

Study of the $C\mathcal{P}$ properties of the top-quark Yukawa interaction in $t\bar{t}H$ and tH events with $H \rightarrow \gamma\gamma$

Marcos Miralles López^{a,*} on behalf of the ATLAS Collaboration

^a*Instituto de Física Corpuscular (CSIC-UV)*

Parque Científico, C/Catedrático José Beltrán, 2 | E-46980 Paterna, Spain

E-mail: marcos.miralles.lopez@cern.ch

These proceedings present the first search for $C\mathcal{P}$ -violation in the top-Yukawa coupling using $t\bar{t}H$ and tH production modes in the diphoton decay channel ($H \rightarrow \gamma\gamma$). The analysis is based on the 139 fb^{-1} of proton-proton collision data at center-of-mass energy $\sqrt{s} = 13 \text{ TeV}$ recorded with the ATLAS detector at the Large Hadron Collider. Two separate set of selections are introduced to capture the full hadronic and leptonic decay of the $t\bar{t}$ system each. This search is performed using a simultaneous fit to the diphoton invariant mass distribution in analysis categories defined to enhance signal over continuum background and also to separate the $C\mathcal{P}$ -even and $C\mathcal{P}$ -odd signal hypotheses.

*** *The European Physical Society Conference on High Energy Physics (EPS-HEP2021)*, ***

*** *26-30 July 2021* ***

*** *Online conference, jointly organized by Universität Hamburg and the research center DESY* ***

*Speaker



1. Introduction

The study of the interaction of the top quark with the Higgs boson is key to probe Higgs properties such as the charge conjugation and parity ($C\mathcal{P}$) of this coupling. For this, the production of the Higgs boson in association with top quarks at the LHC [1, 2] presents the first opportunity to directly study top-Higgs interactions. The Standard Model (SM) predicts a scalar Higgs ($J^{C\mathcal{P}} = 0^{++}$) with a prescribed coupling to the top quark. However, the presence of a $J^{C\mathcal{P}} = 0^{+-}$ pseudoscalar admixture, which introduces a second coupling to the top quark, has not yet been excluded. The observation of this $C\mathcal{P}$ -odd contribution would be a sign of beyond the SM physics.

There are two main processes to explore the $C\mathcal{P}$ properties of the top-Yukawa coupling: the $t\bar{t}H$ and the tH production. The $C\mathcal{P}$ -odd component impacts the production rates and some kinematic distributions. This effect can be seen in Figure 1 for the mentioned processes as well as for the $H \rightarrow \gamma\gamma$ branching ratio and the ggF production.

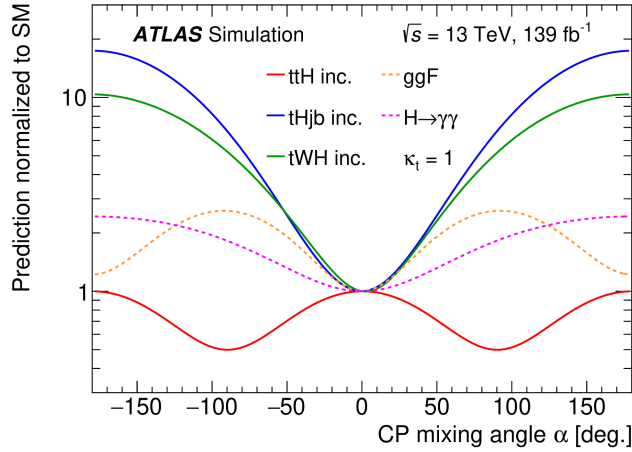


Figure 1: Dependence of the expected rates on the mixing angle for $\kappa_t = 1$. Normalised to the SM expectation. Figure from Ref. [3].

This analysis is performed using 139 fb^{-1} of $\sqrt{s} = 13 \text{ TeV}$ proton–proton (pp) collision data recorded from 2015 to 2018 with the ATLAS detector [4–6]. The effective field theory (EFT) definition to introduce a $C\mathcal{P}$ -odd component in the top-Yukawa coupling is provided by the Higgs characterization model [7], which is implemented in the MADGRAPH5_AMC@NLO generator [8]. Within this model, the term in the effective Lagrangian that describes the top-Yukawa coupling is

$$\mathcal{L} = -\frac{m_t}{v} \{ \bar{\psi}_t \kappa_t [\cos \alpha + i \sin \alpha \gamma_5] \psi_t \} H, \quad (1)$$

where m_t is the top quark mass, v is the Higgs vacuum expectation value, $\kappa_t (> 0)$ is the top-Yukawa coupling parameter, and α is the $C\mathcal{P}$ -mixing angle. The term proportional to $\sin \alpha$ introduces the $C\mathcal{P}$ -odd interaction.

2. Event Selection and Categorisation

Events are required to have two isolated photons (decaying from the Higgs boson) with transverse momenta (p_T) greater than 35 GeV and 25 GeV. Both photons must satisfy the tight identi-

fication requirement [9]. In addition, events are separated into two $t\bar{t}H$ -enriched regions. In first place, the ‘‘Lep’’ region ($\geq 1j, \geq 1b$) includes events where a W boson (from a top quark decay) will decay leptonically. Secondly, the ‘‘Had’’ region ($0l, \geq 3j, \geq 1b$) will include all other events with hadronically decaying W bosons. These two additional jets must have a $p_T > 25$ GeV. An object-level boosted decision tree (BDT) ‘‘Top Reco BDT’’ is trained with the $t\bar{t}H$ sample by using the XGBoost package [10] to extract a top quark candidate. Its goal is to separate between random jet triplets from $t\bar{t}H(\gamma\gamma)$ events from those coming from a top quark decay. The BDT uses information from the decay products of the top quark (W boson and b -jet) from reconstructed jets or the final-state lepton for the ‘‘Had’’ and ‘‘Lep’’ regions respectively. Figure 2 shows the reconstructed top quark candidate mass from the BDT.

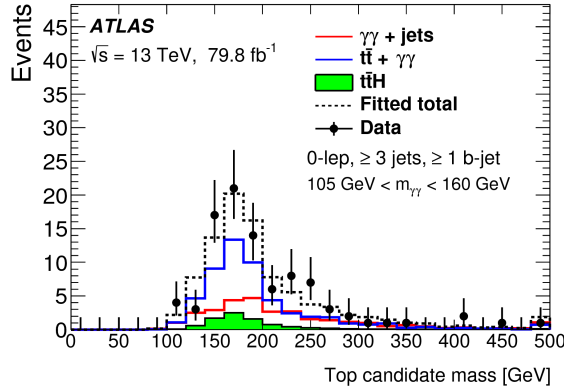


Figure 2: Distribution of the reconstructed top quark-candidate in the diphoton mass window from a triplet of jets selected by the dedicated ‘‘Top Reco BDT’’. There is a good agreement with data. Figure from Ref. [1].

To improve the analysis sensitivity, selected events are categorized using partitions of a two-dimensional BDT space. Two additional BDTs are trained separately in both the ‘‘Lep’’ and ‘‘Had’’ channels:

- **Background BDT** aims to separate between $t\bar{t}H$ signal and background events [1]. It exhibits good background rejection and $t\bar{t}H/tH$ acceptance and weak dependence with $C\mathcal{P}$ -mixing angle.
- **CP BDT** aims to separate $C\mathcal{P}$ -even for $C\mathcal{P}$ -odd events.

Figure 3 shows the BDT discriminant distributions in the data as well as those expected from $C\mathcal{P}$ -even and $C\mathcal{P}$ -odd Higgs boson signals in the ‘‘Had’’ region. There are 20 categories in total: 12 in the ‘‘Had’’ region and 8 in the ‘‘Lep’’ region.

3. Fitting Procedure and Results

A simultaneous maximum-likelihood fit is performed to the diphoton invariant mass ($m_{\gamma\gamma}$) spectra in all the categories. Signal and background shapes are modeled by analytic functions as discussed in Ref. [11]. Figure 4 shows the distributions of the reconstructed masses for the diphoton system and primary top quark. Figure 5 shows the yields for the $C\mathcal{P}$ -even and $C\mathcal{P}$ -odd signals

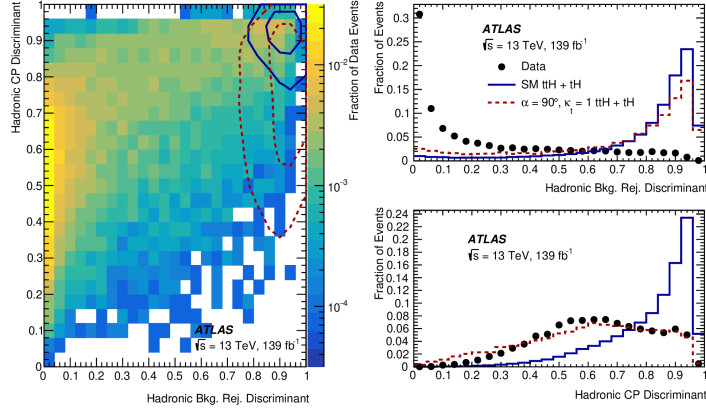


Figure 3: Two dimensional BDT distribution for the hadronic channel (left). Inner (outer) contours capture 25% (50%) of $t\bar{t}H$ and tH signal events. Projections to Background and CP BDT axes normalized to unity (right). Figure from Ref. [3].

as well as the data calculated in the smallest $m_{\gamma\gamma}$ interval containing 90% of the signal shown in three groups of categories. The binning is obtained by the combination of the appropriate event categories defined from the previous BDT categorisation. The last bin shows a clear favour towards the $C\mathcal{P}$ -even hypothesis.

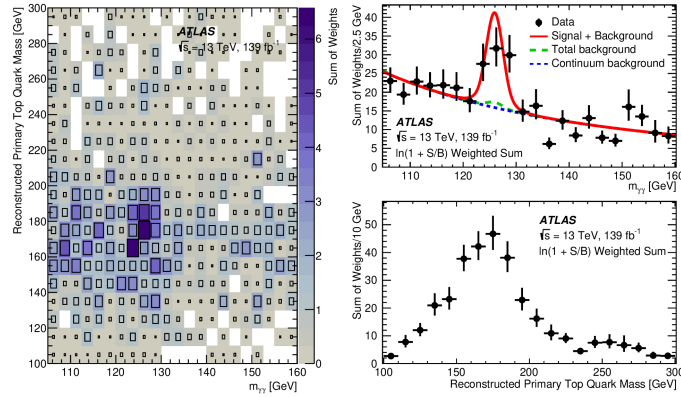


Figure 4: Distribution of reconstructed primary top mass versus Higgs mass in data events (left). Projection to both particle mass axes (right). Data error bars are statistical. The events are weighted by $\ln(1 + S/B)$ with S and B being the fitted signal and background yields in the smallest $m_{\gamma\gamma}$ interval containing 90% of the signal in each category. Figure from Ref. [3]

3.1 $t\bar{t}H$ signal strength

Assuming a $C\mathcal{P}$ -even coupling, and constraining all non- $t\bar{t}H$ Higgs Boson processes to their SM predictions, the measured rate for $t\bar{t}H$ is $\mu_{t\bar{t}H} = 1.43^{+0.33}_{-0.31}(\text{stat}) +^{+0.21}_{-0.15}(\text{syst})$ times the SM expectation. The background-only hypothesis is rejected with an observed (expected) significance of 5.2σ (4.4σ).

3.2 tH cross-section upper limit

Under the same assumptions, the used CLs method [12] yields a 95% confidence level (CL) upper limit of 12 times the SM prediction for tH production cross-section, the same as expected assuming the presence of SM tH signal.

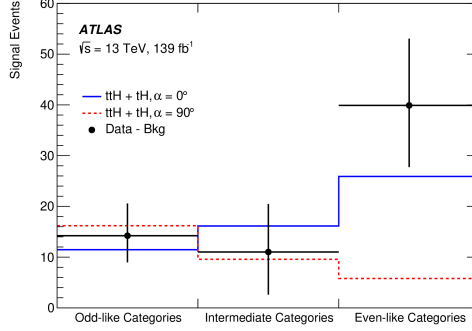


Figure 5: Data comparison with $C\mathcal{P}$ -even and $C\mathcal{P}$ -odd hypothesis. SM $C\mathcal{P}$ -even coupling is favoured. Figure from Ref. [3].

3.3 Limits on $C\mathcal{P}$ -mixing angle α

From the fit, the 1D (mixing angle α) and 2D ($\kappa_t \cos \alpha - \kappa_t \sin \alpha$ contours) limits on the $C\mathcal{P}$ properties of the top-Yukawa coupling can be extracted. These are shown in Figure 6 in the left and right panels respectively.

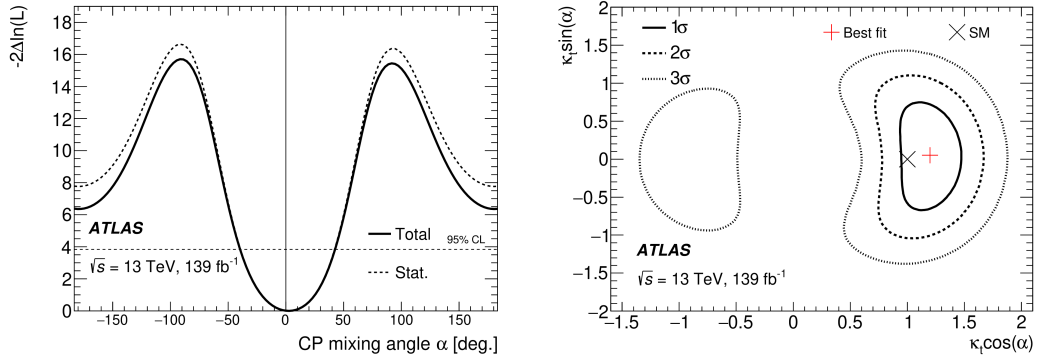


Figure 6: One-dimensional likelihood distribution of the mixing angle α (left), and 2D two-dimensional $\kappa_t \cos \alpha - \kappa_t \sin \alpha$ likelihood contours (right) results for $C\mathcal{P}$ limit. Figures from Ref. [3].

Additional information about the top-Yukawa coupling is required. The corresponding Higgs boson coupling modifiers κ_γ and κ_g are taken from the Run 2 Higgs boson coupling combination [13] (without including the $t\bar{t}H$ process). With these, the expected (observed) exclusion limit on the mixing angle (without a prior constraint on κ_t) is $|\alpha| > 43(63)^\circ$ at 95% CL. The $C\mathcal{P}$ -odd hypothesis is excluded at 3.9σ (2.5σ). In this analysis, the statistical uncertainty is the dominant one. These results are documented in Ref. [3].

References

- [1] ATLAS Collaboration, *Observation of Higgs boson production in association with a top quark pair at the LHC with the ATLAS detector*, *Phys. Lett. B* **784** (2018), 173-191 [arXiv:1806.00425 [hep-ex]].
- [2] CMS Collaboration, *Observation of $t\bar{t}H$ production*, *Phys. Rev. Lett.* **120** (2018), 231801 [arXiv:1804.02610 [hep-ex]].
- [3] ATLAS Collaboration, *$C\mathcal{P}$ Properties of Higgs Boson Interactions with Top Quarks in the $t\bar{t}H$ and tH Processes Using $H \rightarrow \gamma\gamma$ with the ATLAS Detector*, *Phys. Rev. Lett.* **125** (2020), 061802 [arXiv:2004.04545 [hep-ex]].
- [4] ATLAS Collaboration, *The ATLAS Experiment at the CERN Large Hadron Collider*, *JINST* **3** (2008) S08003.
- [5] ATLAS Collaboration, *ATLAS Insertable B-Layer Technical Design Report*, Report No. ATLAS-TDR-19, 2010, <https://cds.cern.ch/record/1291633>; Addendum: Report No. ATLAS-TDR-19-ADD-1, 2012, <https://cds.cern.ch/record/1451888>.
- [6] B. Abbott et al., *Production and integration of the ATLAS Insertable B-Layer*, *JINST* **13** (2018) T05008
- [7] F. Demartin, F. Maltoni, K. Mawatari, B. Page and M. Zaro, *Higgs characterisation at NLO in QCD: CP properties of the top quark Yukawa interaction*, *Eur. Phys. J. C* **74** (2014), 3065 [arXiv:1407.5089 [hep-ph]].
- [8] J. Alwall, R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, O. Mattelaer, H. S. Shao, T. Stelzer, P. Torrielli and M. Zaro, *The automated computation of tree-level and next-to-leading order differential cross sections, and their matching to parton shower simulations*, *JHEP* **07** (2014), 079 [arXiv:1405.0301 [hep-ph]].
- [9] ATLAS Collaboration, *Electron and photon performance measurements with the ATLAS detector using the 2015–2017 LHC proton-proton collision data*, *JINST* **14** (2019) P12006.
- [10] T. Chen and C. Guestrin, *XGBoost: A Scalable Tree Boosting System*, [arXiv:1603.02754 [cs.LG]].
- [11] ATLAS Collaboration, *Measurements of Higgs boson properties in the diphoton decay channel with 36 fb^{-1} of pp collision data at $\sqrt{s} = 13\text{ TeV}$ with the ATLAS detector*, *Phys. Rev. D* **98** (2018), 052005 [arXiv:1802.04146 [hep-ex]].
- [12] A. L. Read, *Presentation of search results: The $CL(s)$ technique*, *J. Phys. G* **28** (2002), 2693-2704
- [13] ATLAS Collaboration, *Combined measurements of Higgs boson production and decay using up to 80 fb^{-1} of proton-proton collision data at $\sqrt{s} = 13\text{ TeV}$ collected with the ATLAS experiment*, *Phys. Rev. D* **101** (2020), 012002 [arXiv:1909.02845 [hep-ex]].

- [14] J. Ellis, D. S. Hwang, K. Sakurai and M. Takeuchi, *Disentangling Higgs-Top Couplings in Associated Production*, [JHEP **04** \(2014\), 004](#) [[arXiv:1312.5736 \[hep-ph\]](#)].