# First observation of the $t\bar{t}H$ process at CMS

# KIRILL SKOVPEN (ON BEHALF OF THE CMS COLLABORATION)

Inter-University Institute for High Energies, Vrije Universiteit Brussel Pleinlaan 2, 1050 Brussels, Belgium

The top quark and the Higgs boson play a special role in the fundamental interactions of the standard model. The observation of the top quark pair production in association with the Higgs boson establishes the first direct measurement of the tree-level coupling of the Higgs boson to the top quark. The analysis of the data collected by the CMS detector results in the overall observed significance of 5.2 standard deviations for this process.

#### PRESENTED AT

11<sup>th</sup> International Workshop on Top Quark Physics Bad Neuenahr, Germany, September 16–21, 2018

#### 1 Introduction

The top quark and the Higgs boson are the two heaviest elementary particles of the standard model (SM). The top quark, the heaviest of all elementary particles, decays before any hadronization takes place and provides an access to a direct study of its intrinsic properties in experiment. The second heaviest elementary particle that was recently discovered, the Higgs boson, gives mass to all SM particles through the omnipresent Higgs field [1, 2, 3]. It is of particular interest to probe the strength of the interaction between these two heavyweight elementary objects: due to the large mass of the top quark and the fact that the strength of Yukawa interaction is proportional to the mass of the quark that interacts with the Higgs field, the top quark is expected to have a large Yukawa coupling,  $y_t \simeq 1$ .

The precise determination of  $y_t$  represents a direct test of the SM predictions, as well as a probe of new physics. Moreover, the properties of this fundamental interaction are potentially connected to the vacuum stability of the Higgs potential. High precision measurements of the masses of the top quark and the Higgs boson, as well as the strength of the top-Higgs interaction, can shed some light on the stability state of our universe at the Planck scale [4, 5].

The direct measurement of  $y_t$  is possible from the study of the process of associated production of the Higgs boson with a pair of top quarks ( $t\bar{t}H$ ), as well as in the associated production of the single top quark with the Higgs boson (tHq). However, the tHq process is strongly suppressed due to destructive interference present in the production diagrams.

#### 2 Search channels

The  $t\bar{t}H$  production is a rare process with the predicted cross section of  $\simeq 0.5$  pb at 13 TeV in the center-of-mass of proton-proton collisions at the LHC. There are several decays of the Higgs boson that define three main search channels considered in analysis of recorded data by the CMS detector [6]. The channel with  $H \to bb$  benefits from the large branching fraction of this decay process but suffers from high background rates. The  $H \to \gamma\gamma$  channels provides a relatively clean experimental signature but is also attributed to rather small branching ratio of this decay channel. Finally, the  $H \to WW^*, ZZ^*, \tau\tau$ , or multilepton channel, represents a compromise between the modest background rates and reasonably high rates of signal events. In the combined analysis of the 8 TeV data collected by ATLAS and CMS detectors, the observed (expected) significance of 4.4 (2.0) standard deviations ( $\sigma$ ) was obtained for the  $t\bar{t}H$  process [7].

The search for  $t\bar{t}H$  is performed for  $H \to bb$  decays with at least one lepton present in the final state [8]. With the requirement of at least three b-tagged jets

to be identified in association with one or two leptons, the dominant background becomes top quark pair production with additional jets ( $t\bar{t}$ +jets). In the first step of the event categorisation procedure the events are classified based on the number of b-tagged jets and leptons to define a multivariate analysis (MVA), deep neural network (DNN) and matrix-element method (MEM) discriminants, with the final choice of the method driven by the obtained highest expected sensitivity for the signal process. In the single lepton case the multi-classification procedure is used to define several physics process categories corresponding to the signal and  $t\bar{t}$ +jets background processes, split by the flavour of additional jets. The DNN discriminant is used as a final variable to perform a maximum likelihood fit. In the dilepton selection, the BDT and MEM discriminants are used. The simultaneous combined likelihood fit is performed over all considered event categories. The final observed (expected) significance is 1.6 (2.2)  $\sigma$ .

The search for  $t\bar{t}H$  in  $H\to bb$  decays is also done in the fully hadronic channel with considering only hadronic decays of the top quarks [9]. Even though the combinatorial background represents an important challenge in this analysis, the full event reconstruction is possible. Events with reconstructed leptons are vetoed, and the dominant background arises from the QCD multijet production. This background is estimated from the dedicated control regions in data corresponding to the low be tag multiplicity. Six exclusive categories are defined based on the number of be and non-b-tagged jets. The final signal contribution is extracted from a maximum likelihood fit to the MEM discriminant which is based on the probabilities of an event to correspond to either the  $t\bar{t}H$  or  $t\bar{t}+b\bar{b}$  hypothesis for the production process, the best fit signal strength is compatible with the background-only hypothesis and therefore only the observed (expected) upper limit at 95% confidence level of 3.8 (3.1) can set on the signal strength.

One of the search channels where one can aim at observing a reconstructed mass peak corresponding to the Higgs boson production is the channel with Higgs decaying to a pair of photons [10]. The background processes consist of irreducible prompt diphoton production and reducible backgrounds associated with  $\gamma$ +jet and di-jet events, in which jets are misidentified as isolated photons. Identification of primary vertex with two photons is performed with boosted decision trees (BDT) and has a strong impact on the reconstructed di-photon invariant mass resolution. The BDT uses several observables including the information on the tracks recoiling against the di-photon system. The background model is defined from fitting data to extract the signal from a maximum likelihood fit which is performed on the di-photon mass distribution. Evidence for the  $t\bar{t}H$  process is obtained with an observed (expected) significance of 3.3 (1.5)  $\sigma$ .

The analysis of the multilepton final states is sensitive to  $H \to WW^*, ZZ^*, \tau\tau$  decays [11]. Six exclusive categories are defined based on the number of reconstructed electrons, muons and hadronic tau decays, starting from the single lepton selection

and going up to the four-lepton requirement. The inclusion of hadronic decays of taus in the analysis allows to significantly enhance the signal sensitivity to  $H \to \tau\tau$ . The dominant background is associated with the  $t\bar{t}+Z$  and  $t\bar{t}+W$  processes, as well as misidentified leptons. The lepton misidentification and charge flip probabilities are estimated with data-driven approaches. A dedicated MVA is used for the optimized lepton selection and is based on several variables associated with the reconstructed lepton and the properties of the closest reconstructed jet. The MEM is additionally used in the event categories with three leptons, and all channels exploit BDT as the final discriminating variable, except for the four lepton selection where the cut-and-count analysis is performed. The extracted signal contribution from the maximum likelihood fit corresponds to the observed (expected) significance of 3.2 (2.8)  $\sigma$ . This result also leads to the observed evidence for the  $t\bar{t}H$  process.

#### 3 Combination

The results obtained in each of the search channels are combined in a maximum likelihood fit. The fitted signal strength values for each channel, as well as for the combination, are represented in Fig. 1. The observed and predicted event yields in all the bins of the final discriminants are ordered by the pre-fit expected signal-to-background ratio. The result distribution is shown in Fig. 2. The overall observed (expected) significance is 5.2 (4.2)  $\sigma$  with the measured production rate consistent with the SM prediction within one standard deviation. This result represents the first experimental observation of the  $t\bar{t}H$  process.

# 4 Search for the tHq process

The single top associated production with the Higgs boson is an important process that is sensitive not only to the absolute value of  $y_t$ , but also to the sign of this parameter. However, this is a very rare process with the predicted cross section that is almost one order of magnitude smaller than that of  $t\bar{t}H$ . The search for tHq process was done with 8 and 13 TeV data in the similar final states as considered in the  $t\bar{t}H$  analysis  $(H \to bb, \gamma\gamma, WW^*, ZZ^*, \tau\tau)$ , and led to stringent contraints set on the anomalous values of  $y_t$  [13, 14]. A positive value of the modifier of the top-Higgs coupling is favoured by about 1.5 standard deviations with excluding values outside the ranges of [-0.9, -0.5] and [1.0, 2.1] at 95% C.L. The standard model-like Higgs couplings to vector bosons are assumed.

### 5 Conclusion

The observation of the ttH process represents an important milestone in our study of the top-Higgs interactions, and it is the first step towards our better understanding of the properties of this fundamental interaction. The analysis of the full data set collected at 13 TeV, as well as the future experiments at the LHC, will allow us to significantly improve the precision of the current measurements of the ttH production cross section and to study differential distributions of the particles involved in this process with a potential probe of new physics.

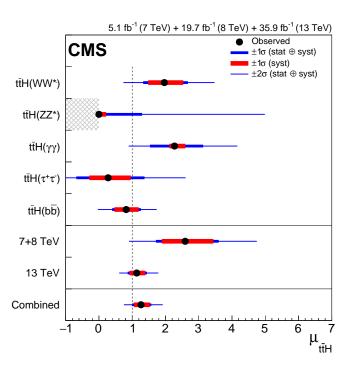


Figure 1: Best fit value for the  $t\bar{t}H$  signal strength modifier [12]. The SM prediction is shown as a dashed vertical line.

### References

- [1] G. Aad, et al. (ATLAS Collaboration), Phys. Lett. B 716 (2012) 1.
- [2] V. Khachatryan, et al. (CMS Collaboration), Phys. Lett. B 716 (2012) 30.
- [3] V. Khachatryan, et al. (CMS Collaboration), JHEP 06 (2013) 081.
- $[4]\ F.$  Bezrukov, M. Yu. Kalmykov, B. A. Kniehl, et al., JHEP  ${\bf 10}$  (2012) 140.

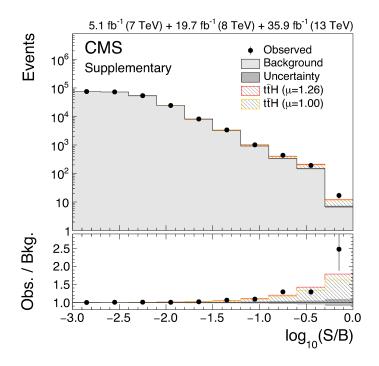


Figure 2: Number of events as a function of  $\log S/B$ , where S and B are the signal and background yields, respectively [12].

- [5] F. Bezrukov and M. Shaposhnikov, J. Exp. Theor. Phys. 120 (2015) 3.
- [6] V. Khachatryan, et al. (CMS Collaboration), JINST3 (2008) S08004.
- [7] G. Aad, et al. (ATLAS and CMS Collaborations), JHEP 08 (2016) 045.
- [8] A. M. Sirunyan, et al. (CMS Collaboration), arXiv:1804.03682 (Submitted to JHEP).
- [9] A. M. Sirunyan, et al. (CMS Collaboration), JHEP 06 (2018) 101.
- [10] A. M. Sirunyan, et al. (CMS Collaboration), arXiv:1804.02716 (Submitted to JHEP).
- [11] A. M. Sirunyan, et al. (CMS Collaboration), JHEP 08 (2018) 066.
- [12] A. M. Sirunyan et al. (CMS Collaboration), Phys. Rev. Lett. 120 (2018) 231801.
- [13] V. Khachatryan et al. (CMS Collaboration), JHEP 06 (2016) 177.
- [14] A. M. Sirunyan et al. (CMS Collaboration), CMS-PAS-HIG-18-009, http://cds.cern.ch/record/2628662.