

1     Measurements of the charge asymmetry in top-quark pair  
 2     production in the dilepton final state at  $\sqrt{s} = 8$  TeV with the  
 3                     ATLAS detector

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7             Measurements of the top–antitop quark pair production charge asym-  
 8     metry in the dilepton channel, characterized by two high- $p_T$  leptons (elec-  
 9     trons or muons), are presented using data corresponding to an integrated  
 10    luminosity of  $20.3 \text{ fb}^{-1}$  from  $pp$  collisions at a center-of-mass energy  $\sqrt{s} =$   
 11    8 TeV collected with the ATLAS detector at the Large Hadron Collider at  
 12    CERN. Inclusive and differential measurements as a function of the invari-  
 13    ant mass, transverse momentum, and longitudinal boost of the  $t\bar{t}$  system  
 14    are performed both in the full phase space and in a fiducial phase space  
 15    closely matching the detector acceptance. Two observables are studied:  
 16     $A_C^{\ell\ell}$  based on the selected leptons and  $A_C^{t\bar{t}}$  based on the reconstructed  $t\bar{t}$  fi-  
 17    nal state. No significant deviation from the Standard Model expectations  
 18    is observed.

19                             PRESENTED AT

20                             9<sup>th</sup> International Workshop on Top Quark Physics  
 21                             Olomouc, Czech Republic, September 19–23, 2016

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<sup>1</sup>Work supported by the Helmholtz Association of German Research Centres.



# 1 Introduction

The measurements of the charge asymmetry provide a good precision test of the Standard Model (SM). In the SM, the asymmetry is produced by interferences between the Born and one-loop diagram of the  $q\bar{q} \rightarrow t\bar{t}$  processes and between  $q\bar{q} \rightarrow t\bar{t}g$  diagrams with initial-state and final-state radiation. In the  $t\bar{t}$  rest frame, this asymmetry causes the top quark to be preferentially emitted in the direction of the initial quark, and causes the antitop quark to be emitted in the direction of the initial antiquark. In the  $pp$  collision at the LHC, valence quarks carry on average a larger fraction of the proton momentum than sea antiquarks, hence top antiquarks produced through  $q\bar{q}$  annihilation are more central than top quarks. In dileptonic events, the charge asymmetry can be measured in two complementary ways: using the pseudorapidity of the charged leptons or using the rapidity of the top quarks. The asymmetry based on the charged leptons uses the difference of the absolute pseudorapidity values of the positively and negatively charged leptons,  $|\eta_{\ell^+}|$  and  $|\eta_{\ell^-}|$ . The leptonic asymmetry is defined as

$$A_C^{\ell\ell} = \frac{N(\Delta|\eta| > 0) - N(\Delta|\eta| < 0)}{N(\Delta|\eta| > 0) + N(\Delta|\eta| < 0)} \quad \text{with } \Delta|\eta| = |\eta_{\ell^+}| - |\eta_{\ell^-}|, \quad (1)$$

where  $N(\Delta|\eta| > 0)$  and  $N(\Delta|\eta| < 0)$  represent the number of events with positive and negative  $\Delta|\eta|$ , respectively. For the  $t\bar{t}$  charge asymmetry the absolute values of the top and antitop quark rapidities ( $|y_t|$  and  $|y_{\bar{t}}|$ , respectively) are used. The  $t\bar{t}$  charge asymmetry is defined as

$$A_C^{t\bar{t}} = \frac{N(\Delta|y| > 0) - N(\Delta|y| < 0)}{N(\Delta|y| > 0) + N(\Delta|y| < 0)} \quad \text{with } \Delta|y| = |y_t| - |y_{\bar{t}}|, \quad (2)$$

where  $N(\Delta|y| > 0)$  and  $N(\Delta|y| < 0)$  represent the number of events with positive and negative  $\Delta|y|$ , respectively.

In these proceedings, the inclusive and differential measurements of the leptonic and  $t\bar{t}$  charge asymmetry in the dilepton channel using data collected by the ATLAS detector [1] corresponding to an integrated luminosity of  $20.3 \text{ fb}^{-1}$  from  $pp$  collisions at a center-of-mass energy  $\sqrt{s} = 8 \text{ TeV}$  are presented [2]. The differential measurements are performed as a function of the mass ( $m_{t\bar{t}}$ ), transverse momentum ( $p_T^{t\bar{t}}$ ) and boost ( $\beta_z^{t\bar{t}}$ ) of the  $t\bar{t}$  system. The measurements are performed in a fiducial region and in the full phase space.

## 2 Event Selection and Reconstruction

Events are required to have exactly two leptons of opposite electric charge and at least two jets with  $p_T > 25 \text{ GeV}$  within  $|\eta| < 2.5$ . In all three final states, exactly

53 two isolated leptons with opposite charge and an invariant mass  $m_{\ell\ell} > 15$  GeV  
 54 are required. In the same-flavor channels ( $ee$  and  $\mu\mu$ ), the invariant mass of the  
 55 two charged leptons is required to be outside of the  $Z$  boson mass window such that  
 56  $|m_{\ell\ell} - m_Z| > 10$  GeV. Furthermore, it is required that missing transverse momentum  
 57 is greater than 30 GeV and at least one of the jets must be  $b$ -tagged. In the  $e\mu$  channel,  
 58 the scalar sum of the  $p_T$  of the two leading jets and leptons is required to be larger  
 59 than 130 GeV.

60 The main background contribution comes from Drell–Yan production of  $Z/\gamma^* \rightarrow$   
 61  $\ell\ell$ , which is estimated by a combination of simulated samples and corrections derived  
 62 from data. The smaller contributions from diboson and single-top-quark production  
 63 are evaluated purely via simulations. Contributions arising from events including a  
 64 jet or a lepton from a semileptonic hadron decay misidentified as an isolated charged  
 65 lepton as well as leptons from photon conversions, are estimated using simulated  
 66 samples, modified with corrections derived from data.

67 The  $t\bar{t}$  system is reconstructed in order to perform the inclusive and differential  
 68 measurements of  $A_C^{t\bar{t}}$ . The system is reconstructed using the KIN method. The KIN  
 69 method assumes the mass of the top quarks (172.5 GeV) and  $W$  mass (80.4 GeV),  
 70 and solve the system of equations obtained from momentum conservation numerically  
 71 using the Newton-Rhapson method. The reconstruction efficiency is above 90%. A  
 72 comparison between observations and expectations is shown in Fig. 1 after event  
 73 reconstruction. A good agreement within the uncertainties is observed.

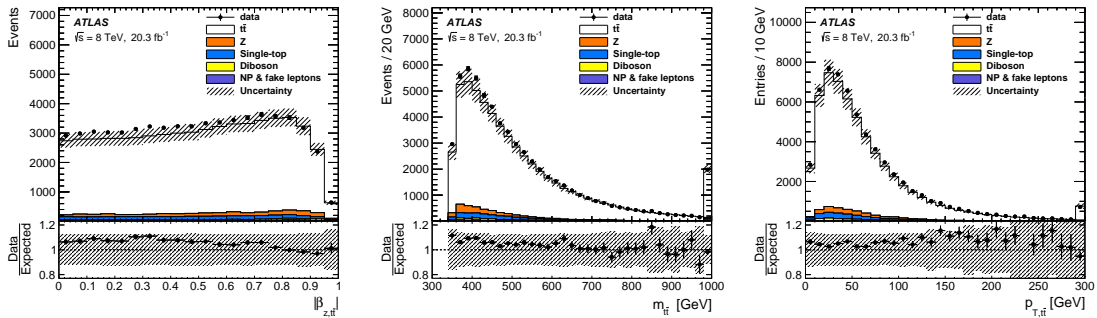


Figure 1: Comparison between observations and expectations for the boost (left) mass (middle), and  $p_T$  of the  $t\bar{t}$  system. The total uncertainty on the distributions are shown [2].

### 74 3 Unfolding

75 The measurements of  $A_C^{\ell\ell}$  and  $A_C^{t\bar{t}}$  are corrected in order to remove the effects in-  
 76 troduced by the detector. This correction is performed by using the Fully Bayesian

77 Unfolding (FBU) method [5]. The measurements are unfolded back to a stable par-  
78 ticle level in a fiducial region closely matching the detector acceptance, and back to  
79 parton level in the full phase space. The combined measurement of the three decay  
80 channels is performed during the unfolding procedure. The asymmetries are com-  
81 puted using the posterior probability density obtained as an output of the unfolding  
82 procedure. Systematic uncertainties related with detector modeling and background  
83 modeling are evaluated during the unfolding by using a marginalization procedure.

## 84 4 Results

85 Figure 2 shows the inclusive and differential measurements performed for  $A_C^{\ell\ell}$  and  $A_C^{t\bar{t}}$   
86 at parton level in the full phase space. The total uncertainty on the measurements  
87 is shown. The main source of uncertainty on the different measurements is the sta-  
88 tistical uncertainty, followed by the signal modeling uncertainty. The measurements  
89 that involve the reconstruction of the  $t\bar{t}$  system are also affected by a reconstruction  
90 uncertainty which is approximately half of the size of the statistical uncertainty. The  
91 uncertainties corresponding to the detector and background modeling do not con-  
92 tribute significantly to the total uncertainty. The results are compatible with the SM  
93 predictions [3]. A similar behavior is observed on the uncertainties in the measure-  
94 ments performed in the fiducial region, however, there is a reduction in the modeling  
95 uncertainties. Figure 3 shows the unfolded distribution for the  $\Delta|y|$  and  $\Delta|\eta|$  observ-  
96 ables in the fiducial region. The distribution is in agreement with SM predictions.  
97 Figure 4 shows the  $A_C^{\ell\ell}$  and  $A_C^{t\bar{t}}$  measurements in comparison with several models be-  
98 yond the SM [4] in the full phase space. In these models, the values of the asymmetry  
99 are expected to be different from the SM expectations. The ellipses correspond to  
100 the  $1\sigma$  and  $2\sigma$  total uncertainty on the measurements. The correlation between  $A_C^{\ell\ell}$   
101 and  $A_C^{t\bar{t}}$  is about 48%. The measurements are compatible with the SM and do not  
102 exclude the two sets of BSM models considered.

## 103 References

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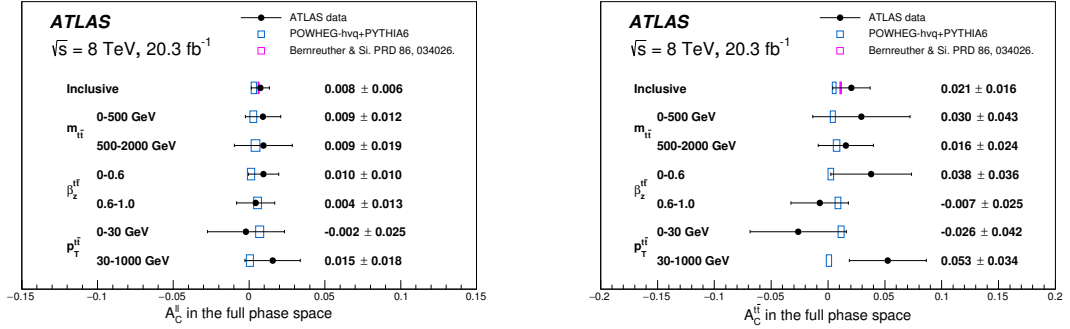


Figure 2: Summary of the inclusive and differential measurements of the  $t\bar{t}$  asymmetry (left) and lepton asymmetry (right) performed in the full phase space. The measurements are compared with theoretical predictions [2].

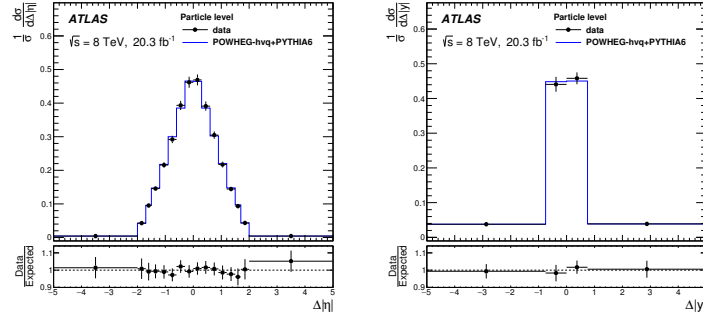


Figure 3: Data distribution after the unfolding procedure compared with the prediction for the inclusive  $\Delta|\eta|$  (left) and  $\Delta|y|$  (right) observables in the fiducial volume. The data/expected ratio is also shown [2].

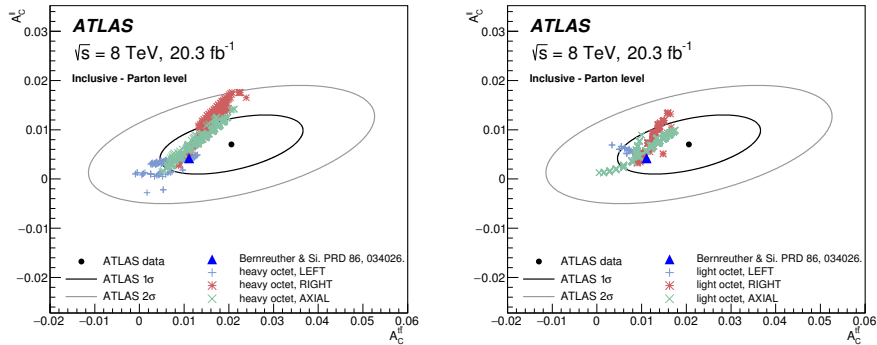


Figure 4: Comparison of the inclusive  $A_C^{\ell\bar{\ell}}$  and  $A_C^{t\bar{t}}$  measurement values in the full phase space to the SM and to two benchmark BSM models [2].