline{\text{fixed}}$ installation you will be using	\Box To be modified						
[C6D6 detectors]							
If relevant, write here the name of the	\boxtimes To be used without any modification						
fixed installation you will be using [Si-	\Box To be modified						
Mon]							
If relevant, write here the name of	\boxtimes To be used without any modification						
the $\underline{\text{fixed}}$ installation you will be using	\Box To be modified						
[sTED]							
If relevant, describe here the name	□ Standard equipment supplied by a manufacturer						
of the flexible/transported equipment	\Box CERN/collaboration responsible for the design						
you will bring to CERN from your In-	and/or manufacturing						
stitute							
Small spare parts, such as detector	\boxtimes Standard equipment supplied by a manufacturer						
holders, cabling, spare detectors etc.	\boxtimes CERN/collaboration responsible for the design						
	and/or manufacturing						
[insert lines if needed]							

HAZARDS GENERATED BY THE EXPERIMENT

Additional hazard from flexible or transported equipment to the CERN site:

Domain	Hazards/Hazardous Activities	Description	
	Pressure		[pressure] [bar], [volume][l]
	Vacuum		
Mechanical Safety	Machine tools		
	Mechanical energy (moving parts)		
	Hot/Cold surfaces		
Cryogenic Safety	Cryogenic fluid		[fluid] [m3]
Electrical Safety	Electrical equipment and installations		[voltage] [V], [current] [A]
Electrical Safety	High Voltage equipment		[voltage] [V]
	CMR (carcinogens, mutagens and toxic		[fluid], [quantity]
to reproduction)			
	Toxic/Irritant		[fluid], [quantity]
Chemical Safety	Corrosive		[fluid], [quantity]
	Oxidizing		[fluid], [quantity]

	Flammable/Potentially explosive atmospheres		[fluid], [quantity]
	Dangerous for the environment		[fluid], [quantity]
Non-ionizing	Laser		[laser], [class]
radiation Safety	UV light		
radiation Salety	Magnetic field		[magnetic field] [T]
	Excessive noise		
Worlenlago	Working outside normal working hours		
Workplace	Working at height (climbing platforms, etc.)		
	Outdoor activities		
	Ignition sources		
Fire Safety	Combustible Materials		
	Hot Work (e.g. welding, grinding)		
Other hazards			
Other nazarus			