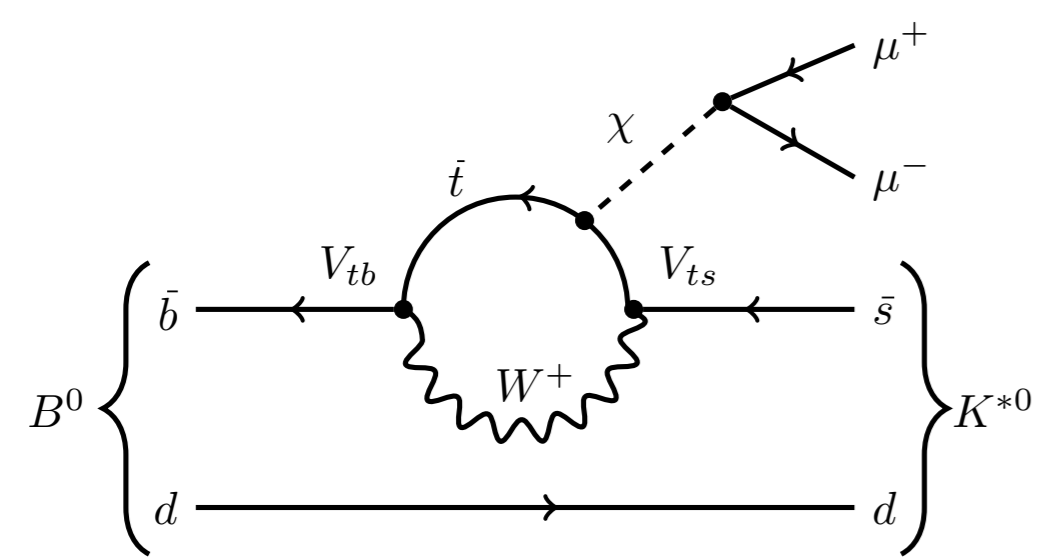


Motivation

- Many theoretical models predict the existence of new particles that couple to the SM by **mixing with the Higgs**.
- Inflaton, axion-like, dark matter mediator also predict the new boson to be **light** [1].

If kinematically allowed,
main production via B meson:



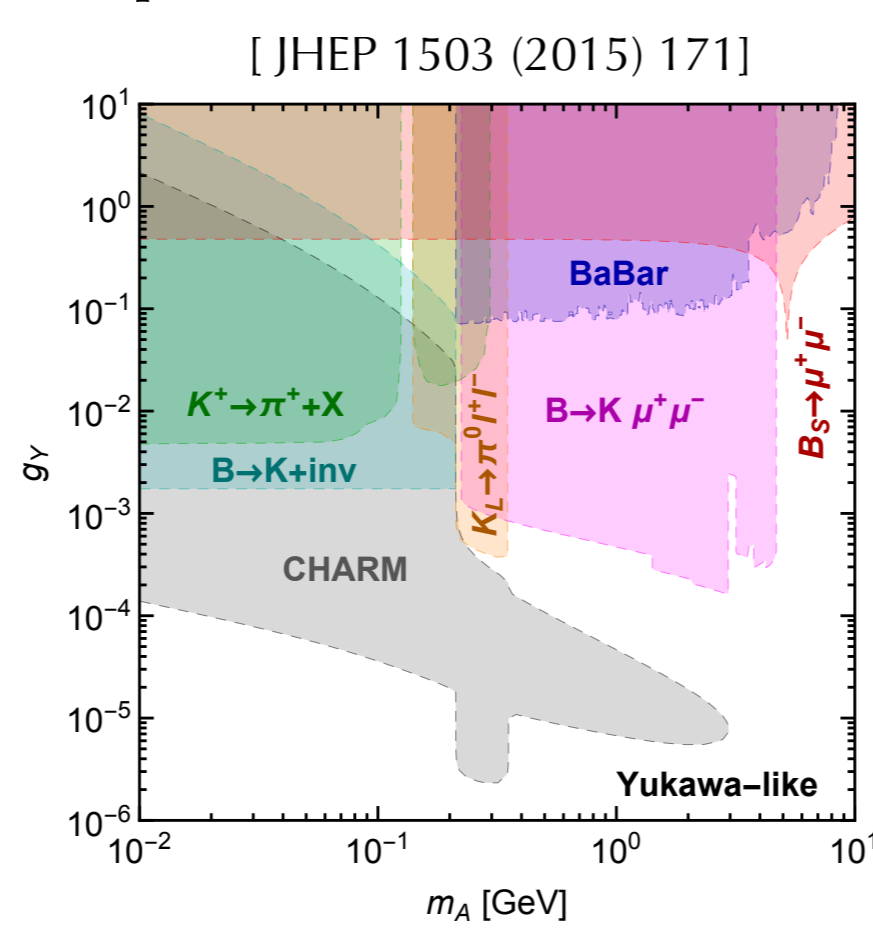
$$\Gamma(B^0 \rightarrow K^{*0} \chi) \propto (m_t^2 |V_{ts}^* V_{tb}|)^2 \propto m_t^4 \lambda^2$$

Branching fraction into muons

$$\mathcal{B}(\chi \rightarrow \mu^+ \mu^-):$$

- dominant till the hadronic threshold ($\chi \rightarrow 2h$, $\chi \rightarrow 3h$)
- always significant $\mathcal{O}(10^{-2})$ in the full mass range.

Current experimental limits: [2]

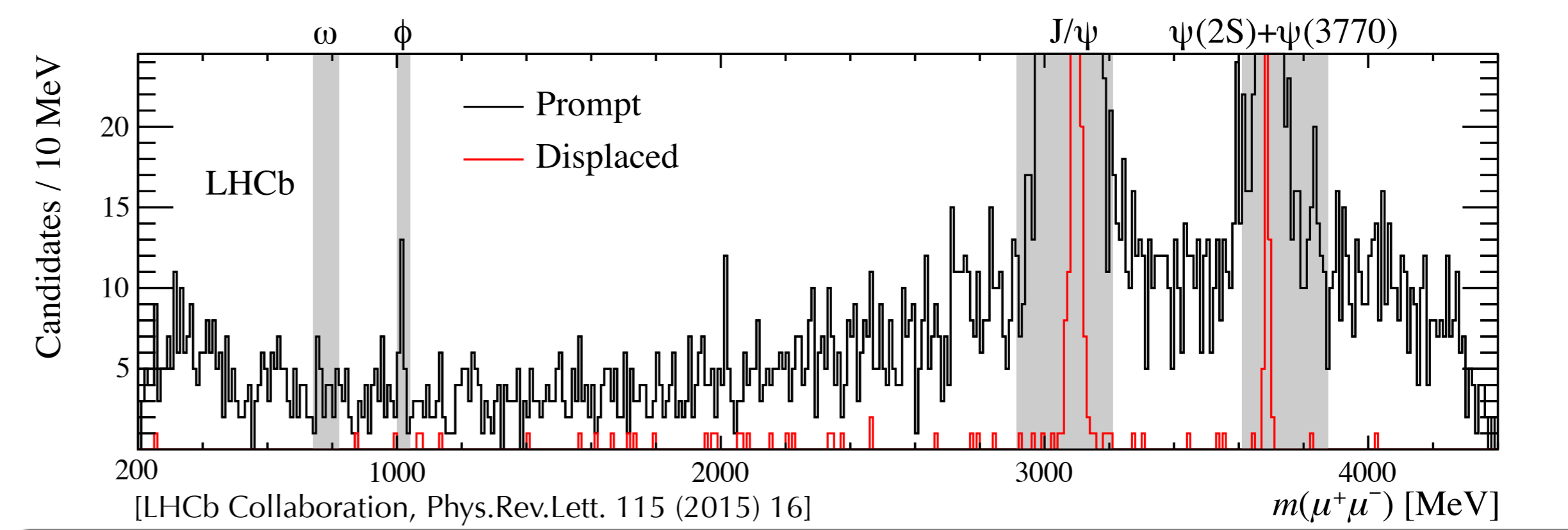


$g_Y = \sin \theta$, θ mixing angle with the Higgs

Main current constraints come from B factories and CHARM experiment (at low mass).

Result

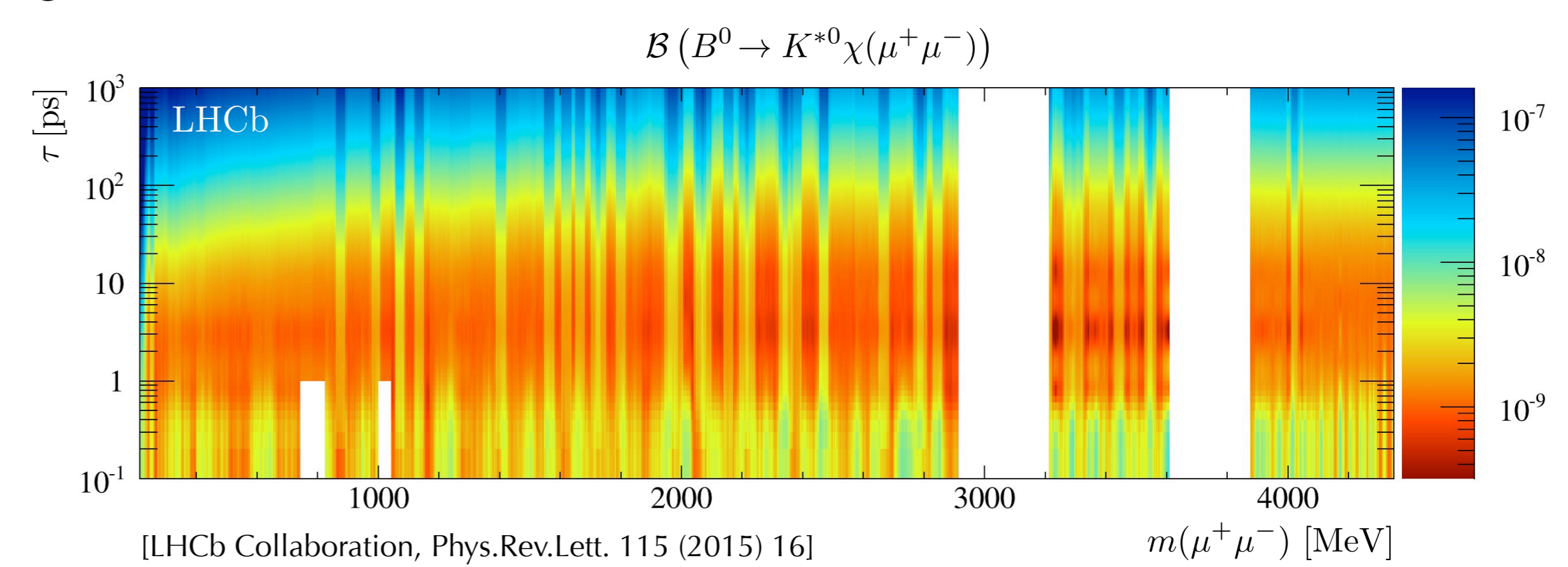
- Look Elsewhere Effect taken into account.
- Largest deviation at $m = 253$ MeV, local p -value = 0.02.
- No deviation from the background-only hypothesis is observed** [5].



Grey regions are vetoed in the analysis.

Exclusion limit

- 95% CL upper limit on the branching fraction $\mathcal{B}(B^0 \rightarrow K^{*0} \chi(\rightarrow \mu^+ \mu^-))$ [5]
- function of mass and lifetime
- assumed spin 0
- long lifetimes have weaker limit due to the loss of reconstruction efficiency



- Model-independent limit: can be interpreted in your favorite BSM theory.

Signal properties

- Depending on the coupling SM/hidden sector we can identify two **lifetime** regimes:

Short lifetime

Pros:

- high reconstruction efficiency

Cons:

- prompt decay
 - large SM background contamination

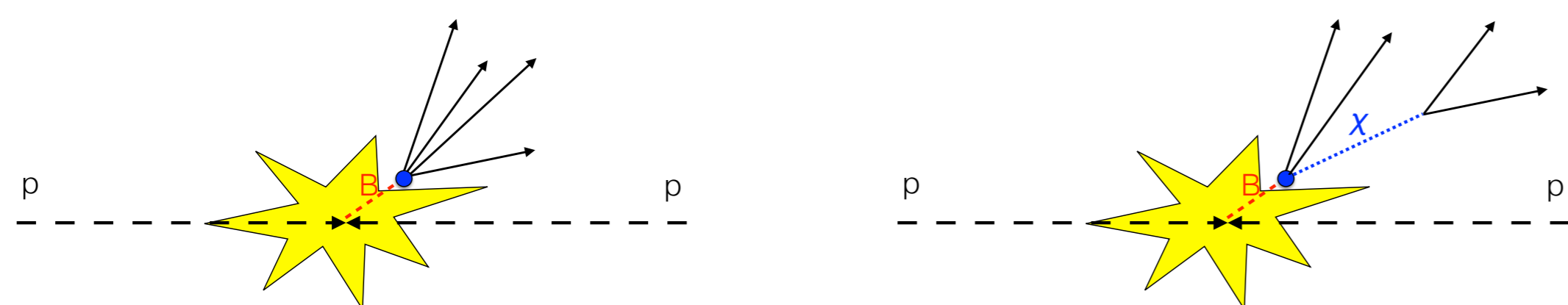
Long lifetime

Pros:

- displaced vertex
 - very clear signature

Cons:

- signal events start to decay after the vertex detector
 - lower efficiency



- A di-muon displaced vertex is allowed, but not required.

Selection & Strategy

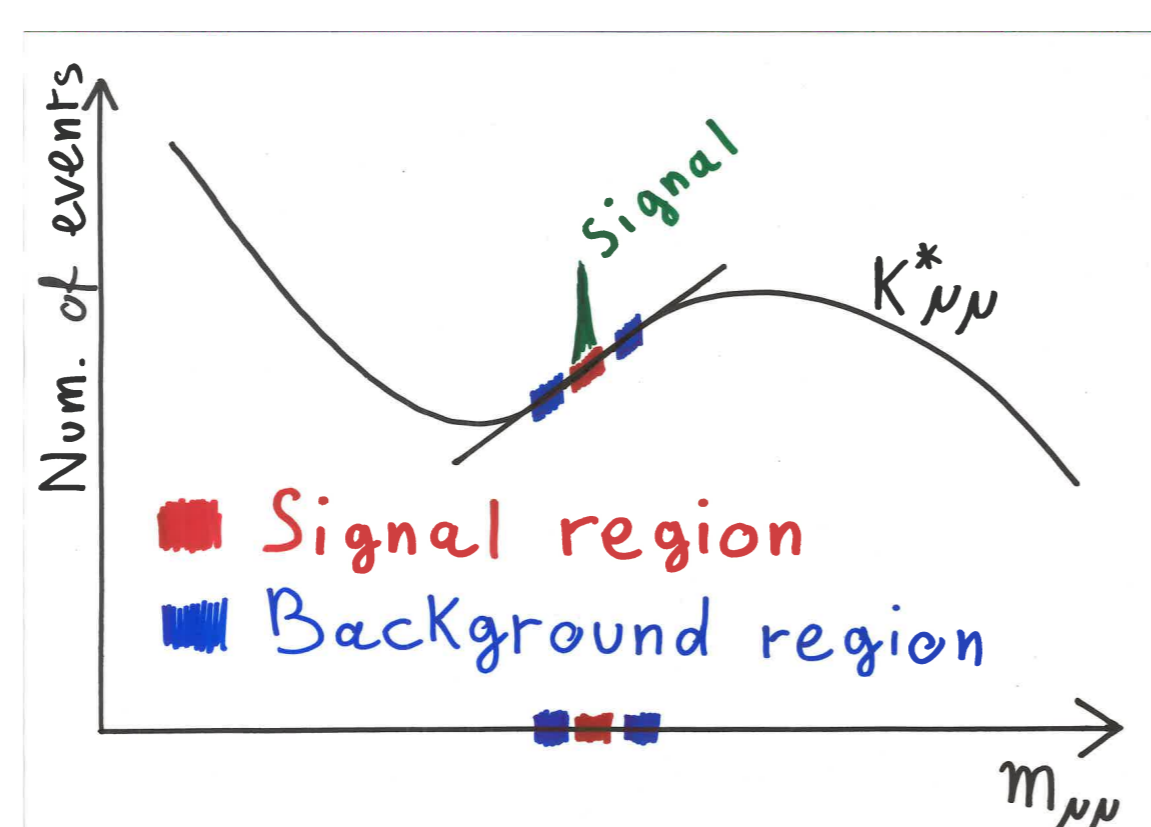
- Multivariate selection: **uBDT** - uniform Boosted Decision Tree [3]
 - performance independent of mass and lifetime of χ
- Optimized maximizing Punzi figure-of-merit $P_{f.o.m.} = \frac{S}{\frac{5}{2} + \sqrt{B}}$
- Factorize lifetime into two components:
 - prompt**, $\tau < 3\sigma_\tau$ (irreducible SM $B^0 \rightarrow K^{*0} \mu^+ \mu^-$)
 - displaced**, $\tau > 3\sigma_\tau$ (only combinatorial background)

lifetime resolution: $\sigma_\tau \sim 0.2$ ps

$$\mathcal{L} = \mathcal{L}^{prompt} \cdot \mathcal{L}^{displaced}$$

- Looking for a di-muon excess:

- $\sigma_m = 2 \div 8$ MeV.
- B^0 mass is constrained to its known value to improve to mass resolution.
- scan the di-muon mass in step of $\frac{1}{2}\sigma_m$ [4].
- di-muon sidebands used to extract the expected number of background events assuming local linearity.



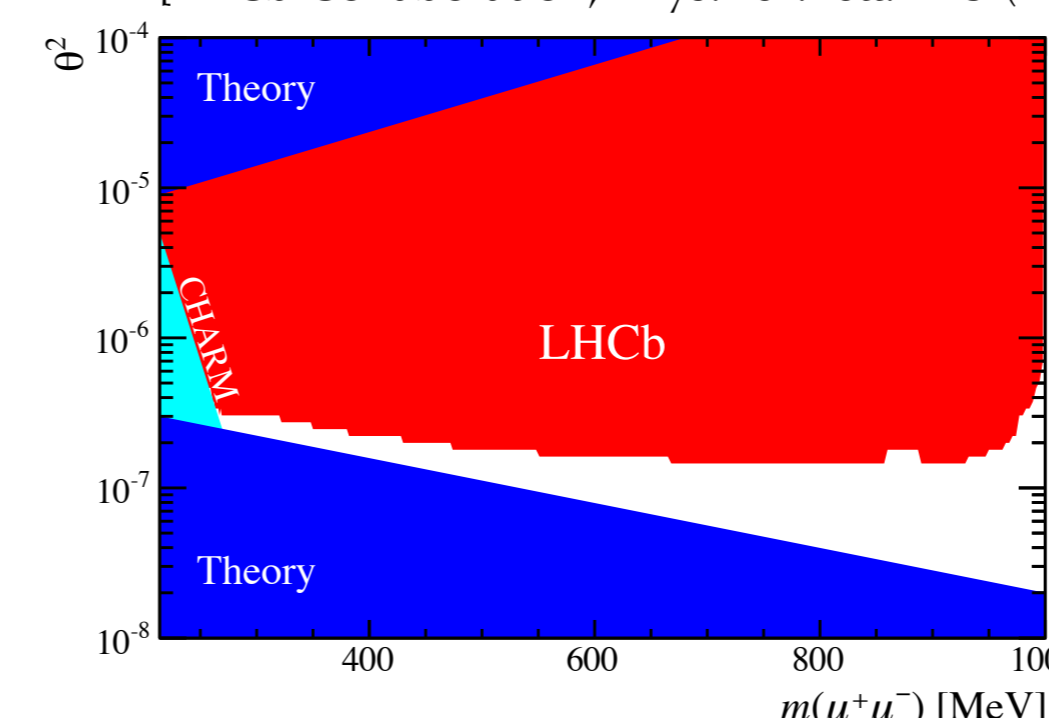
Benchmark models

- Two model-dependent exclusion limits:

Inflaton [6]

- $0.1 < m_\chi < 1$ MeV
- $\tau_\chi \sim 10^{-5} \div 10^{-12}$ s

[LHCb Collaboration, Phys.Rev.Lett. 115 (2015) 16]

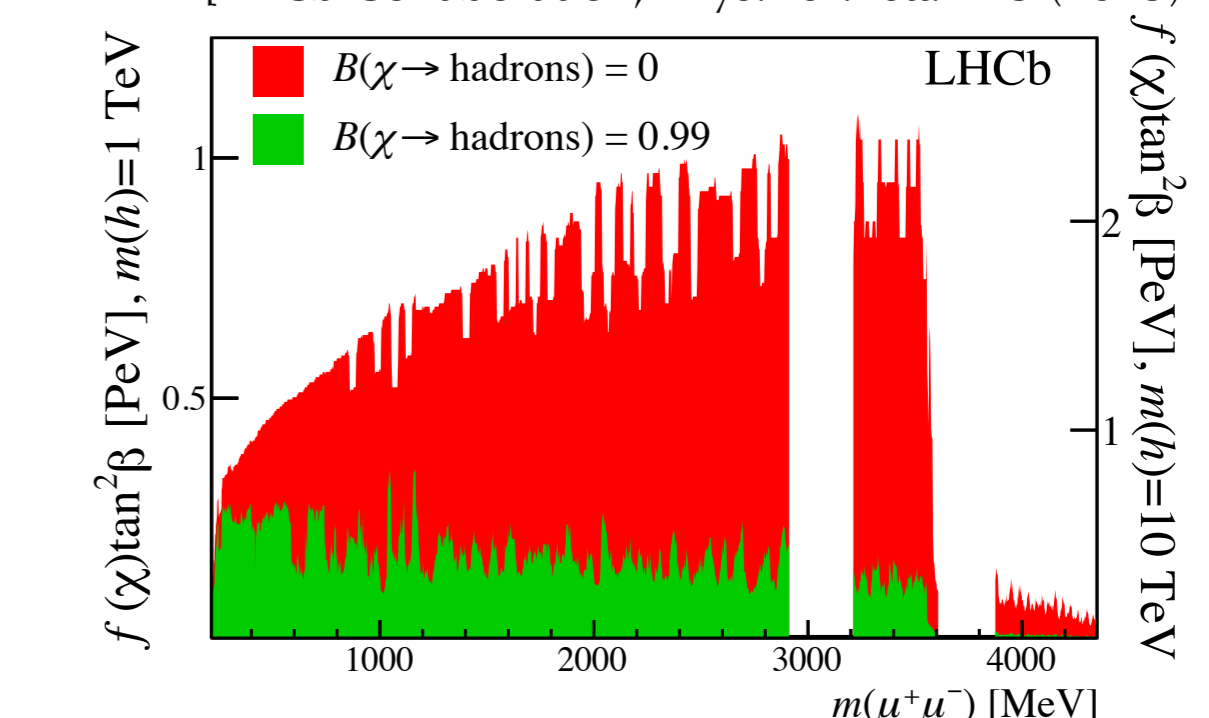


Cosmological and CHARM experiment excluded regions are also shown.

Axion-portal [7]

- larger allowed mass range
- short lifetime

[LHCb Collaboration, Phys.Rev.Lett. 115 (2015) 16]



Excluded region by the analysis for two extreme branching fractions of the axion-like particle.

Complementary searches

- A second very similar search is ongoing at the LHCb experiment:

$$B^+ \rightarrow K^+ \chi(\rightarrow \mu^+ \mu^-)$$

Differences:

- Only 3 tracks in the final state
 - less precise B decay vertex reconstruction.
- More sensitive to (pseudo)scalar particle
 - angular momentum conservation:
 - $J^P(K^{*0}) = 1^-$, $J^P(K^+) = 0^-$
 - effect significant only at high mass.

References

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- J. Stevens and M. Williams, JINST 8, P12013 (2013)
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