

New Physics with missing energy at future lepton colliders - Snowmass White Paper

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Abstract

Two models that extend the particle content of the SM and provide dark matter candidates, namely the Inert Doublet Model and the Two-Higgs Doublet model with additional pseudoscalar, are confronted with current experimental and theoretical constraints and predictions for production cross sections for various standard pair-production modes within these models at future lepton colliders are presented.

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I. INTRODUCTION

We discuss two new physics models that extend the Standard Model (SM) particle sector by additional scalars and provide a dark matter candidate, namely the Inert Doublet Model (IDM) and the Two-Higgs Doublet Model with additional pseudoscalar (THDMa). Both models are confronted with current theoretical and experimental constraints. From the theoretical side, these include the minimization of the vacuum as well as the requirement of vacuum stability and positivity. We also require perturbative unitarity to hold, and perturbativity of the couplings at the electroweak scale.

Experimental bounds include the agreement with current measurements of the properties of the 125 GeV resonance discovered by the LHC experiments, as well as agreement with the null-results from searches for additional particles at current or past colliders. We also confront the models with bounds from electroweak precision observables (via S , T , U parameters), B-physics observables ($B \rightarrow X_s \gamma$, $B_s \rightarrow \mu^+ \mu^-$, ΔM_s), as well as agreement with astrophysical observables (relic density and direct detection bounds). We use a combination of private and public tools in these analyses, where the latter include HiggsBounds [1], HiggsSignals [2], 2HDMC [3], SPheno [4], Sarah [5], micrOMEGAs [6, 7], and MadDM [8]. Experimental numbers are taken from [9, 10] for electroweak precision observables, [11] for $B_s \rightarrow \mu^+ \mu^-$, [12] for ΔM_s and [13] and [14] for relic density and direct detection, respectively. Bounds from $B \rightarrow X_s \gamma$ are implemented using a fit function from [15, 16]. Predictions for production cross sections shown here have been obtained using Madgraph5 [17].

II. THE INERT DOUBLET MODEL

A. The model

The Inert Doublet Model (IDM) [18–20] is an intriguing new physics model. In this model, the SM scalar sector is enhanced by an additional $SU(2) \times U(1)$ gauge doublet ϕ_D . As in the SM, the first doublet ϕ_S contains the SM-like Higgs boson h , while the inert one ϕ_D contains four scalar states H , A and H^\pm . A discrete exact \mathbb{Z}_2 symmetry is introduced with the following transformation properties

$$\phi_S \rightarrow \phi_S, \phi_D \rightarrow -\phi_D, \text{ SM} \rightarrow \text{SM}. \quad (1)$$

The additional doublet does not acquire a vacuum expectation value (vev) and does not couple to fermions. Electroweak symmetry breaking is as in the SM. The symmetry also insures that the lightest particle of ϕ_D is stable, making this a good dark matter candidate.

The scalar potential of the model is given by

$$V = -\frac{1}{2} \left[m_{11}^2 (\phi_S^\dagger \phi_S) + m_{22}^2 (\phi_D^\dagger \phi_D) \right] + \frac{\lambda_1}{2} (\phi_S^\dagger \phi_S)^2 + \frac{\lambda_2}{2} (\phi_D^\dagger \phi_D)^2 + \lambda_3 (\phi_S^\dagger \phi_S) (\phi_D^\dagger \phi_D) + \lambda_4 (\phi_S^\dagger \phi_D) (\phi_D^\dagger \phi_S) + \frac{\lambda_5}{2} \left[(\phi_S^\dagger \phi_D)^2 + (\phi_D^\dagger \phi_S)^2 \right]. \quad (2)$$

The model features 7 free parameters, which we chose in the so-called physical basis [21]

$$v, M_h, M_H, M_A, M_{H^\pm}, \lambda_2, \lambda_{345}, \quad (3)$$

with $\lambda_{345} \equiv \lambda_3 + \lambda_4 + \lambda_5$. As two parameters (vev v and $M_h \sim 125$ GeV) are fixed by experimental measurements, we end up with a total number of 5 free parameters. Here, we consider the case where H is the dark matter candidate, which implies $M_{A, H^\pm} \geq M_H$.¹

The model is subject to a large number of theoretical and experimental constraints, discussed at length e.g. in [21–25]. In the scan, we make use of the publicly available tools `2HDMC` [3], `HiggsBounds-5.10.1` [1, 26–29], `HiggsSignals-2.6.2` [2, 30], as well as `micrOMEGAs_5.2.4` [7]. Cross sections are calculated using `Madgraph5` [17] with a UFO input file from [31]². Experimental values are taken from `GFitter` [10, 34], and the Planck [35] and XENON1T [14] experiments. Direct collider searches as well as agreement with the 125 GeV coupling strength measurements are implemented via `HiggsBounds` and `HiggsSignals`, where we also compare to the total width upper limit [36] and invisible branching ratio [37] of h . Finally, recast results from a LEP-SUSY search [38] were included.

B. Current Status

The experimental and theoretical constraints lead to a large reduction of the allowed parameter space. As an example, the masses are usually quite degenerate, as can be seen from figure 1. This is caused by an interplay of electroweak constraints and theoretical requirements on the potential. We also consider the case when $M_H \leq M_h/2$, where constraints from $h \rightarrow$ invisible start to play an important role and an interesting interplay arises, between bounds from signal strength measurements, that require $|\lambda_{345}|$ to be rather small $\lesssim 0.3$, and bounds from dark matter relic density, where too low values of that parameter lead to small annihilation cross sections and therefore too large relic density values. The resulting parameter space is shown in figure 2. In [21], it was found that this in general leads to a lower bound of $M_H \sim 50$ GeV, with exceptions presented in [25].

¹ Note that the new scalars in the IDM do not have CP quantum numbers, as they do not couple to fermions. In the subsequent discussion, we can replace $H \longleftrightarrow A$ if we simultaneously use $\lambda_5 \longleftrightarrow -\lambda_5$. All phenomenological considerations are identical for these cases.

² Note the official version available at [?] exhibits a wrong CKM structure, leading to false results for processes involving electroweak gauge bosons radiated off quark lines. In our implementation, we corrected for this. Our implementation corresponds to the expressions available from [33].

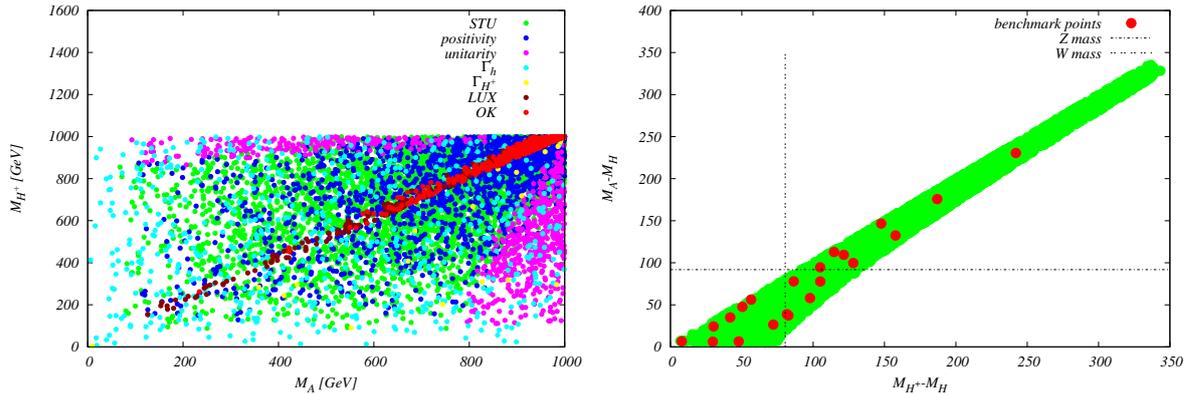


FIG. 1. Masses are requested to be quite degenerate after all constraints have been taken into account. *Left*: In the (M_A, M_{H^\pm}) plane (taken from [21]). *Right*: In the $(M_{H^\pm} - M_H, M_A - M_H)$ plane (taken from [24]).

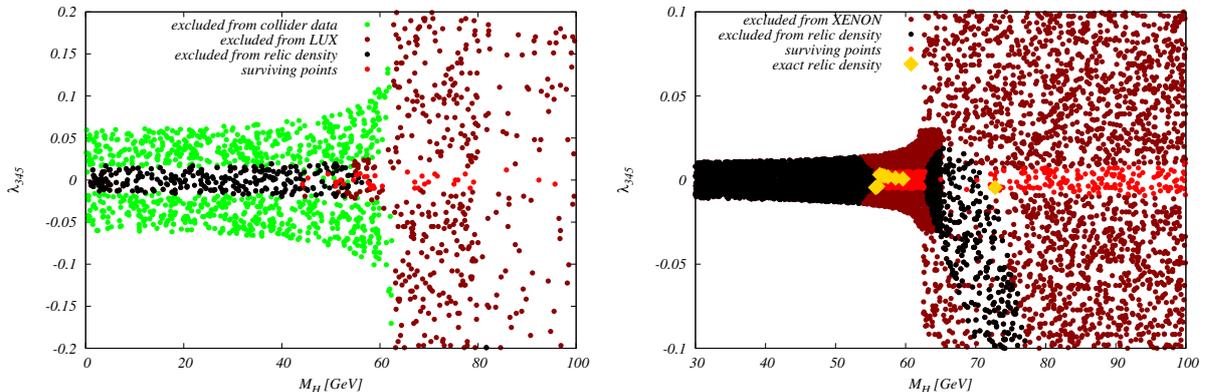


FIG. 2. Interplay of signal strength and relic density constraints in the (M_H, λ_{345}) plane. *Left*: Using LUX constraints [39], bounds labelled "excluded from collider data" have been tested using HiggsBounds and HiggsSignals (taken from [21]). *Right*: Using XENON1T results, with golden points labelling those points that produce exact relic density (taken from [22]).

C. Discovery prospects at ILC and CLIC

So far, no publicly available search exists that investigates the IDM parameter space with actual collider data. The discovery potential of ILC and CLIC was investigated in [40–45] for several benchmark points proposed in [24], for varying center-of-mass energies from 250 GeV up to 3 TeV. We concentrated on AH and H^+H^- production with $A \rightarrow ZH$ and $H^\pm \rightarrow W^\pm H$, where the electroweak gauge bosons subsequently decay leptonically. For event generation, we used WHizard 2.2.8 [46, 47], with an interface via SARAH [48] and SPheno 4.0.3 [4, 49] for model implementation. For CLIC results energy spectra [50] were

also taken into account.

For the production modes above, we considered leptonic decays of the electroweak gauge bosons. The investigated final states were

$$e^+ e^- \rightarrow \mu^+ \mu^- + \cancel{E}, \quad e^+ e^- \rightarrow \mu^\pm e^\mp + \cancel{E}$$

for HA and $H^+ H^-$ production, respectively. For a more accurate description, we did not specify the intermediate states in the event generation, which in turn means all processes leading to the above signatures were taken into account, including interference between the contributing diagrams. This includes final states, where the missing energy stems from additional neutrinos, *e.g.* from τ^\pm decays. Event selection was performed using a set of preselection cuts as well as boosted decision trees, as implemented in the TMVA toolkit [51]. Results for the ILC running at 250 GeV and 500 GeV, and the first CLIC stage at 380 GeV are shown in figure 3. The expected discovery reach of 500 GeV ILC extends up to neutral scalar mass sum of 330 GeV up to charged scalar masses of 200 GeV. Results for the

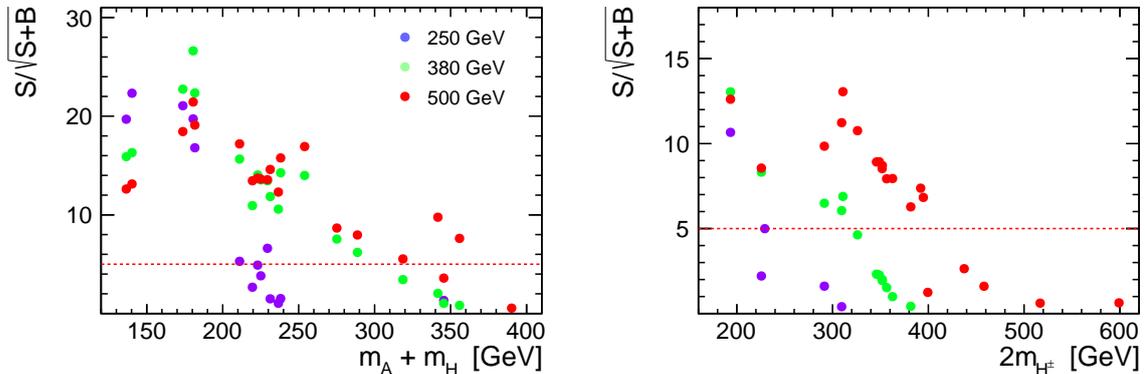


FIG. 3. Significance of the deviations from the Standard Model predictions expected for 1 ab^{-1} of data collected at centre-of-mass energy of 250 GeV, 380 GeV and 500 GeV, for: (left) events with two muons in the final state ($\mu^+ \mu^-$) as a function of the sum of neutral inert scalar masses and (right) events with an electron and a muon in the final state ($e^+ \mu^-$ or $e^- \mu^+$) as a function of twice the charged scalar mass.

discovery reach of CLIC, including center-of-mass energies of 1.5 TeV and 3 TeV, are shown in figure 4. In general, production cross sections $\gtrsim 0.5 \text{ fb}$ seem to be accessible, where best prospects for the considered benchmark points are given for 380 GeV or 1.5 TeV center-of-mass energies. Along similar lines, mass sums up to 1 TeV seem accessible, where the $\mu^\pm e^\mp$ channel seems to provide a larger discovery range in general.

For leptonic signatures, the sensitivity of high energy $e^+ e^-$ colliders to pair-production of IDM scalars is limited by the production cross section in the considered channel. The signal cross section for the charged scalar pair-production increases by about an order of magnitude when the semi-leptonic final state is considered, *i.e.* when the hadronic decay of one of the

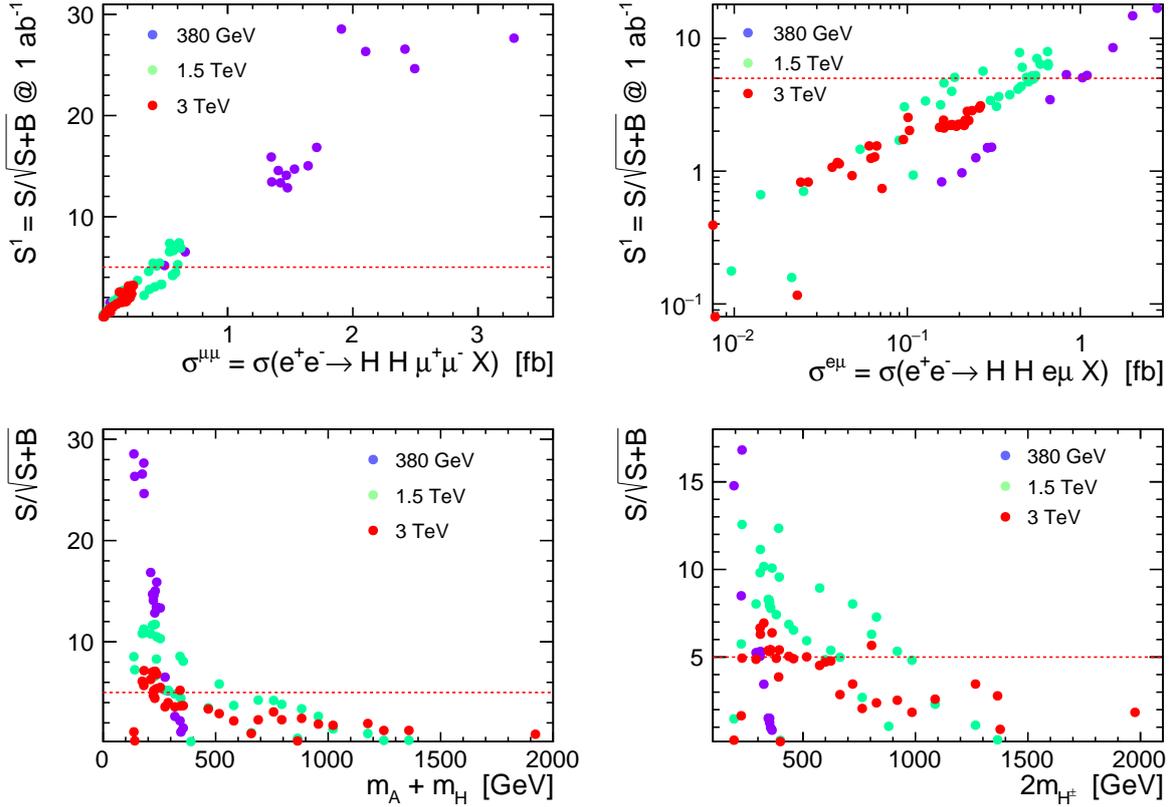


FIG. 4. Discovery prospects at CLIC for the IDM in $\mu^+\mu^- + \cancel{E}$ (left) and $\mu^\pm e^\mp + \cancel{E}$ (right) final states, as a function of the respective production cross-sections (top) and mass sum of the produced particles (bottom). Taken from [40].

W bosons is considered. The sensitivity of high energy CLIC to H^+H^- production with the semi-leptonic final state is shown in Fig. 5. The accessible scalar mass range increases by about a factor of two, to about 2 TeV [43–45]. It is also interesting to note that highest signal observation significance is obtained for scenarios with scalar mass difference of 20 to 60 GeV.

D. Sensitivity comparison at future colliders

After a dedicated analysis of the IDM benchmarks in the high energy CLIC environment, an important question is whether other current or future collider options provide similar or better discovery prospects. For the same set of benchmarks[24, 40], production cross sections for a variety of processes have been presented in [25], including VBF-type topologies. Cross sections were calculated using `Madgraph5`. We concentrate on production modes at muon colliders and refer to [25] regarding processes at proton-proton machines. We

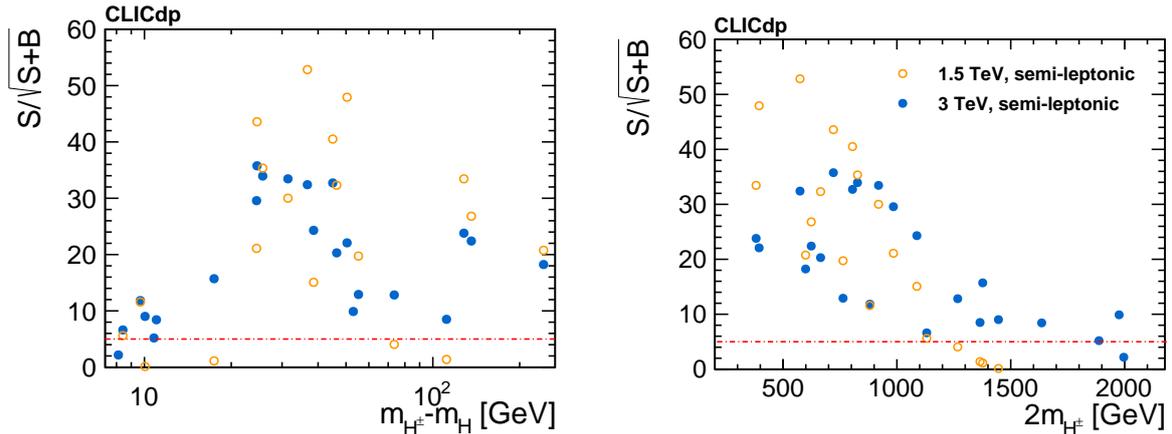


FIG. 5. Expected statistical significance of IDM charged scalar pair-production observation as a function of the IDM scalar mass difference (left) and of the total mass of the produced IDM scalars (right). Results are presented for CLIC running at 1.5 TeV (orange circles) and 3 TeV (blue points). The red horizontal lines indicate the 5σ threshold. Figure taken from [45].

collider	cm energy [TeV]	$\int \mathcal{L}$	σ_{1000} [fb]
ee	3	5 ab^{-1}	0.2
$\mu\mu$	10	10 ab^{-1}	0.1
$\mu\mu$	30	90 ab^{-1}	0.01

TABLE I. Collider parameters used in the discovery reach study performed in [25]. Collider specifications have been taken from [52] for the muon collider. The last column denotes the minimal cross section required to produce 1000 events using full target luminosity.

list the considered collider types and nominal center-of-mass energies as well as integrated luminosities in table I.

A scenario is called "realistic" when 1000 events can be produced using target luminosity and center-of-mass energies as specified above. Obviously, more detailed studies, including both background contribution and detector response simulation, are necessary to assess the actual collider reach.

We here concentrate on production at future muon colliders:

$$\mu^+ \mu^- \rightarrow \nu_\mu \bar{\nu}_\mu AA, \quad \mu^+ \mu^- \rightarrow \nu_\mu \bar{\nu}_\mu H^+ H^-.$$

which corresponds to VBF-like production modes. Here again intermediate states are not specified, so in fact several diagrams contribute which not all have a typical VBF topology. See appendix B and C of [25] for details.

collider	all others	AA	$AA + VBF$
HL-LHC	1 TeV	200-600 GeV	500-600 GeV
HE-LHC	2 TeV	400-1400 GeV	800-1400 GeV
FCC-hh	2 TeV	600-2000 GeV	1600-2000 GeV
CLIC, 3 TeV	2 TeV	-	300-600 GeV
$\mu\mu$, 10 TeV	2 TeV	-	400-1400 GeV
$\mu\mu$, 30 TeV	2 TeV	-	1800-2000 GeV

TABLE II. Sensitivity of different collider options specified in table I, using the "realistic" criterium of 1000 generated events in the specific channel. Shown are minimal and maximal mass scales that are reachable. Numbers for CLIC correspond to results from detailed investigations [40, 41]. Table taken from [53].

Figure 6 shows the production cross sections as a function of the mass sum of produced particles for various center-of-mass energies. Understanding the behaviour of the VBF-induced channels is non-trivial; this can be attributed to the fact that more diagrams, apart from the naive pair-production process, contribute. Large jumps between cross-section predictions for scenarios with similar mass scales can be traced back mainly to a fine-tuned cancellation of various contributing diagrams, as discussed in greater detail in [25].

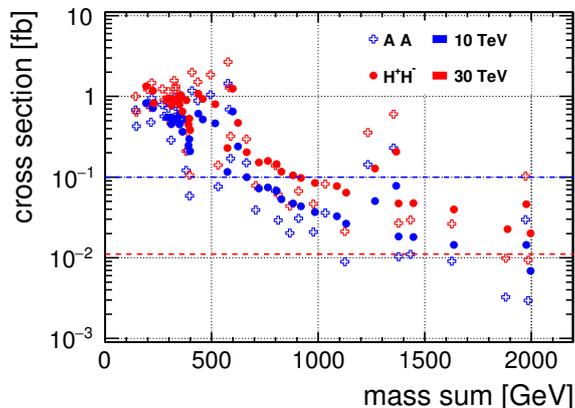


FIG. 6. Cross sections as a function of the mass sum at muon colliders in the VBF-type production mode. Taken from [25].

The summary of sensitivities in terms of mass scales is given in table II, where for completeness the reach for proton colliders as estimated in the above work has been added. It is seen that especially for AA production the VBF mode at both proton and muon colliders serves to significantly increase the discovery reach of the respective machine. Using the simple

counting criterium above, we can furthermore state that a 27 TeV proton-proton machine has a similar reach as a 10 TeV muon collider, while 100 TeV FCC-hh would correspond to a 30 TeV muon-muon machine. Obviously, detailed investigations including SM background are needed to give a more realistic estimate of the respective collider reach.

III. THDMA

The THDMA is a type II two-Higgs-doublet model that is extended by an additional pseudoscalar a mixing with the "standard" pseudoscalar A of the THDM. In the gauge-eigenbasis, the additional scalar serves as a portal to the dark sector, with a fermionic dark matter candidate, denoted by χ . More details can e.g. be found in [54–60].

The model contains the following particles in the scalar and dark matter sector: h, H, H^\pm, a, A, χ . It depends on 12 additional new physics parameters

$$v, m_h, m_H, m_a, m_A, m_{H^\pm}, m_\chi; \cos(\beta - \alpha), \tan\beta, \sin\theta; y_\chi, \lambda_3, \lambda_{P_1}, \lambda_{P_2},$$

where v and either m_h or m_H are fixed by current measurements in the electroweak sector.

We here report on results of a scan that allows all of the above novel parameters float in specific predefined ranges [60]. In such a scenario, it is not always straightforward to display bounds from specific constraints in 2-dimensional planes. Two examples for scenarios where this is possible are shown in figure 8. The first plot shows bounds in the $(m_{H^\pm}, \tan\beta)$ plane from B-physics observables. The result is similar to a simple THDM, and shows that in general low masses $m_{H^\pm} \lesssim 800$ GeV as well as values $\tan\beta \lesssim 1$ are excluded. The second plot displays the relic density as a function of the mass difference $m_a - 2m_\chi$. Here, a behaviour can be observed that is typical in many models with dark matter candidates: in the region where this mass difference remains small, relic density annihilates sufficiently to stay below the observed relic density bound. On the other hand, too large differences lead to values $\Omega h_c \gtrsim 0.12$ and therefore are forbidden from dark matter considerations.

Finally, it was investigated which cross-section values would still be feasible for points that fulfill all constraints [60] at e^+e^- colliders. We here concentrate on signatures that include missing energy and therefore do not exist in a THDM without a portal to the dark sector. Processes like $e^+e^- \rightarrow hA, ha$ are suppressed due to alignment, which makes $e^+e^- \rightarrow HA, Ha$ the most interesting channel that contains novel signatures. Due to the interplay of B-physics and electroweak constraints, such points typically have mass scales $\gtrsim 1$ TeV. Therefore production cross sections for an e^+e^- collider with a center-of-mass energy of 3 TeV are of interest. The corresponding production cross sections are shown in figure 9, which displays predictions for $t\bar{t}t\bar{t}$ and $t\bar{t} + \cancel{E}$ final states using a factorized approach. There is a non-negligible number of points where the second channel is dominant. A "best" point with a large rate for $t\bar{t} + \cancel{E}_\perp$ has been presented in [60] and is repeated here for completeness

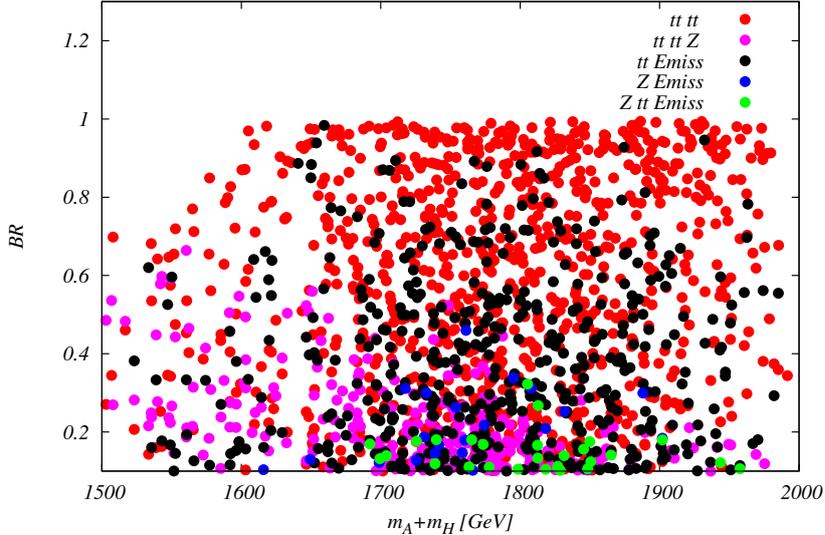


FIG. 7. Branching ratios into various final states for AH production, as a function of the mass sum.

$$\begin{aligned}
 \sin \theta &= -0.626, \quad \cos(\beta - \alpha) = 0.0027, \quad \tan \beta = 3.55 \\
 m_H &= 643 \text{ GeV}, \quad m_A = 907 \text{ GeV}, \quad m_{H^\pm} = 814 \text{ GeV}, \\
 m_a &= 653 \text{ GeV}, \quad m_\chi = 277 \text{ GeV}, \\
 y_\chi &= -1.73, \quad \lambda_{P_1} = 0.18, \quad \lambda_{P_2} = 2.98, \quad \lambda_3 = 8.63.
 \end{aligned}
 \tag{4}$$

For this point, all width/ mass ratios are $\lesssim 6\%$. In addition, branching ratios for various final states as a function of the mass sum for the HA channel are given in figure 7.

IV. CONCLUSION AND OUTLOOK

In this work, we presented two models that extend the particle content of the SM and also provide at least one dark matter candidate. We have presented production cross sections for various standard pair-production modes within these models; for the IDM, we have given an estimate of mass range that can be reached based on a simple counting criterium. A more dedicated investigation of the corresponding signatures, including background simulation and cut optimization, is in the line of future work.

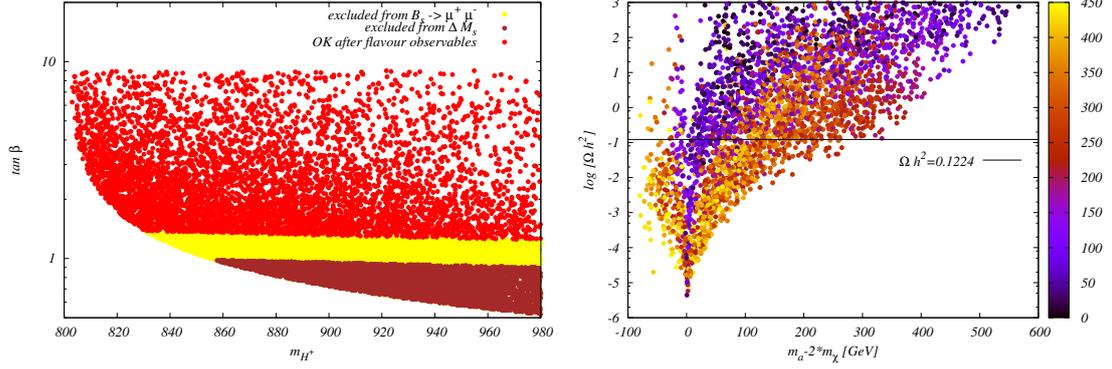


FIG. 8. *Left:* Bounds on the $(m_{H^\pm}, \tan \beta)$ plane from B-physics observables, implemented via the SPheno [4]/ Sarah [5] interface, and compared to experimental bounds [11, 12]. The contour for low $(m_{H^\pm}, \tan \beta)$ values stems from [15, 16]. *Right:* Dark matter constraints in the THDMA model. *Right:* Dark matter relic density as a function of $m_a - 2m_\chi$, with m_χ defining the color coding. The typical resonance-enhanced relic density annihilation is clearly visible. Figures taken from [60].

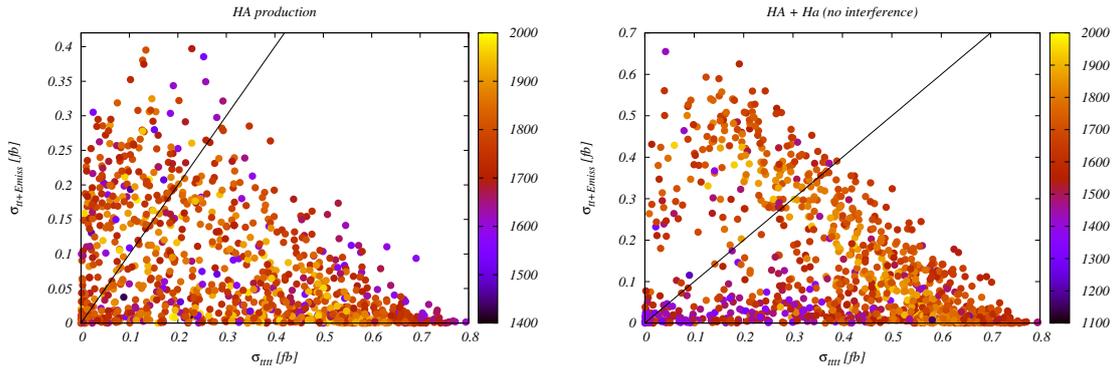


FIG. 9. Production cross sections for $t\bar{t}t\bar{t}$ (x-axis) and $t\bar{t} + E\bar{t}$ (y-axis) final state in a factorized approach, for an e^+e^- collider with a 3 TeV center-of-mass energy. *Left:* mediated via HA , *right:* mediated via HA and Ha intermediate states. Color coding refers to $m_H + m_A$ (left) and $M_H + 0.5 \times (m_A + m_a)$ (right). Figures taken from [60].

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