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Top EFT summary plots November 2023

The ATLAS and CMS Collaborations

This note presents figures that summarise the limits on effective field theory (EFT) operators derived from measurements of the ATLAS top working group and LHC Top working group (ATLAS+CMS). Measurements of top quark pair production, single top production and associated production processes are interpreted within the SMEFT framework. FCNC processes are also included. Individual and marginalised bounds on Wilson coefficients are derived at the 68% and 95% CL.

© 2023 CERN for the benefit of the ATLAS Collaboration. Reproduction of this article or parts of it is allowed as specified in the CC-BY-4.0 license. The Standard Model of Particle Physics (SM) has proven to be a succesful theory both in describing observed phenomena as well as making predictions which have later been confirmed by experiments. However, there remain unaswered questions that the SM fails to answer such as the asymmetry between matter and anti-matter, the hierarchy between the Planck mass scale and the electroweak scale set by the vacuum expectation value of the Higgs field or the nature of the dark matter and dark energy present in our Universe. This makes mandatory the searches for new physics at the LHC. New physics can be probed directly by searching for new states. However, if the mass of such particles lies outside the direct reach at LHC energies, it is still possible to infer their existence by indirect means.

The Effective Field Theory (EFT) [1] provides a model-independent framework for such indirect searches. Within this framework, the SM is regarded as a low-energy approximation of a more fundamental theory involving interactions at an energy scale Λ . New physics is then parametrised in terms of higher dimension operators which only include SM fields. The effective Lagrangian then becomes

$$\mathcal{L}_{\text{Eff}} = \mathcal{L}_{\text{SM}} + \sum_{d,i} \frac{c_i^{(d)}}{\Lambda^{d-4}} O_i^{(d)},\tag{1}$$

where \mathcal{L}_{SM} is the SM Lagrangian, $O_i^{(d)}$ are the effective operators of dimension *d* and the complex coefficients $c_i^{(d)}$ are the Wilson coefficients that parameterise the strength of the interaction.

The measurements included in these summary figures use the Warsaw basis [2] as recommended by the LHC Top Working Group [3]. For all measurements the series is truncated at dimension-6 and the value $\Lambda = 1$ TeV is used.

The bounds in ATLAS-only plots are reported taking into account only the linear term of the SMEFT operator (red lines) and/or the linear and the quadratic terms (blue line), except for Figure 5, where only the quadratic term is reported. The dashed line represents the 95% CL limits while the solid one accounts for the 68% CL limits.

Figure 1 shows the derived bounds on two-quark operators obtained from ATLAS measurements only. Limits on each individual operator are derived by fixing the rest to the SM value. The measurements include:

- ATLAS, *ttZ* [4]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- ATLAS, *tī* rapidity asymmetry [6]

Figure 2 shows the derived bounds on two-quark operators obtained from ATLAS measurements only. In this plot, all other operators are profiled. The measurements include:

- ATLAS, Single top t-channel, top polarization [7]
- ATLAS, $t\bar{t}Z$ [4]
- ATLAS, *tī* l+jets boosted [5]

Figure 3 shows the derived bounds on four-quark operators obtained from ATLAS measurements only. In this plot, all other operators are fixed to zero (their SM expectation). The measurements include:

- ATLAS, $t\bar{t}$ l+jets boosted [5]
- ATLAS, $t\bar{t}$ all-hadronic boosted [8]
- ATLAS, *tī* rapidity asymmetry [6]
- ATLAS, $t\bar{t}$ +jet energy asymmetry [9]
- ATLAS, *ttZ* [4]

Figure 4 shows the derived bounds on four-quark operators obtained from ATLAS measurements only. In this plot, all other operators are profiled. The measurements include:

- ATLAS, $t\bar{t}Z$ [4]
- ATLAS, $t\bar{t}$ l+jets boosted [5]

Figure 5 shows the derived bounds on dimension-6 FCNC EFT operators in vertical formats. In this plot, all other operators are fixed to zero (their SM expectation). The measurements included:

- ATLAS, FCNC tqZ [10]
- ATLAS, FCNC *tqH* [11]
- ATLAS, FCNC tqg [12]
- ATLAS, FCNC $tq\gamma$ [13]

Figures 6 and 7 show the derived bounds on dimension-6 top quark EFT operators (excluding FCNCs) in vertical and horizontal formats respectively. In these plots, all other operators are fixed to zero (their SM expectation). The measurements include:

- CMS, $t\bar{t}$ dilepton [14]
- ATLAS, $t\bar{t}Z$ [4]
- CMS, $t\bar{t}$ spin correlations [15]
- CMS, 4 top quarks [16]
- CMS, $t\bar{t}$ and tW, BSM search [17]
- CMS, *ttZ* [18]
- ATLAS+CMS, *W* helicity [19]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- CMS, *tZq/tīZ* [21]
- ATLAS, Single top t-channel, top polarization [7]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- ATLAS, $t\bar{t}$ +jet energy asymmetry [9]
- ATLAS, $t\bar{t}$ all-hadronic boosted [8]
- CMS, *tt*γ [22]

- ATLAS, *tī* rapidity asymmetry [6]
- CMS, $t\bar{t}$ +boosted Z/H [23]
- CMS, tīH, tīlv, tīll, tllq, tHq, tītī [24]

Figures 8 and 9 show the derived bounds on dimension-6 top quark EFT operators (excluding FCNCs) in vertical and horizontal formats respectively. In these plots, all other operators are profiled. The measurements include:

- CMS, 4 top quarks [16]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- CMS, *tZq/ttZ* [21]
- ATLAS, Single top t-channel, top polarization [7]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- CMS, *tt*γ [22]
- CMS, $t\bar{t}$ +boosted Z/H [23]
- CMS, $t\bar{t}H$, $t\bar{t}\ell\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq, $t\bar{t}t\bar{t}$ [24]

Figure 10 shows the derived bounds on dimension-6 FCNC EFT operators in vertical formats. In this plot, all other operators are fixed to zero (their SM expectation). The measurements include:

- CMS, $t\bar{t}$ and tW, BSM search [17]
- ATLAS, FCNC tqZ [10]
- ATLAS, FCNC *tqH* [11]
- ATLAS, FCNC *tqg* [12]
- ATLAS, FCNC $tq\gamma$ [13]

Figures 11 and 13 show the derived bounds on dimension-6 operators related to interactions with scalar bosons (excluding FCNCs). In these plots, all other operators are fixed to zero (their SM expectation). The measurements include:

- ATLAS, $t\bar{t}Z$ [4]
- CMS, $t\bar{t}$ and tW, BSM search [17]
- CMS, *ttZ* [18]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- CMS, *tZq/ttZ* [21]
- CMS, $t\bar{t}$ +boosted Z/H [23]
- CMS, *tīH*, *tīlv*, *tīll*, *tllq*, *tHq*, *tītī* [24]

Figures 12 and 14 show the derived bounds on dimension-6 operators related to interactions with scalar bosons (excluding FCNCs). In these plots, all other operators are profiled. The measurements include:

- ATLAS, $t\bar{t}Z$ [4]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- CMS, *tZq/ttZ* [21]
- CMS, $t\bar{t}$ +boosted Z/H [23]
- CMS, *tīH*, *tīlv*, *tīll*, *tllq*, *tHq*, *tītī* [24]

Figures 15 and 16 show the derived bounds on dimension-6 operators related to interactions with vector bosons (excluding FCNCs). In these plots, all other operators are fixed to zero (their SM expectation). The measurements include:

- CMS, tī dilepton [14]
- ATLAS, $t\bar{t}Z$ [4]
- CMS, *tī* spin correlations [15]
- CMS, $t\bar{t}$ and tW, BSM search [17]
- CMS, *ttZ* [18]
- ATLAS+CMS, W helicity [19]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- CMS, *tZq/t*tZ [21]
- ATLAS, Single top t-channel, top polarization [7]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- CMS, *tt*γ [22]
- ATLAS, *tī* rapidity asymmetry [6]
- CMS, $t\bar{t}$ +boosted Z/H [23]
- CMS, tīH, tīlv, tīll, tllq, tHq, tītī [24]

Figures 17 and 18 show the derived bounds on dimension-6 operators related to interactions with vector bosons (excluding FCNCs). In these plots, all other operators are profiled. The measurements include:

- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- CMS, *tZq/ttZ* [21]
- ATLAS, Single top t-channel, top polarization [7]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- CMS, *tt*γ [22]
- CMS, $t\bar{t}$ +boosted Z/H [23]

• CMS, tīH, tīlv, tīll, tllq, tHq, tītī [24]

Figures 19 and 20 show the derived bounds on dimension-6 operators related to four-fermion interactions (excluding FCNCs). In these plots, all other operators are fixed to zero (their SM expectation). The measurements include:

- CMS, 4 top quarks [16]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- ATLAS, $t\bar{t}$ +jet energy asymmetry [9]
- ATLAS, $t\bar{t}$ all-hadronic boosted [8]
- ATLAS, $t\bar{t}$ rapidity asymmetry [6]
- CMS, $t\bar{t}H$, $t\bar{t}\ell\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq, $t\bar{t}t\bar{t}$ [24]

Figures 21 and 22 show the derived bounds on dimension-6 operators related to four-fermion interactions (excluding FCNCs). In these plots, all other operators are fixed to zero (their SM expectation). The measurements include:

- CMS, 4 top quarks [16]
- CMS, $t\bar{t} + Z/W/H$, tZq, tHq [20]
- ATLAS, $t\bar{t}$ l+jets boosted [5]
- CMS, *tīH*, *tīlv*, *tīll*, *tllq*, *tHq*, *tītī* [24]

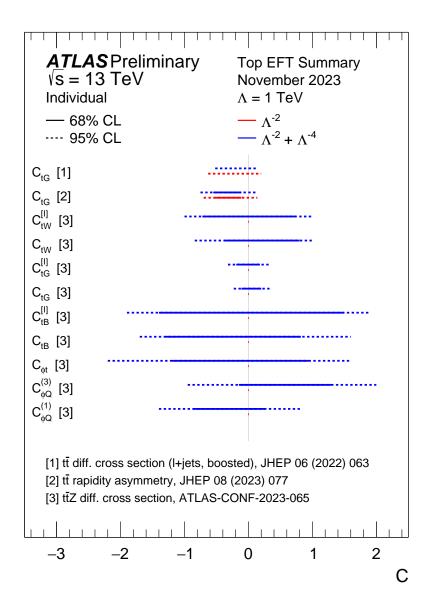


Figure 1: Summary of constraints on two-fermion SMEFT operators from top-quark measurements at the ATLAS experiment. The bounds on the Wilson coefficients are reported at the 68% CL (solid) and/or 95% CL (dashed) depending on the availability in the corresponding measurement. The bounds are reported without (red) and/or with (blue) taking into account the quadratic term of the SMEFT operator, depending on the availability in the corresponding measurement. Limits on each individual operator are derived by fixing the rest to the SM value. Interpretations use the SMEFT framework and the Warsaw basis. The vertical bar represents the SM prediction.

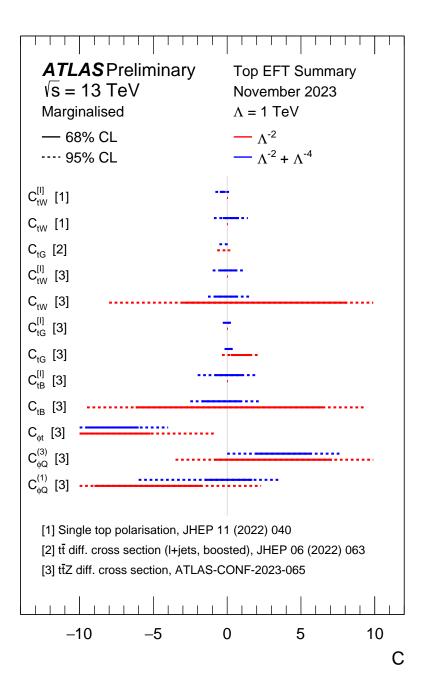


Figure 2: Summary of constraints on two-fermion SMEFT operators from top-quark measurements at the ATLAS experiment. The bounds on the Wilson coefficients are reported at the 68% CL (solid) and/or 95% CL (dashed) depending on the availability in the corresponding measurement. The bounds are reported without (red) and/or with (blue) taking into account the quadratic term of the SMEFT operator, depending on the availability in the corresponding measurement. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. Interpretations use the SMEFT framework and the Warsaw basis. The vertical bar represents the SM prediction.

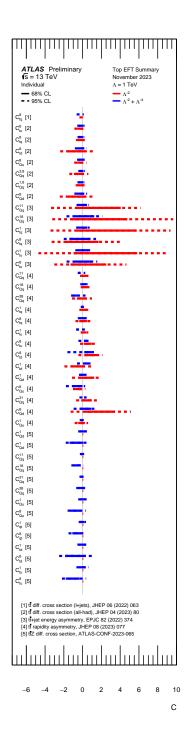


Figure 3: Summary of constraints on four-fermion SMEFT operators from top-quark measurements at the ATLAS experiment. The bounds on the Wilson coefficients are reported at the 68% CL (solid) and/or 95% CL (dashed) depending on the availability in the corresponding measurement. The bounds are reported without (red) and/or with (blue) taking into account the quadratic term of the SMEFT operator, depending on the availability in the corresponding measurement. Limits on each individual operator are derived by fixing the rest to the SM value. Interpretations use the SMEFT framework and the Warsaw basis. The vertical bar represents the SM prediction. Only the most stringent limit from Ref. [6] is quoted.

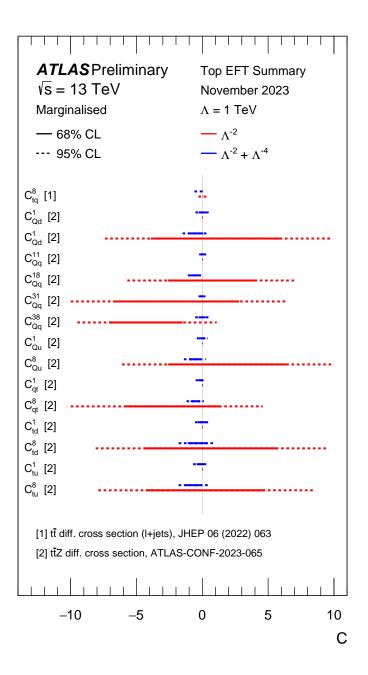


Figure 4: Summary of constraints on four-fermion SMEFT operators from top-quark measurements at the ATLAS experiment. The bounds on the Wilson coefficients are reported at the 68% CL (solid) and/or 95% CL (dashed) depending on the availability in the corresponding measurement. The bounds are reported without (red) and/or with (blue) taking into account the quadratic term of the SMEFT operator, depending on the availability in the corresponding measurement. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. Interpretations use the SMEFT framework and the Warsaw basis. The vertical bar represents the SM prediction. Only the most stringent limit from Ref. [6] is quoted.

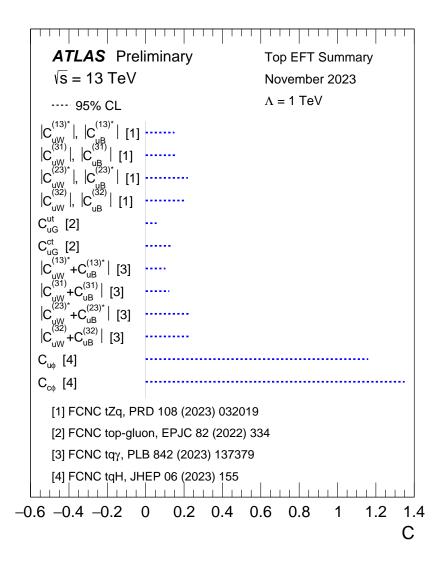


Figure 5: Summary of constraints on FCNC SMEFT operators from top-quark measurements at the ATLAS experiment. The bounds on the Wilson coefficients are reported at the 95% CL. The quadratic term of the SMEFT operator is taken into account for the interpretation (the linear term is negligible). Limits on each individual operator are derived by fixing the rest to the SM value. Interpretations use the SMEFT framework and the Warsaw basis. The vertical bar represents the SM prediction.

HCtopWG ATLAS	p quark EFT operators - Individ — ATLAS+CMS	ual limits CM	Following arXiv:1802.07237 Dimension 6 operators $\hat{C}_i \equiv C_i$	$/\Lambda^2$
Č ¹ _{ao}			CMS, <i>titi</i> [3]	36 fb
			CMS, tÎH, tÎlv, tÎtt, têtq, tHq, tÎtî [14	
Č ¹ _{at}			CMS, tĨtĨ [3] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [14	36 fb ⁻ 138 fb ⁻
Č ⁸ αt			CMS, tĨtĨ [3] CMS, tĨH, tĨlv, tĨtl, tElq, tHq, tĨtĨ [14	36 fb ⁻ 138 fb ⁻
Č ¹ tt			CMS. titi 131	36 fb
0 ₁₁			CMS, tÎH, tÎlv, tÎtt, tttq, tHq, tÎtÎ [14	
			ATLAS, tī + jet energy asymmetry [9] ATLAS, tī l + jets boosted [11]	139 fb 139 fb
\tilde{C}_{tq}^{8}			ATLAS, t [†] all-hadronic boosted [13] CMS, t [†] H, t [†] Iν, t [†] tℓ, tℓℓq, tHq, t [†] t [†] [14	139 fb 138 fb
	-		ATLAS, tt rapidity asymmetry [16]	139 fb
Č _{to}			CMS, tĨ + Z/W/H, tZq, tHq [7] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [14	42 fb
			CMS, tł + boosted Z/H [15]	138 fb
			CMS, tÎZ [5] CMS, tÎ + Z/W/H, tZq, tHq [7]	78 fb 42 fb
$\tilde{C}^{-}_{\phi Q}$	±		CMS, tZq / tĨZ [8]	138 fb
			CMS, tĨH, tĨlv, tĨtl, tℓℓq, tHq, tĨtĨ [14 CMS, tĨ + boosted Z/H [15]	138 fb 138 fb
	-		CMS, II and IW, BSM search [4]	36 fb
č3			CMS, tī + Z/W/H, tZq, tHq [7] CMS, tZq / tīZ [8]	42 fb 138 fb
$\hat{C}^3_{\phi Q}$	<u> </u>		CMS, tĨH, tĨlv, tĨtt, tttq, tHq, tĨtĨ [14 CMS, tĨ + boosted Z/H [15]	138 fb 138 fb
	-		CMS, If + boosted Z/H [15] ATLAS, ITZ diff. cross section [17]	140 fb
			CMS, <i>tiZ</i> [5]	78 fb
Č _{el}			CMS, tī + Z/W/H, tZq, tHq [7] CMS, tZq / tīZ [8]	42 fb 138 fb
-01			CMS, tîH, tîlv, tîll, têlq, tHq, tîtî [14 CMS, tî + boosted Z/H [15]	138 fb 138 fb
			ATLAS, <i>tīZ</i> diff. cross section [17]	140 fb
Ē _{otb}			CMS, tĨ + Z/W/H, tZq, tHq [7] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [14	
			CMS, tł + boosted Z/H [15]	138 fb
	<u> </u>		CMS, tt and tW, BSM search [4] ATLAS+CMS, W helicity [6] 2/	36 fb +20 fb
			CMS, $t\tilde{t} + Z/W/H$, tZq , tHq [7]	42 fb
Č _{tW}	Į.		CMS, tZq / tĪZ [8] ATLAS, Top polarization [12]	138 fb 139 fb
	±		CMS, tÎH, tÎlv, tÎtt, tttq, tHq, tÎtÎ [14 CMS, tÎ + boosted Z/H [15]	138 fb 138 fb
	+		ATLAS, <i>tIZ</i> diff. cross section [17]	140 fb
$\tilde{C}_{tW}^{[I]}$	<u>-</u>		ATLAS, Top polarization [12] ATLAS, <i>tlZ</i> diff. cross section [17]	139 fb 140 fb
	+		CMS, tīZ [5]	78 fb
à	Ŧ		CMS, tł + Z/W/H, tZq, tHq [7] CMS, tZq / tłZ [8]	42 fb 138 fb
Č _{1Z}			CMS, tÎ ₂ [10] CMS, tÎH, tÎlv, tĨtt, tttq, tHq, tĨtĨ [14	137 fb
	+		CMS, <i>tī</i> + boosted Z/H [15]	138 fb
$\tilde{C}_{tZ}^{(I)}$	÷		CMS, tĨZ [5] CMS, tĨγ [10]	78 fb 137 fb
a.			CMS, tī + Z/W/H, tZq, tHq [7]	42 fb
Ĉ _{bW}			CMS, tĨH, tĨlv, tĨtl, tťtℓq, tHq, tĨtĨ [14 CMS, tĨ + boosted Z/H [15]	138 fb
	1		CMS, tr dilepton [1]	36 fb 36 fb
			CMS, tt spin correlations [2] CMS, tt and tW, BSM search [4]	36 fb
Č _{1G}	- †		CMS, tÎ + Z/W/H, tZq, tHq [7] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [14	42 fb 138 fb
	1		ATLAS, <i>tī</i> rapidity asymmetry [16] ATLAS, <i>tī</i> Z diff. cross section [17]	139 fb 140 fb
Ĉic/ac]		ATLAS, <i>ti</i> 2 diff. cross section [17]	140 fb
С _{tG} /gs z3(t) -			CMS, tī + Z/W/H, tZq, tHq [7]	42 fb
ai			CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [14	138 fb
Č_(ℓ)			CMS, t1 + Z/W/H, tZq, tHq [7] CMS, t1H, t11v, t1et, teeq, tHq, t1t1 [14	138 fb
$\tilde{C}_{Qe}^{(l)}$			CMS, tī + Z/W/H, tZq, tHq [7] CMS, tīH, tīlv, tlℓℓ, tℓℓq, tHq, tītī [14	42 fb 138 fb
$\tilde{C}_{tl}^{(t)}$			CMS, tĨ + Z/W/H, tZq, tHq [7] CMS, tĨH, tĨIv, tĨtℓ, tℓℓq, tHq, tĨtĨ [14	42 fb 138 fb
$\tilde{C}_{te}^{(l)}$			CMS, tī + Z/W/H, tZq, tHq [7] CMS, tīH, tīlv, tītt, tttq, tHq, tītī [14	42 fb
$\tilde{C}_{t}^{S(\ell)}$			CMS, tł + Z/W/H, tZq, tHq [7]	42 fb
			CMS, tĨH, tĨlv, tĨtt, tttq, tHq, tĨtĨ [14	138 fb 42 fb
Č ^{T(l)}		NO HED A passa on	CMS, tī + Z/W/H, tZq, tHq [7] CMS, tīH, tīlv, tītt, ttttq, tHq, tītī [14	
 JHEP 02 (2019) 14 PRD 100 (2019) 07 JHEP 11 (2019) 08 JHEP 11 (2019) 08 JHEP 03 (2020) 05 JHEP 08 (2020) 05 	0 [7] JHEP 03 (2021) 095 2002 [8] JHEP 12 (2021) 045 2 [9] EPU5 82 (2022) 374 5 [10] JHEP 05 (2022) 001 5 [11] JHEP 05 (2022) 003 1 [12] JHEP 12 (2022) 040	[13] JHEP 04 (2023) 80 [14] arXiv:2307.15781 * [15] PRD 108 052008 [16] JHEP 08 (2023) 077 [17] ATLAS-CONF-2023-065 * * Preliminary	EFT formatism is employed at different levels of experimental analysies	
[v] vi en, no (2020) 02	(14) VI 12 ¹ 11 (2022) 040	·,	1	

Figure 6: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to interactions involving top quarks, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. In JHEP 06 (2022) 063, the limit is derived for the coefficient c_{tG} normalised with the strong coupling, g_S , as implemented in SMEFT@NLO. Vertical format.

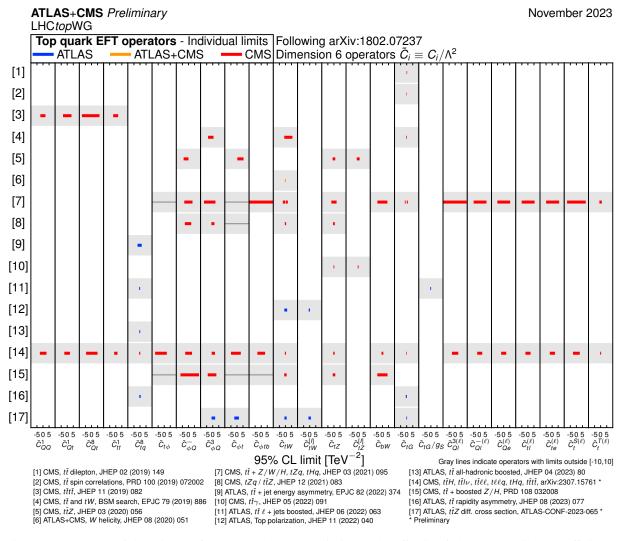


Figure 7: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to interactions involving top quarks, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. In JHEP 06 (2022) 063, the limit is derived for the coefficient c_{tG} normalised with the strong coupling, g_S , as implemented in SMEFT@NLO. Horizontal format.

ATLAS	Top quark EFT operators - Margi	nalised limits Following arXiv:1802.07237 — CMS Dimension 6 operators $\tilde{C}_i \equiv C_i/$	A2
Č ¹ _{QQ}	· · · ·	CMS, tilt [1] CMS, tilt, tilv, till, till, till, till (17]	36 fb-
Ĉ₁		CMS, tỉtỉ [1] CMS, tỉH, tỉlư, tỉtl, tlưq, tHq, tỉtỉ [7]	36 fb 138 fb
Ĉ ⁸ ₀≀		CMS, títí [1] CMS, títl, tílv, títl, téla, tHq, títí [7]	36 fb ⁻ 138 fb ⁻
Ĉ ¹ t		CMS, tỉtỉ [1] CMS, tỉtỉ, tỉtν, tỉtℓ, tℓℓq, tHq, tỉtỉ [7]	36 fb ⁻ 138 fb ⁻
Ĉ ⁸ tq	<u>.</u>		139 fb-
Ĉ _{to}		CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tĨH, tĨlv, tĨℓ, tℓℓq, tHq, tĨtĨ [7]	42 fb ⁻ 138 fb ⁻
σιφ		CMS, tt + boosted Z/H [8]	138 fb
$\tilde{C}^{-}_{\phi Q}$			42 fb ⁻ 138 fb ⁻
çu .		CMS, tÎH, tÎlv, tÎtℓ, tℓℓq, tHq, tÎtÎ [7] CMS, tÎ + boosted Z/H [8]	138 fb 138 fb
		CMS, tī + Z/W/H, tZq, tHq [2] CMS, tZq / tīZ [3]	42 fb ⁻ 138 fb ⁻
$\tilde{C}^{3}_{\phi Q}$		CMS, tĨH, tĨlv, tĨℓℓ, tℓℓq, tHq, tĨtĨ [7]	138 fb 138 fb
		ATLAS, <i>tTZ</i> diff. cross section [9]	140 fb
		CMS, $t\bar{t} + Z/W/H$, tZq , tHq [2] CMS, $tZq / t\bar{t}Z$ [3]	42 fb 138 fb
Ĉ _{¢t}			138 fb 138 fb
			140 fb
Ĉ _{¢tb}		CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tĨH, tĨlv, tĨtℓ, ttℓq, tHq, tĨtĨ [7] CMS, tĨ + boosted Z/H [8]	42 fb 138 fb 138 fb
	<u> </u>	CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tZq / tĨZ [3]	42 fb 138 fb
Ĉ _{tW}	<u>+</u>		139 fb 138 fb
	±	CMS, tt + boosted Z/H [8]	138 fb 140 fb
Ĉ ^{[ŋ} rw	1		139 fb 140 fb
	<u>+</u>	CMS, $t\bar{t} + Z/W/H$, tZq , tHq [2] CMS, $tZq/t\bar{t}Z$ [3]	42 fb 138 fb
Ĉ _{tZ}	<u>-</u>		137 fb 138 fb
			138 fb
			137 fb
Ĉ _{bW}	<u> </u>	CMS, tî + Z/W/H, tZq, tHq [2] CMS, tîH, tîlv, tîlê, têlq, tHq, tîtî [7] CMS, tî + boosted Z/H [8]	42 fb 138 fb 138 fb
Ĉ₁ _G	-	CMS, tī + Z/W/H, tZq, tHq [2] CMS, tī H, tī \u03c6, ti tdq, tHq, tī tī [7] ATLAS, tī Z diff. cross section [9]	42 fb 138 fb 140 fb
Ĉ₁ _G /g _S	•	ATLAS, $t\bar{t} \ell$ + jets boosted [5]	139 fb
$\tilde{C}^{3(\ell)}_{QI}$		CMS, tī + Z/W/H, tZq, tHq [2] CMS, tīH, tīlv, tītēt, tēta, tHq, tītī [7]	42 fb 138 fb
$\tilde{C}_{QI}^{-(\ell)}$		CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [7]	42 fb 138 fb
$\tilde{C}^{(\ell)}_{Qe}$		CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [7]	42 fb 138 fb
$\tilde{C}_{tl}^{(\ell)}$		CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tĨH, tĨlv, tĨtℓ, ttℓq, tHq, tĨtĨ [7]	42 fb 138 fb
$\tilde{C}_{te}^{(\ell)}$		CMS, tĨ + Z/W/H, tZq, tHq [2] CMS, tĨH, tÎlv, tĨtℓ, ttℓq, tHq, tĨtĨ [7]	42 fb 138 fb
$\tilde{C}_{l}^{S(\ell)}$		CMS, tỉ + Z/W/H, tZq, tHq [2] CMS, tỉ H, tỉlv, tỉ ℓℓ, titq, tỉtĩ [7]	42 fb 138 fb
$\tilde{C}_{t}^{T(\ell)}$	Ŧ	CMS, tī + Z/W/H, tZq, tHq [2] CMS, tīH, tīlv, tītē, tēta, tHq, tītī [7]	42 fb 138 fb
[1] JHEP 11 [2] JHEP 03 [3] JHEP 12 [4] JHEP 05	(2019) 082 [5] JHEP 06 (2022) 063 (2021) 095 [6] JHEP 11 (2022) 040 (2021) 083 [7] arXiv:2307.15761 *	[8] PRD 108 032008 EFT formalism is employed at different levels of patrus_CONF-2023-065 * * Preliminary	

Figure 8: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to interactions involving top quarks, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. In JHEP 06 (2022) 063, the limit is derived for the coefficient c_{tG} normalised with the strong coupling, g_S , as implemented in SMEFT@NLO. Vertical format.

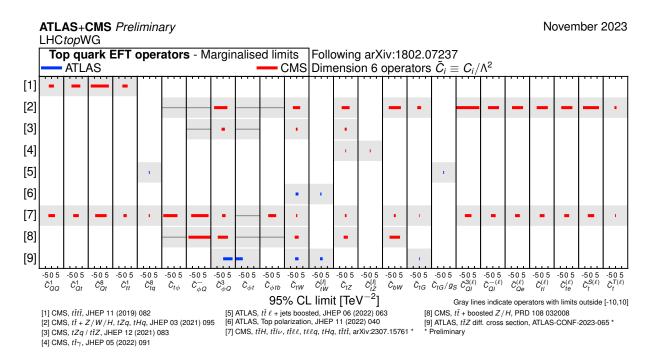


Figure 9: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to interactions involving top quarks, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. In JHEP 06 (2022) 063, the limit is derived for the coefficient c_{tG} normalised with the strong coupling, g_S , as implemented in SMEFT@NLO. Horizontal format.

FCNC operators - Individua	Following arXiv:1802.07237 CMS Dimension 6 operators $\tilde{C}_i \equiv C_i$	/Λ ²
$\left \tilde{C}^{(32)}_{UW} + \tilde{C}^{(32)}_{UB} \right $	ATLAS, FCNC $tq\gamma$ [4]	139 fb ⁻
$\tilde{C}_{UW}^{(23)*} + \tilde{C}_{UB}^{(23)*}$	ATLAS, FCNC $tq\gamma$ [4]	139 fb ⁻
$\tilde{C}_{uW}^{(31)} + \tilde{C}_{uB}^{(31)}$	ATLAS, FCNC $tq\gamma$ [4]	139 fb ⁻
$\left. \tilde{C}_{uW}^{(13)*} + \tilde{C}_{uB}^{(13)*} \right $	ATLAS, FCNC $tq\gamma$ [4]	139 fb ⁻
$ar{C}^{32}_{uW}$	ATLAS, FCNC tZq [5]	139 fb ⁻
	ATLAS, FCNC tZq [5]	139 fb-
$ar{C}^{23}_{\nu W}*$	ATLAS, FCNC tZq [5]	139 fb ⁻
	ATLAS, FCNC tZq [5]	139 fb ⁻
$ar{C}^{31}_{uW}$	ATLAS, FCNC tZq [5]	139 fb ⁻
\tilde{C}^{31}_{uB}	ATLAS, FCNC tZq [5]	139 fb
$ar{C}^{13}_{uW}*$	ATLAS, FCNC tZq [5]	139 fb ⁻
$\left. \tilde{C}^{13}_{uB} * \right $	ATLAS, FCNC tZq [5]	139 fb ⁻
$ ilde{C}_{UG} $	CMS, $t\bar{t}$ and tW , BSM search [1] ATLAS, FCNC tqg [2]	36 fb 139 fb
	CMS, $t\bar{t}$ and tW , BSM search [1] ATLAS, FCNC tqg [2]	36 fb 139 fb
$ ilde{\mathcal{C}}_{u\phi}$	ATLAS, FCNC $tqH (H \rightarrow \tau \tau)$ [3]	139 fb
	ATLAS, FCNC $tqH (H \rightarrow \tau \tau)$ [3]	139 fb ⁻
[1] EPJC 79 (2019) 886 [3] JHEP 06 (2023) 155 [2] EPJC 82 (2022) 334 [4] PLB 842 (2023) 137379	[5] PRD 108 (2023) 032019 EFT formalism is employed at different levels of experimental analyses	

Figure 10: Summary of constraints on FCNC SMEFT operators from top-quark measurements, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. The definitions of operators O_{uG} and O_{cG} are different in the two searches from references EPJC 82 (2022) 334 and EPJC 79 (2019) 886. For comparison purposes in this plot, the limits for coefficients C_{uG} and C_{cG} from Ref. arXiv:2112.01302 are multiplied by factor $g_S(m_{top}^2)/2 = 0.57$ to be compatible with the definition used in Ref. EPJC 79 (2019) 886 as it is recommended in Ref. arXiv:1802.07237. Vertical format.

ATLAS+CMS Preliminal _HCtopWG	ry	November 202
	calar boson operators - Individual limits	Following arXiv:1802.07237 — CMS Dimension 6 operators $\tilde{C}_i \equiv C_i / \Lambda^2$
$ ilde{C}_{t\phi}$		CMS, $t\bar{t} + Z/W/H$, tZq , tHq [3] 42 fb ⁻ CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$ [5] 138 fb ⁻ CMS, $t\bar{t}$ + boosted Z/H [6] 138 fb ⁻
$ ilde{C}_{\phi t}$		CMS, $t\bar{t}Z$ [2] 78 fb ⁻¹ CMS, $t\bar{t} + Z/W/H$, tZq , tHq [3] 42 fb ⁻¹ CMS, $tZq / t\bar{t}Z$ [4] 138 fb ⁻¹
<i>σ</i> _φ ι		CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$ [5]138 fb ⁻ CMS, $t\bar{t}$ + boosted Z/H [6]138 fb ⁻ ATLAS, $t\bar{t}Z$ diff. cross section [7]140 fb ⁻
$ ilde{C}_{\phi tb}$ -		CMS, $t\bar{t} + Z/W/H$, tZq , tHq [3] 42 fb ⁻ CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$ [5] 138 fb ⁻ CMS, $t\bar{t}$ + boosted Z/H [6] 138 fb ⁻
$ ilde{C}^{-}_{\phi Q}$		CMS, $t\bar{t}Z$ [2] 78 fb ⁻ CMS, $t\bar{t} + Z/W/H$, tZq , tHq [3] 42 fb ⁻ CMS, $tZq / t\bar{t}Z$ [4] 138 fb ⁻ CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$ [5] 138 fb ⁻ CMS, $t\bar{t}$ + boosted Z/H [6] 138 fb ⁻
$\tilde{C}^3_{\phi a}$		CMS, $t\bar{t}$ and tW , BSM search [1] 36 fb ⁻ CMS, $t\bar{t} + Z/W/H$, tZq , tHq [3] 42 fb ⁻ CMS, $tZq / t\bar{t}Z$ [4] 138 fb ⁻ CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$ [5] 138 fb ⁻ CMS, $t\bar{t} + boosted Z/H$ [6] 138 fb ⁻ ATLAS, $t\bar{t}Z$ diff. cross section [7] 140 fb ⁻
[1] EPJC 79 (2019) 886 [2] JHEP 03 (2020) 056 [3] JHEP 03 (2021) 095	[4] JHEP 12 (2021) 083 [7] ATLAS-CONF-2023- [5] arXiv:2307.15761 * * Preliminary [6] PRD 108 032008	065 * EFT formalism is employed at different levels of experimental analyses

Figure 11: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with scalar bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. Vertical format.

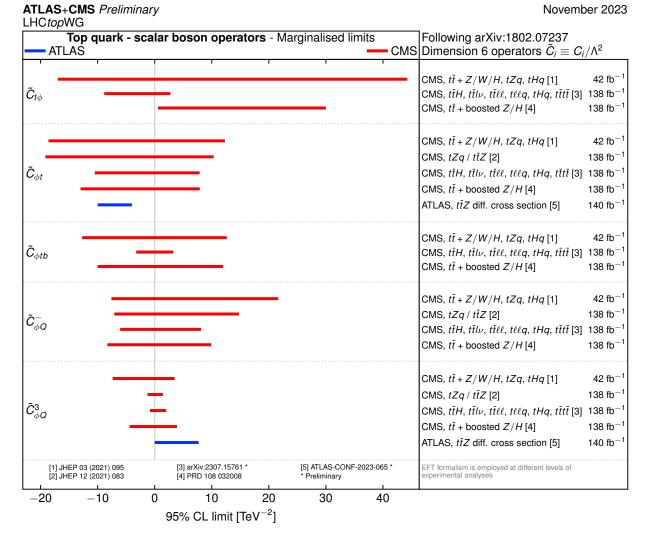


Figure 12: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with scalar bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. Vertical format.

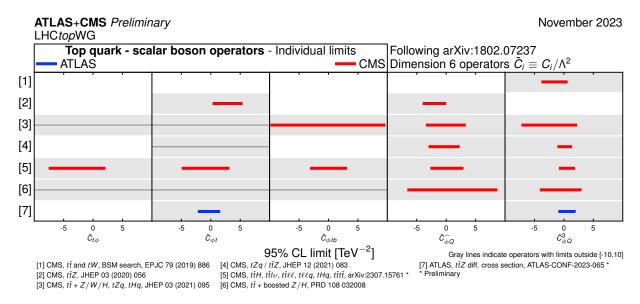


Figure 13: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with scalar bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. Horizontal format.

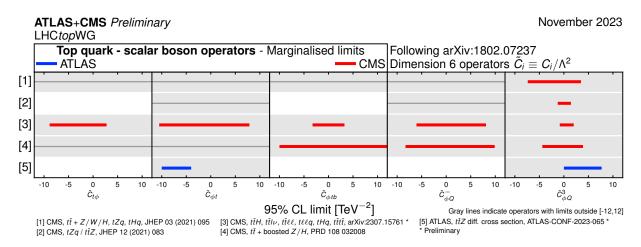


Figure 14: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with scalar bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. Horizontal format.

	ector boson operators		Following arXiv:1802.07237	
- ATLAS	ATLAS+CMS	CN	IS Dimension 6 operators $\hat{C}_i \equiv$	C_i/Λ^2
			CMS, <i>tīZ</i> [4]	78 fb ⁻¹
			CMS, $t\bar{t} + Z/W/H$, tZq , tHq [6]	42 fb
			CMS, $tZq / t\bar{t}Z$ [7]	138 fb
\tilde{C}_{tZ}			CMS, $t\bar{t}\gamma$ [8]	137 fb ⁻¹
			CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$	
			CMS, $t\bar{t}$ + boosted Z/H [12]	138 fb ⁻¹
ãИ			CMS, <i>tīZ</i> [4]	78 fb ⁻¹
$ ilde{C}_{tZ}^{[I]}$			CMS, $t\bar{t}\gamma$ [8]	137 fb ⁻¹
\tilde{C}_{tB} —	·····		ATLAS, $t\bar{t}Z$ diff. cross section [14]	140 fb ⁻¹
			CMS, $t\bar{t}$ and tW , BSM search [3]	36 fb ⁻¹
			ATLAS+CMS, W helicity [5]	20+20 fb ⁻¹
			CMS, $t\bar{t} + Z/W/H$, tZq , tHq [6]	42 fb ⁻¹
$ ilde{C}_{tW}$			CMS, <i>tZq / ttZ</i> [7]	138 fb ⁻¹
			ATLAS, Top polarization [10]	139 fb ⁻¹
			CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$	[11]138 fb ⁻¹
			CMS, $t\bar{t}$ + boosted Z/H [12]	138 fb ⁻¹
			ATLAS, $t\bar{t}Z$ diff. cross section [14]	140 fb ⁻¹
$ ilde{C}^{[I]}_{tW}$			ATLAS, Top polarization [10] ATLAS, <i>tīZ</i> diff. cross section [14]	139 fb ⁻¹ 140 fb ⁻¹
				14010
ã			CMS, $t\bar{t} + Z/W/H$, tZq , tHq [6]	42 fb ⁻¹
<i>С_{ьw}</i>			CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$ CMS, $t\bar{t}$ + boosted Z/H [12]	138 fb ⁻¹ 138 fb ⁻¹
Ĉ _{tG} ∕g _S			ATLAS, <i>tt ℓ</i> + jets boosted [9]	139 fb ⁻¹
	<u> </u>		CMS, tt dilepton [1]	36 fb ⁻¹
	-		CMS, $t\bar{t}$ spin correlations [2]	36 fb ⁻¹
			CMS, $t\bar{t}$ and tW , BSM search [3]	36 fb ⁻¹
\tilde{C}_{tG}			CMS, $t\bar{t} + Z/W/H$, tZq , tHq [6]	42 fb ⁻¹
	<u> </u>		CMS, $t\bar{t}H$, $t\bar{t}l\nu$, $t\bar{t}\ell\ell$, $t\ell\ell q$, tHq , $t\bar{t}t\bar{t}$	[11]138 fb ⁻¹
			ATLAS, <i>tī</i> rapidity asymmetry [13]	139 fb ⁻¹
	+		ATLAS, $t\bar{t}Z$ diff. cross section [14]	140 fb ⁻¹
$ ilde{C}_{tG}^{[I]}$			CMS, $t\bar{t}$ spin correlations [2]	36 fb ⁻¹
[1] JHEP 02 (2019) 149 [2] PRD 100 (2019) 072002 [3] EPJC 79 (2019) 886 [4] JHEP 03 (2020) 056 [5] JHEP 08 (2020) 051	[6] JHEP 03 (2021) 095 [7] JHEP 12 (2021) 083 [8] JHEP 05 (2022) 091 [9] JHEP 06 (2022) 063 [10] JHEP 11 (2022) 040	[11] arXiv:2307.15761 * [12] PRD 108 032008 [13] JHEP 08 (2023) 077 [14] ATLAS-CONF-2023-065 * * Preliminary	EFT formalism is employed at different levels of experimental analyses	
_4 _2	0 2	4	6	

Figure 15: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with vector bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237.

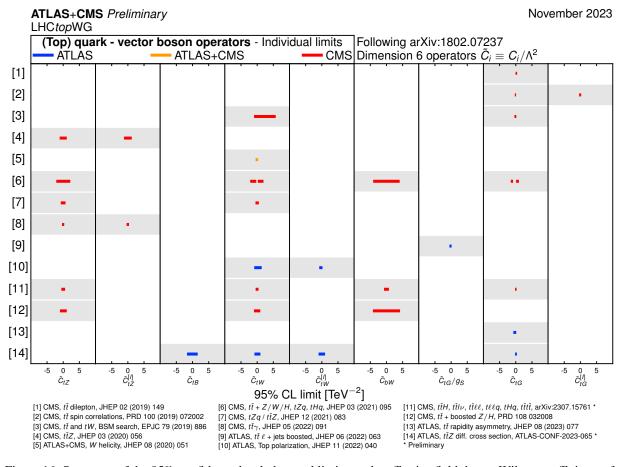
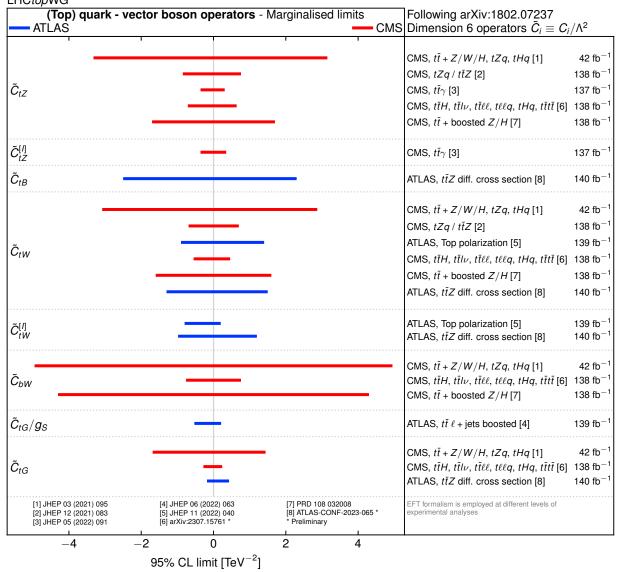


Figure 16: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with vector bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. Horizontal format.



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Figure 17: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with vector bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. Vertical format.

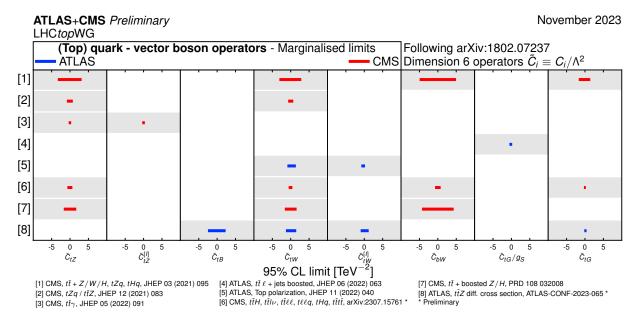


Figure 18: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to (top) quark interaction with vector bosons, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. Horizontal format.

ATLAS	r-fermion operators - Individual limits	Following arXiv:1802.07237 CMS Dimension 6 operators $\tilde{C}_i \equiv C_i / \Lambda^2$
	· · · · · · · · · · · · · · · · · · ·	CMS, $t\bar{t}t\bar{t}$ [1] 36 fb ⁻
\tilde{C}_{tt}^{1}		CMS, tĨH, tĨ/v, tĨℓℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻
Ĉ₁ αι		CMS, tītī [1] 36 fb ⁻ CMS, tīH, tīlv, tītt, tttq, tHq, tītī [6] 138 fb ⁻
Ĉ ¹ aa		CMS, tĨtĨ [1] 36 fb ⁻ CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻
\tilde{C}^8_{Ql}		CMS, tÎtî [1] 36 fb ⁻ CMS, tÎH, tÎlv, tÎtê, têlq, tHq, tÎtî [6] 138 fb ⁻
$\bar{C}_{Ql}^{3(\ell)}$		CMS, tĨ + Z/W/H, tZq, tHq [2] 42 fb- CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb-
$\tilde{C}_{Ql}^{-(\ell)}$		CMS, tĨ + Z/W/H, tZq, tHq [2] 42 fb ⁻ CMS, tĨH, tĨlv, tĨℓℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻
$\tilde{C}_{Qe}^{(\ell)}$		CMS, tĨ + Z/W/H, tZq, tHq [2] 42 fb ⁻ CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻
$\tilde{C}_{tl}^{(\ell)}$		CMS, tī + Z/W/H, tZq, tHq [2] 42 fb ⁻ CMS, tī H, tī lv, tī tt, tī tt, tt (6] 138 fb ⁻
$\tilde{C}_{te}^{(\ell)}$		CMS, tī + Z/W/H, tZq, tHq [2] 42 fb ⁻ CMS, tĨH, tĨlv, tĨℓℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻
C ^{S(ℓ)}		CMS, tī + Z/W/H, tZq, tHq [2] 42 fb ⁻
$\tilde{C}_{t}^{T(\ell)}$		CMS, tÎH, tÎlv, tÎtê, têlq, tHq, tÎtî [6] 138 fb CMS, tÎ + Z/W/H, tZq, tHq [2] 42 fb
ь _t .,		CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻ ATLAS, tĨ + jet energy asymmetry [3] 139 fb ⁻
Ĉ ¹¹ Qq	Ŧ	ATLAS, it + jet energy asymmetry [5] 139 fb CMS, třH, třl/v, třt4, tt42, tH4, třtří [6] 138 fb ⁻ ATLAS, tř rapidity asymmetry [7] 139 fb ⁻ ATLAS, třZ diff. cross section [8] 140 fb ⁻
	<u> </u>	ATLAS, tt + jet energy asymmetry [3] 139 fb ⁻ ATLAS, tt all-hadronic boosted [5] 139 fb ⁻
5 ¹⁸ Og	<u> </u>	CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb- ATLAS, tĨ rapidity asymmetry [7] 139 fb-
		ATLAS, <i>ti</i> Z diff. cross section [8] 140 fb ⁻
Ē ³¹		CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [6] 138 fb ⁻ ATLAS, tĨ rapidity asymmetry [7] 139 fb ⁻ ATLAS, tĨZ diff. cross section [8] 140 fb ⁻
Ĉ ¹ tq	Ŧ	ATLAS, tĨ + jet energy asymmetry [3] 139 fb CMS, tĨ H, tĨ Iν, tĨ LI, tếL (tếL q, tHq, tĨ tĨ [6] 138 fb ATLAS, tĨ rapidity asymmetry [7] 139 fb
		ATLAS, tt + jet energy asymmetry [3] 139 fb- ATLAS, tt + jets boosted [4] 139 fb-
\tilde{C}_{tq}^{8}	<u>_</u>	ATLAS, tt all-hadronic boosted [5] 139 fb-
	—	CMS, tÎH, tÎlv, tĨlt, tĨtt, tĨtt, tĨtt, [6] 138 fb ⁻ ATLAS, tĪ rapidity asymmetry [7] 139 fb ⁻
Ĉ _{tu}	Ξ	ATLAS, t ⁷ + jet energy asymmetry [3] 139 fb ⁻ ATLAS, t ⁷ rapidity asymmetry [7] 139 fb ⁻ ATLAS, t ⁷ Z diff. cross section [8] 140 fb ⁻
Ĉ ¹ td		ATLAS, tł̃ rapidity asymmetry [7] 139 fb ⁻ ATLAS, tł̃Z diff. cross section [8] 140 fb ⁻
		ATLAS, tł + jet energy asymmetry [3] 139 fb-
28 Tu		ATLAS, tt all-hadronic boosted [5] 139 fb ⁻ ATLAS, tt rapidity asymmetry [7] 139 fb ⁻ ATLAS, ttZ diff. cross section [8] 140 fb ⁻
		ATLAS, tt all-hadronic boosted [5] 139 fb-
Ĉ ⁸ td		ATLAS, tī rapidity asymmetry [7] 139 fb ⁻ ATLAS, tī Z diff. cross section [8] 140 fb ⁻
õe a	<u> </u>	ATLAS, <i>t</i> t all-hadronic boosted [5] 139 fb ⁻ ATLAS, <i>t</i> t rapidity asymmetry [7] 139 fb ⁻
Ĉ ⁸ od		ATLAS, t ² rapidity asymmetry [7] 139 fb ⁻ ATLAS, t ² Z diff. cross section [8] 140 fb ⁻
Ĉ ⁸ αυ	=	ATLAS, tř all-hadronic boosted [5] 139 fb- ATLAS, tř rapidity asymmetry [7] 139 fb- ATLAS, třZ diff. cross section [8] 140 fb-
Ĉ1 _{au}	+	ATLAS, <i>t</i> ⁷ rapidity asymmetry [7] 139 fb ⁻ ATLAS, <i>t</i> ⁷ Z diff. cross section [8] 140 fb ⁻
Ĉ¹ Qd	=	ATLAS, <i>tī</i> rapidity asymmetry [7] 139 fb ⁻ ATLAS, <i>tī</i> Z diff. cross section [8] 140 fb ⁻
	+	ATLAS, <i>ti</i> all-hadronic boosted [5] 139 fb ⁻
\tilde{C}^{38}_{Qq}	<u> </u>	CMS, tî <i>H</i> , tî <i>lv</i> , tî <i>lt</i> , t <i>ît</i> , t <i>ît</i> , tî <i>t</i> , [6] 138 fb ⁻ ATLAS, tî rapidity asymmetry [7] 139 fb ⁻ ATLAS, tî <i>Z</i> diff. cross section [8] 140 fb ⁻
[1] JHEP 11 (2019) 082 [2] JHEP 03 (2021) 095 [3] EPJC 82 (2022) 374	[4] JHEP 06 (2022) 063 [7] JHEP 08 (2023) 073 [5] JHEP 04 (2023) 80 [8] ATLAS-CONF-2023 [6] aX/k-2307.15761 * Poliminary	
	-5.0 -2.5 0.0 2.5 5.0 7.5	

Figure 19: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to four-fermion interactions, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. Vertical format.

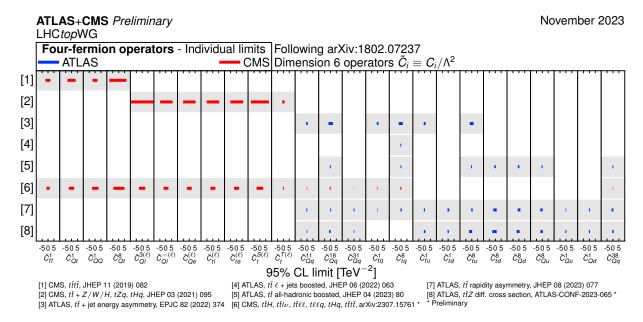


Figure 20: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to four-fermion interactions, as obtained by the ATLAS and CMS Collaborations. The results are reported as individual constraints assuming new physics contributions from one specific operator at a time. Interpretations use the SMEFT framework and the Warsaw basis. The formalism is employed at different levels of the experimental analyses, from the interpretation of measured observables to a comparison of the data to simulations at the detector level. Most interpretations follow the LHCtopWG recommendations from arXiv:1802.07237. Horizontal format.

ATLAS+CMS	Preliminarv

November 2023

ATLAS	Four-fermion operators - Marginalised limits	Following arXiv:1802.07237 Dimension 6 operators $\tilde{C}_i \equiv C_{i/i}$	/Λ ²
\tilde{C}_{tt}^{1}		CMS, tītī [1] CMS, tīH, tīlv, tītl, tllq, tHq, tītī [4]	36 fb⁻ 138 fb⁻
\tilde{C}^{1}_{Qt}		CMS, tĨtĨ [1] CMS, tĨH, tĨlv, tĨtℓ, tℓℓq, tHq, tĨtĨ [4]	36 fb⁻ 138 fb⁻
\tilde{C}^{1}_{QQ}		CMS, tītī [1] CMS, tīH, tīlv, tītl, ttlq, tHq, tītī [4]	36 fb⁻ 138 fb⁻
\tilde{C}^8_{Qt}		CMS, tītī [1] CMS, tīH, tīlv, tītl, ttlq, tHq, tītī [4]	36 fb⁻ 138 fb⁻
$\tilde{C}^{3(\ell)}_{Ql}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī̄ [4]	42 fb⁻ 138 fb⁻
$ ilde{C}_{Ql}^{-(\ell)}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	42 fb 138 fb
$ ilde{C}^{(\ell)}_{Qe}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	42 fb 138 fb
$ ilde{C}_{tl}^{(\ell)}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	42 fb 138 fb
$ ilde{C}_{te}^{(\ell)}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	42 fb 138 fb
$ ilde{C}_t^{S(\ell)}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	42 fb 138 fb
$ ilde{C}_t^{T(\ell)}$		CMS, tī̄ + Z/W/H, tZq, tHq [2] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	42 fb ⁻ 138 fb ⁻
\tilde{C}^{11}_{Qq}	ŧ	CMS, tĪH, tĪlv, tĪtl, ttlq, tHq, tĪtĪ [4] ATLAS, tĪZ diff. cross section [5]	138 fb 140 fb
\tilde{C}^{18}_{Qq}	_	CMS, tīH, tīlv, tīll, tllq, tHq, tītī [4] ATLAS, tīZ diff. cross section [5]	138 fb⁻ 140 fb⁻
\tilde{C}^{31}_{Qq}	<u>+</u>	CMS, tīH, tīlv, tītl, ttlq, tHq, tītī [4] ATLAS, tīZ diff. cross section [5]	138 fb ⁻ 140 fb ⁻
\tilde{C}_{tq}^1	+	CMS, tĪH, tĪlv, tĪtl, ttlq, tHq, tĪtĪ [4]	138 fb⁻
\tilde{C}^8_{tq}	=	ATLAS, tī ℓ + jets boosted [3] CMS, tīH, tīlν, tīℓℓ, tℓℓq, tHq, tītī [4]	139 fb⁻ 138 fb⁻
\tilde{C}_{tu}^1		ATLAS, tTZ diff. cross section [5]	140 fb ⁻
\tilde{C}^{1}_{td}		ATLAS, tTZ diff. cross section [5]	140 fb ⁻
\tilde{C}^8_{tu}		ATLAS, tTZ diff. cross section [5]	140 fb-
\tilde{C}^8_{td}		ATLAS, tTZ diff. cross section [5]	140 fb ⁻
\tilde{C}^8_{Qd}		ATLAS, tTZ diff. cross section [5]	140 fb-
\tilde{C}^8_{Qu}		ATLAS, <i>tTZ</i> diff. cross section [5]	140 fb-
\tilde{C}^{1}_{Qu}	<u> </u>	ATLAS, <i>t_tZ</i> diff. cross section [5]	140 fb ⁻
\tilde{C}^{1}_{Qd}	+	ATLAS, <i>tīZ</i> diff. cross section [5]	140 fb ⁻
\tilde{C}^{38}_{Qq}	+	CMS, tīH, tīlv, tītl, ttlq, tHq, tītī [4] ATLAS, tīZ diff. cross section [5]	138 fb 140 fb
[1] JHEP 11 (2 [2] JHEP 03 (2	2019) 082 [3] JHEP 06 (2022) 063 [5] ATLAS-CONF-2023-065 * 2021) 095 [4] arXiv:2307.15761 * * Preliminary	EFT formalism is employed at different levels of experimental analyses	

Figure 21: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to four-fermion interactions, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. Vertical format.

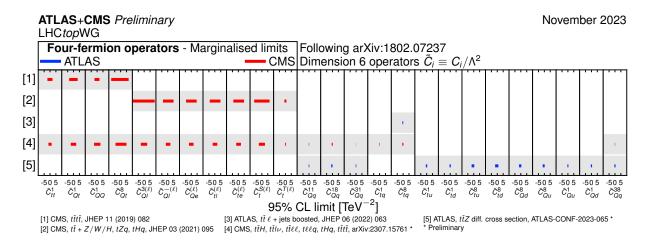


Figure 22: Summary of the 95% confidence level observed limits on the effective field theory Wilson coefficients of the dimension-6 operators related to four-fermion interactions, as obtained by the ATLAS and CMS Collaborations. The results are reported as marginalised constraints, treating all Wilson coefficients contributing to a given process as free parameters. The effect of a given Wilson coefficient is considered in multiple processes, where indicated in the references, and across multiple bins of differential measurements. Each row presents all the marginalised constraints obtained from a single fit. Horizontal format.

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