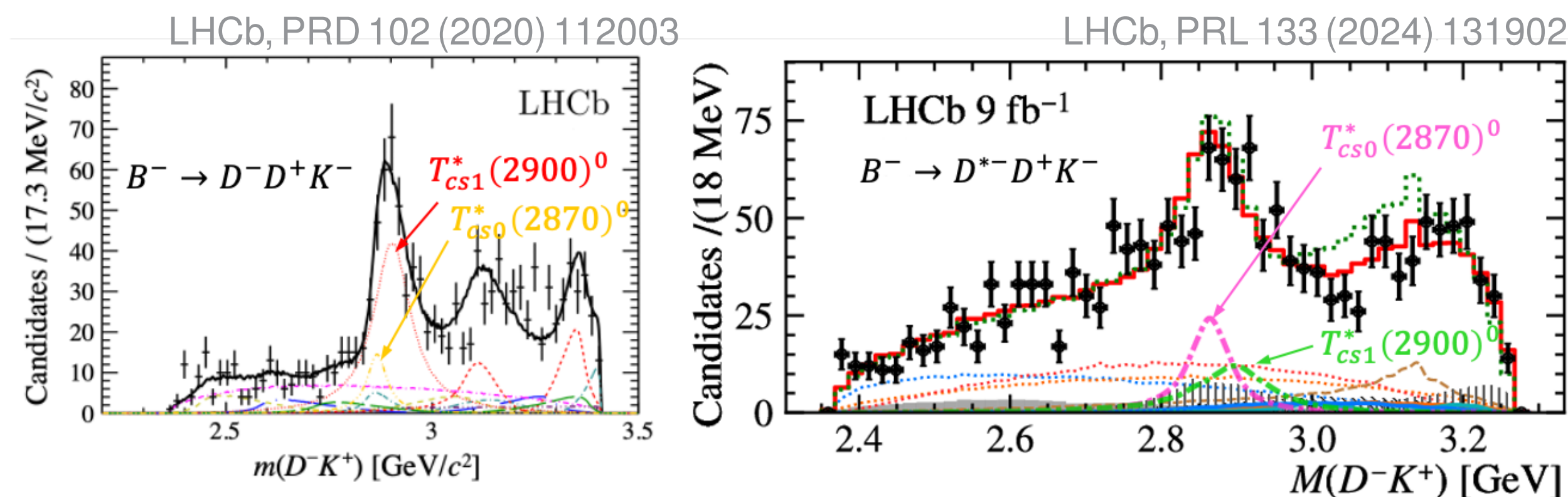


1. Introduction

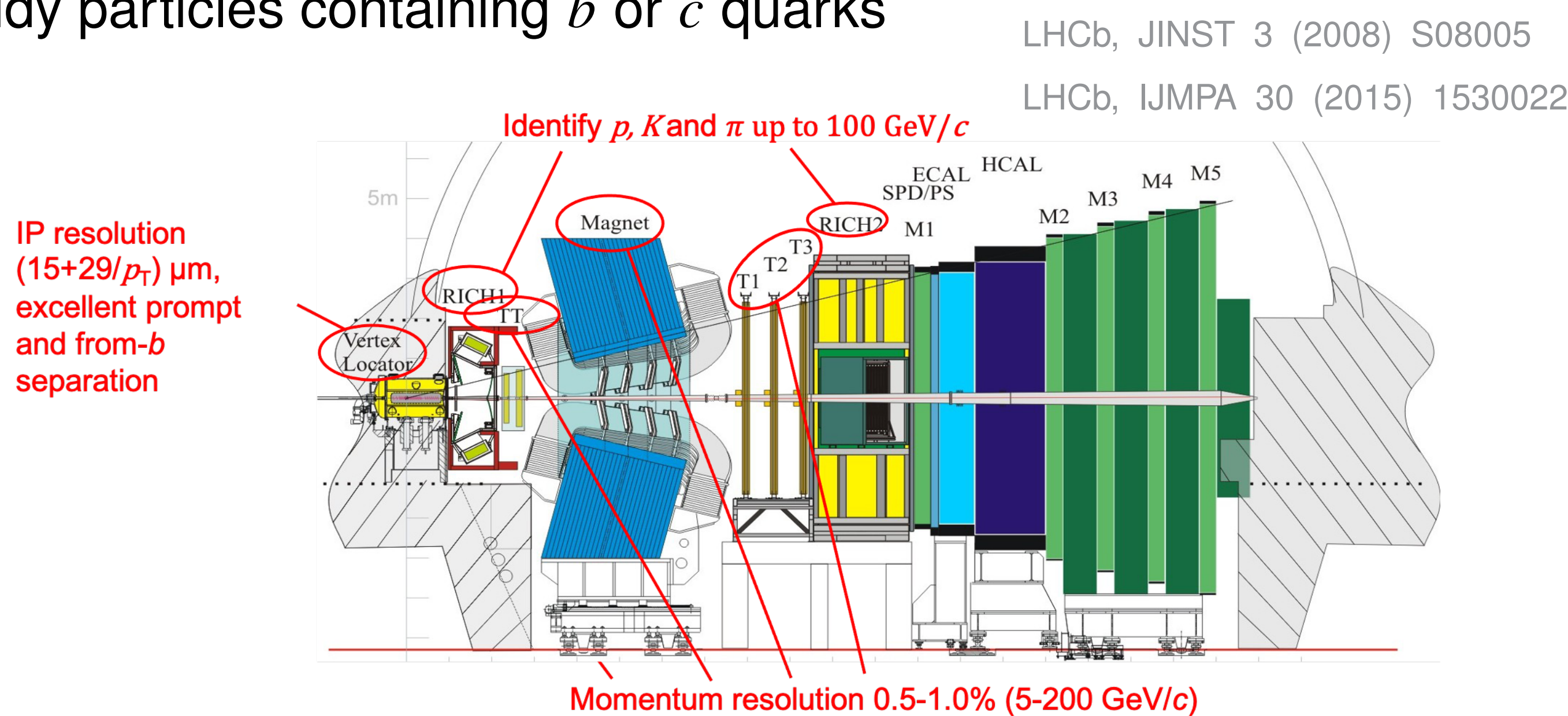
- Discoveries of open-charm tetraquark states $T_{cs0}^*(2870)^0$ and $T_{cs1}^*(2900)^0$ ($cs\bar{u}\bar{d}$)
 - Firstly observed in the $B^- \rightarrow D^- D^+ K^-$ decay LHCb, PRD 102 (2020) 112003 LHCb, PRL 125 (2020) 242001
 - Confirmed in the same channel of the $B^- \rightarrow D^{*-} D^+ K^-$ decay



- There are multiple interpretations for their internal structures.
- $B^- \rightarrow D^- D^0 K_S^0$ decay: an ideal process to search for the T_{cs}^* exotic states in the $D^0 K_S^0$ channel.

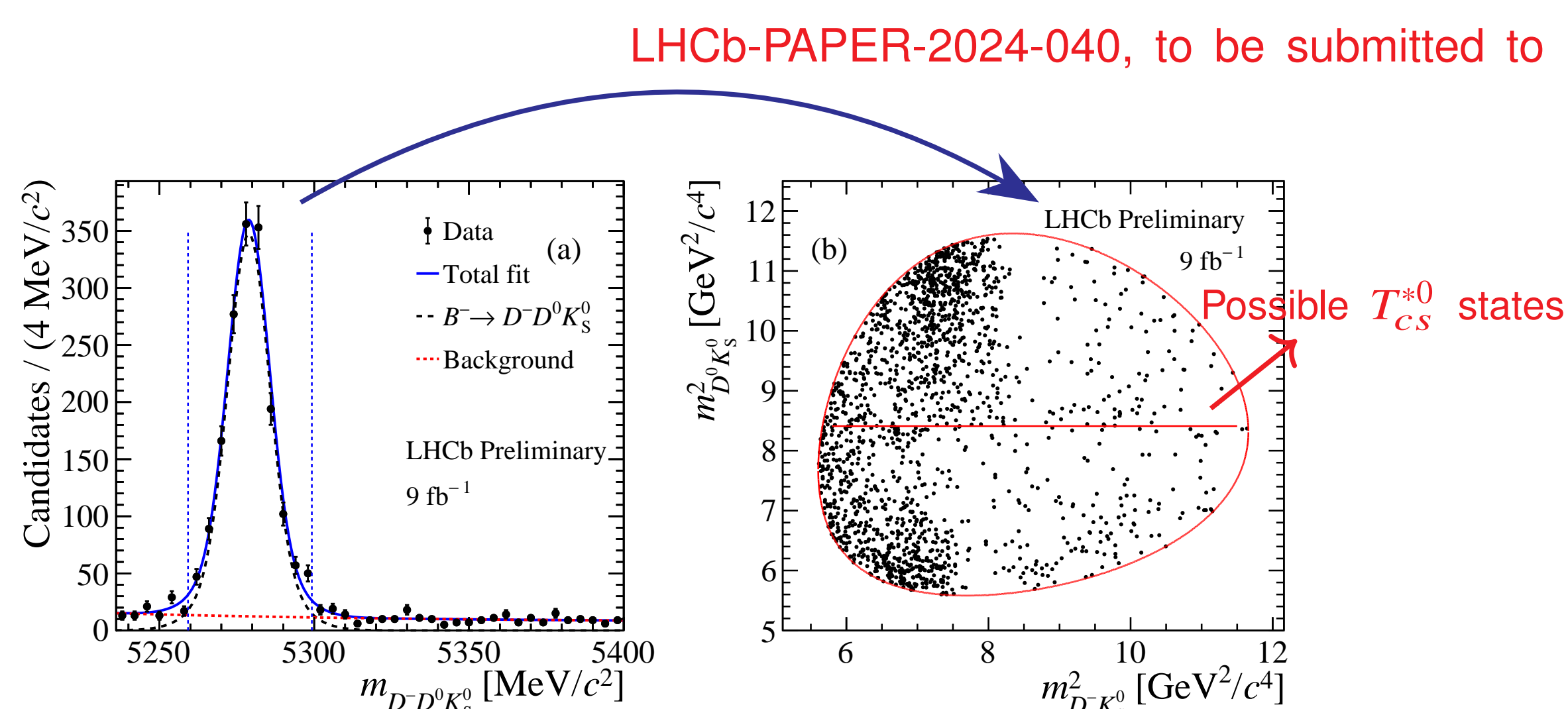
2. LHCb detector

- A single-arm forward spectrometer, covering the pseudorapidity range of $2 < \eta < 5$
- Study particles containing b or c quarks



3. Signal extraction

- Dataset: LHCb Run1+Run2 data
- Selections: cut-based selections + MVA
- Mass fit to extract the signal yields:
 - Signal shape: $f \times \text{Gaus}(m, \sigma_1) + (1 - f) \times \text{Gaus}(m, \sigma_2)$
 - Combinatorial background shape: exponential



Signal yield $N_s = 1544 \pm 42$, signal purity $f_s = (92.6 \pm 0.6)\%$

4. Amplitude model

- For the $B \rightarrow R(\rightarrow ab)c$ decay, helicity amplitude is $\mathcal{M}_R(m_{ab}, \theta_{ab}|\vec{\omega})$
- Signal PDF:

$$\mathcal{P}_s(m_{ab}, \theta_{ab}|\vec{\omega}) = \frac{\epsilon(m_{ab}, \theta_{ab})}{I(\vec{\omega})} \left| \sum_R \mathcal{M}_R(m_{ab}, \theta_{ab}|\vec{\omega}) \right|^2$$

- Efficiency map $\epsilon(m_{ab}, \theta_{ab})$: derived from simulation samples
- Background PDF \mathcal{P}_b : derived from the high B mass region of data
- Log-likelihood for the amplitude fit:

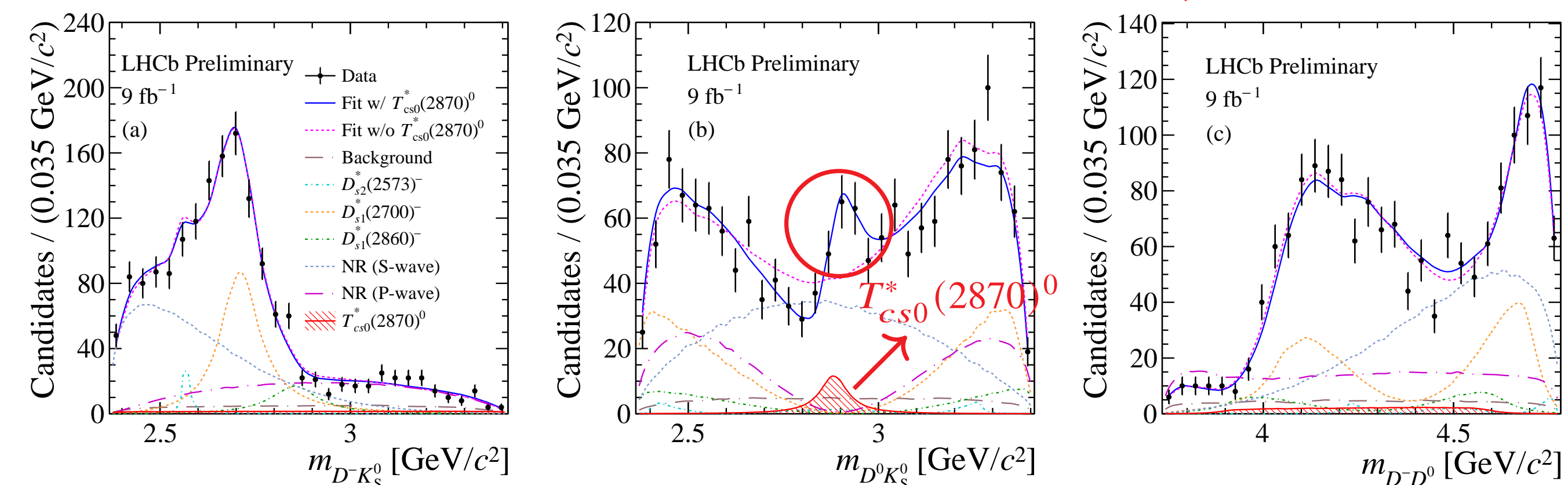
$$\ln \mathcal{L} = \sum_j \ln \left[f_s \mathcal{P}_s(m_{ab}^j, \theta_{ab}^j|\vec{\omega}) + (1 - f_s) \mathcal{P}_b(m_{ab}^j, \theta_{ab}^j) \right] \quad f_s \text{ is fixed}$$

5. Fit strategy

- Null hypothesis model (H_0): only known D_{sJ}^* resonances + NRs
- Add new T_{cs}^* states in the model to find the nominal one (H_1) resulting in the largest $\Delta \ln \mathcal{L} \equiv \ln \mathcal{L}(H_1) - \ln \mathcal{L}(H_0)$.
- Significance of T_{cs}^* is estimated using pseudo-experiments.

6. Results

- The nominal model: known D_{sJ}^* resonances + NRs + $T_{cs0}^*(J^P = 0^+)$



- $M(T_{cs0}^*) = 2883 \pm 11 \pm 7 \text{ MeV}/c^2$, $\Gamma(T_{cs0}^*) = 87_{-47}^{+22} \pm 6 \text{ MeV}$.
- The significance of the $T_{cs0}^*(J^P = 0^+)$ state is 5.3σ after accounting for systematic uncertainties.
- The T_{cs0}^* state is identical to the $T_{cs0}^*(2870)^0$ state observed in the $D^+ K^-$ mass spectrum of the $B^- \rightarrow D^- D^+ K^-$ decay.
- No significant $T_{cs1}^*(J^P = 1^-)$ state is found.

7. Check of isospin invariance

Quantities used to check the isospin invariance between $T_{cs}^* \rightarrow D^0 \bar{K}^0$ and $T_{cs}^* \rightarrow D^+ K^-$ decays:

- Relative decay width $R_I(T_{cs}^*) \equiv \Gamma(T_{cs}^* \rightarrow D^0 \bar{K}^0) / \Gamma(T_{cs}^* \rightarrow D^+ K^-) \approx 1$

$$R_I(T_{cs}^*) = \frac{\mathcal{B}(B^- \rightarrow D^- D^0 \bar{K}^0) \text{FF}(T_{cs}^* \rightarrow D^0 \bar{K}^0)}{\mathcal{B}(B^- \rightarrow D^- D^+ K^-) \text{FF}(T_{cs}^* \rightarrow D^+ K^-)}$$

- Ratio of the fit fractions

$$R_{\text{FF}}(D^0 K_S^0) \equiv \text{FF}(T_{cs1}^* \rightarrow D^0 K_S^0) / \text{FF}(T_{cs0}^* \rightarrow D^0 K_S^0)$$

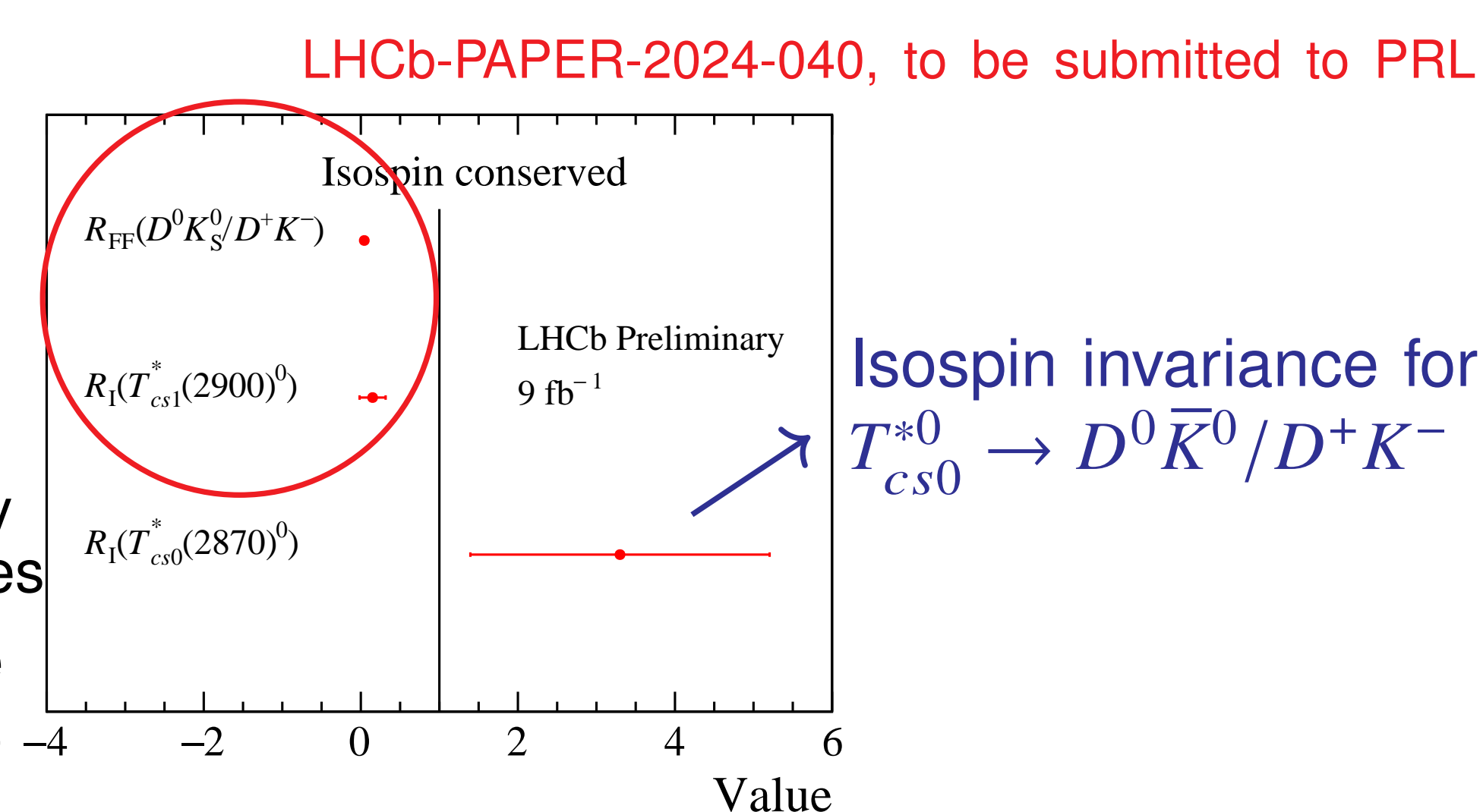
$$R_{\text{FF}}(D^+ K^-) \equiv \text{FF}(T_{cs1}^* \rightarrow D^+ K^-) / \text{FF}(T_{cs0}^* \rightarrow D^+ K^-)$$

$$R_{\text{FF}}(D^0 K_S^0 / D^+ K^-) \equiv R_{\text{FF}}(D^0 K_S^0) / R_{\text{FF}}(D^+ K^-) \approx 1$$

Results:

Indicate a violation of isospin invariance for $T_{cs1}^* \rightarrow D^0 \bar{K}^0 / D^+ K^-$

- May be caused by triangle singularities
- May be a genuine state not having a definite isospin



8. Summary

- The $T_{cs0}^*(2870)^0$ state is observed in a new decay mode $D^0 K_S^0$ in the $B^- \rightarrow D^- D^0 K_S^0$ decay.
- The mass and width of the $T_{cs0}^*(2870)^0$ state are measured.
- No significant T_{cs1}^* state with $J^P = 1^-$ is observed in the $B^- \rightarrow D^- D^0 K_S^0$ decay.
- Isospin invariance is checked, and an isospin violation between the $T_{cs1}^*(2900)^0 \rightarrow D^0 \bar{K}^0$ and $T_{cs1}^*(2900)^0 \rightarrow D^+ K^-$ decays is indicated.