



LAPP-EXP-94.06 March 1994

Sw 9419

A FIRST ELECTROMAGNETIC CALORIMETER PROTOTYPE OF PbWO₄ CRYSTALS

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ABSTRACT

The main properties of PbWO₄ crystals have been measured. They are of particular interest for electromagnetic calorimetry. Beam tests of 21-cells matrix prototype are reported. An energy resolution better than 1% has been achieved for electrons energies larger than 25 GeV.

(Submitted to Nucl. Instr. and Methods)

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1. INTRODUCTION

The lead tungstate scintillating crystal PbWO₄ [1-6] (hereafter designated as PWO) has recently been shown to be amongst the most promising candidates for future electromagnetic (EM) calorimetry, together with CeF₃ [7] and PbF₂, a Cherenkov light radiator [8]. CeF₃ provides a high light yield and PbF₂ has fast timing properties with a high stopping power. But PWO crystals present a whole set of physical and chemical parameters (Table 1) which are attractive for large EM calorimeters in the extreme conditions that are foreseen for example at LHC. Despite a moderate mechanical hardness (Moh = 4), it is easy to machine and polish. PWO is grown from a mixture of PbO and WO₃ oxides by the Czochralski method at a rather moderate temperature. The recent study of the light emission mechanism of PWO samples grown in various conditions allows to control the growing technology in order to optimise desirable characteristics [9-10].

This work presents the first beam tests of electromagnetic calorimeter prototypes made of large size crystals ($\geq 20 \text{ X}_0$) at IHEP Protvino, with 4 cells grown in different conditions, and at CERN, with a matrix of 21 cells grown in optimised conditions as much as possible.

2. SCINTILLATION AND OPTICAL PROPERTIES OF PWO CRYSTAL

The PWO crystals are grown in platinum crucibles in an atmosphere similar to air in composition. A reference sample $(1 \times 1 \times 1 \text{ cm}^3)$ and a long bar have been cut out from a series of ingots (~ 240 mm long, \emptyset ~ 32 mm).

The luminescence characteristics of a reference crystal are presented in Fig. 1. A large energy gap separates the maximum of luminescence emission around 480 nm and the strong absorption edge below 350 nm. The scintillation light is a superposition of several bands between $\lambda = 420$ nm and 650 nm [9]. Crystals have been grown in conditions that favour green luminescence.

The scintillation kinetics (or decay time) of PWO was measured by the start-stop method with a 22 Na source and a Lecroy 3001 QVT analyser. Measurements shown in Fig. 2 have been fitted by a sum of 3 exponential functions and a constant background. Decay times are $\tau_1 = (2 \pm 0.1)$ ns, $\tau_2 = (7.4 \pm 0.3)$ ns and $\tau_3 = (26 \pm 3)$ ns respectively. Their relative intensities

are 41%, 39% and 20% after background subtraction. In Fig. 3a, the observed scintillation yield for PWO and BGO samples of the same dimensions ($20 \times 20 \times 40 \text{ mm}^3$) wrapped in millipore (a diffusing material) are shown for different radioactive excitation sources. The light was detected by a photomultiplier tube (PM) Philips XP2262 coupled to the crystal with the Dow Corning Q₂-3067 optical compound (refractive index n = 1.48). All spectra have been measured with a Lecroy QVT pulse charge analyser. The gate widths were 350 ns for BGO and 35 ns for PWO, respectively. In the case of BGO irradiated by 137 Cs, the signal was attenuated by 24 db. The mean value of the pedestals corresponds to the same channel number, in both Fig. 3a and 3b.

The measured scintillation efficiency of PWO is 5% of that of BGO when the resolution of the 137 Cs photopeak (0.660 MeV) is used to estimate the number of photoelectrons. Fig. 3a also shows the 60 keV photopeak obtained with BGO and 241 Am source. Fig. 3b represents the 137 Cs γ -peak and the 4.4 MeV Am-Be source γ -peak detected by a (2.2 x 2.2 x 2.4 cm³) PWO crystal. Here high voltage settings allow a direct comparison of spectra without attenuation. The ratios of photopeak positions and of energies are in agreement.

3. RADIATION HARDNESS

The γ -radiation hardness of PWO is related to charge exchange processes in Pb³⁺ defect centers [10]. It has been checked with a 500 R/s intensity 60 Co source at room temperature (300°K). Crystals with a large number of Pb³⁺ centers are initially slightly yellowish. Their transparency increases under small irradiation doses and decreases above a 105 rad absorbed dose, but it recovers very fast. Fig. 4a shows such a recovery after 10 days without any special UV or visible light treatment.

Negligible changes (Fig. 4b) of transparency are observed up to 5 Mrad in a more transparent crystal grown with less Pb³⁺ centers. Further production of crystals for the calorimeter prototype has been made along this line.

A sample (2 x 2 x 2.4 cm³) has been irradiated with neutrons from the Saclay reactor ULYSSE. Experimental conditions of irradiation are described elsewhere [11]. Fast neutrons cumulative fluxes range from 10^{12} to 10^{14} n/cm². The γ contamination amounts to 1.2 Mrad

for 10¹⁴ fast n/cm². Accompanying thermal and epithermal neutron fluxes are three times larger than the fast neutron flux. Due to activation, the transmission has been measured after a delay from 1 to 2 weeks according to the initial dose, meanwhile samples were kept in darkness to prevent any annealing effect.

The irradiation effect on light attenuation is shown in Fig. 4c. The corresponding absorption coefficient variation Δk (defined by Tirrad/Tinit = $e^{-\Delta KL}$) is displayed in Fig. 4d. Almost no influence is observed up to 10^{13} n/cm². The variation becomes significant above this fluence but remains very low. After 10^{14} n/cm² the radiation induced attenuation length is about 1 m.

4. BEAM TESTS

4.1 Tests Conditions

a) at IHEP Protvino in an electron beam from 4 to 26 GeV/c (Spring 93) with a momentum spread $\sigma p/p \sim 2\%$ [12]:

Four PWO crystals grown in different conditions have been studied. Their cross-section was $22 \times 22 \text{ mm}^2$ with lengths ranging from 150 to 180 mm (17-20 X_0). They were assembled inside a 4×4 matrix, as shown in Fig. 5, where they occupy positions PWO1, PWO2, PWO5, PWO6. The others cells were made of Cherenkov radiators like Gd₃ Ga₅ O₁₂ alias (GGG) or NaBi(WO₄)₂ alias (NBW) with close refractive index and radiation length (Table 1). The photomultipliers were Philips XP1911 or FEU-147. The crystals were polished and wrapped in aluminised mylar. The scintillation yield of each cell had been previously measured with a Philips XP2020Q and with a multialcali photocathode FEU-84. These PMs have a photocathode larger than the crystal output area. The light collection geometrical factor is $\varepsilon_g = 1$. Optical coupling was done with DowCorning O2-3067.

b) at CERN in an electron beam of 10, 25 and 100 GeV/c available in the H8 beam of the SPS North Area (Autumn 93), with a momentum spread $\sigma p/p \sim 0.4\%$ [9]:

A matrix detector of 21 crystals has been assembled. A Philips XP1911 phototube was attached to each crystal. The matrix cells were separated by 50 µm of iron sheets. They

4.2 Single Cell Tests

At IHEP, the spectra produced by 26 GeV/c electrons in GGG and in PWO associated to the same XP2020Q PM have been compared. The σ A/A resolution of the PWC $\frac{1}{12}$ $\frac{1}{1$

The collected charge as a function of the gate width for various PWO and GGG cells have been measured after 95 m of 50Ω cable. Results for the four PWO cells and one GGG cell are shown in Fig. 7a, taking as a reference the charge amplitude obtained with a 400 ns gate. PWO6 is almost as fast as the Cherenkov radiator GGG. It has been selected as a reference for further crystal production. The pulse width (recorded with a XP2020Q PM) is 15 ns (FWHM) and 25 ns (at 10% of the maximum amplitude).

At CERN, 100 GeV/c electron signals measured with a XP1911 PM are shown in Fig. 7b. They were observed with a Tektronix 2467B oscilloscope after a 4 m long 50 Ω cable (30 ns separate the two vertical bars). The rise time is ~ 5 ns. 85% of the signal lies inside a ~ 35 ns window.

4.3 Coordinate Resolution and Shower Profile Measurements

The matrix structure shown in Fig. 5 has also been used to determine the profile of the shower of 26 GeV electrons (IHEP). The beam was scanned across the PWO6 and PWO5

cells, as indicated by the arrow. Fig. 8a displays the horizontal position of the center of gravity of the calibrated pulse amplitudes in each matrix cell as a function of the impact of the incident electron as determined with a 200 μ m pitch silicon microstrip detector. The origin of the coordinates was just at the interface between PWO5 and PWO6. After applying a correction procedure described in [14] real impact coordinates were obtained. The position resolution σ_x of the center of gravity, that is displayed in Fig. 8b with the origin at the center of the cell, is less than 400 μ m near the edge and it rises to 900 μ m in the center.

The measured 26 GeV electron shower profile is shown in Fig. 9a as the fraction A₆/A_{tot} of the charge amplitude that is collected in PWO6 as a function of beam position. The origin is the center of PWO6. Up to 86% of the electromagnetic shower energy is absorbed in the central 22 x 22 mm² cell. For comparison, Fig. 9b displays the shower profiles for GGG (1), NBW (2), PWO (3) and TF1(SF2) (4) lead glass. The profile is twice narrower in heavy crystals than in TF1.

4.4 Muon Beam Tests and Attenuation Length

The attenuation length of the scintillation light through PWO has been measured with muons crossing cells perpendicularly to their axis. Fig. 10a displays scans of three crystals (20 X_0 , CERN matrix) as a function of the distance of the muon impact to the PM photocathode. The vertical scale is logarithmic and represents the light intensity in arbitrary units before cell calibration. The curves exhibit the same behaviour. Neglecting the small edge effect far from the PM (right side of the curves), the attenuation length is of the order of 2 m. Crossing muons that release about 33 MeV produce an average number of the order of 100 photoelectrons (spectrum of Fig. 10b).

4.5 Electron Spectra and Energy Resolution

10, 25 and 100 GeV data have been recorded at two high voltage settings. Fig. 11 presents the calibrated 25 GeV spectra obtained within the 3 x 3 central core of the detector installed at CERN (Fig. 6), when the beam hits a 4 x 4 mm² area around the axis of the central cell. Two different abscissa scales are used in order to make the energy deposit in external cells

more visible. Both represent the calibrated charge. 80% of the total energy is deposited in the central crystal. Similar results are obtained at other beam energies as it is illustrated on Fig. 12 which displays the normalized central counter spectra at 10, 25 and 100 GeV.

The total energy summed over the 21 calibrated counters of the matrix is shown in Fig. 13 for electrons traversing the central crystal near its axis, in a 6 x 6 mm² area.

The shape of the spectrum is nearly gaussian except for a small tail at the left side. The resolution has been determined by fitting the spectrum with a gaussian between 50% of the maximum on the left wing and 10% of the maximum on the right wing of the peak.

The result of the fit includes the beam momentum spread $\sigma p/p = 0.4\%$ and an electronic noise component which can be split in two contributions: the first one is correlated to pedestal fluctuations; the second is due to the minimum detection threshold of 3 in ADC units which increases the fluctuation of the number of unrecorded small amplitudes. The corrected resolution is presented in Fig. 14 with error bars. The prototype matrix resolution is better than 1% above 25 GeV. It is fitted by the following relation:

$$\frac{\sigma E}{E} = \frac{3.25 \pm 0.6}{\sqrt{E}} \oplus 0.8 \pm 0.06 \tag{1}$$

Both errors are statistical*. These results are in agreement with a preliminary ones [9] that was quoted without errors bars. The difference between the present constant term value and the one quoted in [9] reflects the systematic uncertainty due to experimental conditions of the test.

Pure photoelectron statistics should contribute to the energy dependent term as $1.7\%/\sqrt{E}$ only, if one extrapolates from muons results (§ 4.4). The residual contribution comes from the structure of multi-cell calorimeters [15].

4.6 Photodiode Tests

A single 24 X_0 crystal (2 x 2 cm²) has been tested with a photodiode (Hamamatsu S3204-03, 300 μ m, 1 x 1 cm² area, fully depleted by a -60 volt bias). The signal was fed to a charge preamplifier designed for L3 experiment by Oakridge Laboratory, followed by a negative output shaping amplifier designed by the Electronic Department of IPN Lyon. Both

^{*} The symbol \oplus implies the square root of the quadratic sum of the two terms.

have a 2 µs shaping time. The shaping amplifier was connected to a Lecroy 2259 peak sensitive ADC.

Fig. 15 shows the spectra of 25 GeV (a) and 100 GeV/c (b) electrons hitting the crystal in a 5 x 5 mm² area around its axis. The same optical coupling as for the PM test was used with a collection geometrical factor $\varepsilon_g = 0.25$. To keep the measurements in the same ADC range at both energies, a factor 4 amplification has been applied at 25 GeV. Respective resolutions are roughly 2.7 and 2.3%. The leakage of the 100 GeV shower through the diode is clearly observed.

Fig. 16 displays the spectrum of 25 GeV muons crossing the diode itself after blinding scintillation light by insertion of a piece of black paper. An amplification factor of 2 was applied for ease of comparison with the 25 GeV spectrum of Fig. 15a. A mean number of 32 000 electron pairs are produced in a 300 μ m silicium diode crossed by a minimum ionising particle. Using the energy scale set by the 25 GeV spectrum (Fig. 15a) and an overall correction factor, that takes into account a geometrical collection efficiency $\epsilon_g = 0.25$ and the nominal 55% mean quantum efficiency of the diode, one gets a mean photons number of \sim 120 photons/MeV seen by the diode without Fresnel correction (of the order of 4%).

CONCLUSIONS

For the first time, after optimisation of the industrial growing process for the production of large optically homogenous PWO heavy crystals with a good radiation hardness, an electromagnetic calorimeter prototype made of a matrix of 21 cells with a length of at least 20 X_0 has been tested. An energy resolution of $\sigma_{E/E} = 3.25/\sqrt{E} \oplus 0.8$ has been achieved using XP1911 photomultipliers with a light collection efficiency of only 38% and a non optimised optical coupling.

ACKNOWLEDGEMENTS

The authors are grateful to B. Ille from IPN Lyon who has kindly provided the basic electronics and advices for the photodiode tests and to J.P. Vialle for improvements of laboratory tests software. They are indebted to A. Bazan and T. Le Flour who have modified

and adapted the NA12 acquisition system to the needs of these tests.

They also thank the entire team operating the ULYSSE reactor and more particularly G. Dauphin and J. Safieh from INSTN/SENE CEA Saclay. They are grateful to C. Jeanney from CEA/DAPNIA/SED for his valuable help for optical machining and measurements.

Some participants are indebted to J. Schukraft and to the ALICE collaboration for their financial support.

Finally, they would like to thank M. Kobayashi and T. Tsuru from KEK for fruitful discussions and help during part of the tests and K. Takamatsu and J.P. Stroot for their constant interest and active support.

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FIGURES CAPTIONS

- Fig. 1 1-Absorption coefficient K (left scale) and 2-the γ -ray excited (Co⁵⁷, 122 KeV) luminescence of a reference crystal at T = 300°K (right scale).
- Fig. 2 PWO scintillation decay time spectrum under γ ray excitation (Na²²), at T = 300°K.
- Fig. 3 The amplitude spectra of PWO et BGO scintillators excited by different energy gamma sources. $T = 300^{\circ}$ K. [137 Cs, 241 Am] and Am-Be.
- Fig. 4 a) PWO recovery of transparency after γ irradiation (transmission through 1 cm) for enriched Pb³⁺ sample.
 - b) PWO transparency versus γ absorbed dose (through 1 cm) for depleted Pb³⁺ sample. Bottom curve remind the shape of scintillation spectrum.
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- Fig. 6 The 21 PWO crystal matrix at CERN.
- Fig. 7 a) Fraction of the collected charge amplitude A/A(400) versus the ADC gate width

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- b) Oscilloscope traces of 100 GeV electrons after 4m of 50 Ω cable at CERN.
- Fig. 8 a) The horizontal x position of the center of gravity of the calorimeter signals versus the beam coordinate.
 - b) Position resolution σ_x versus the beam coordinate x(beam).
- Fig. 9 a) The 26 GeV electron shower profile A₆/A_{tot}.
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- Fig. 16 Muons spectrum through the photodiode itself.

Table 1: Basic characteristics of some heavy crystals

	Density	Radiation	Moliere	Light yield	Decay time	Peak	Refractive
Crystal		length	radius			emission	index
	(g/cm ³)	(cm)	(cm)	(% of BGO)	(ns)_	(nm)	
BGO	7.13	1.12	2.25	100	300	480	2.15
CeF ₃	6.16	1.68	2.63	45	5/20	300/340	1.62
BaF ₂	4.89	2.05	4.4	40/160	0.6/630	210/320	1.49
PWO	8.28	0.87	2.19	5	2.1/7.5/26	480	2.16
NaBi(WO ₄) ₂	7.57	0.98	2.66	č	-	-	2.05
Gd ₃ Ga ₅ O ₁₂	7.02	1.45	2.4	č	-	-	2.0
PbF ₂	7.56	0.95	2.2	Č	-	-	1.82

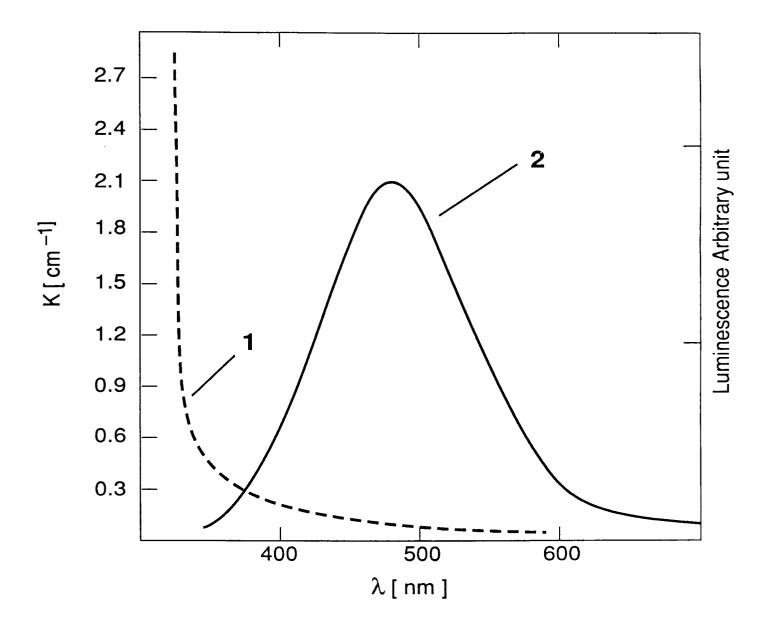
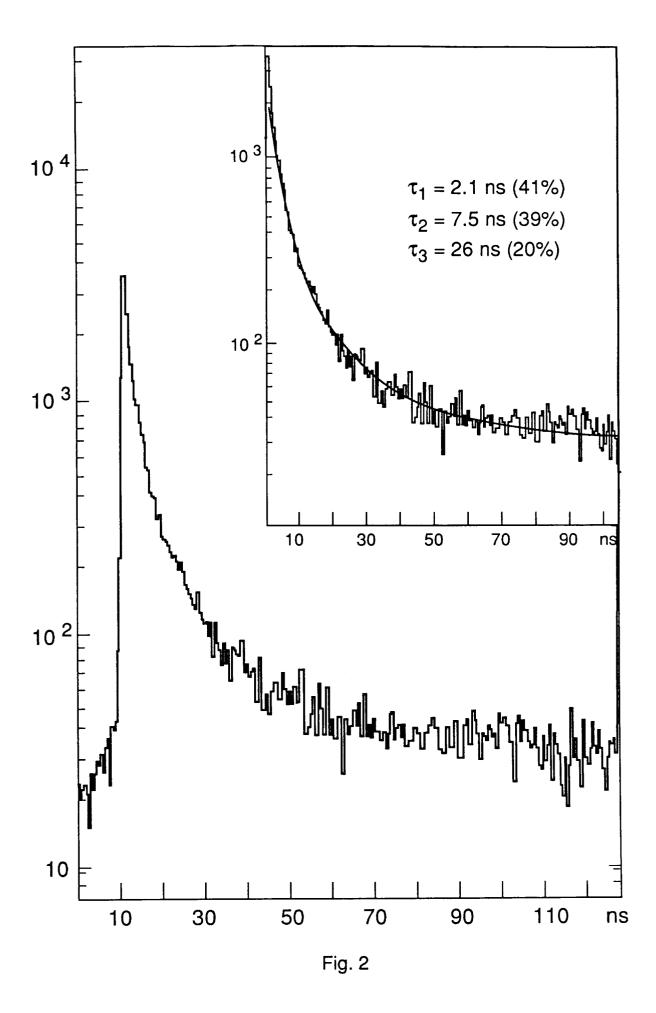
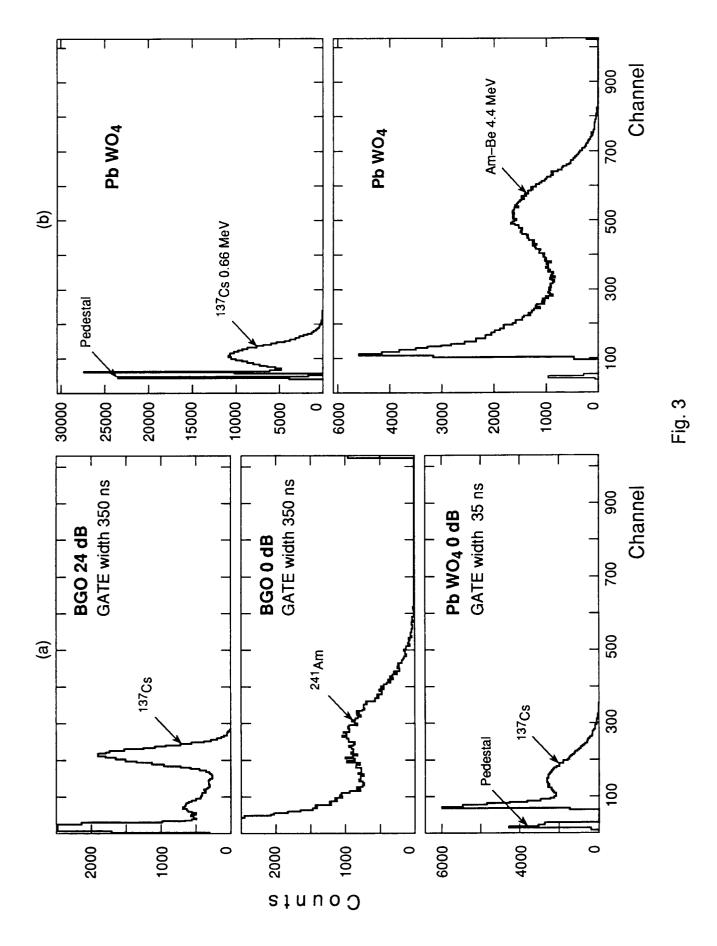


Fig. 1





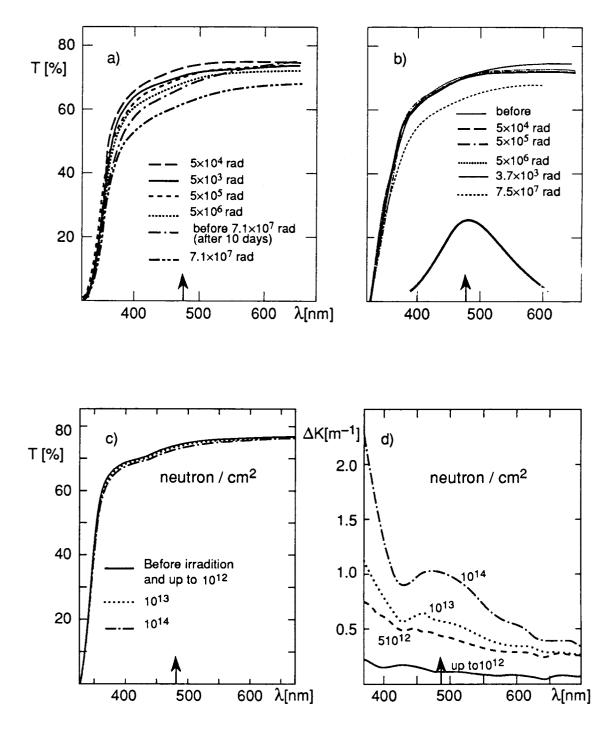


Fig. 4

GGG	PW01	NBW	GGG	
17X ₀	17X ₀	18X ₀	17X ₀	
PW02	PW06	PW05	NBW	
17X ₀	20X ₀	18X ₀	19X ₀	
GGG	NBW	NBW	GGG	
17X ₀	18X ₀	18X ₀	17X ₀	
GGG	GGG	GGG	GGG	
17X ₀	17X ₀	17X ₀	17X ₀	

Fig. 5

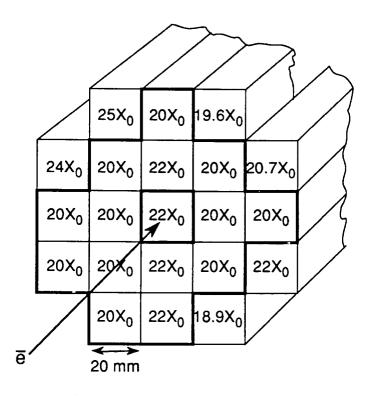


Fig. 6

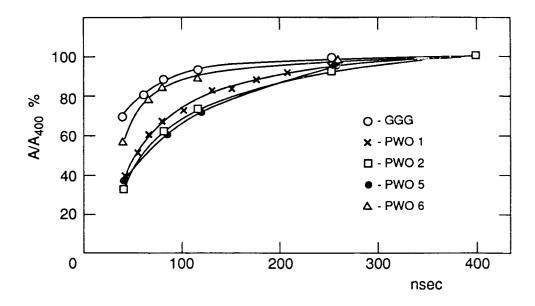


Fig. 7a

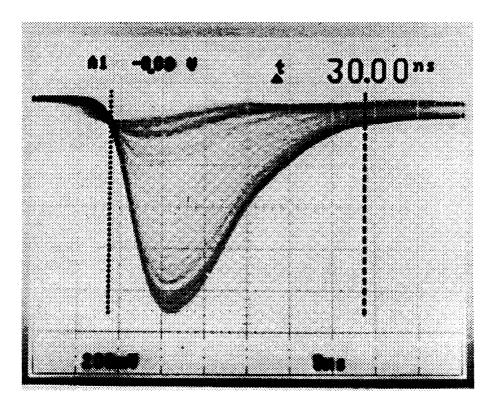


Fig. 7b

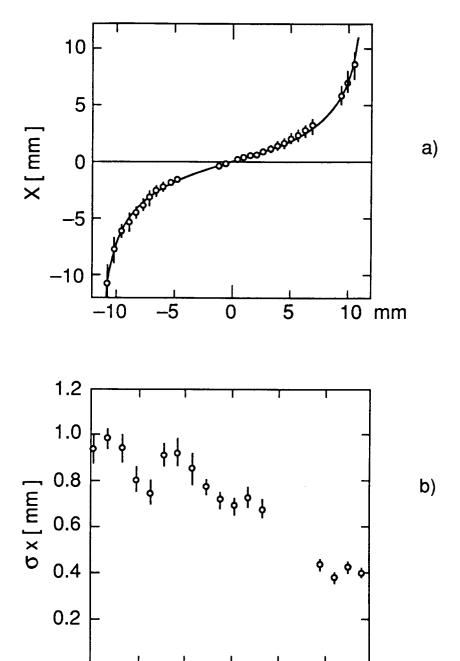
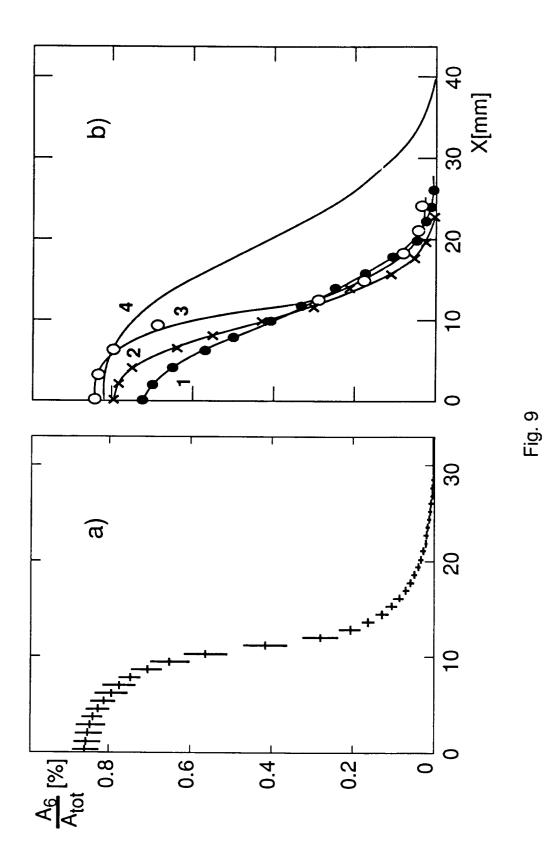


Fig. 8

X [mm]



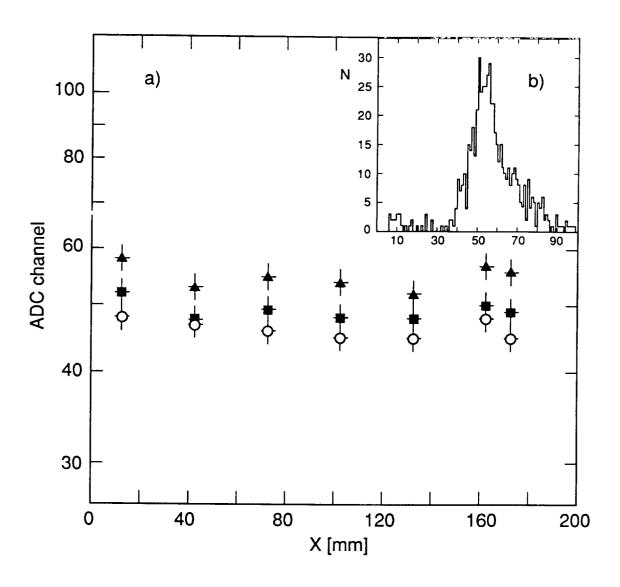


Fig. 10

CENTRAL PART of Pb WO₄ MATRIX e⁻ 25 GeV

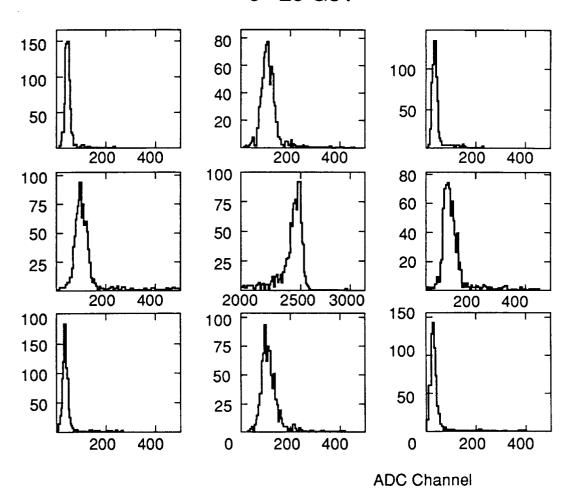


Fig. 11

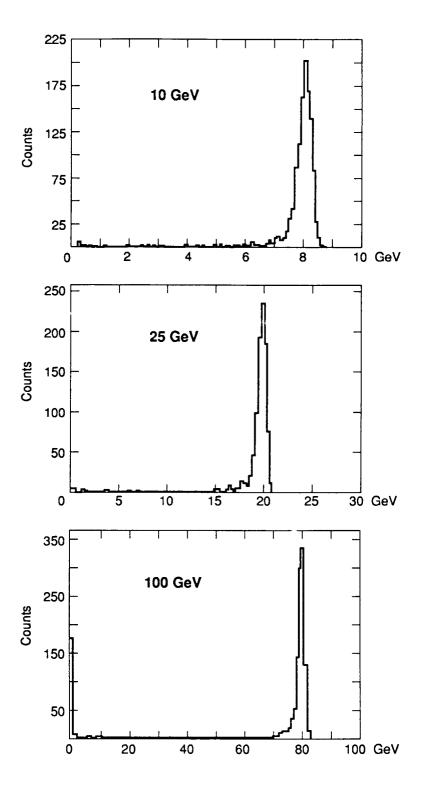


Fig. 12

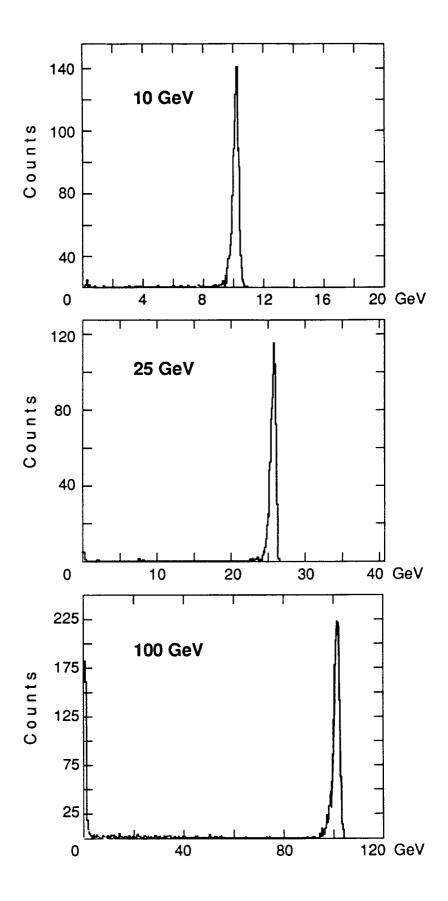


Fig. 13

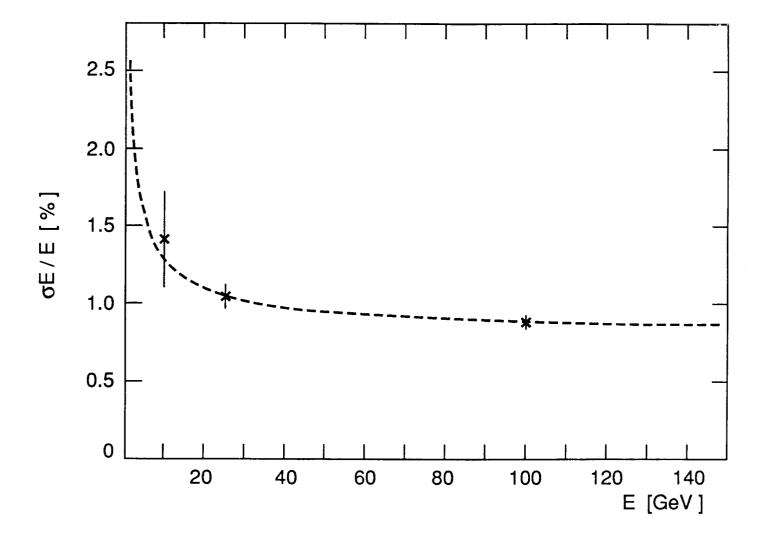


Fig. 14

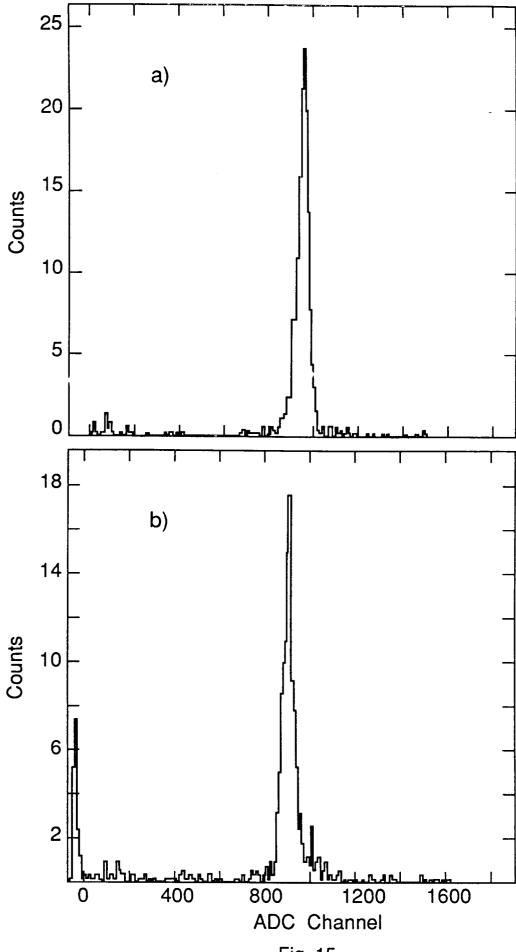


Fig. 15

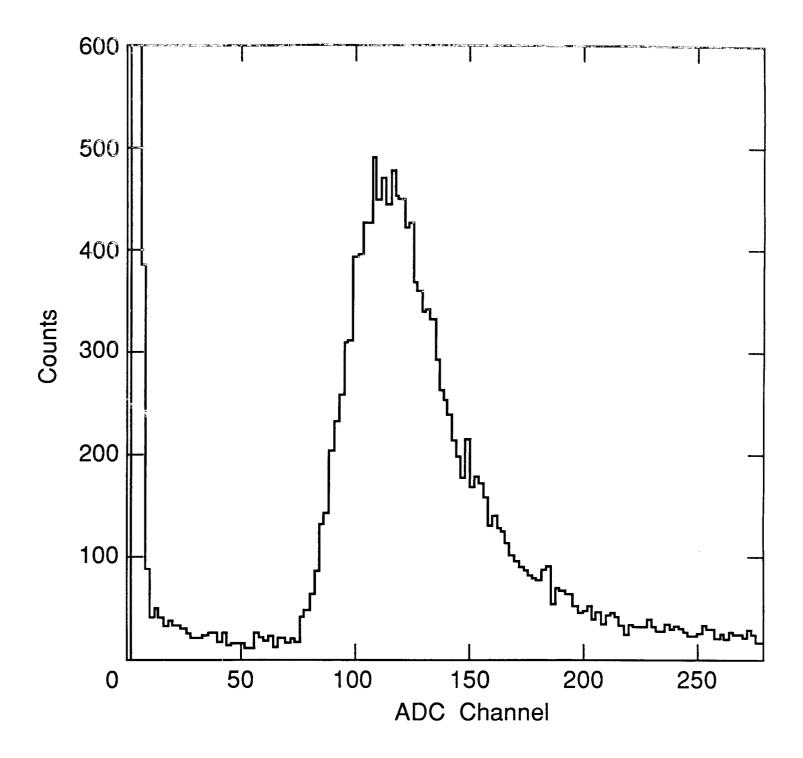


Fig. 16