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Measurements with the PTB Bonner Sphere Spectrometer in High-Energy Neutron Calibration Fields at CERN (CERN-CEC Experiment H6J93, July 1993)

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(CERN-CEC Experiment H6J93, July 1993) in High-Energy Neutron Calibration Fields at CERN Measurements with the PTB Bonner Sphere Spectrometer

by

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The results are compared with calculations. calibration fields produced at the CERN Super Proton Synchrotron. neutron spectral fluences contributing to four mixed high-energy Abstract —— The PTB Bonner sphere spectrometer was used to measure

de CERN. Les résultats sont comparés avec des calculs. d'éta1onnage haute énergie produits au Super Synchrotron de Protons mesurer les spectres des neutrons contribuant à quatre champs mixtes Resume —— Le spectrométre multisphéres du PTB a été utilisé pour

Die Meßergebnisse werden mit Rechnungen verglichen. Proton Synchrotron fiir Kalibrierzwecke erzeugt wurden, zu bestimmen. verschiedenen hochenergetischen gemischten Feldern, die am Super CERN eingesetzt, um die spektrale Neutronenfluenz in vier Zusammenfassung —— Das Bonner—Kugel Spektrometer der PTB wurde am

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1. INTRODUCTION

aircraft personnel during flights at high altitude [ll. the ICRP's acknowledgement of the occupational exposure of civil high-energy particles in cancer therapy and, last but not least, to become established, but also to the increasing use of the protection practice at high-energy (GeV+TeV) accelerators, which has metry, of high-energy particle fields is due not only to radiation Increased interest in the dosimetry, and by implication the spectro

for other instruments. equivalent with acceptable accuracy, and so to establish a reference the neutron spectra, spectrometry is applied to determine the dose As the dose-equivalent responses of the present dosemeters depend on particular attention is paid to neutron dosimetry in these fields. major contributor to the dose equivalent, and for this reason shields of high-energy accelerators, the neutron component is the In the complex stray radiation fields encountered outside the

participated with the PTB "C" Bonner sphere spectrometer. these series which was carried out in July 1993, the PTB with spectrometers and dosemeters in these investigations. In one of at CERN, and several European laboratories were invited to take part measurements was performed in various high-energy radiation fields Within the framework of a CERN — CEC project a series of

the PTB Bonner spheres are reported. published elsewhere. In the present paper the results obtained with A comparison of the results obtained by all the participants will be

2. THE FIELDS INVESTIGATED

 $1/3$ pions $[2, 3]$. charged particles estimated to be a mixture of about 2/3 protons and North Experimental Hall. The beam consisted of 205 GeV/c positively of the CERN Super Proton Synchroton (SPS) in the H6 beam line of the The mixed high-energy radiation fields were produced using the beam shown in Fig. 1. dead—time losses. The time structure of the CERN SPS "spills" is pile—up and to properly correct the measured count rates for structure which must be taken into account in order to avoid pulse As is usual for high—energy accelerators, the beam has a time

Measurements were performed during a gate period $T_m = 1.7$ s. length was T_{0} = 2.3 s and the beam pulse period was T_{n} = 14.4 s. Fig.1 The time structure of the CERN — SPS beam. The beam pulse

measurement positions chosen for the PTB Bonner spheres. reference [2] and slightly modified in order to indicate the and the measurement geometry are shown in Fig. 2 as taken from specially prepared for this series of measurements. The shieldings was located alternatively in one of two shielding configurations The beam collided with a thick copper target (7 cm \emptyset x 50 cm), which

total thickness of 40 cm were used. of 160 cm concrete, while for the top shielding iron plates with a In the other configuration (see cut BB) the side shielding consisted cm concrete both at the side and on the top (see cut AA in Fig. 2). ln one configuration the shielding around the target consisted of 80

shielding was defined. The four measurement positions of the PTB same size and placed at the height of the beam close to the side similar way a grid of 8 cubic cells (numbered from l to 8) of the are the points of measurement for the various instruments. In a ding. The centres of these cells, placed 25 cm above the shielding, 16) each with a side length of 50 cm was defined on the top shiel respect to this point a grid of 16 cubic cells (numbered from 1 to guration the centre of the target was taken as reference and with In order to establish fixed points for measurements, in each confi

Fig. 2 The two shielding configurations and the measurement positions of the PTB Bonner spheres during the July 1993 series of measurements at CERN - SPS. The "iron side" position (IS#4) is also called "thick concrete" (TC#4).

Bonner spheres which coincide with those of other participants are indicated in Fig. 2 by squares with a point in their centre and listed in Table 1.

Position	Cell	Short Names		Distance to Downstream Shielding from Target
CONCRETE TOP $ Nr.6 CT#6, CT $ CONCRETE SIDE $Nr.2 CS#2, CS$ IRON TOP IRON SIDE		$Nr.6$ IT#6, IT Nr.4 IS # 4, IS	25 cm 25 cm 25 cm 25 cm	125 cm 125 cm 125 cm 25 cm

Table 1. The measurement positions of the PTB Bonner spheres.

The neutron fields encountered in practice are usually accompanied by other radiations. It is typical of high energy accelerators that even in places protected by heavy shielding, not only neutrons and concrete shieldings, the lowest ones behind iron. 3% to 8% in dose equivalent. The highest values are obtained behind points of measurement is estimated to be 1% to 2.5% in fluence and neutron component, the mixed proton—pion contribution at various code, the results being reported in Ref. 3. In relation to the calculated by Roesler and Stevenson using the FLUKA Monte Carlo (protons and pions) components of the radiation fields have been measurement geometry shown in Fig. 2, the neutron and charged hadron photons but also secondary charged particles can be found. For the

3. THE PTB "C" BONNER SPHERE SPECTROMETER

are also used as parts of the spectrometer. bare counter and occasionally the bare counter under a cadmium cover SP90, produced by Centronic Ltd, UK, filled with 200 kPa of ³He. The at their centre a spherical proportional counter (Ø 3.2 cm) of type (3", Z3.5", 4", 4.5", 5", 6", 7", 8", lO", l2", 15" and l8") having polyethylene spheres with diameters ranging from 7.62 cm to 45.72 cm The spectrometer consists of twelve high-density (0.946 g/cm^3)

equivalent. differential or integral, expressed in terms of fluence or dose correlated part must be quadratically added to all final results, measuring method applied for the experimental calibrations. This to the data basis used for the calculations and the fluence diameters in the entire energy range, reflecting the limitations due various spheres, and a correlated part common to all sphere internal shape or height inconsistencies between the responses of the table. They consist of an uncorrelated part which accounts for deviation, as everywhere throughout this report) are also given in uncertainties of the response values (in terms of one standard given numerically in Table 3 of Appendix A. The estimated energies from l meV to 25.12 MeV (C—D_90.RES) is shown in Fig. 3 and extrapolation purposes. The response matrix resulting for neutron responses calculated with the ANISN code [6] for interpolation and [4] and with thermal neutrons [5], using the energy-dependent calibrations with monoenergetic neutrons from l.l7 keV to 14.8 MeV The fluence responses of the spheres are based on experimental

 $\overline{4}$

energy spectrum produced in the present experiment, the response Although inadequate for the investigation of the whole neutron

thermal energies to 25.12 MeV (C—D_90.RES). Fig. 3. The fluence responses of the PTB "C" Bonner spheres from

spectra. below about 10 MeV without using any a priori information on the with the purpose of obtaining the partial neutron spectral fluence matrix from Fig. 3 was used in the first series of data unfolding

functions. Finally the responses obtained were empirically account the regularities of shape of the other IAR response the l8" sphere, also not found in Ref. 7, was guessed taking into sphere diameters have been performed. The high-energy response of our BS set but not considered in Ref. 7, spline interpolations on tions. For sphere diameters of 3.5", 4", 4.5" and 7", contained in up to 387.5 MeV, proved to be compatible with the 55.15 MeV calibra energy shapes of the IAR Lausanne response matrix [7], which extends the most appropriate extension to higher energies. Only the high neutrons [10] were performed at PSI Villigen / Switzerland to select from 1MeV to 10 MeV, and calibration measurements with 55.15 MeV literature [7-9] were fitted to our responses in the energy range the high—energy part of different calculated responses found in the matrix up to a few hundred MeV became necessary. For this purpose, range covered in this experiment, the extension of the PTB response In order to make possible the investigation in the whole energy

is shown in Fig. 4 and given numerically in Table 4 of Appendix A. extrapolated up to 1 GeV. The resulting extended "C" response matrix

thermal energies to 1 GeV $(C-1)$ 94.A61). Fig. 4. The fluence responses of the PTB "C" Bonner spheres from

quadratically added to the final results. a supplementary correlated uncertainty of $\pm 10\%$ should be fluences up to about 20 MeV. For higher-energy parts of a spectrum, same as for the C-D_90.RES response and are still valid for spectral included in this result. The uncertainties given in Table 3 are the proton-recoil telescope $(E > 48 \text{ MeV})$. The guessed 18" response is in absolute fluence with reference measurements using a solid-state $(E > 5$ MeV, NE 213 scintillation detector) and agrees within \pm 8% shows full agreement in shape with time-of-flight measurements field of 55.15 MeV peak energy determined using this response matrix The spectral fluence of the PSI Villigen quasi-monoenergetic neutron

4. THE MEASUREMENTS

were normalized to one PIC count. incident on the target. All the results presented in this report this instrument being estimated to correspond to $2 \cdot 10^4$ particles Chamber (PIC), used as transmission chamber in beam, one count of The beam monitoring was realized by means of a Precision Ionisation

the same during the BS measurements at a certain point. count rates were not corrected for dead—time, but they were almost position, where the dead-time correction amounted to 2.7%. The PIC was attained in this measurement series by the 5" sphere in the IT#6 to the PIC counter. A maximum count rate of the order of $7{\cdot}10^3$ ${\rm s}^{-1}$ Fig. 1) for the Bonner sphere counting and simultaneously applied time window of 1.7 s has been set in the beam—on periods (see constant during the measurement. In order to fulfil this condition a per event. A dead-time correction is simple if the count rate is scaler, the entire system having a non-extending dead-time of 4 μ s preamplifier, an amplifier and a discriminator connected to a central detector were counted using electronics consisting of a In the Bonner sphere measurements, the pulses produced by the

induced by neutrons from those due to noise and photon events. In least as far as it is possible to clearly distinguish the pulses spectra from all other measurements reported here are similar, at sphere in the CT#6 position is shown as an example in Fig. 5. The pulse-height spectrum taken during the measurement with the 8" from the central counter were measured simultaneously. The events were not included in the reading, the pulse—height spectra Nevertheless, in order to ensure that noise or gamma—ray induced have the advantage of being sensitive practically only to neutrons. The Bonner spheres equipped with a 3 He-filled proportional counter

position. detector during the measurement with the $8''$ sphere in the $CT#6$ Fig. 5. The pulse—height spectrum obtained from the central

readings. about 1800 pulses per second) and they were included in the chiefly due to the pile-up of neutron events (the count rate was events and noise. Pulses with amplitudes higher than channel 750 are attributed to neutrons, those of lower amplitude to photon-induced Fig. 5 the pulses with amplitudes higher than channel 130 were

proton—pion fluence. (from a few tens of MeV) a certain fraction of the 1% to 2.5% mixed fluence should be interpreted as including in their high-energy part this reason the results expressed in terms of spectral neutron neutrons that have a chance of being thermalized and detected. For neutrons, at least with regard to the production of secondary polyethylene of the spheres are similar to those of the high-energy compared with the neutrons (see section 2 of this report), in the high-energy protons and pions, which amount to 1% to 2.5% in fluence the high-energy neutrons. The nuclear interactions of the evaluated and was therefore not separated from the contribution of the investigated mixed—fields on the BS count rates was not The influence of the secondary charged hadron components present in

this table. Leake—type REM counter normalized in the same way are also given in attributed to the fluence responses. The readings obtained with a statistical uncertainties and the uncorrelated uncertainties PIC count are listed in Table 5 of Appendix A, together with their The readings corrected for dead-time losses and normalized to one

the iron shield (curve c) these contributions appear to be lower. are expected in both the high-energy and thermal regions. But above rather large readings, considerable contributions to the spectrum similar. As large-diameter spheres and the bare counter as well show those obtained near the concrete shielding (a and b) are very measurements. The curves presented generally differ in shape, but the good stability of all instruments involved in the four series of positions in Fig. 6 (a to d). The smoothness of the curves indicates sphere diameter. These are shown for all four PTB measurement BS readings (relative to the 5" sphere reading) as a function of the period of measurement can be judged from the representation of the The stability of the monitor and the Bonner spheres during the whole

function of the sphere diameter for the four measurement positions. Fig. 6. The relative readings of the PTB Bonner spheres as a

5. RESULTS

the guess spectrum should be modified. observed. These changes are used as indications of the way in which allowed and the changes in shape due to the first few iterations are unfolding a small number of iterations (usually no more than 10) are the shape of the guess spectrum given in the input. During the this kind of unfolding, the solution is not unique and depends on program based on the SAND—II algorithm [11] was used. As usual in For the few-channel unfolding of the measured BS data a home-made

for this spectrum calculating Assuming a solution spectrum, Φ^s , we can simulate our BS readings

$$
C_{\mathbf{d}}^{s} = \sum_{i=1}^{n} R_{\mathbf{d},i} \cdot \Phi_{i}^{s} \qquad (d = 1, 2, \dots, n_{D})
$$
 (1)

solution fluence in the same bin. response of the d-th BS in the i-th energy bin, and Φ_i^s is the spectrum, n_p is the number of BS's used, R_{d, i} is the fluence where $n_{\rm e}$ is the number of energy bins in which we describe the Using these calculated readings and the measured readings

$$
M_d
$$
 (d = 1, 2, ..., n_n)

we can construct the reading ratios

$$
r_d^s = \frac{M_d}{C_d^s} \quad (d = 1, 2, \dots, n_p)
$$
 (2)

and obtain their relative uncertainties from
\n
$$
\sigma_{rel}^{2}(r_{d}^{s}) = \sigma_{rel}^{2}(M_{d}) + \sigma_{rel}^{2}(C_{d}^{s})
$$
\n(3)

Table 5. uncorrelated relative uncertainty of the response, also given in of Appendix A, and $\sigma_{rel}(C_d^S)$ is considered to consist only of the reading (statistical of BS and monitor readings) as given in Table 5 where $\sigma_{rel}^{}(M_{d}^{})$ is the relative statistical uncertainty of the

The weighted mean value of the ratios r_{d}^{s} is obtained from

$$
\langle r_{d}^{s} \rangle = \frac{\sum_{d=1}^{n} r_{d}^{s} \cdot w_{d}}{\sum_{d=1}^{n} w_{d}} = \frac{\sum_{d=1}^{n} r_{d}^{s} / \sigma^{2} (r_{d}^{s})}{\sum_{d=1}^{n} w_{d}}
$$
(4)

 $w_{\perp} = 1/\sigma^2(r^s)$ (5) the weights being the inverses of the variances of r_A^s

The absolute standard deviation of $\langle r \rangle^s$ is obtained from where $\sigma(r_d^s)$ is the absolute standard deviations of r_d^s .

$$
\sigma^{2}(\langle r_{d}^{s}\rangle) = \frac{1}{\sum_{d=1}^{n}1/\sigma^{2}(r_{d}^{s})}
$$
(6)

described by the value of the reduced chi-squared: The spread of the individual r_d^s values about their mean $\langle r_d^s \rangle$ is well

$$
\chi_r^2 = \frac{1}{n_D - 1} \sum_{d=1}^{n_D} \frac{\left[\langle r_d^s \rangle - r_d^s \right]^2}{\sigma^2(r_d^s)}
$$
(7)

confidence interval and we have used $n_{\rm n}$ = 13 detectors in this work, statistics. As one standard deviation, $\sigma(r_{\rm d}^{\rm s})$, corresponds to a 68% unity and their spread about unity is compatible with the appropriate and produces r_d^s values so that $\langle r_d^s \rangle$ is very close to A solution spectrum, Φ^s , is considered valid if it is physically

lie close to unity or even lower. than $2 \cdot \sigma(r^2)$. In this case the value of the *reduced chi-squared* will than one standard deviation, and the others will not deviate by more solution at least nine r_d^s values will deviate from unity by less i.e. 12 spheres and the bare counter, this means that for a valid

the bar on the right indicating the uncertainty $\sigma(r_{\text{a}}^s)$. indicates the relative uncertainty of the measured reading, $\sigma(M_{_2})$, represented. The uncertainty bar on the left side of a point measured / calculated readings of the various BS's, r_a^s , are spectrum is shown as a histogram, while in the lower part the ratios Appendices B to D. In the upper part of each figure the solution Various results are presented numerically and graphicaly in

5.1 The results obtained using the response matrix C—D_90.RES

uncertainties of the response matrix at high energies. reality, the quality of the result being unaffected by the larger 12.59 MeV can be done almost independently of the high-energy On the other hand, the unfolding for neutron energies below "high—energy" fluence which is obtained underestimates the reality. but decrease with increasing energy, it is obvious that the 31.62 MeV. As beyond this interval the BS responses are not constant in the last two bins covering the energy interval from 12.59 MeV to the neutrons with energies higher than 12.59 MeV were concentrated an adequate shape was found. In constructing the guess input spectra spectral shapes were produced and used as guess input spectra until use of a priori information. This means that a large variety of The unfolding of the measured data was done in this case without any

ted in Appendix B and also shown in Fig. 7 as dot-dot-dashed curves. the extension "53E") for all four measurement positions are presen The solutions obtained from this kind of unfolding (file names with

C-l__94.A61 and a limited a priori information 5.2 The results obtained using the extended response matrix

to 1 GeV can be produced. Unfortunately, for neutron energies higher With the response matrix C-1_94.A61, solution spectra extending up obligatory. extremely poor. In this case the use of some a priori information is account of that the spectrometric ability of the BS's becomes functions of various spheres are very similar in shape and on than about 50 MeV, although different in magnitude the response

unfolding. this contribution was carefully determined in each case by about 80 MeV was introduced in all guess spectra and the amount of spallation spectrum, a Maxwellian distribution with the maximum at of neutrons produced by spallation reactions. In order to approach a expected that the high—energy part of the spectrum mainly consists shielding configurations in this series of measurements, it can be Taking into account the way the neutrons are produced and the

Fig. 7 as smooth full curves. the extension "6lE") are presented in Appendix C and also shown in The solutions obtained from this kind of unfolding (file names with

a priori information from FLUKA calculations 5.3 The results obtained using the response matrix C—1_94.A6l and

calculations. bars indicate the statistical uncertainties from the Monte Carlo binning, are shown in Fig. 7 as discrete points whose uncertainty renormalized to one PIC count but still in their original energy 1/3 pion-induced calculated spectra. These mixed calculated spectra, constructing the guess spectra we took a mixture of 2/3 proton- and Appendix A and numerical data listed in Appendices Cand D). In logarithmically equidistant, 5 per decade (see matrix C-l_94.A61 in match the description of our data in 61 energy groups multiplied by $2 \cdot 10^4$ incident particles / PIC count) and rebinned to particle (proton or pion), were renormalized to one PIC count (i.e. the calculated spectra, which are normalized to one incident charged as a priori information in the input guess spectra. For this purpose investigated with our BS's became available [3], these could be used As neutron spectra calculated for the same geometries as

completely the low-energy part of the spectra (the lowest energy in From Fig. 7 it can be seen that the calculations fail to describe

Fig. 7. The various BS solution spectra (53E: dot-dot-dashed curve; 61E: smooth full curve; 61X: histogram) compared with calculations (points). See more explanations in text.

solutions. input guess spectra, the rest being taken from our "53E"-type neutron energies higher than about 10 keV were introduced in the calculations is 0.5 eV). For this reason only calculated results for

presented in Appendix D and also shown in Fig. 7 as histograms. the extension "61X") for all four measurement positions are The solutions obtained from this kind of unfolding (file names with

of BS solutions together with the calculations. Fig. 7 allow a qualitative comparison to be made of different types

0.4 eV, 10 keV, 10 MeV and 1 GeV. 1 GeV, or over four adjacent energy groups with the limits at 1 meV, dose equivalent integrated over the whole energy range from 1 meV to in Appendix E. The reported values refer to the neutron fluence and B to D, or systematized for all solution spectra in the tables given calculated and are given together with the spectra in Appendices For quantitative comparisons the integral data of interest were

The dose equivalent values reported in this work are:

function according to ICRP 21 [12]. $H21 = \sum h21(E_i) \cdot \Phi_i$, where h21 is the fluence-to-dose conversion

conversion function (i.e. $H^*(10)/\Phi$) according to ICRP 60 [1,15]. $H60G = \sum h60G(E_i) \cdot \Phi_i$, where h60G is the fluence-to-ambient-dose
proposion function (i.e. $H^*(10)/\Phi_i$) according to ICPP 60 [1.15] conversion function (i.e. $H^*(10)/\Phi$) according to ICRU 39 [13,14]. $H39 = \sum h39(E_i) \cdot \Phi_i$, where h39 is the fluence-to-ambient-dose
nversion function (i.e. $H^*(10)/\Phi$) according to ICRU 39 [13.14].

the TEPC measurements in "SIDE" positions are spatially slightly measurement positions for HANDI are indicated in column four) that of the table, relative to the BS results. lt should be noted (the listed in Table 2, in absolute values and, in the last two columns obtained with the CERN—HANDl TEPC and reported in Ref. 16 are also counter (calibrated with a bare ²⁵²Cf neutron source) and those neutrons. For comparison, the results we obtained with the Leake—REM 35 Z uncertainty of the dose equivalent produced by higher-energy equivalent produced by neutrons with energies below 10 MeV, and a obtained taking into account a 15 % uncertainty of the dose measurement positions. The uncertainties of these results were solutions of type "61X" are given in Table 2 for all four The integral values H21, H39 and H60G resulting from the BS spectral

Table 2. The dose equivalent values obtained in the four measurement positions with the PTB Bonner spheres, with the Leake-type REM counter and with the CERN HANDI-type TEPC*).

The dose equivalent is given in pSv/PIC-count :

1) C93-CT.DOS ---> CERN $(24.07.93)$ - CONCRETE TOP POS#6

2) C93-CS.DOS \longrightarrow CERN (24.07.93) - CONCRETE SIDE POS#2

	$PTB-C-BS$	PTB-7.22-LEAKE	AHS-HANDI-POS#3	L/BS	H/BS
H ₂₁	611 ± 28	317 ± 75	587	0.52	0.96
H39	714 ± 29 %	$321 \pm 75%$		0.45	
H60G	734 ± 27 %	404 ± 75	679	0.55	0.93

3) C93-IT.DOS -- > CERN (25.07.93) - IRON TOP POS#6

4) C93-IS.DOS \longrightarrow CERN (25.07.93) - IRON SIDE POS#4 (THICK CONCRETE)

*'The AHS-HANDI results are taken from Ref. 16.

of the shieldings, the results still bear comparison. downstream compared with the BS's. Taking into account the thickness HANDI was placed SO cm upstream and in the "IS" measurement 50 cm shifted in relation to the Bonner spheres. In the "CS" measurement

6. DISCUSSION

in the following. Carlo simulations of the experiment [3]. They are taken as reference energies, namely fluence spectral distributions obtained from Monte were produced using the best available a priori information at high experiment are considered to be the "61X" type spectra because they (measurement positions) investigated in the H6J93-CERN-CEC The PTB Bonner sphere best solutions for the four fields

the FLUKA Monte Carlo calculations. spectrum was placed higher in energy than that which resulted from 22% to 35%. This is obviously due to the fact that the spallation the contrary, overestimate the fluence in the same energy region by measurement positions by 21% to 33%, but solutions of type "61E", on solutions of type "53E" underestimate the fluence in all four for neutron energies below 10 MeV. Above 1O MeV, as expected, the "61E" solution types agree within a few percent with the reference Fig. 7 and Table 10a (Appendix E) indicate that both the "53E" and

that the FLUKA code encounters some difficulties in such cases. 37 MeV, 31 MeV and 15 MeV in the other positions), and it could be softer than all the others (mean energy 5.3 MeV compared with result is lower by 34%. lt turns out that this spectrum is much discrepancy is seen for the IRON TOP position (IT#6) were the FLUKA lower by 14% in absolute fluence than our reference. A large discrepancy). For the IRON SIDE position (lS#4) the FLUKA result is agreement is good in both CONCRETE positions (less than 8% by the BS measurements and even for the absolute fluence values the IO keV. In this case the shapes of the spectra are fully confirmed reference one should rather consider the neutron energies above Comparing the spectra from Monte Carlo calculations with the

measurements with the Bonner spheres in both CONCRETE geometries, a Table 2 indicates very good agreement of the TEPC dose equivalent

agree with the BS's only in the soft field IT#6. underestimate the dose equivalent in all hard-spectra fields and source (mean energy 2.13 MeV), are not surprising. They clearly Leake REM counter, which was calibrated with a bare californium the softest spectrum was obtained. The results obtained with the underestimation of the dose equivalent in the IT#6 position, where tolerable overestimation in the IS#4 position, but about 30%

7. ACKNOWLEDGEMENTS

neutron spectra reported in Ref. 3. the numerical results from the Monte Carlo calculations of the comparison exercise, and to Dr. G.R. Stevenson who made available the excellent working conditions and efficient organisation of the The authors are indebted to Dr. M. Höfert and his staff at CERN for

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APPENDIX A

THE BONNER SPHERE READINGS IN THE FOUR MEASUREMENT POSITIONS THE FLUENCE RESPONSE MATRICES USED IN THE DATA ANALYSIS AND

Explanations to Tables 3 and 4:

Table 3. The response matrix C-D 90.RES (1 meV \div 25 MeV).

Table 4. The extended response matrix $C-1$ 94.A61 (1 meV ÷ 1 GeV).

- The neutron energies, E_n , are given in eV.

- The fluence responses are given in cm^2 [counts/(neutron/cm²)].

 $cRAT(E_n) = R_{c0CO}(E_n) / R_{0CO}(E_n).$ cadmium—covered bare counter, while cRAT indicates the ratio the 3.5" sphere, the 12" sphere, etc. cOCO indicates the - OCO, 3CO, 3C5, 12C indicate the bare counter, the 3" sphere,

applied for the experimental calibrations. basis used for the calculations and the fluence measuring method entire energy range, reflecting the limitations due to the data $\small{\mathsf{correlated}}$ part $(\sigma_{\small{\mathsf{corr}}}),$ common to all sphere diameters in the height inconsistencies between responses of various spheres, and a uncorrelated part ($\sigma_\textrm{uncorr}$) which accounts for internal shape or estimated in terms of one standard deviation. They consist of an - The uncertainties given at the lower part of the tables are

Explanations to Table 5:

uncertainties in the four points of measurement. Table 5. The readings of the PTB-"C" Bonner spheres and their

(Precision Ionisation Chamber) count. - The dead-time corrected readings are given relative to one PIC

LEAKE rem counter readings have the same meaning. counts and of the PIC counts. The uncertainties of the deviation, contain the statisical uncertainties of the Bonner sphere - The uncertainties of the readings, always in terms of one standard

shape from the field in which the instrument was calibrated. unknown neutron field wich might differ considerably in spectral but to its capability to predict the dose equivalent in an counter does not refer to the fluence response of this instrument, The large uncertainty of 75% given in this column for the Leake rem the "C" responses, i.e. the same as the $\sigma_{\rm uncorr}^{\rm in}$ from Tables 3 and 4. - The last column in Table 5 gives the uncorrelated uncertainties of

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Table 3. The response matrix $C-D_90.RES$ (1 meV ÷ 25 MeV).

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Table 4. The extended response matrix $C-1$ 94.A61 (1 meV ÷ 1 GeV).

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Table 5. The readings of the PTB-"C" Bonner spheres and their uncertainties in the four points of measurement.

APPENDIX B

All the data are given per PIC count ! POSITIONS USING THE FLUENCE RESPONSE MATRIX C—D 90.RES. THE PTB-BS SOLUTIONS (TYPE "53E") FOR ALL FOUR MEASUREMENT

12.59 MeV to 31.62 MeV. concentrated in the last two bins covering the energy interval from was used. The neutrons with energies higher than 12.59 MeV were In constructing the input spectra no external a priori information

Explanations to the upper part of a figure:

the figure. must be multiplied by the factor given in the upper right part of the spectrum in absolute values per PIC count, the plotted values - The solution spectrum is plotted as a histogram. In order to get

- F is the integral fluence in cm^{-2} .

ICRP21, ICRU39 and lCRP60, respectively, in Sv (see Section 5.3). - H21, H39 and H6OG are the integral dose equivalent according to

i-th energy bin. $\mathbf{E_m}$ is the mean energy $\mathbf{E_m} = \Sigma \mathbf{E_i} \cdot \Phi_i / \Sigma \Phi_i$, Φ_i being the fluence in the

F and integral dose equivalent H21. EH39 and EH6OG are similar. - EH21 is the energy of a monoenergetic field with integral fluence

E 6OG are similar. equivalent according to ICRP21 in the i-th energy bin. E 39 and E_m 21 is the mean energy E_m 21= ΣE_i ·H21₁/ Σ H21₁, H21₁ being the dose

Explanations to the lower part of a figure:

r_, obtained according to the Eqs. l and 2. - The points represent the ratios measured/calculated BS readings,

of the fluence response. relative uncertainty of the calculated reading, which is the $\sigma_{{}_{\rm uncorr}}$ the measured reading. The right uncertainty bar includes the the point size, indicates the relative statistical uncertainty of - The left uncertainty bar of a point, in many cases smaller than

- The reduced chi-squared, X_r^2 , given in the upper part of the figure

is obtained according to the Eq. 7.

Explanations to the table:

bin, respectively; $E_i = (EL_i \cdot EU_i)^2$. All energies are given in eV. - EL, and EU, are the lower and upper margins of the i-th energy

The solutions are given as spectral fluence $\Phi_E(E_i)=d\Phi_i/dE_i$ (cm⁻²eV⁻¹)

Fig. 8 The PTB-BS solution "53E" for the position CT#6.

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Fig. 9 The PTB-BS solution "53E" for the position CS#2.

Fig. 10 The PTB-BS solution "53E" for the position IT#6.

Fig. 11 The PTB-BS solution "53E" for the position IS#4.

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Table 6. The PTB-BS solutions "53E" for all four positions.

File C93L4S.53E (ALEVRA - PTB 7.22 * 11.02.94 / 12:00:00)

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APPENDIX C

All the data are given per PIC count ! POSITIONS USING THE FLUENCE RESPONSE MATRIX C-1_94.A61. THE PTB—BS SOLUTIONS (TYPE "6lE") FOR ALL FOUR MEASUREMENT

lation reaction assumed to be the main contributor at high energies. the maximum at about 80 MeV was introduced to account for a spalenergies beyond this limit a Maxwellian spectral distribution with "53E" solutions were used for neutron energies below 12.59 MeV. For In constructing the input spectra the same parameters as for the

Explanations to the upper part of a figure:

the figure. must be multiplied by the factor given in the upper right part of the spectrum in absolute values per PIC count, the plotted values - The solution spectrum is plotted as a histogram. In order to get

- F is the integral fluence in cm^{-2} .

ICRP21, ICRU39 and ICRP60, respectively, in Sv (see Section 5.3). - H21, H39 and H6OG are the integral dose equivalent according to

i—th energy bin. E is the mean energy E = $\Sigma \mathbb{E}\cdot \Phi$ / $\Sigma \Phi$, Φ being the fluence in the $m = \sum_{i} \phi_i / \sum_{i} \phi_i$

F and integral dose equivalent H21. EH39 and EH6OG are similar. - EH21 is the energy of a monoenergetic field with integral fluence

E 6OG are similar. equivalent according to ICRP21 in the i-th energy bin. E_39 and \sum_{m} 21 is the mean energy E_{m} 21= ΣE_{i} ·H21₁/ Σ H21₁, H21₁ being the dose

Explanations to the lower part of a figure:

r_, obtained according to the Eqs. 1 and 2. - The points represent the ratios measured/calculated BS readings,

relative uncertainty of the calculated reading, which is the $\sigma_{_{\rm uncorr}}$ the measured reading. The right uncertainty bar includes the the point size, indicates the relative statistical uncertainty of - The left uncertainty bar of a point, in many cases smaller than

of the fluence response.

is obtained according to the Eq. 7. - The reduced chi-squared, X_r^2 , given in the upper part of the figure

Explanations to the table:

bin, respectively; E=(EL·EU_) 2 . All energies are given in eV. $i = (EL_i \cdot EU_i$ - EL_i and EU_i are the lower and upper margins of the i-th energy

The solutions are given as spectral fluence $\Phi_E(E_i)=d\Phi_i/dE_i$ (cm⁻²eV⁻¹)

Fig. 12 The PTB-BS solution "61E" for the position CT#6.

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Fig. 13 The PTB-BS solution "61E" for the position CS#2.

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Fig. 14 The PTB-BS solution "61E" for the position IT#6.

Fig. 15 The PTB-BS solution "61E" for the position IS#4.

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Table 7. The PTB-BS solutions "61E" for all four positions.

ALEVRA - PTB 7.22 * 11.02.94 / 12:00:00) File C93-4S.61E C				
$\frac{2}{3}$ Unfolding BS-measurements! (code = SANDA7 ; response = C-1_94.A61)				
Data from C93-??.DAT (CERN - JULY '93 : 4 MEASUREMENT POSITIONS) $\frac{1}{2}$				
13 (11) detectors used: 000,300,(305,)400,(405,)500,600,700,800,100,120,150,180. 5				
6	(CT#6)	(CS#2)	(11#6)	(15#4) C93-1S.61E
$\overline{7}$ (Fluence/PIC-count)	C93-CT.61E	C93-CS.61E	C93-IT.61E ∯Ei/cm ² /eV	⊕Ei/cm ² /eV →
$I = Ei(eV)$ $I = Eli$ $I = Eli$ ا 8 nE	Φ Ei/cm ² /eV	≹Ei/cm'/eV		
61 1 1.0000E-03 7.9433E-04 1.2589E-03	1.5850E+00	2.5420E+00	2.4810E+00	9.4820E-01
2 1.5849E-03 1.2589E-03 1.9953E-03	1.7490E+00	2.8930E+00	2.5200E+00	9.5190E-01
3 2.5119E-03 1.9953E-03 3.1623E-03	1.9520E+00	3.3430E+00	2.5340E+00	9.4360E-01
4 3.9811E-03 3.1623E-03 5.0119E-03	2.2070E+00	3.9280E+00	2.5040E+00	9.1780E-01
5 6.3096E-03 5.0119E-03 7.9433E-03	2.5170E+00	4.6720E+00	2.4010E+00	8.6640E-01
6 1.0000E-02 7.9433E-03 1.2589E-02	2.8570E+00	5.5330E+00	2.1900E+00 1.8430E+00	7.7960E-01 6.5030E-01
7 1.5849E-02 1.2589E-02 1.9953E-02 8 2.5119E-02 1.9953E-02 3.1623E-02	3.1320E+00 3.1450E+00	6.3120E+00 6.5600E+00	1.3610E+00	4.8140E-01
9 3.9811E-02 3.1623E-02 5.0119E-02	2.7220E+00	5.8070E+00	9.1920E-01	3.2470E-01
10 6.3096E-02 5.0119E-02 7.9433E-02	1.8210E+00	3.8980E+00	6.2390E-01	2.0690E-01
11 1.0000E-01 7.9433E-02 1.2589E-01	8.1930E-01	1.7100E+00	4.2570E-01	1.2100E-01
12 1.5849E-01 1.2589E-01 1.9953E-01	2.6940E-01	5.1370E-01	2.9190E-01	6.9320E-02
13 2.5119E-01 1.9953E-01 3.1623E-01	1.2090E-01	2.1060E-01	2.0120E-01	4.1940E-02 2.5990E-02
14 3.9811E-01 3.1623E-01 5.0119E-01	7.5210E-02	1.2970E-01 8.2280E-02	1.3940E-01 9.7120E-02	1.6050E-02
15 6.3096E-01 5.0119E-01 7.9433E-01	4.7800E-02 3.0130E-02	5.1740E-02	$6.8080E - 02$	9.8680E-03
16 1.0000E+00 7.9433E-01 1.2589E+00 17 1.5849E+00 1.2589E+00 1.9953E+00	1.8860E-02	3.2310E-02	4.8040E-02	$6.0410E-03$
18 2.5119E+00 1.9953E+00 3.1623E+00	1.1770E-02	2.0110E-02	3.4060E-02	3.6940E-03
19 3.9811E+00 3.1623E+00 5.0119E+00	7.3460E-03	1.2510E-02	2.4250E-02	2.2590E-03
20 6.3096E+00 5.0119E+00 7.9433E+00	4.5780E-03	7.7770E-03	1.7330E-02	1.3810E-03
21 1.0000E+01 7.9433E+00 1.2589E+01	2.8530E-03	4.8340E-03	1.2410E-02	8.4480E-04 5.1690E-04
22 1.5849E+01 1.2589E+01 1.9953E+01	1.7780E-03	$3.0050E - 03$ 1.8690E-03	8.9140E-03 6.4120E-03	3.1630E-04
23 2.5119E+01 1.9953E+01 3.1623E+01 24 3.9811E+01 3.1623E+01 5.0119E+01	1.1090E-03 6.9190E-04	1.1640E-03	4.6180E-03	1.9370E-04
25 6.3096E+01 5.0119E+01 7.9433E+01	4.3200E-04	7.2460E-04	3.3300E-03	1.1870E-04
26 1.0000E+02 7.9433E+01 1.2589E+02	2.6970E-04	$4.5140E - 04$	2.4040E-03	7.2750E-05
27 1.5849E+02 1.2589E+02 1.9953E+02	1.6850E-04	2.8130E-04	1.7370E-03	4.4600E-05
28 2.5119E+02 1.9953E+02 3.1623E+02	1.0530E-04	1.7530E-04	1.2570E-03	2.7350E-05
29 3.9811E+02 3.1623E+02 5.0119E+02	6.5800E-05	1.0930E-04	9.0990E-04	1.6780E-05 1.0300E-05
30 6.3096E+02 5.0119E+02 7.9433E+02	4.1160E-05 2.5770E-05	6.8230E-05 4.2610E-05	6.5900E-04 4.7750E-04	6.3280E-06
31 1.0000E+03 7.9433E+02 1.2589E+03 32 1.5849E+03 1.2589E+03 1.9953E+03	1.6140E-05	2.6630E-05	3.4610E-04	3.8900E-06
33 2.5119E+03 1.9953E+03 3.1623E+03	1.0120E-05	1.6660E-05	2.5100E-04	2.3930E-06
34 3.9811E+03 3.1623E+03 5.0119E+03	6.3580E-06	1.0440E-05	1.8220E-04	1.4740E-06
35 6.3096E+03 5.0119E+03 7.9433E+03	$4.0030E - 06$	6.5590E-06	1.3260E-04	$9.1030E - 07$
36 1.0000E+04 7.9433E+03 1.2589E+04	2.5300E-06	4.1370E-06	$9.6720E - 05$	5.6430E-07
37 1.5849E+04 1.2589E+04 1.9953E+04	1.6110E-06	2.6280E-06	7.0960E-05 $5.2510E-05$	3.5250E-07 2.2340E-07
38 2.5119E+04 1.9953E+04 3.1623E+04 39 3.9811E+04 3.1623E+04 5.0119E+04	1.0410E-06 6.8900E-07	1.6930E-06 1.1180E-06	3.9370E-05	1.4520E-07
40 6.3096E+04 5.0119E+04 7.9433E+04	4.7590E-07	7.6970E-07	3.0060E-05	9.8530E-08
41 1.0000E+05 7.9433E+04 1.2589E+05	$3.5030E - 07$	5.6460E-07	2.3420E-05	7.1070E-08
42 1.5849E+05 1.2589E+05 1.9953E+05	2.7930E-07	$4.4840E - 07$	1.8530E-05	5.4960E-08
43 2.5119E+05 1.9953E+05 3.1623E+05	2.4060E-07	3.8450E-07	1.4550E-05	4.4800E-08
44 3.9811E+05 3.1623E+05 5.0119E+05	2.1680E-07	3.4420E-07	1.0870E-05	3.6450E-08
45 6.3096E+05 5.0119E+05 7.9433E+05	1.9260E-07	3.0270E-07	5.9730E-06 2.3590E-06	2.6840E-08 1.6130E-08
46 1.0000E+06 7.9433E+05 1.2589E+06	1.5550E-07 1.0250E-07	2.4060E-07 1.5440E-07	5.6710E-07	7.1300E-09
47 1.5849E+06 1.2589E+06 1.9953E+06 48 2.5119E+06 1.9953E+06 3.1623E+06	4.9740E-08	7.2260E-08	6.9270E-08	2.4460E-09
49 3.9811E+06 3.1623E+06 5.0119E+06	1.7250E-08	2.4190E-08	1.0150E-08	1.0570E-09
50 6.3096E+06 5.0119E+06 7.9433E+06	6.9050E-09	9.4440E-09	7.3300E-09	6.8140E-10
51 1.0000E+07 7.9433E+06 1.2589E+07	5.2950E-09	7.0770E-09	6.9870E-09	5.0090E-10
52 1.5849E+07 1.2589E+07 1.9953E+07	4.9220E-09	6.6320E-09	7.1200E-09	3.9270E-10
53 2.5119E+07 1.9953E+07 3.1623E+07	4.8120E-09	6.3940E-09	7.5600E-09	3.3960E-10 $3.0160E-10$
54 3.9811E+07 3.1623E+07 5.0119E+07	4.7770E-09 4.4050E-09	6.2560E-09 5.7100E-09	7.5270E-09 6.8560E-09	2.5840E-10
55 6.3096E+07 5.0119E+07 7.9433E+07 56 1.0000E+08 7.9433E+07 1.2589E+08	3.5040E-09	4.4800E-09	5.3420E-09	1.9700E-10
57 1.5849E+08 1.2589E+08 1.9953E+08	2.0770E-09	2.6420E-09	3.1990E-09	1.1240E-10
58 2.5119E+08 1.9953E+08 3.1623E+08	8.0900E-10	1.0250E-09	1.2540E-09	4.2670E-11
59 3.9811E+08 3.1623E+08 5.0119E+08	1.6090E-10	2.0410E-10	2.5000E-10	8.3400E-12
60 6.3096E+08 5.0119E+08 7.9433E+08	1.0940E-11	1.3900E-11	1.7090E-11	5.6220E-13
61 1.0000E+09 7.9433E+08 1.2589E+09	1.3640E-13	1.7320E-13	2.1260E-13	$6.9510E - 15$
			1.6974E+01	$2.3433E-01$
ětotal (n/cm²PIC-count) Ξ	1.8059E+00	2.7586E+00		

 $\hat{\mathcal{E}}$

APPENDIX D

POSITIONS USING THE FLUENCE RESPONSE MATRIX C—1 94.A6l. THE PTB—BS SOLUTIONS (TYPE "61X") FOR ALL FOUR MEASUREMENT

All the data are given per PIC count !

to match our responses and normalized to one PIC count. presented in Ref. 3 were used. The calculated spectra were rebinned beyond this limit the numerical results of the calculations Appendix B were used for neutron energies below IO keV. For energies In constructing the input spectra the "53E" solutions from

Explanations to the upper part of a figure:

the figure. must be multiplied by the factor given in the upper right part of the spectrum in absolute values per PIC count, the plotted values - The solution spectrum is plotted as a histogram. In order to get

- **F** is the integral fluence in cm^{-2} .

ICRP21, ICRU39 and ICRP60, respectively, in Sv (see Section 5.3). - H21, H39 and H6OG are the integral dose equivalent according to

i—th energy bin. $\mathbf{E_m}$ is the mean energy $\mathbf{E_m} = \mathbf{\Sigma E_i \cdot \Phi_i / \Sigma \Phi_i}$, Φ_i being the fluence in the

F and integral dose equivalent H21. EH39 and EH6OG are similar. - EH21 is the energy of a monoenergetic field with integral fluence

E 60G are similar. equivalent according to ICRP21 in the i-th energy bin. E 39 and $E_{\rm m}$ 21 is the mean energy $E_{\rm m}$ 21= $\Sigma E_{\rm i}$ •H2l₁/ Σ H2l₁, H2l₁ being the dose

Explanations to the lower part of a figure:

r_, obtained according to the Eqs. I and 2. - The points represent the ratios measured/calculated BS readings,

of the fluence response. relative uncertainty of the calculated reading, which is the $\sigma_{{\rm uncorr}}$ the measured reading. The right uncertainty bar includes the the point size, indicates the relative statistical uncertainty of - The left uncertainty bar of a point, in many cases smaller than

is obtained according to the Eq. 7. - The reduced chi-squared, X^2 , given in the upper part of the figure

Explanations to the table:

bin, respectively; E=(EL·EU_)². All energies are given in eV. $i=[EL_i \cdot EU_i]$ - EL, and EU, are the lower and upper margins of the i-th energy

The solutions are given as spectral fluence $\Phi_E(E_i) = d\Phi_i/dE_i$ (cm⁻²eV⁻¹)

Fig. 16 The PTB-BS solution "61X" for the position CT#6.

Fig. 17 The PTB-BS solution "6XE" for the position CS#2.

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Fig. 18 The PTB-BS solution "61X" for the position IT#6.

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Fig. 19 The PTB-BS solution "61X" for the position IS#4.

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APPENDIX E

All the data are given per PIC count ! FOR ALL FOUR MEASUREMENT POSITIONS AND ALL SOLUTION TYPES THE GROUP VALUES OF THE FLUENCE AND THE DOSE EQUIVQLENT

investigations from 1 meV to 1 GeV was divided into 4 energy groups: The neutron energy range covered by the present spectrometric

> Group 4: from IO MeV to l GeV Group 3: from 10 keV to 10 MeV Group 2: from 0.4 eV to IO KeV Group I: from 1 meV to 0.4 eV

Type "6lX" (spectra *.6lX) as presented in Appendix D, - Type "61E" (spectra *.61E) as presented in Appendix C, Type "S3E" (spectra *.53E) as presented in Appendix B, the rest of our data and normalized to one PIC count, - Type "61N" (spectra *.61N) calculated in Ref. 3, rebinned to match The spectral solutions of different types:

were considered for each of the four measurement positions:

- CT#6 = CONCRETE TOP POSITION #6
- CS#2 = CONCRETE SIDE POSITION #2
- $-$ IT#6 = IRON TOP POSITION #6
- IS#4 = IRON SIDE POSITION #4 (THICK CONCRETE)

(a and b). equivalent H21, H39 and H6OG were calculated and listed in Table 9 and the corresponding group values for neutron fluence and dose

relative to the "61X" values and listed in Table 10 (a and b). group values for the other types were recalculated in percent As the solutions of type "6lX" are considered to be the best, the

Table 9a. The group values, absolute and relative to total, of the neutron fluence and dose equivalent H21, for all types of solutions in all four measurement positions.

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Table 9b. The group values, absolute and relative to total, of the dose equivalent H39 and H60G, for all types of solutions in all four measurement positions.

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Percentual group fluence relative to *.61X

Table 10a. The group values of the neutron fluence and dose equivalent H2l, for all types of solutions relative to the solution of type "61X", in all four measurement positions.

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Table 10b. The group values of the dose equivalent H39 and H60G, for all types of solutions relative to the solution of type "61X", in all four measurement positions.

 $\sim 10^{11}$ km s $^{-1}$

PTB-Berichte der Serie N (Neutronenphysik)

120 S., 6 Abb., 57 Tab., ISBN 3-88314-906-3, 1989. DM 32,50 Regard to the Dose from Induced Photons. Sections of Oxygen between 6 and 15 MeV. Equivalent, H_e, from Neutron Irradiation with Special Elastic and Inelastic Differential Neutron Scattering Cross Comparison of the Effective Dose, E, and the Effective Dose N-1: G. Börker, R. Böttger, H. J. Brede, H. Klein u. a.: N-12: R. A. Hollnagel:

B. R. L. Sicbert: N-13: H. Schuhmacher, U. J. Schrewe: N—2: B. W. Bauer, W. G. Alberts, M. Luszig-Bhardra.

at the PTB. Monte Carlo Simulation of Fast Neutron Scattering Experi Analysis of Results for the Albedo Neutron Dosemeter in Use N-14: D. Schmidt, B. R. L. Siebert: Response of Individual Neutron Dosemeters: Detailed Energy, Angle of Incidence and Phantom Shape on the 42 S., 14 Abb., 5 Tab., ISBN 3-89429-306-3, 1993, DM 20,00

54 S.. 21 Abb., ISBN 3-89429-005—6. 1990. DM 20,50 ments Including DD-Breakup Neutrons.

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