

Bottomonia and B_c^+ Production at LHCb

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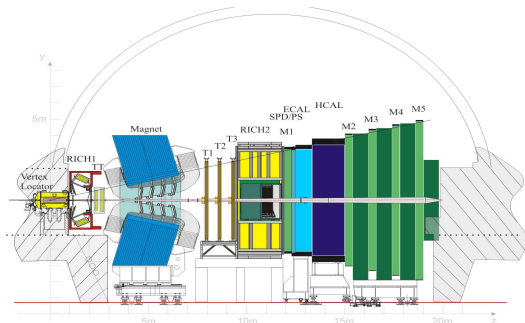


Outline

- LHCb detector
- Bottomonium studies
 - $\Upsilon(nS)$ production
 - χ_b production
- B_c^+ studies
 - B_c^+ mass and production using $B_c^+ \rightarrow J/\psi\pi^+$
 - Observation of $B_c^+ \rightarrow J/\psi\pi^+\pi^+\pi^-$



LHCb Detector



Smooth running so far:

- ✓ 2010: 37 pb^{-1}
- ✓ 2011: 1 fb^{-1}
- ✓ 2012: 2 fb^{-1}
delivered by last
Thursday!

Daria Savrina's talk this morning: *LHCb status and overview*.

- Unique rapidity range: $2 < \eta < 5$;
- Precision vertexing and tracking;
- Excellent particle identification;
- Precise and efficient muon reconstruction: $\sim 90\%$ trigger rate for dimuon channels!



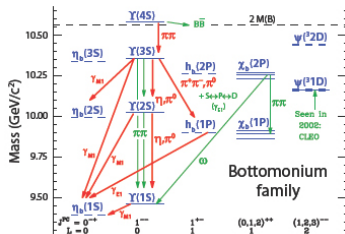
Bottomonia Productions

Bottomonia production mechanism is not yet well understood, several models exist:

- Colour Singlet Model (CSM, NLO CSM)
- Nonrelativistic QCD expansion (NRQCD) with contributions from Colour Octet Mechanism
- Colour Evaporation Model

Experimental inputs from LHC are needed.

- ✓ $\Upsilon(nS)$ inclusive production ($n = 1, 2, 3$)
- ✓ Υ from χ_b feed-down

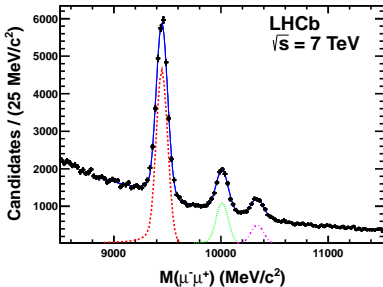




$\Upsilon(nS)$ Production

[Eur.Phys.J.C\(2012\)72:2025](#)

- 25 pb⁻¹ of 2010 data, $\sqrt{s} = 7$ TeV
- $\Upsilon(nS)$ reconstructed from di-muon final states:
 - 1 μ $p_T > 1.4$ GeV/c; or
2 μ $p_T(\mu_1) > 0.56$ GeV/c,
 $p_T(\mu_2) > 0.48$ GeV/c
typical trigger requirement
 - $M(\mu^+\mu^-) > 2.9$ GeV/c
 - good fit quality of track & vertex
- The fitting is repeated for each p_T and y bin to obtain differential cross-section.



Fitting functions:
Crystal Ball for $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$;
Exponential background.

$$N(1S) = 26410 \pm 212$$

$$N(2S) = 6726 \pm 142$$

$$N(3S) = 3260 \pm 112$$



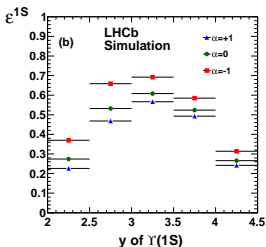
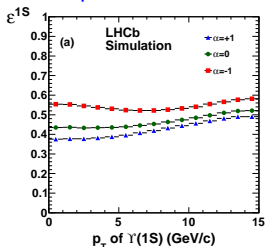
Integrated Cross-sections

$$\sigma(pp \rightarrow \Upsilon(1S)X) \times \mathcal{B}^{1S} = 2.29 \pm 0.01 \pm 0.10_{-0.37}^{+0.19} \text{ nb},$$

$$\sigma(pp \rightarrow \Upsilon(2S)X) \times \mathcal{B}^{2S} = 0.562 \pm 0.007 \pm 0.023_{-0.092}^{+0.048} \text{ nb},$$

$$\sigma(pp \rightarrow \Upsilon(3S)X) \times \mathcal{B}^{3S} = 0.283 \pm 0.005 \pm 0.012_{-0.048}^{+0.025} \text{ nb}$$

- Uncertainties: statistical, systematics and from the **unknown Υ polarisation**.



α : polarisation parameter
+1 fully transverse
-1 fully longitudinal
0 unpolarised

- A factor of 3 smaller than the ATLAS and CMS results, though different rapidity ranges are covered (ATLAS: $|y| < 2.4$; CMS: $|y| < 2.0$; LHCb: $2.0 < y < 4.5$).



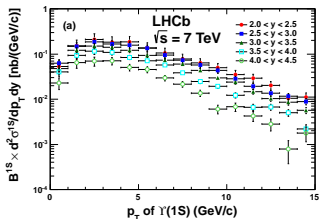
Determination of Differential Cross-section

$$\frac{d^2\sigma^{nS}}{dp_T dy} \times \mathcal{B}^{nS} = \frac{N^{nS}}{\mathcal{L} \times \varepsilon^{nS} \times \Delta y \times \Delta p_T}, n = 1, 2, 3.$$

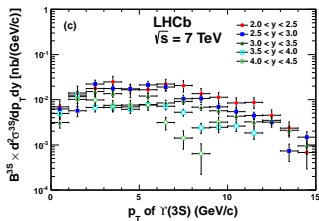
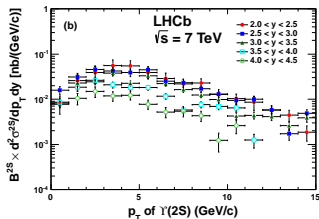
- N^{nS} : signal yield from fits
- \mathcal{L} : luminosity with uncertainty of 3.5%
- ε^{nS} : efficiencies estimated from Monte Carlo
- $\Delta y, \Delta p_T$: $\Delta y = 0.5, \Delta p_T = 1 \text{ GeV}/c$ for $2.0 < y < 4.5, p_T < 15 \text{ GeV}/c$



Results - Differential Cross-sections

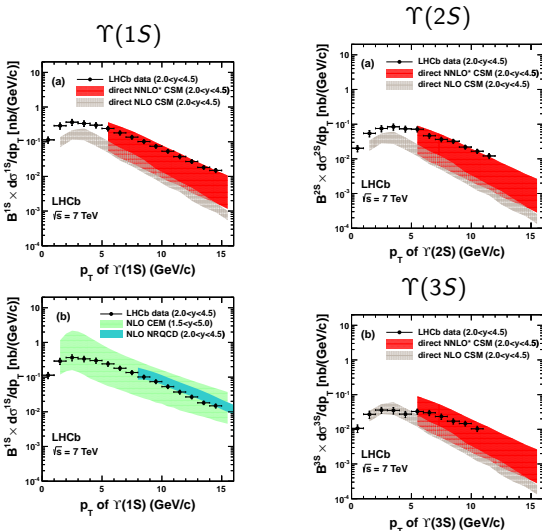


The first measurement of Υ production in the forward region at $\sqrt{s} = 7$ TeV!





Comparison with Theory



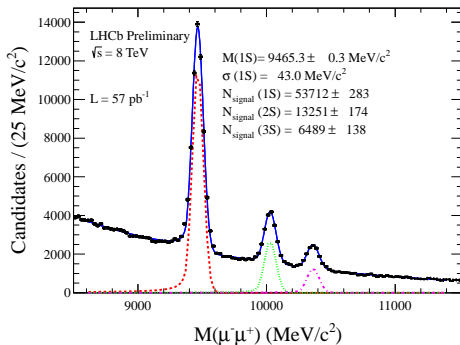
- General agreement with NNLO CSM, NLO NRQCD and NLO CEM
- NNLO CSM only calculate the direct production - feed-down information is important

NNLO CSM: [Artoisenet et al, Phys.Rev.Lett.101,152001\(2008\)](#)
 NRQCD: [Ma et al, Phys.Rev.Lett.106,042002\(2011\)](#)
 NLO CEM: [Frawley et al, Phys.Rep.462,125\(2008\)](#)



$\Upsilon(nS)$ with 8 TeV Data

[LHCb-CONF-2012-025](#)



- 57 pb^{-1} of early 2012, $\sqrt{s} = 8 \text{ TeV}$
- Fitted resolution: $43 \text{ MeV}/c^2$ ($47 \text{ MeV}/c^2$ in 2011)

Study on $\Upsilon(nS)$ polarisation is ongoing with the full 2011+2012 datasets.

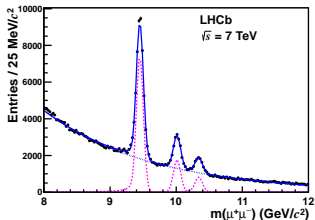
$\chi_b(1P)$ Production

[JHEP 11\(2012\),31](#)

In order to interpretate the cross-section and polarisation of bottomonia data, we need to better understand the Υ feed-down from higher states.

- $\chi_b(1P) \rightarrow \Upsilon(1S)\gamma \rightarrow \mu^+\mu^-\gamma$
- 32 pb⁻¹ data from 2010, $\sqrt{s} = 7$ TeV
- $2.0 < y^{1S} < 4.5$, $6 < p_T < 15$ GeV/c

$\Upsilon(1S)$ candidates:



$$N_{\Upsilon(1S)} = 39635 \pm 252$$

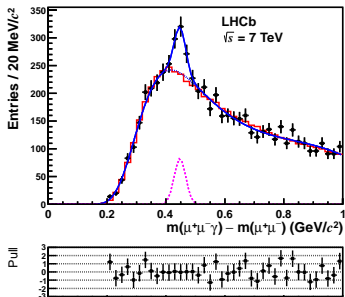
$$M_{\Upsilon(1S)} = 9449.2 \pm 0.4 \text{ MeV}/c^2$$

$$\sigma_{\Upsilon(1S)} = 51.7 \pm 0.4 \text{ MeV}/c^2$$

Reconstruction of χ_b from $\Upsilon(1S)$:

- $p_T^{\Upsilon(1S)} > 6$ GeV/c
- $9.36 < M(\mu\mu) < 9.56$ GeV/c²
- $p_T^\gamma > 0.6$ GeV/c²
- $\cos\theta_\gamma^* > 0$, θ_γ^* is polar angle of γ in χ_b rest frame

Reconstructed $\chi_b(1P)$ Signal



$$\frac{dN}{dx} = A_1 \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-\Delta M)^2}{2\sigma^2}} + A_2(x-x_0)^\alpha e^{-(c_1x+c_2x^2+c_3x^3)},$$

Gaussian signal + background

$$N_{\chi_b(1P)} = 201 \pm 55$$

$$\Delta M = 447 \pm 4 \text{ MeV}/c^2$$

$$\sigma = 19.0 \pm 4.2 \text{ MeV}/c^2$$

- Clear $\chi_b(1P)$ signal
- Yet no hint for higher state ($\Delta M(\chi_b(2P)) \sim 800 \text{ MeV}/c^2$)
- $\chi_{bJ}(J = 0, 1, 2)$ mass difference is comparable with resolution, so the observed signal is a mixture ($\Delta M_{2-1} \sim 20 \text{ MeV}/c^2$, $\Delta M_{1-0} \sim 33 \text{ MeV}/c^2$).



Determination of $\chi_b(1P)$ Yield

The fraction of $\Upsilon(1S)$ originating from $\chi_b(1P)$:

$$f_{\chi_b \rightarrow \Upsilon} = \frac{N_{\text{rec}}(\chi_b)}{N_{\text{rec}}(\Upsilon)} \cdot \varepsilon_{\text{cond}}(\chi_b)$$

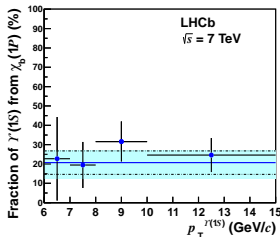
where $\varepsilon_{\text{cond}}(\chi_b)$ is the reconstruction efficiency after $\Upsilon(1S)$ is selected
 $\Rightarrow \varepsilon_{\text{cond}}(\chi_b) = (9.4 \pm 0.1)\%$

$$f_{\chi_b \rightarrow \Upsilon} = (20.7 \pm 5.7(\text{stat.}) \pm 2.1(\text{syst.})_{-5.4}^{+2.7}(\text{pol.}))\%$$

Comparing with CDF result
 $(\sqrt{s} = 1.8 \text{ TeV}, p_T > 8 \text{ GeV}/c)$:

$$f_{\chi_b \rightarrow \Upsilon} = (27.1 \pm 6.9 \pm 4.4)\%$$

CDF, Phys.Rev.Lett.84(2000)2094



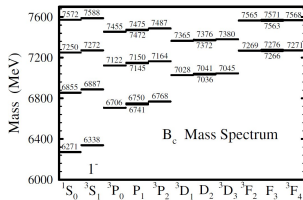
The ratio shows no dependence on p_T



B_c^+ Studies

- B_c : unique meson family formed by two different heavy flavoured quarks
- Mass prediction: $6.2 - 6.4 \text{ GeV}/c^2$ by potential models and Lattice QCD
- Produced mainly via gluon-gluon fusion at LHC, $\sigma \sim 0.4 \mu\text{b}$, one order of magnitude higher than Tevatron ($\sqrt{s} = 1.96 \text{ TeV}$)
- Decay modes: \bar{b} or c decay, annihilation; yet only a few channels observed experimentally.

- ✓ B_c^+ Production and Mass Measurement with $B_c^+ \rightarrow J/\psi \pi^+$
- ✓ Observation of $B_c^+ \rightarrow J/\psi \pi^+ \pi^+ \pi^-$ decay



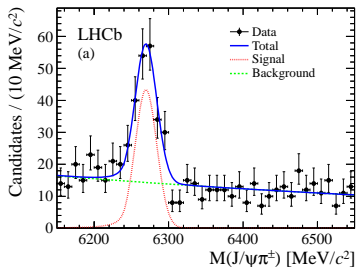
Godfrey, Phys.Rev.D 70,054017(2004)



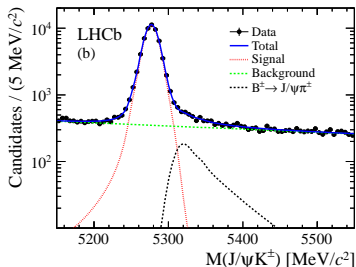
B_c^+ Production and Mass Measurement

[arXiv:1209.5634](https://arxiv.org/abs/1209.5634)

- $B_c^+ \rightarrow J/\psi\pi^+$, $J/\psi \rightarrow \mu^+\mu^-$, control channel $B^+ \rightarrow J/\psi K^+$
- Based on 0.37 fb^{-1} 2011 data ($\sqrt{s} = 7 \text{ TeV}$), $p_T > 4 \text{ GeV}/c$, $2.5 < \eta < 4.5$



$$N = 162 \pm 18$$



$$N = 56243 \pm 256$$

- Signal fitted with double-sided Crystal Ball function
- Cabbibo suppressed background $B^+ \rightarrow J/\psi\pi^+$ for $B^+ \rightarrow J/\psi K^+$ is considered



B_c Production Cross-section

$$\frac{\sigma(pp \rightarrow B_c^+ X) \cdot \mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)}{\sigma(pp \rightarrow B^+ X) \cdot \mathcal{B}(B^+ \rightarrow J/\psi K^+)}$$

$$= (0.68 \pm 0.10(\text{stat.}) \pm 0.03(\text{syst.}) \pm 0.05(\text{lifetime}))\%$$

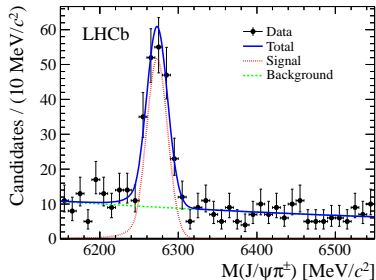
- Efficiencies are calculated in individual (p_T, η) bin to avoid bias
- Systematics uncertainty is dominated by the contribution from B_c^+ lifetime (0.453 ± 0.041) ps, reducible with a better lifetime measurement



B_c Mass Measurement

To measure $M(B_c^+) - M(B^+)$, invariant masses of $J/\psi\pi^+$ and $J/\psi K^+$ are fitted simultaneously

⇒ systematic uncertainty is reduced.



$$M(B_c^+) = 6273.7 \pm 1.3(\text{stat.}) \pm 1.6(\text{syst.}) \text{ MeV}/c^2$$

$$M(B_c^+) - M(B^+) = 994.6 \pm 1.3(\text{stat.}) \pm 0.6(\text{syst.}) \text{ MeV}/c^2$$

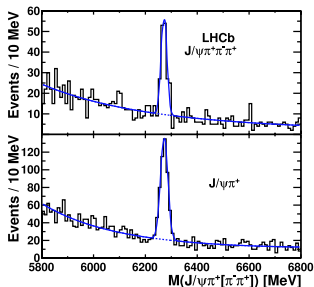
World best measurement!



Observation of $B_c^+ \rightarrow J/\psi \pi^+ \pi^+ \pi^-$

[Phys.Rev.Lett.108,251802\(2012\)](#)

- 0.8 fb^{-1} data from 2011
- $N(B_c^+ \rightarrow J/\psi \pi^+ \pi^+ \pi^-) = 135 \pm 14$
- **First observation!**
- $N(B_c^+ \rightarrow J/\psi \pi^+) = 414 \pm 25$



$$\frac{\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+ \pi^+ \pi^-)}{\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)} = 2.41 \pm 0.30(\text{stat.}) \pm 0.33(\text{syst.})$$

The main systematic uncertainty: model dependence of efficiency

Consistent with theoretical predictions 1.9 – 2.3 [Likhoded and Luchinsky, Phys.Rev.D](#)

[81,014015\(2010\)](#)



Ongoing Studies on B_c^+

- Lifetime measurement
 - $B_c^+ \rightarrow J/\psi(\mu\mu)\mu^+\nu$ and $B_c^+ \rightarrow J/\psi\pi^+$
- Search for new decay channels:
 - $B_c^+ \rightarrow B_s\pi^+, J/\psi K^+, \psi(2S)\pi^+, J/\psi D_s^+, \dots$
- Search for excited states:
 - $B_c^{*+} \rightarrow B_c^+\gamma, B_c(2S) \rightarrow B_c^{*+}\pi^+\pi^-$

...



Summary

- LHCb brings great opportunities for bottonomia and B_c studies
 - Inclusive $\Upsilon(nS)$ differential cross-section measured, and updated with $\sqrt{s} = 8$ TeV data;
 - $\Upsilon(1S)$ feed-down from $\chi_b(1P)$ measured; New state of $\chi_b(3P)$ observed;
 - B_c^+ mass measured to the world best precision using $B_c^+ \rightarrow J/\psi\pi^+$, and its production measured w.r.t $B^+ \rightarrow J/\psi K^+$;
 - $B_c^+ \rightarrow J/\psi\pi^+\pi^+\pi^-$ first observed;
- ... and many more going on!