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Letter of Intent

Study of the  $(\bar{K}^0_{\pi^+ \pi^-})$ -system produced in the  
reaction  $K^- p \rightarrow \bar{K}^0_{\pi^+ \pi^-} n$  at 12 GeV/c using the  
Omega Prime Spectrometer and the RF-separated S1 beam.

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## 1. PHYSICS MOTIVATION

This experiment is to be regarded as the high statistics continuation of experiment S116 which used the  $\Omega$  - spectrometer and a PS beam. The statistics can be increased by almost two orders of magnitude by the use of the more intense SPS beam, the RF - separator and the MWPC's of the  $\Omega$  prime - detector.

These last years the main interest in the physics of hadrons has been the discovery of new quarks and their mesonic bound states. New theoretical ideas, like QCD, have some promise of predictive power. In this situation we feel that a detailed study of "old" mesons, formed by the u-, d- and s-quarks, can be extremely useful. In fact, the properties of many mesons, which are required to exist in quark-models, are not well established. Any meson with well established properties, or even the definite non-existence of such a particle, can be a severe test of theoretical models.

The proposed experiment

$$K^- p \rightarrow \bar{K}^0 \pi^+ \pi^- n \quad (1)$$

at  $\sim 12$  GeV/c is able to see all low mass  $K^*$ - resonances decaying into  $K \pi \pi$ . It is free of complications due to diffractive production. A detailed partial wave analysis (PWA) of the three-meson-system will give a precise determination of spin-parity  $J^P$  of these resonances. This is necessary since only the  $J^P$  of the  $2^+$  - state  $K^*(1440)$  is well established. All the other states, predicted by the quark model, have either not yet been observed or their resonance character and their quantum numbers are not well known up to now.

We believe that this experiment can answer all these questions for the  $K\pi\pi$ -mass region below 2 GeV, if it contains a sample of more than  $10^5$  events.

Apart from surprises, which are always possible in experiments that have a new sensitivity orders of magnitude higher, we intend to study the following subjects:

- a) From a study of diffractive ( $K\pi\pi$ )-systems, various groups find good evidence for the production of the two  $Q$ -resonances with  $J^P = 1^+$  at 1290 and 1400 MeV [1]. However due to the diffractive Deck background below the resonances the interpretation of the results is complicated. To establish these two resonances without any doubt, a study of the  $Q$ 's in non-diffractive channels is necessary. Figs. 1 and 2 show the results of a PWA of reaction (1) from the S116 experiment [2]. It is seen, that  $Q_2$  and probably also  $Q_1$  are produced in reaction (1) and can therefore be well studied in a high statistics experiment.
- b) Several groups report evidence for structure in the ( $K\pi$ ) and ( $K\pi\pi$ )-system in the mass region 1600 - 2000 MeV. The situation concerning this structure is still confused since the results differ in the values of mass and width of this enhancement (for a review see [3]) and it might be possible that more than one  $K^{*}$ -enhancement is responsible for this complex situation. In our PWA of reaction (1) we find one clear peak in the  $3^-$  state (with  $M = 1812 \pm 28$  MeV,  $\Gamma = 181 \pm 24$  MeV) which was interpreted as  $K^*(1780)$ .

A high statistics experiment of reaction (1) could clarify the situation concerning several  $K^{*}$ -resonances in this high ( $K\pi\pi$ ) mass region. In addition a reliable measurement of the relative phase of this state should be possible in order to establish the resonance nature of the peak.

- c) Meson resonances belonging to the natural spin-parity series ( $0^+, 1^-, 2^+, \dots$ ) can decay into two pseudoscalar particles and are therefore relatively easy to detect and to study. Resonances of the unnatural series ( $0^-, 1^+, \dots$ ) can only be found in a system of at

least three pseudoscalar mesons and are much harder to investigate. This explains the fact that  $0^+$ ,  $1^-$ ,  $2^+$ ,  $3^-$  and partly also  $4^+$  resonances predicted by the quark model are already well known while only some of the states of the unnatural series have been observed. Some of the  $1^+$  states were found recently, but none of the  $2^-$  resonances are clearly established. According to the quark model we expect the  $2^-$  states in the  $(K\pi\pi)$ -system to be in the region 1600 - 1800 MeV. As seen from fig. 1, they are certainly not produced with large cross-section. Therefore only high statistics at high energy can solve the problem. An experiment at low energy is less well suited to find these states since overlaps with other reaction mechanisms complicate the study of these small effects.

- d) A spectrometer experiment at SLAC investigated the  $(K\pi)$ -system and found evidence (in two out of four possible partial wave solutions) for a broad  $J^P = 1^-$  resonant-like structure at 1650 MeV [4]. This resonance, if it exists at all, is rather inelastic (elasticity of 0.3) and should therefore be found in reaction (1). Indeed a broad enhancement can be seen in fig. 1 for the  $1^-$ -wave around 1600 MeV which is compatible with the SLAC result. However, with the limited statistics of this experiment, no conclusive answer concerning the existence of  $K^*(1650)$  can be given.

Another interesting question concerning this  $K^*(1650)$  is, how to classify this resonance within the quark model. One possibility is that it is a radial excitation of  $K^*(890)$  (that means the S-wave  $(q\bar{q})$ -triplet) and the other possibility is that  $K^*(1650)$  is a D-wave  $(q\bar{q})$ -triplet. From the interference between the  $J^P = 3^-$   $K^*(1780)$  and  $K^*(1650)$ , it should be possible to solve this problem [5].

- e) In fig. 3 we show the  $(K\pi\pi)$ -system of reaction (1) produced by unnatural spin-parity exchange, decomposed in natural and unnatural spin-parity states. It is seen that it occurs mainly in the natural

series ( $1^-$ ,  $2^+$ ,  $3^-$ ) which points strongly to a  $\pi$ -exchange production mechanism <sup>+) . Reaction (1) is therefore well suited for a study of inelastic  $K\pi$ -scattering:  $K^- \pi^+ \rightarrow K^0 \pi^+ \pi^-$ . It is complementary to elastic  $K\pi$ -scattering and should help to solve the ambiguity problem in the elastic  $K\pi$ -scattering. In addition good determination of the branching ratios of  $K^*(1440)$  and  $K^*(1780)$  into  $K^* \pi$  and  $K\rho$  should be possible.</sup>

In conclusion, a high statistics experiment of reaction (1) could bring much progress in the  $K^*$ -spectroscopy and could solve many of the waiting problems in this field.

## 2. APPARATUS AND TRIGGER

We propose to study reaction (1) in the CERN- $\Omega$  prime-spectrometer, which will be equipped with wire chambers instead of the optical spark chambers used in the  $\Omega$ -detector. A version of  $\Omega$  prime with only forward chambers is envisaged, which is expected to be available by mid 1979.

The trigger will be a (2prong +  $V^0$ )-trigger, similar to the one used in the previous S116 experiment. It will consist of two multiwire proportional chambers, one of which has an active surface of  $19 \times 19 \text{ cm}^2$  and a wire spacing of 1 mm (MWPC1). This chamber will be placed about 5 cm downstream of the  $H_2$ -target. The second chamber (MWPC2) has a sensitive surface of  $96 \times 147 \text{ cm}^2$ , a wire spacing of 2 mm, and it will be placed about 55 cm downstream of the target. Both chambers exist and will be available from the current  $\Omega$ -equipment. The trigger condition will be: 2 charged tracks in MWPC1 and 4 charged tracks in MWPC2.

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<sup>+) We would like to remind, that elastic ( $\pi\pi$ ) and ( $K\pi$ ) scattering were investigated in reactions similar to reaction (1), i.e.  $\pi p \rightarrow (\pi\pi)n$  resp.  $Kp \rightarrow (K\pi)n$  [4,6].</sup>

In addition to this topology trigger a cylindrical veto counter (TS) surrounding the target will be used to further eliminate events not belonging to the (2prong +  $V^0$ )-topology. This target hodoscope is especially important to veto events from the 4 prong channels, which give the most serious background to reaction (1). We emphasize that the TS affects the (2prong +  $V^0$ )-events such that only events with the  $V^0$  decaying before MWPC1 are rejected, which should not fulfill the topology trigger condition anyway. The trigger set up is schematically drawn in fig. 4.

### 3. BEAM AND TARGET

We plan to use the RF-separated S1 beam with a flux of  $\sim 3 \cdot 10^5$   $K^-$ /burst, which means a total flux of  $\sim 1 \cdot 10^6$  particles/burst at the  $\Omega$  prime target. The beam momentum should be fixed at 12 GeV. The target will be a 30 cm long cylindrical  $H_2$ -target.

### 4. ACCEPTANCE AND RATES

For the above described trigger set up we have calculated the geometrical acceptance for reaction (1) as well as for the most important background channels which also fulfill the trigger condition. This was done by subjecting events from a  $K^-p$  bubble chamber experiment at 10 GeV/c to the simulated trigger system. Table 1 shows the results for geometrical acceptance and the expected trigger rates assuming a beam flux of  $3 \cdot 10^5$   $K^-$ /burst and the use of i) the (2prong +  $V^0$ )-topology trigger only, ii) the topology trigger and the TS veto counter.

For the following calculations we assume that a trigger rate of  $\leq 80$  trigger/burst will be realistic, which should be well manageable by

the  $\Omega$  prime detector. \*)

Assuming  $9 \cdot 10^3$  good bursts/day, a  $\Omega$  prime running efficiency of 75% and 6 days of running time, we shall record  $\sim 3,2 \cdot 10^6$  triggers, out of which  $\sim 2,5 \cdot 10^5$  events belong to reaction (1). This corresponds to a sensitivity of  $\sim 800$  events/ $\mu\text{b}$  for this channel.

Assuming a data reconstruction efficiency of 80%, we shall obtain  $\sim 200.000$  good  $n\bar{K}^0\pi^+\pi^-$  events after geometrical reconstruction. About 130.000 of these events will have a  $(K\pi\pi)$ -mass between 1.0 and 2.0 GeV and therefore will be useful for a partial wave analysis.

## 5. COMPUTING LOAD

The whole statistics will be analysed at the PDP10 computer at Aachen. However, testing and calibration runs will be necessary to be done at CERN, for which we request  $\sim 20$  hours of CDC 7600 C.P.time.

Assuming a total statistics of  $3,2 \cdot 10^6$  triggers, the estimated requirement for pattern recognition and geometrical reconstruction is the equivalent of  $\sim 130$  days of PDP10 computer time. The time available at this machine is such that data reconstruction will be performed within  $\sim 50$  weeks.

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\*) The background can be further reduced by using a  $\pi^0$ -veto counting system.

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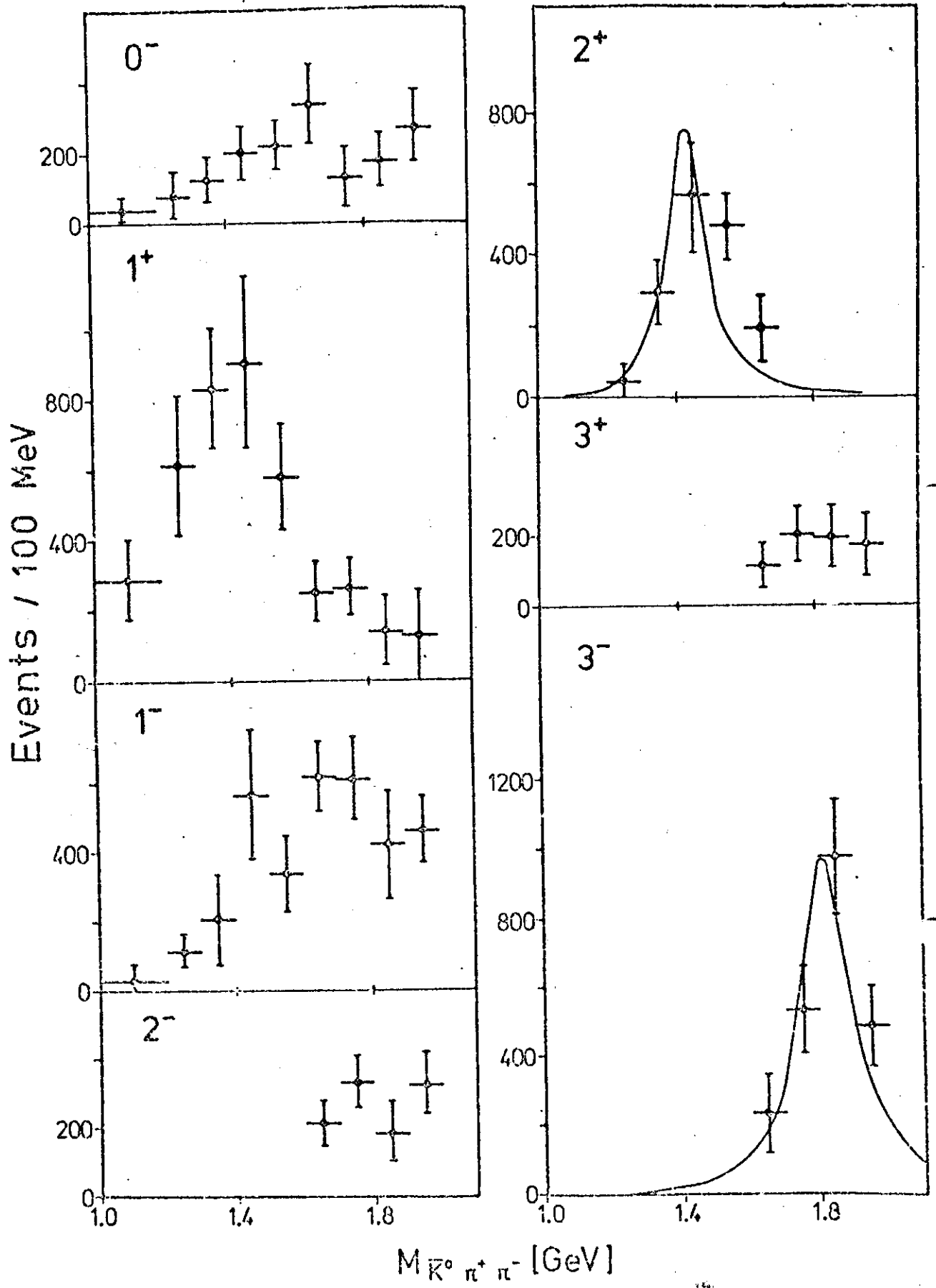


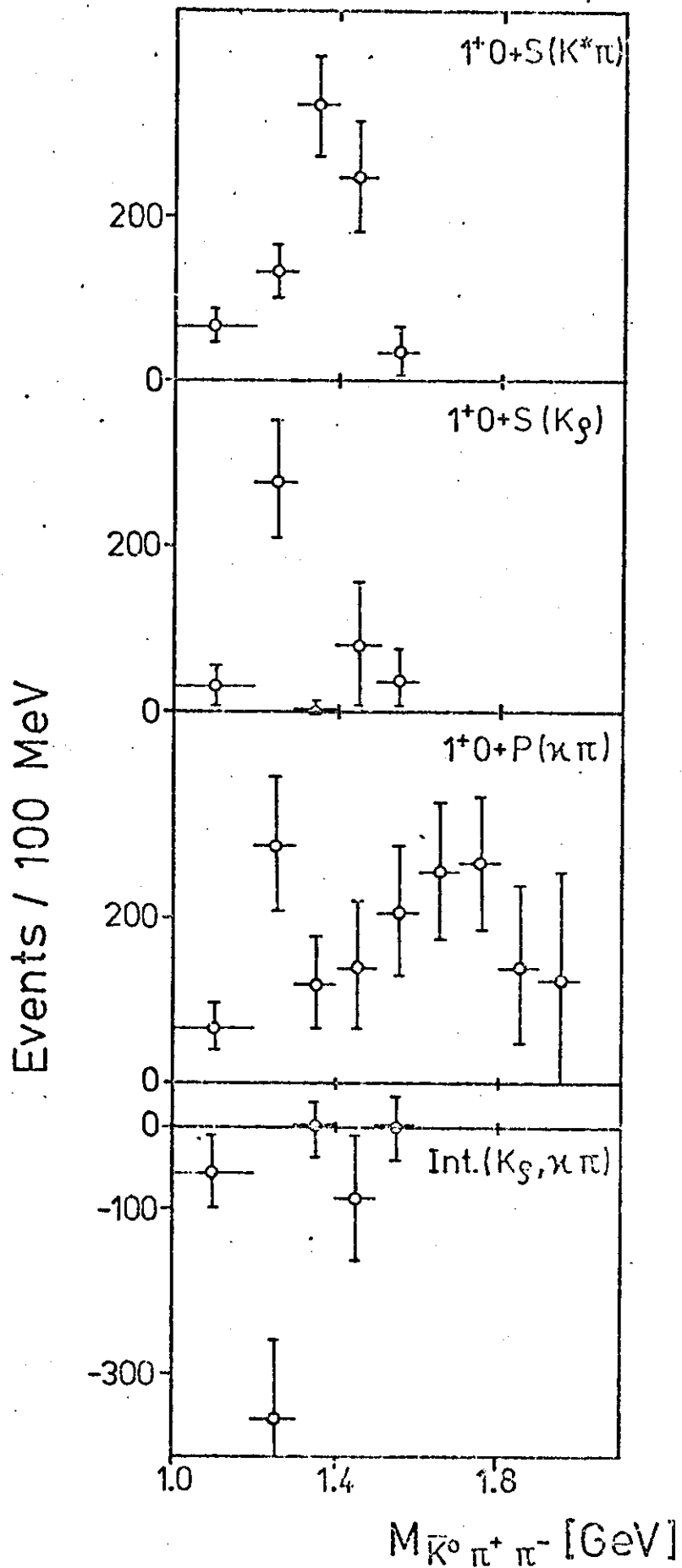
Table I : Acceptance and trigger rates

	i) topology trigger only		ii) total trigger	
	acceptance	trigger/burst	acceptance	trigger/burst
$K^- p \rightarrow \bar{K}^0 \pi^+ \pi^- n$ $\quad \quad \quad \downarrow$ $\quad \quad \quad \rightarrow \pi^+ \pi^-$	$\sim 17 \%$	$\sim 7$	$\sim 15 \%$	$\sim 6$
$\rightarrow 2$ prong + $V^0$ ( $\bar{K}^0 \pi^+ \pi^- n$ excluded)	$\sim 10 \%$	$\sim 65$	$\sim 9 \%$	$\sim 59$
$\rightarrow 4$ prong	$\sim 3,5 \%$	$\sim 83$	}	$\leq 15$
$\rightarrow 4$ prong + $V^0$	$\sim 4,6 \%$	$\sim 26$		
$\rightarrow 6$ prong	$\sim 2,5 \%$	$\sim 18$		
$\rightarrow$ rest		$\sim 2$		
total		$\sim 201$		$\leq 80$

FIGURE CAPTIONS

- Fig. 1 Results of a PWA from experiment S116. The intensities of the various  $J^P$  - states as function of the  $(K^0 \pi^+ \pi^-)$ -mass are given.
- Fig. 2 Intensities of the  $J^P = 1^+$  wave, taken from experiment S116.
- Fig. 3 Contribution of the natural  $J^P$  - series ( $1^-, 2^+, 3^-$ ) and of the unnatural  $J^P$  - series ( $0^-, 1^+, 2^-, 3^+$ ) produced by unnatural spin-parity exchange.
- Fig. 4 Experimental set-up. A forward version of the  $\Omega$  prime detector with A- and B-type chambers is envisaged.





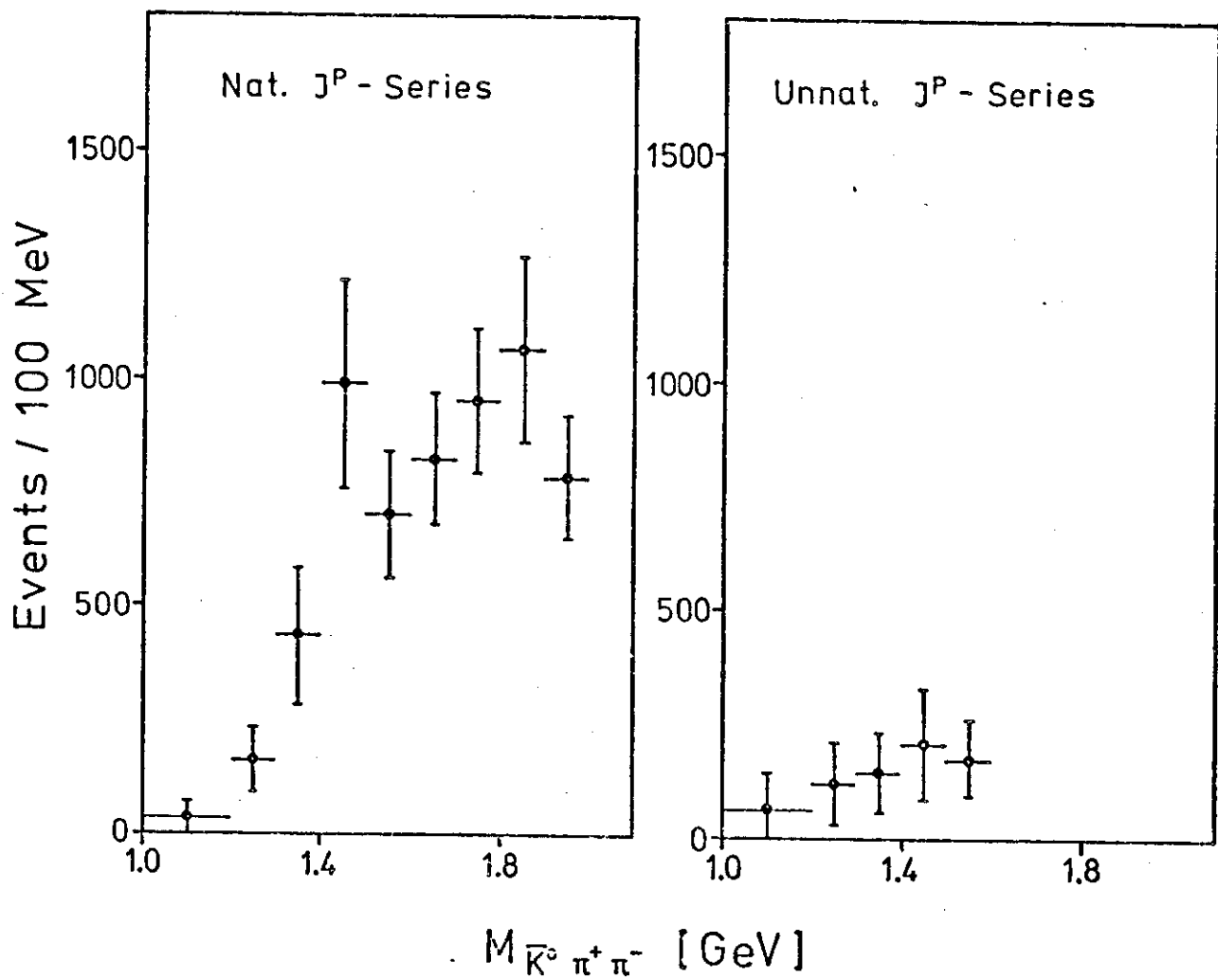


Fig.3 Unnatural Spin-Parity Exchange

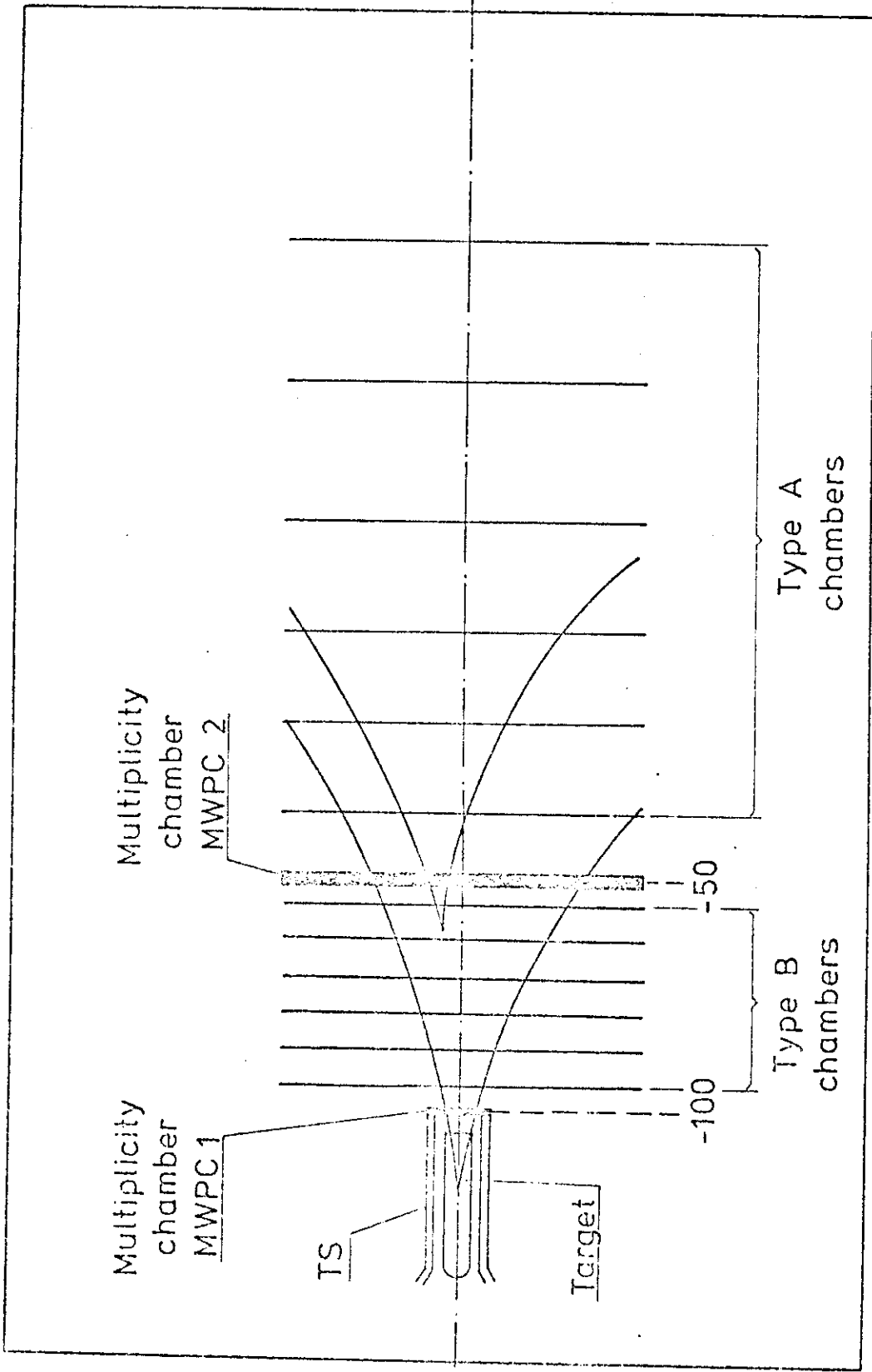


Fig.4 Trigger set up