

The pixel tracking telescope at the Fermilab Test Beam Facility

Simon Kwan^a, CM Lei^a, Dario Menasce^b, Luigi Moroni^b, Jennifer Ngadiuba^b, Alan Prosser^a, Ryan Rivera^a, Stefano Terzo^b, Marcos Turqueti¹^a, Lorenzo Uplegger^a, Luigi Vigani^b, Mauro E. Dinardo^b,

a Fermi National Accelerator Laboratory, Batavia, IL, USA

b Istituto Nazionale di Fisica Nucleare, Sezione di Milano Bicocca, and Università degli Studi di Milano Bicocca, Piazza della Scienza 3, 20126 Milano, Italy

E-mail: uplegger@fnal.gov

Abstract

An all silicon pixel telescope has been assembled and used at the Fermilab Test Beam Facility (FTBF) since 2009 to provide precise tracking information for different test beam experiments with a wide range of Detectors Under Test (DUTs) requiring high resolution measurement of the track impact point. The telescope is based on CMS pixel modules left over from the CMS forward pixel production. Eight planes are arranged to achieve a resolution of less than 8 μm on the 120 GeV proton beam transverse coordinate at the DUT position. In order to achieve such resolution with $100 \times 150 \mu\text{m}^2$ pixel cells, the planes were tilted to 25 degrees to maximize charge sharing between pixels. Crucial for obtaining this performance is the alignment software, called Monicelli, specifically designed and optimized for this system. This paper will describe the telescope hardware, the data acquisition system and the alignment software constituting this particle tracking system for test beam users.

1 Introduction

The Fermilab Test Beam Facility (FTBF) at the Fermi National Accelerator Laboratory provides beam in a multitude of particle types and a range of energies with which users can test their detectors [1]. The beam is resonantly extracted in a slow spill for each Main Injector cycle delivering a single 4.2 second long spill per minute. The primary beam (bunched at 53 MHz) consists of high energy protons (120 GeV) at variable intensities between 1 and 300 kHz. This beam can also be targeted to create secondary particle beams of pions, muons or electrons with energies down to about 1 GeV. Users have access to the facility instrumentation to measure the position and energy of the incident beam. Four pre-installed scintillation counters give rough beam position, a lead glass calorimeters measure the beam energy to a precision of $\sim 3\%$, two time-of-flight detectors can be set up for particle identification, and finally a silicon pixel telescope provides a precision position measurement of less than 8 μm using the primary 120 GeV proton beam. The silicon pixel telescope will

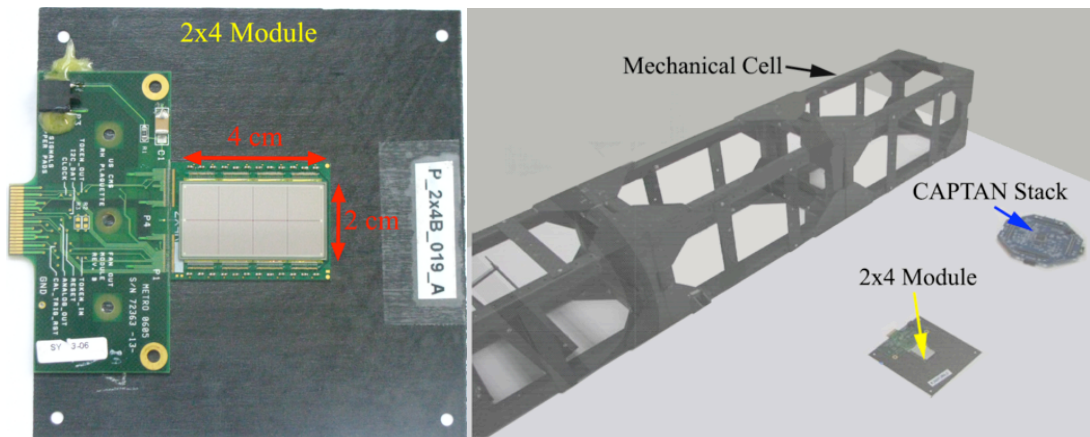
¹ *Now at Lawrence Berkeley National Laboratory, Berkeley, CA, USA*

43 be described in this paper along with its data acquisition (DAQ) system and alignment
44 software.

45 This facility adds to other operating test-beam facilities equipped with high precision
46 tracking telescopes as those of the AIDA project in Europe [2] [3]. The main features
47 of the present one are: a high-energy bunched proton beam (limited multiple Coulomb
48 scattering effects at 120 GeV and precise time localization of the incoming protons), a
49 narrow event time-window (within ~ 38 nsec), fast readout system capable to
50 accumulate up to 600,000 beam tracks per minute during the 4.2 second spill, and
51 beam tracking with better than $8 \mu\text{m}$ precision at the DUT.

52 2 The Pixel Telescope

53 A CMS pixel-based telescope has been built to provide precision tracking information
54 to any test beam experiment at the FTBF that requires the particle impact point on
55 their detector under tests (DUTs) with accurate precision. The telescope is placed
56 along the FTBF beam line and consists of eight detector planes - each made of
57 modules left over from the CMS Forward Pixel detector production - mounted on a
58 carbon fiber frame. Four of the eight telescope planes use modules composed of six
59 (2×3) PSI46V2 Read Out Chips (ROCs) [4], while the remaining four planes are
60 equipped with 2×4 modules. One of the 2×4 modules is shown in Figure 1, left side,
61 along with a picture, on the right, of the carbon fiber mechanical structure.
62



63
64 **Figure 1. Left: CMS PSI46V2 module. Right: telescope mechanical structure together with its**
65 **components.**
66

67 Each ROC reads an array of 52×80 pixel cells, where each pixel cell size is $100 \times$
68 $150 \times 285 \mu\text{m}^3$ (except for the pixels at the edge columns and upper row which have
69 sizes $100 \times 300 \times 285 \mu\text{m}^3$ and $200 \times 150 \times 285 \mu\text{m}^3$, respectively), for a total square
70 active area of $0.81 \times 0.81 \text{ cm}^2$. The total active area for a 2×4 module is then $1.62 \times$
71 3.24 cm^2 while the 2×3 module total active area is $1.62 \times 2.43 \text{ cm}^2$. The eight planes
72 are arranged in two stations with a section available in between for the DUT. The
73 DUT is considered as a separate third station itself.

74 A three-dimensional schematic view of the telescope detector planes is presented in
75 Figure 2, where the laboratory coordinate system is also indicated. In this reference
76 frame the Z axis is along the beam direction with +Z pointing downstream.

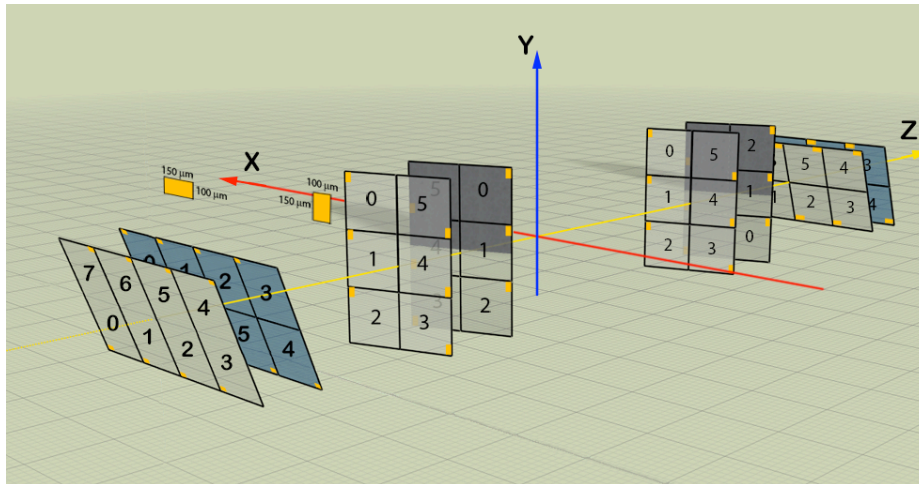
77 In order to improve the spatial resolution the planes are tilted by 25 degrees to
78 maximize the clusters of two adjacent hits, in particular:

79

- 80 • Four planes (2×4 detectors) tilted around the X axis with the long pixel side
- 81 oriented in the X direction;
- 82 • Four planes (2×3 detectors) tilted around the Y axis with the long pixel side
- 83 oriented in the Y direction.

84
85
86
87
88

In the first case, the most precise measurement is in the Y coordinate, while in the second case the most precise measurement is in the X coordinate. This geometry gives a total overlap active area of $\sim 1.6 \times 1.6 \text{ cm}^2$.

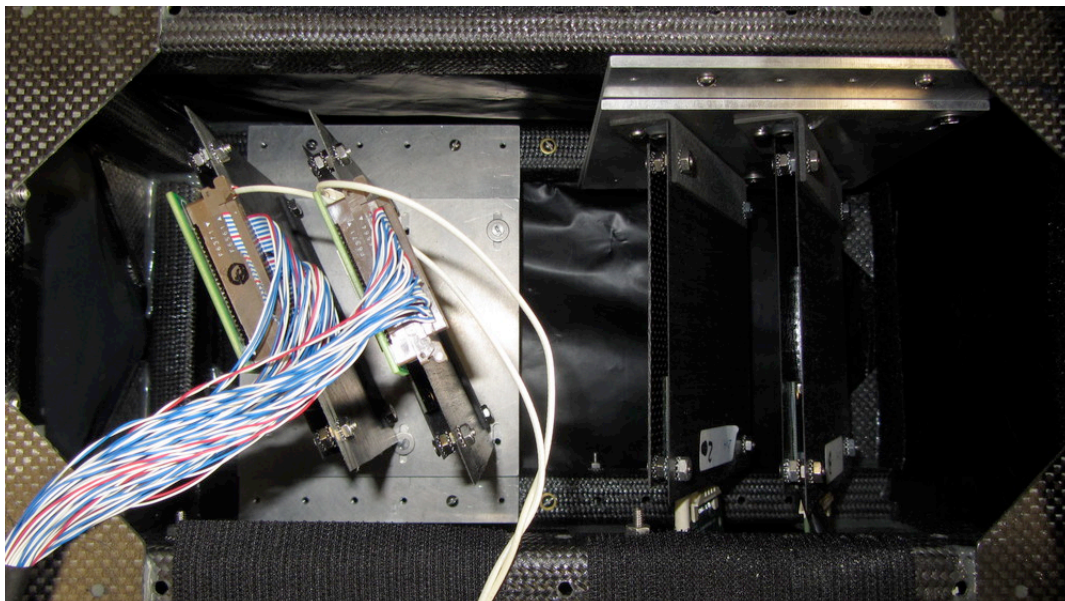


89
90

Figure 2. Three-dimensional schematic view of the pixel telescope geometry.

91
92
93
94
95
96
97

Each detector has a carbon fiber layer with the module glued on it. Heat is passively dissipated through the carbon fiber and hence no active cooling is required. For each station, the detectors are grouped into two pairs and each pair is attached on a light aluminum mechanical support mounted on the mechanical structure of the telescope, as shown in Figure 3.

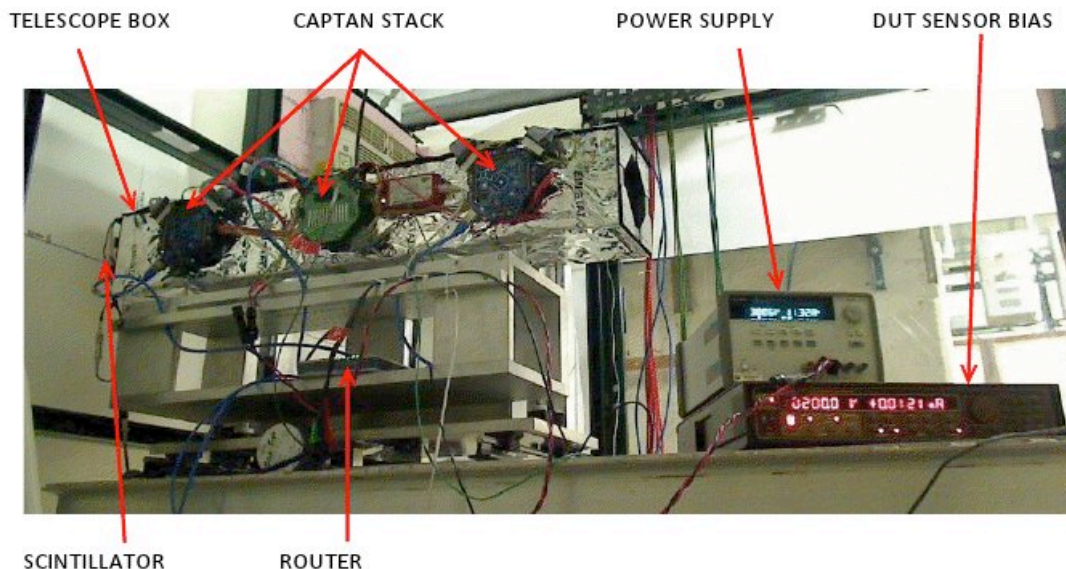


98
99
100

Figure 3. Modules mounted on the aluminum support at 25 degrees inside the carbon fiber frame.

101 The mechanical structure has a modular design: the frame consists of three basic cells,

102 one for each station, built with carbon fiber tubes of dimension $17.0 \times 17.0 \times 34.0$
103 cm^3 . The frame is covered by a Mylar anti-static layer, as can be seen in Figure 4,
104 which also serves the secondary purpose of keeping the detectors dark.
105



106
107 **Figure 4. The CAPTAN pixel telescope complete assembly in situ at the FTBF.**
108

109 The beam particle is triggered by a coincidence signal generated from three
110 scintillation counters placed behind the telescope. The trigger signal opens a small
111 time window in which the data acquisition system collects and sends data from the
112 detectors to the computer. The data from each ROC are tagged with the trigger count,
113 so that time correlated data recorded from all the detectors have the same tag. The hits
114 that share the same trigger count are identified as an event. An event may contain hit
115 data associated with one or more particle tracks passing through the telescope.
116 To keep the ROCs synchronized with the particle beam, the accelerator clock signal is
117 fed into one of the stations and then redistributed to the other stations through SATA
118 cables. Since the CMS ROC has been designed to run at a maximum frequency of 40
119 MHz while the Main Injector accelerator frequency is 53 MHz, the clock that is
120 distributed to the stations runs at 26.5 MHz, half of the Main Injector frequency,
121 allowing the detector to work properly and remain synchronized with the beam.

122 **3 The Data Acquisition System**

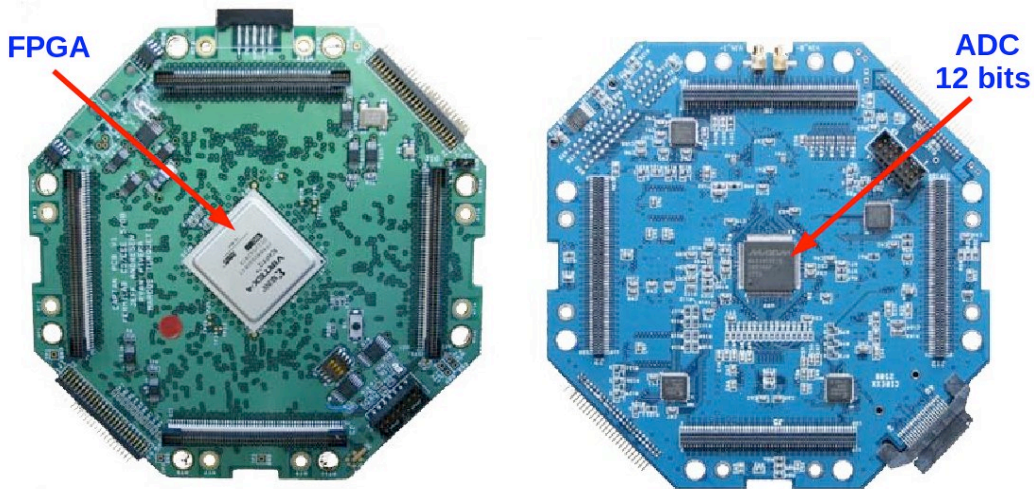
123 **3.1 The DAQ: hardware**

124 The Data Acquisition (DAQ) hardware is based on the CAPTAN system developed at
125 Fermilab [5]. The CAPTAN (Compact And Programmable daTa Acquisition Node) is
126 a flexible and versatile data acquisition system designed to meet the readout and
127 control demands of a variety of pixel and strip detectors for high energy physics
128 applications.

129 The system consists of three CAPTAN nodes, one for each station. The node is a
130 stack of different function boards with a vertical bus for high-speed data exchange
131 and features a Gigabit Ethernet Link (GEL) for high-speed communication through
132 the network.

133 In the present case, two boards are stacked in each CAPTAN node (Figure 5):

- 134 • The Node Processing and Control Board (NPCB) provided with a Virtex-4 FPGA.
- 135 • The Data Conversion Board (DCB) provided with a 12-bit MAX1438 Analogue
- 136 to Digital Converter.
- 137



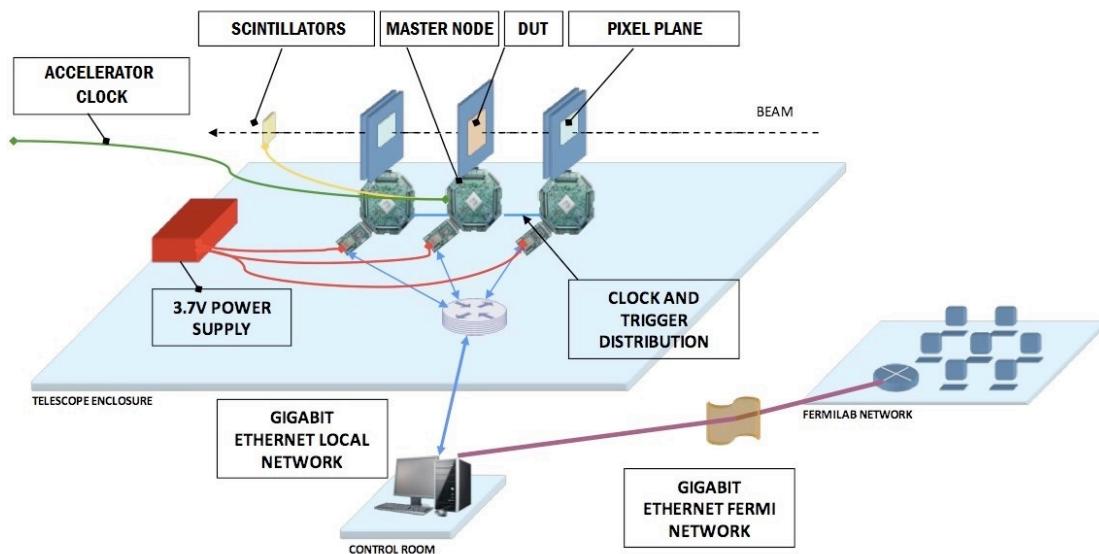
138
139 **Figure 5. Pictures of the primary boards of the CAPTAN stack used for telescope data**
140 **acquisition. Left: Node Processing and Control Board. Right: Data Conversion Board.**

141
142 Each NPCB is connected to a gigabit Ethernet router and the router is, in turn,
143 connected to the computer placed in the FTBF control room through an Ethernet
144 cable. This computer runs the Graphical User Interface, which controls the entire
145 telescope system.

146 The master CAPTAN node, which connects to the DUTs, receives the accelerator
147 clock and the trigger from the scintillators and redistributes them to the other two
148 nodes after having reduced the clock to half the frequency (26.5 MHz).

149 In each node, the data from ROCs are received by the DCB, digitalized by the ADC,
150 and then sent to the FPGA through the vertical bus. The formatted data are then
151 transferred to the control room PC via gigabit Ethernet. The diagram in Figure 6
152 shows the full readout system schematically.

153

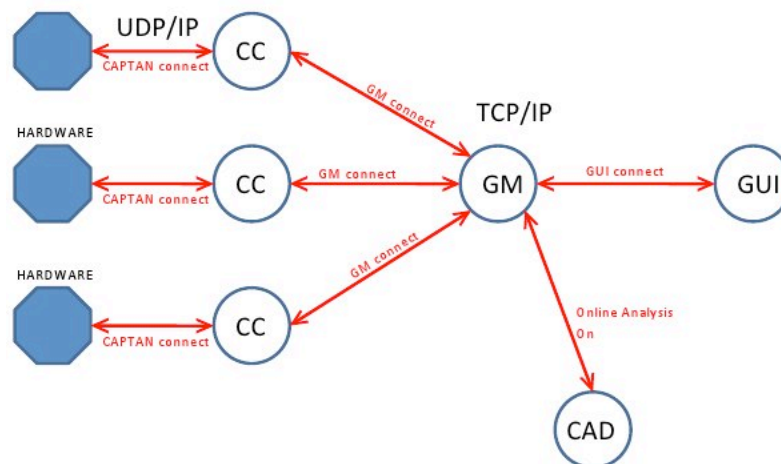


154
155 **Figure 6. Schematic of the telescope readout system based on CAPTAN hardware.**
156

157 At the control room PC, the data from each CAPTAN node are saved in a binary file
 158 for each Run. At the end of a Run, the CAPTAN Analysis and Display software,
 159 described in the next paragraph, scans the files from the different nodes and merges
 160 all temporally correlated data from every pixel plane (marked with the same trigger
 161 count) to form a representation of an event. A single binary file containing the merged
 162 events is then created for each Run and saved on disk. Each pixel data are stored in a
 163 64 bit word (8 bytes), so the typical size of a raw data event for a single telescope
 164 track with clusters made of two-pixel hits on all eight planes is only about 128 bytes.
 165 A typical Run with 400000 events at a beam-particle rate of ~25 kHz has an average
 166 ~0.9 reconstructed telescope-tracks per event and is of the order of 50 Mbytes
 167 including the hits on the DUT. Even at this low beam-particle rate, it requires only ~5
 168 spills (~5 minutes) to be accumulated. The system has been tested to work up to beam
 169 rates of ~150 kHz without any performance degradation, approaching the maximum
 170 trigger frequency that the readout chips can tolerate with low hit occupancy.

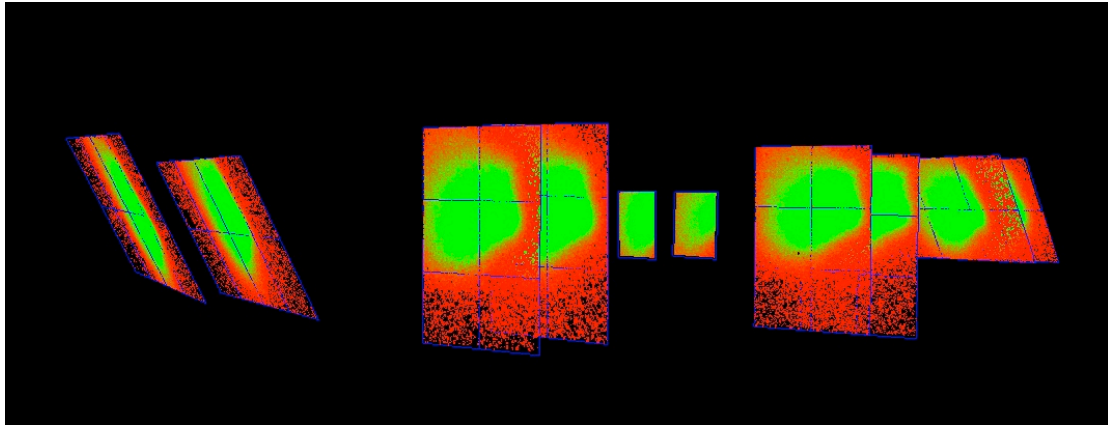
171 3.2 The DAQ: software

172 A complete software solution for interfacing with the CAPTAN system has been
 173 designed for Microsoft Windows using Microsoft Visual C++ 2008 [6].
 174 The diagram in Figure 7 shows the CAPTAN software topology conceptually. The
 175 building blocks of the software are the Global Master (GM), the CAPTAN Controller
 176 (CC), a Graphical User Interface (GUI) and the CAPTAN Analysis and Display
 177 (CAD).
 178



179
 180 **Figure 7. Schematic of the telescope readout system based on CAPTAN hardware.**
 181

182 The GM is the server for the entire system. It handles the communications among the
 183 CCs, the GUI and the CAD. The GUI allows the user to set up the readout chips,
 184 trigger and clock system, run calibration procedures, as well as to start or stop the data
 185 acquisition. It also controls when the CC stores data to disk. Another crucial block of
 186 the software is represented by the CAD, which allows the user to immediately
 187 visualize the merged telescope data in three dimensions as shown in Figure 8. This is
 188 critical for the data acquisition process since it allows the user to quickly check the
 189 quality of the data. For instance, this feature makes it possible to assess whether the
 190 beam is properly positioned on the detectors.
 191



192
193

Figure 8. Three-dimensional data visualization with the CAPTAN software.

194

4 Telescope Tracking and Alignment

195

196

197

198

199

200

201

202

203

204

205

206

207

208

209

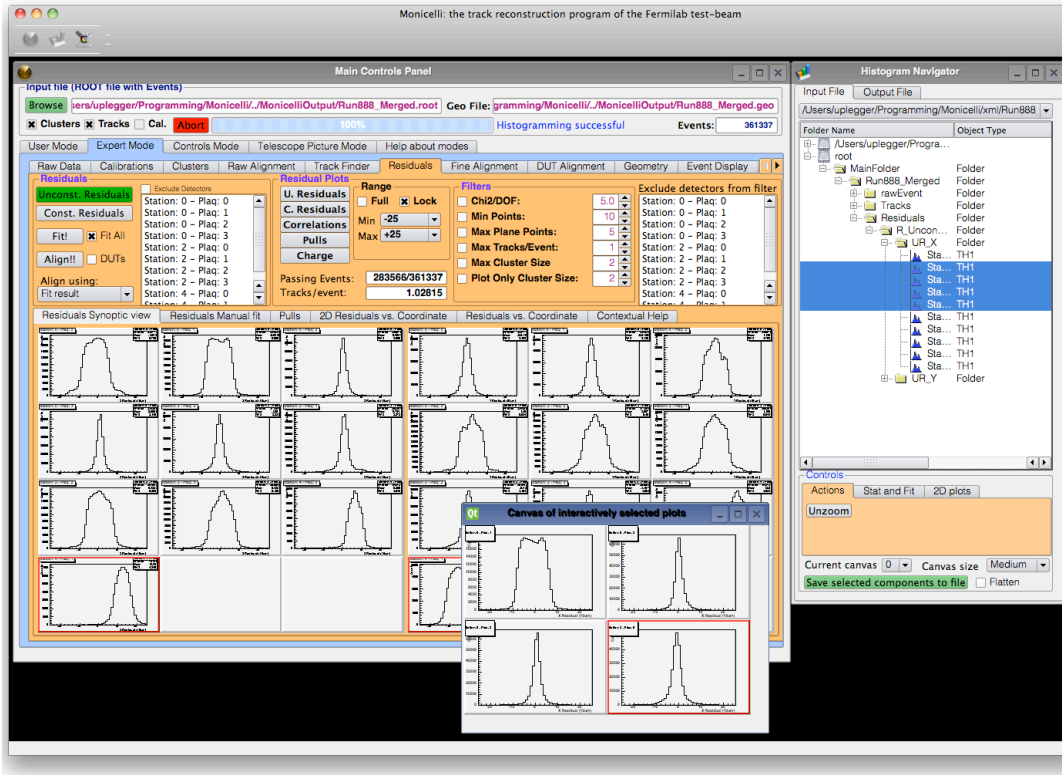
210

211

212

The main goal of the test beam experiments is to probe the performance of the DUT with a set of well reconstructed beam tracks. For this purpose the telescope has been designed to achieve an optimal resolution on the coordinates of the track impact point on the DUT, which is placed at the center of the telescope. The track reconstruction and the telescope alignment are performed by means of the C++ Monicelli application, developed by the INFN Milano-Bicocca group. This software provides the user with an appropriate iterative procedure to converge toward the optimal alignment of the telescope. All the operations can be accomplished in steps through a graphical interface (GUI) that makes the software user-friendly. For debugging, the GUI allows the user to perform the individual steps of the alignment, change critical parameters interactively and finally monitor the partial results by means of specific histograms and distributions. The software is also equipped with efficient interactive tools allowing the user to browse, examine, print and save these distributions in real-time. Figure 9 shows a snapshot of GUI and its components.

The next paragraphs give a quick overview of the main features of the track reconstruction and alignment procedure. It is worth noting that all the quoted results are for track having hits on all eight planes.



213
214
215
216
217
218

Figure 9. Monicelli Graphical User Interface. Inside the desktop style GUI, the main panel on the top left shows controls for all alignment operations and on the right the histogram browser used to navigate across all control histograms created during the alignment phases. Histograms can be viewed inside the control panel and also in separate canvases, on the bottom, opened by the histogram browser.

219 4.1 Track Reconstruction

220 The track-reconstruction code implemented in Monicelli was developed for the
221 purpose of this pixel telescope. It performs a simple and fast χ^2 straight line fit to the
222 coordinates of the arrays of hits, which were preselected by a pattern recognition
223 based on their alignment. In case of clusters of adjacent hits, the coordinates are
224 linearly interpolated over the expected charge-sharing width. The errors attributed to
225 the coordinates were calibrated during the first commissioning phase of the telescope.
226 They were initially estimated on the basis of what one would a priori expect for the
227 different cluster topologies. Once the alignment was performed they have been
228 refined, by an iterative procedure, which continued until the unconstrained pulls on
229 each telescope plane became Gaussian with nearly unitary RMS. The unconstrained
230 (or unbiased) residuals and pulls of a track on a plane are calculated excluding the hits
231 on that particular plane from the track fit.

232 To perform this track reconstruction process, Monicelli reads the merged binary data
233 file together with an XML geometry file describing the overall configuration and
234 geometrical details of the telescope detectors for that particular data set. The user can
235 change this file by means of a graphical XML editor included in the package. The
236 geometrical details of each plane are set specifying its space coordinates and rotation
237 angles in the laboratory frame, as well as the number and orientation of ROCs and, for
238 each ROC, the number of columns and rows together with their pitch.

239 The Monicelli output file consists in a ROOT TTree containing, for each event, the
240 reconstructed telescope tracks together with the associated clusters of hits and the raw
241 data, including those of the DUT. This file provides the user with all the information

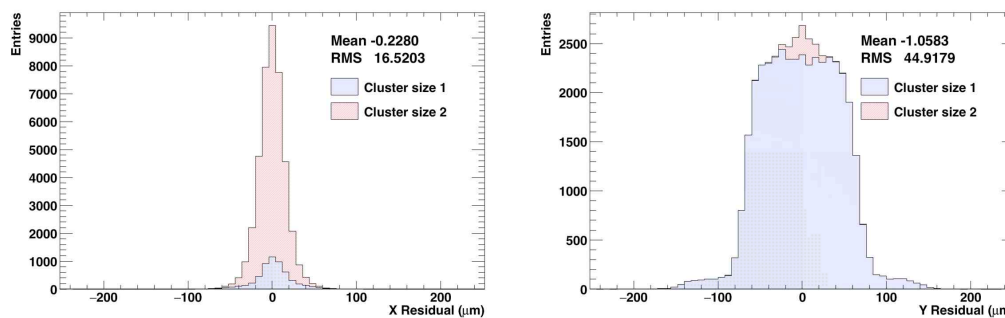
242 needed for the analysis of the test-beam data. A general-purpose analysis program is
243 also available and can be easily adapted to the particular features of the DUT.

244 4.2 Alignment

245 The alignment code was expressly developed for this telescope and it is based on a
246 standard χ^2 successive-approximation minimization of the unconstrained residuals on
247 each plane. It returns all the six geometrical degrees of freedom of each plane. Only
248 the tracks having hits on all the planes are considered in this process. Typically, no
249 more than 10 iterations with 40000 events are required to converge to a reliable
250 solution.

251 Figure 10 shows the two unconstrained residuals on a detector plane after a complete
252 alignment. The narrow residual on the left is dominated by charge sharing between
253 adjacent pixels along the tilted coordinate that is measured with the best resolution
254 (pitch = 100 μm), resulting in distributions with RMS of $\sim 16.5 \mu\text{m}$. The other is
255 mainly driven by single-hit events along the non-tilted coordinate (pitch = 150 μm)
256 and, therefore, presents a RMS of about $150/\sqrt{12} \mu\text{m}$.

257

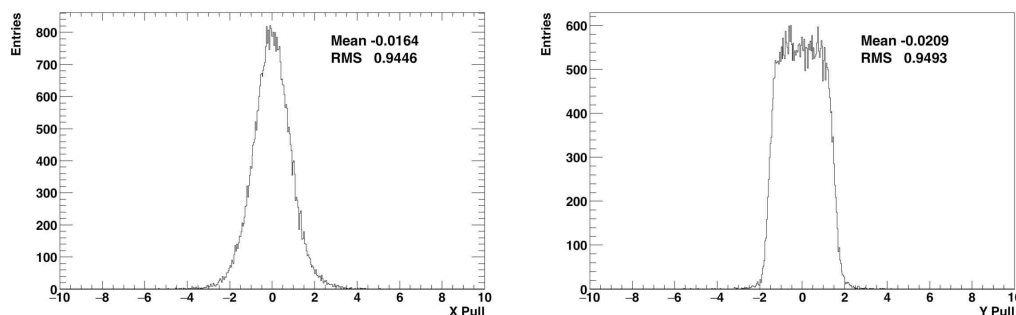


258
259
260

Figure 10. The x and y unconstrained residual distributions for one telescope detector after a complete alignment. x and y are the local coordinates of the plane.

261 Figure 11 shows the two unconstrained pull distributions on a detector plane after a
262 complete alignment of the telescope. The quasi unitarity of the x and y pull
263 distributions confirms both the precision of the alignment and the correct evaluation
264 of the errors associated to the different clusters of adjacent hits.

265

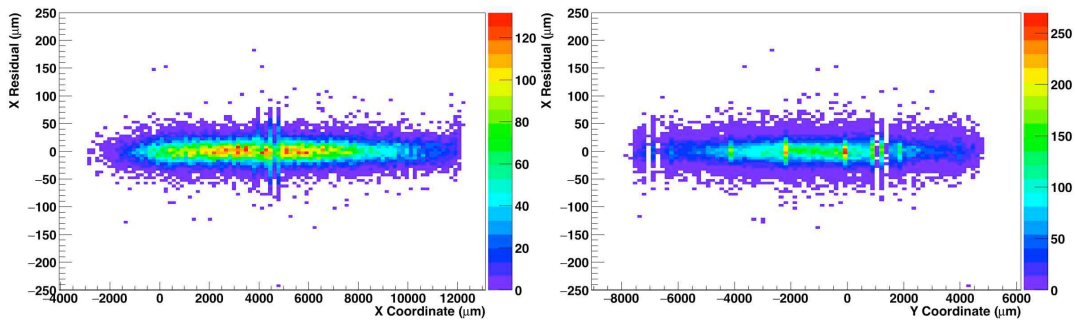


266
267

Figure 11. The x and y unconstrained pull distributions for one of the telescope detectors

268 The accuracy of the alignment is further investigated studying the plots of Figure 12,
269 showing the correlations between the unconstrained residuals and the impact point
270 coordinates on the detector. Any deviation from a flat distribution would signal a
271 residual problem in the alignment. As the alignment procedure goes ahead, the
272 correlation plots are flattened and resemble those shown in Figure 12.

273

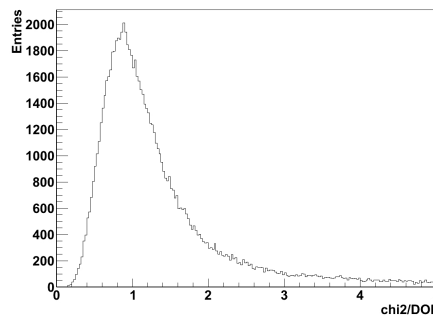


274
275
276

Figure 12. Plots of the correlation between unconstrained residuals and impact point coordinate for one of the telescope detectors.

277

The shape of the fitted track χ^2/DoF distribution is shown in Figure 13.

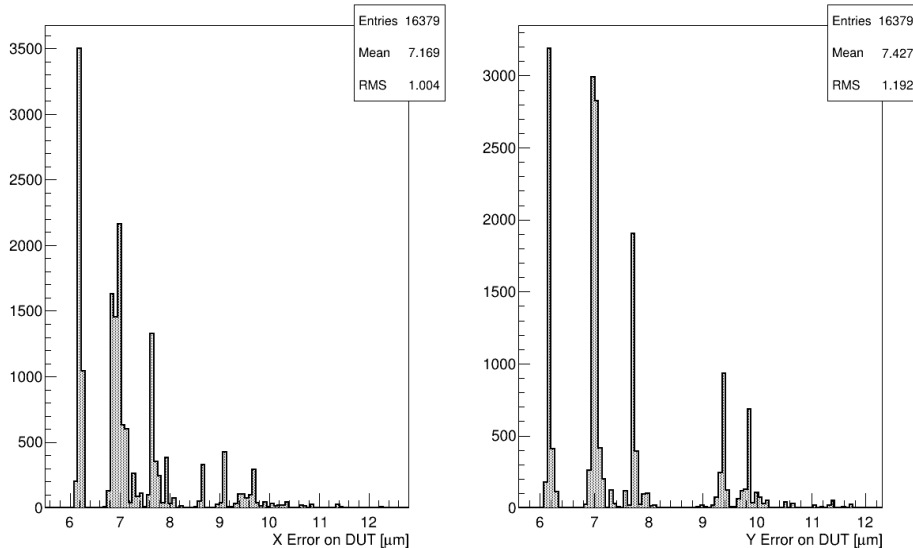


278
279
280

Figure 13. The track χ^2/DoF distribution after a complete alignment.

281
282

Finally, the transverse error on track fit extrapolation at DUT Z position ($Z \approx 0$) is plotted in Figure 14 (for tracks having hits on all eight planes).



283
284
285
286
287

Figure 14. Track fit extrapolated error distributions at the DUT Z position ($Z \approx 0$) after a complete alignment. The discrete peaks result from the combinations of different hit resolution that can be associated to the tracks.

288
289

It turns out that the best achievable telescope resolutions on the DUT is as small as $6.2 \mu\text{m}$ in both X and Y coordinates, and that the bulk of tracks give resolutions better

290 than 8 μm . Excluding tracks with extrapolation error larger than 8 μm , as it is
 291 typically done for data analysis, the average error at DUT is approximately 6.9 μm on
 292 both projections (6.87 μm in X and 6.94 μm in Y). This error will be called Pure
 293 Telescope extrapolation error.

294 To obtain the Total Expected extrapolation error (at 120 GeV proton energy), one has
 295 to add in quadrature the error due to the Multiple Coulomb Scattering (MCS) in the
 296 DUT. For three or less DUTs placed in the standard configuration around the
 297 telescope center at $Z_1 = +5$ cm, $Z_2 = -5$ cm and $Z_3 = -6$ cm, the Total Expected
 298 extrapolation error is given by
 299

$$300 \quad \sigma_X = \sqrt{6.87^2 + 1.37^2 \times L_{DUT_1} + 1.95^2 \times L_{DUT_2} + 2.01^2 \times L_{DUT_3}} \quad \mu\text{m}$$

301

$$302 \quad \sigma_Y = \sqrt{6.94^2 + 2.10^2 \times L_{DUT_1} + 2.68^2 \times L_{DUT_2} + 2.74^2 \times L_{DUT_3}} \quad \mu\text{m}$$

303

304

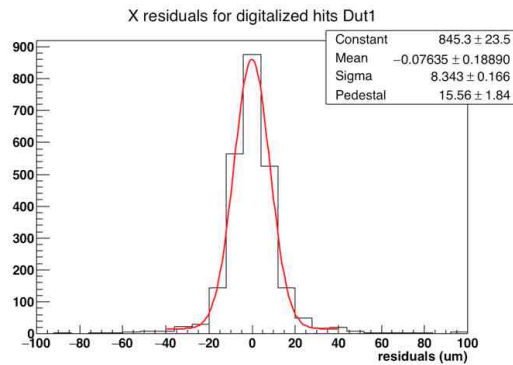
305 L_{DUT_n} is the thickness of the nth DUT in percent of a radiation length. The errors in
 306 the previous formulae are negligible ($< 0.1 \mu\text{m}$) and hence omitted.

307 **5 Performance in Beam Tests**

308 The test-beam performance of the telescope in testing a DUT and resolving its details
 309 was studied on several occasions during the numerous test-beam campaigns that were
 310 conducted on different pixel sensors. Here, only the most critical figure is reported
 311 and, namely, the actual measured track extrapolation precision in real experimental
 312 conditions.

313 As a general criterion for the analysis of the beam test results, any event must
 314 simultaneously fulfill the following requirements: only one telescope track is
 315 reconstructed; the track must have hits on all the planes; no hits other than those
 316 associated to the track can be present on the telescope planes; and the Pure Telescope
 317 extrapolation error at the DUT position is less than 8 μm in both views.

318 In Figure 15 the measured residual of a normally incident telescope track with respect
 319 to a doublet of adjacent hits on a silicon 3D pixel detector is reported. The coordinate
 320 attributed to the doublet was exactly the divide between the two adjacent pixels.
 321



322

323

324

Figure 15. Track residuals of two adjacent hits along X on the 3D Si-pixel detector.

325 The silicon 3D pixel sensor (Sintef 2E) had a $100 \times 150 \mu\text{m}^2$ pixel cell with two
326 central columnar signal-electrodes ($10 \mu\text{m}$ radius) and six columnar field-electrodes
327 ($7 \mu\text{m}$ radius) on the periphery. The sensor thickness was approximately $220 \mu\text{m}$.
328 The measured RMS of the residual, $8.34 \pm 0.17 \mu\text{m}$, results from the sum in quadrature
329 of the actual track extrapolation error and that of the coordinate measured on the
330 DUT. From a study of the correlation between the telescope track impact point and
331 the charge asymmetry of the two pixels of the doublet on the DUT, it is found that the
332 effective region of sharing is concentrated around the divide of the doublet with RMS
333 $= 2.38 \pm 0.60 \mu\text{m}$. Unfolding this error from the measured residual, the actual track
334 extrapolation error results $7.99 \pm 0.25 \mu\text{m}$ and it is fully consistent with a Total
335 Expected value of $7.86 \mu\text{m}$, resulting from the sum in quadrature of the Pure
336 Telescope resolution, $6.87 \mu\text{m}$, and a multiple Coulomb scattering error of $3.82 \mu\text{m}$,
337 due to the simultaneous presence of three DUTs ($\sim 1.5\%$ of a radiation length each).
338 On this basis one can conclude that the actual track extrapolation error at the DUT is
339 well matched by the Total Expected extrapolation error resulting from the formulae in
340 the previous section. For instance, the Total Expected extrapolation error for a typical
341 test of a single DUT (1.5% of a radiation length) is $7.07 \mu\text{m}$ in X and $7.40 \mu\text{m}$ in Y.
342 Obviously, this performance could be further improved by selecting only the best
343 resolution tracks and, specifically, those belonging to the first peak of Fig. 14 at 6.2
344 μm .
345 Summarizing, the present telescope allows for a real track extrapolation error at the
346 DUT typically better than $8 \mu\text{m}$ both in X and Y. This figure includes all experimental
347 effects that can affect the precision during a long Run of 24 hours at least such as
348 thermal excursion, vibrations and all the other possible phenomena.
349

350 **6 Conclusions**

351 This paper describes the pixel silicon telescope installed at the Fermilab Test
352 Beam Facility (FTBF) that provides high resolution tracking of the beam particles.
353 The impact point of the available 120 GeV beam protons at the center of the telescope
354 can be reconstructed with an accuracy of $8 \mu\text{m}$, at least, on both the transversal
355 coordinates. The detector hardware and control software is complemented by the
356 Monicelli software, which provides to the users the reconstructed tracks necessary for
357 their data analysis. The telescope has already been used by several experiments and,
358 in particular, has been extensively used by the CMS forward pixel community to test
359 sensor candidates for future CMS pixel upgrades [7-15].
360

361 **Acknowledgements**

362 We wish to thank the Fermilab Test Beam Facility personnel, and in particular Aria
363 Soha, Mandy Rominsky and Eugene Schmidt, for the continuous support they provide
364 us. This research was supported, in part, by the U.S. Department of Energy and the
365 Italian Istituto Nazionale di Fisica Nucleare and Ministero della Ricerca Scientifica e
366 Tecnologica.

367 We also wish to thank the LHC Physics Center (LPC) for the support given to our
368 students Stefano Terzo, Jennifer Ngadiuba and Luigi Vigani.

369 **References**

- 370 [1] Fermilab Test Beam Facility web site: <http://www-ppd.fnal.gov/FTBF>
- 371 [2] I.Rubinskiy: An EUDET/AIDA pixel beam telescope for detector development.
- 372 Physics Procedia 37 (2012) 923 – 931
- 373 [3] I.Rubinskiy: A EUDET/AIDA pixel beam telescope as a tool for testing tracking
- 374 detectors. Instrumentation Seminar at DESY (March 15, 2013)
- 375 [4] H. Chr. Kaestli et al.: Design and Performance of the CMS Pixel Detector
- 376 Readout Chip. Nucl.Instrum.Meth. A565 (2006) 188-194
- 377 [5] Marcos Turqueti, Ryan A. Rivera, Alan Prosser, Jeffry Andresen, and John
- 378 Chramowicz. CAPTAN: A hardware architecture for integrated data acquisition,
- 379 control, and analysis for detector development. In Nuclear Science Symposium
- 380 Conference Record, 2008. NSS '08. IEEE, pages 3546-3552.
- 381 [6] Ryan A. Rivera, Marcos Turqueti, and Alan Prosser. A software solution for the
- 382 control, acquisition, and storage of CAPTAN network topologies. In Nuclear
- 383 Science Symposium Conference Record, 2008. NSS '08. IEEE, pages 805-808.
- 384 [7] Mayur Bubna et al.: Laboratory and testbeam results for thin and epitaxial planar
- 385 sensors for HL-LHC JINST 10.1088/1748-0221/10/08/C08002 C08002 10 08
- 386 2015 2015-08-04.
- 387 [8] Mayur Bubna,et al.: Testbeam and Laboratory Characterization of CMS 3D Pixel
- 388 Sensors JINST 10.1088/1748-0221/9/07/C07019 C07019 9 2014 2014
- 389 [9] Luigi Moroni et al.: Tracking performance of a single-crystal and a polycrystalline
- 390 diamond pixel-detector JINST 10.1088/1748-0221/8/06/P06006 P06006 8 Solid
- 391 state detectors Diamond Detectors Radiation-hard detectors Particle tracking
- 392 detectors (Solid-state detectors) 2013 2013-06-01
- 393 [10] Alexander Krzywda et al.: Pre- and post-irradiation performance of FBK 3D
- 394 silicon pixel detectors for CMS Nucl.Instrum.Meth. 10.1016/j.nima.2014.06.029
- 395 404-411 A763.
- 396 [11] Enver Alagoz et al.: Pre- and post-irradiation performance of FBK 3D silicon
- 397 pixel detectors for CMS Nucl.Instrum.Meth. 10.1016/j.nima.2014.06.029 404-411
- 398 A763.
- 399 [12] Enver Alagoz et al. Testbeam and laboratory test results of irradiated 3D CMS
- 400 pixel detectors Nucl.Instrum.Meth. 10.1016/j.nima.2013.07.042 52-56 A732.
- 401 [13] Ada Solano et al.: Performance of CMS 3D silicon pixel detectors before and
- 402 after irradiation Nucl.Instrum.Meth. 10.1016/j.nima.2013.04.048 33-37 A730.
- 403 [14] Obertino Maria et al.: 3D-FBK pixel sensors with CMS readout: First test
- 404 results Nucl.Instrum.Meth. 10.1016/j.nima.2012.11.076 342-344 A718.
- 405 [15] Lorenzo Uplegger et al.: Test-beam studies of diamond sensors for SLHC
- 406 Nucl.Instrum.Meth. 10.1016/j.nima.2012.10.011 376-379 A718.