a Congress Parts unthelra on

EUR OPEAN ORGANIZATION FOR NUCLEAR RESEARCH

NOT TO BE CIRCULATED

ONLY FOR EEC

PROPOSAL FOR AN EXPERIMENT TO DETECT THE POSSIBLE EXISTENCE OF FRACTIONALLY CHARGED PARTICLES

 $\label{eq:1} \mathcal{L}=\frac{1}{2}\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right)\mathcal{L}^{2}+\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right)\mathcal{L}^{2}+\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right)\mathcal{L}^{2}+\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right)\mathcal{L}^{2}+\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right)\mathcal{L}^{2}+\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right)\mathcal{L}^{2}+\left(\mathcal{L}^{2}+\mathcal{L}^{2}\right$

A. Buhler, G. Finocchiaro, T. Massam, A. Michelini Th. Muller, G. Petrucci, M. Schneegans and A. Zichichi

sk.

Recently much interest has been generated around the possible existence of fractionally charged particles").

In what follows we propose two possible experiments to be performed at the Proton Synchrotron; the first before the shut-down and the second immediately after. $\mathcal{O}(\mathcal{E}(\mathcal{E}_\mathcal{A})\cap \mathcal{E}(\mathcal{E}))$

USING M3 BEAM Ι.

 $\mathcal{C} \subset \mathcal{C}(\mathcal{E})$. The principle of the method is simultaneous magnetic and electrostatio analysis with high-velocity threshold rejection to reduce background and dE/dx measurements to provide supporting evidence.

Figure 1 shows the experimental set-up. The following formulae are relevant: $\rightarrow \pm \frac{1}{2}$

bending magnet

 $\Theta_{\rm M} \sim \frac{2}{5}$

 $\Theta_{\rm s} \sim \frac{2}{\rm p_{\beta}}$.

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{A}}_{\mathcal{A}}(\mathcal{L}^{\mathcal{A}}_{\mathcal{A}})) = \mathcal{L}(\mathcal{L}^{\mathcal{A}}_{\mathcal{A}}(\mathcal{L}^{\mathcal{A}}_{\mathcal{A}}))$

 $\label{eq:2} \mathcal{L} \left(\mathcal{L} \right) = \frac{1}{2} \sum_{i=1}^n \mathbf{E}_i \left(\mathcal{L} \right) \left(\mathcal{L} \right) \left(\mathcal{L} \right)$

electrostatic separator

If we set the magnetic deflection to accept momentum P_1 for single-charged particles we have $P = qP_1$. Therefore, for charge q we have:

$$
\beta = \frac{qP_1}{\sqrt{q^2P_1^2 + m^2}} \quad ,
$$

*) Gell-Mann and Zweig.

the captured

 $-2 -$

Mass resolution

Image_width $a)$

> The image as shown in Fig. 2 corresponds to $\Delta\beta \simeq 0.01$. a Tanahin Maséhi Tan

b) Dispersion in the separator

$$
\frac{\mathrm{d}\beta}{\beta} = \frac{\mathrm{d}\mathrm{P}}{\mathrm{P}} \times \frac{\mathrm{m}^2 \mathrm{d}\mathrm{P}_1}{\left(\mathrm{q}^2 \mathrm{P}_1^2 + \mathrm{m}^2\right)^{3/2}},
$$

which has maximum value of 0.39 (dP/P) at $m = \sqrt{2} qP$.

' Hence, the dispersion in the separators for 0.5% momentum $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L$ acceptance is $d\beta/\beta \leq 0.002\%$.

Limits of cross-section and time needed

We have 1000 $\frac{5}{\sqrt{2}}$ (with 6×10^{11} circulating protons) with good image conditions.

 \mathcal{L}^{max}

When the third world we We have observed that outside the \bar{P} peak in the separator curve the background counting rate is less than 3% of the P rate. We have not, in fact, looked further and it can be that the tail $\mathbf{P}^{\mathcal{A}}(t)$, where $\mathcal{A}^{\mathcal{A}}(t)$, $\mathcal{A}^{\mathcal{A}}(t)$ $\mathcal{L}^{\mathcal{L}}$ and $\mathcal{L}^{\mathcal{L}}$ are the positive of the contribution of the contribution of $\mathcal{L}^{\mathcal{L}}$ decreases more.

We will perform pulse-height analysis on this 3% background and we expect a factor 10 rejection in favour of the fractionally charged particles. and the company of the state

. In order to scan mass values from 0.4 GeV up to 1.8 GeV (see Fig. 2) we need 40 points, each one with 10^4 events, i.e. 18 hours of data taking. This allows us to detect an event rate of the order of 10⁻⁵ the antiproton rate. 음식 아이는 그 아니?

Conclusion

We will be able to scan a mass range for the fractionally charged particles, which goes from 0.4 GeV up to 1.8 GeV in one day, with a limit on the cross-section of the order of 10^{-3} 2 cm². The method is not limited in the mass value, and, in fact, by reducing the separator

an Kabupatèn
Kabupatèn K

→ 相互の (の) (の) (の) () () (

voltage we can scan in another 24 hours from 1.8 up to 3.0 GeV masses. Clearly the highest possible mass (\sim 5.7 GeV) compatible with the PS energy can be scanned with the same limit for the cross-section.

Contractor

 $\mathcal{A}=\{1,2,3,4,5,6\}$

Time needed

One week of PS as main users.

 $\tau \rightarrow \tau_{\rm eff}$

II. USING THE d15 BEAM

We have also studied the possibility of transforming the d_{15} beam into a 40 GeV/c beam. If the fractionally charged particles exist we should be able to observe them by using a beam transport system capable of deflecting these small charges at this high apparent momenta.

The modifications required for the d15 beam can be performed during the next shut-down. It is essentially a question of displacing a bending magnet and a few quadrupoles. The characteristics of the dis beam will not be destroyed.

With this 40 GeV/c beam we will perform pulse-height analysis in scintillation counters. The lower limit for the detectable crosssection is going to be given by the muon background. We are studying the problem of reducing this source of background by using a highly directional telescope and large anticoincidence counters. In this way we can increase our previous limit by at least two orders of magnitude.

Conclusion

We propose to make the necessary transformations in the dis beam as soon as possible. If this proposal is accepted, a collaboration with the CERN-ETH Group is envisaged.

- 3 -

Figure captions

Fig. 1 Shows the experimental set-up. The momentum is determined by the beam optics, with momentum acceptance of 0.5%. D is a four-metre long ethylene threshold Cerenkov counter to veto pions, electrons and muons. A and B are two plastic scintillation counters, 5 mm thick and 8 cm in diameter, to record the beam intensity. Counters M_1 and M_2 are 1 cm thick, 28 cm diameter counters, which give two independent dE/dx measurements.

 $- 4 - 4$

- Fig. 2 Shows the spacial resolution which can be obtained with the separator for a momentum spread of 1% at 2.5 GeV/c. To give an idea of the scale, the abscissa is also marked with $\sim 10^6$ the values of the mass of a particle 'A' of charge = $\frac{1}{3}$ \ldots which would pass through the mass slit at three settings of tho separator oompensating magnet. $\mathbb{C}^{\mathbb{Z}}$.
- Fig. 3 Shows the variation of the velocity of the $'A'$ as a function of its mass for fixed values of the nominal beam momontum:

$$
\beta = \frac{P}{\sqrt{P^2 + m^2/q^2}} \;
$$

where P is the nominal beam momentum, m is the mass, and *q* is the charge of the 'A'. The dotted lines on the 2.5 GeV/o ourve show the resolution in mass, whioh corresponds ta the width of the image at the mass slit.

 $\mathcal{L}^{\text{max}}_{\text{max}}$

than a

•

N
250

Fig &2

Velocity of as a function of mass for nominal beam momenta of 1.0 and 2.5 Gev/c. The dotted lines on the 2.5 Gev/c curve gwe an estimate of the mass resolution. The points marked p and π are plotted at the values of ρ corresponding to these particles and show where they will coincide with the \AA . $Fig 3 3$