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A laser–plasma platform for photon–photon physics: the two photon Breit–Wheeler process

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Abstract

We describe a laser–plasma platform for photon–photon collision experiments to measure fundamental quantum electrodynamic processes. As an example we describe using this platform to attempt to observe the linear Breit–Wheeler process. The platform has been developed using the Gemini laser facility at the Rutherford Appleton Laboratory. A laser Wakefield accelerator and a bremsstrahlung convertor are used to generate a collimated beam of photons with energies of hundreds of MeV, that collide with keV x-ray photons generated by a laser heated plasma target. To detect the pairs generated by the photon–photon collisions, a magnetic transport system has been developed which directs the pairs onto scintillation-based and hybrid silicon pixel single particle detectors (SPDs). We present commissioning results from an experimental campaign using this laser–plasma platform for photon–photon physics, demonstrating successful generation of both photon sources, characterisation of the magnetic transport system and calibration of the SPDs, and discuss the feasibility of this platform for the observation of the Breit–Wheeler process. The design of the platform will also serve as the basis for the investigation of strong-field quantum electrodynamic processes such as the nonlinear Breit–Wheeler and the Trident process, or eventually, photon–photon scattering.

1. Introduction

1.1. Background

Gregory Breit and John A Wheeler predicted in 1934 that when two photons collide, they can produce an electron and a positron in a process called 'inelastic photon–photon scattering' [1]. This Breit–Wheeler

(BW) process $\gamma + \gamma \rightarrow e^- + e^+$ has been studied in detail using the theoretical framework of quantum electrodynamics (QED) [2]. In the linear BW process, two photons must have sufficient energy to create an electron and a positron, i.e. the invariant mass of the collision must satisfy $\sqrt{s} > 2m_e c^2$, where $\sqrt{s}$ is the total energy measured in the centre of momentum frame. For two photons of energy E_1 and E_2 interacting at angle θ, $s = 2[1 - \cos(\theta)]E_1 E_2$.

To date, the linear BW process has not been observed in the laboratory with real photons due to the lack of suitably bright sources of photons with sufficient energy. Indeed, in their original paper, Breit and Wheeler even stated: 'it is hopeless to try to observe the pair formation in the laboratory' as in 1934 there was no conceivable way of generating such photon sources.

Recent results from the STAR collaboration report the observation of the BW process using quasi-real photons in the peripheral collisions of high energy ions. The high energy of the quasi-real photons involved in those collisions ($\sqrt{s} \approx 2$ GeV) means the virtuality of the interaction, $Q^2 \approx (30 \text{ MeV})^2$, is sufficiently small to make the measurements indistinguishable from collisions between real photons. There have been proposals for γ–γ colliders using conventional particle accelerators [3, 4], but none have yet been realised and a number of schemes involving multi-PW lasers to produce sufficiently bright γ sources have been proposed [5–7].

Today bright electron beams with GeV energies can be produced over a few centimetres by laser Wakefield acceleration (LWFA) [8, 9]. These electron beams can be used to produce brilliant γ ray beams by bremsstrahlung [10–13] and inverse Compton scattering [14–16]. Colliding such brilliant γ ray photons with a suitable x-ray photon field could be a viable route to study two-photon physics in the laboratory for the first time [17, 18].

The BW process plays a pivotal role in a range of astrophysical phenomena [19]. The linear BW process is the key process in pair-instability supernovae (PIS) and the closely related pulsational pair-instability supernova, the fate of massive stars where the BW process softens the equation-of-state at the core, leading to collapse and ignition [20, 21]. PIS are of particular interest as they might be the only stars observable in the epoch of reionization when the first stars are formed, meaning that they may be the only view into star formation in the metal-poor early Universe possible with the James Webb Space Telescope [22, 23]. The BW process also occurs in the coronae of accretion disks in active galactic nuclei [24], shaping particle spectra [25]. Finally, the BW process occurs when γ rays interact with the extragalactic background light (EBL), producing a cut-off of γ rays at $\sim 10^{13}$ eV [26]. There has been some suggestion of the possibility of non-standard model physics in the spectra of blazars, which some studies suggest have anomalously hard high-energy γ ray spectra after the BW process is accounted for, implying less pair-production from QED than anticipated [27, 28]. This problem could well be solved with conventional physics (e.g. corrections to models of the EBL [29]) as happened for the '*compactness problem*', a contradiction that arose from γ ray burst light curve data which was eventually solved by including relativistic motion in the calculations [30]. Nonetheless this illustrates the importance of confirming that QED physics behaves as expected in this regime for this important process.

The existence of a photon–photon experimental platform studying the BW effect will serve as a basis for the investigation of other QED processes including elastic photon–photon scattering $\gamma\gamma \rightarrow \gamma\gamma$ and the nonlinear BW process [31]. The nonlinear BW process is closely related to the BW process but involves the interaction of one high-energy photon with many (n) low energy ones, i.e. $\gamma + n\omega \rightarrow e^- + e^+$. The generation of matter via the nonlinear BW process plays a large role at the poles of rotating black holes [32]. In 1997, an experiment at SLAC collided a 48 GeV electron beam with a then state-of-the art laser pulse [33, 34], demonstrating the nonlinear BW process, thus studying the perturbative regime of strong-field QED [35].

Elastic photon–photon scattering, $\gamma\gamma \rightarrow \gamma\gamma$, is important in models of primordial abundances and may also affect the observed spectra from distant γ ray bursts [36, 37]. Photon–photon scattering involving virtual photons has previously been observed. Delbrück scattering, $\gamma\gamma^* \rightarrow \gamma\gamma^*$, where a real photon scatters from a virtual photon in the electric field of an ion, was observed in 1975 [38] and scattering between two virtual photons in the Coulomb field of two high energy ion beams $\gamma^*\gamma^* \rightarrow \gamma\gamma$ was observed at the large hadron collider in 2017 and 2019 [39–41]. However, photon–photon scattering is most relevant in astrophysics with real photons close to threshold for pair production energy ($\sqrt{s} \lesssim 2m_e c^2$) as at higher energies pair production dominates. Photon–photon scattering with real photons has yet to be observed in the laboratory.

1.2. Laser–plasma approach

In 2014, Pike *et al* [17] proposed an ambitious scheme that could enable the observation of the linear BW process by combining a laser-driven γ beam source with a laser-driven x-ray source. They proposed using a

megajoule laser pulse, e.g. the National Ignition Facility [42], coupled to a petawatt short pulse laser to produce 10 000 pairs in a single shot. However, no such combination of high-power lasers exists worldwide. Ribeyre *et al* (2016) [5] proposed the interaction between two high-power laser generated γ-ray beams. However this method relies on extremely precise spatial and temporal overlap of the two beams, and due to the approximately equal incoming photon energies, would generate pair particles into a wide range of angles, making detection difficult. In this article we present an alternative laser-based approach that can operate at existing laser facilities, benefits from directed pair production, and which can make use of higher shot repetition rates to provide the statistics needed for any given measurement.

Our approach is based on the interaction of a single γ beam colliding with an x-ray bath, similar in nature to that of Pike *et al* [17], but with a few hundred terawatt laser system providing tens of joules (as opposed to a megajoule) at a repetition rate of over 100 shots an hour (as opposed to a few shots a day). A major difficulty that a laser based approach overcomes is the spatial and temporal overlap that goes hand-in-hand with colliding intense photon sources. Using the same laser system to drive the photon sources means that they can be inherently synchronised in time (within tens of femtoseconds) and do not suffer from significant temporal jitter. Making use of a large x-ray radiation field (>100 μm $\times$ 100 μm and tens of picoseconds in duration) means that the spatial alignment of the two photon sources is also relatively straightforward. In addition to the overlap, the strong asymmetry of the photon energy of the two sources also results in a strongly directional emission of the created electron–positron pairs. This is particularly important considering that for many experimental configurations the BW process can be obscured by other more efficient pair creation processes, such as the Bethe–Heitler process, when high energy photons interact with the matter in the vicinity of the photon–photon collision. A strategic and tailored detection system that can decouple the BW signal from background pairs from other processes is required, and a clean particle-free interaction volume is desirable. Two decoupled photon sources with a spatially directed signal is a solution to this problem and presents a major advantage over other suggested schemes that collide two sources of photons that are equal in energy [43].
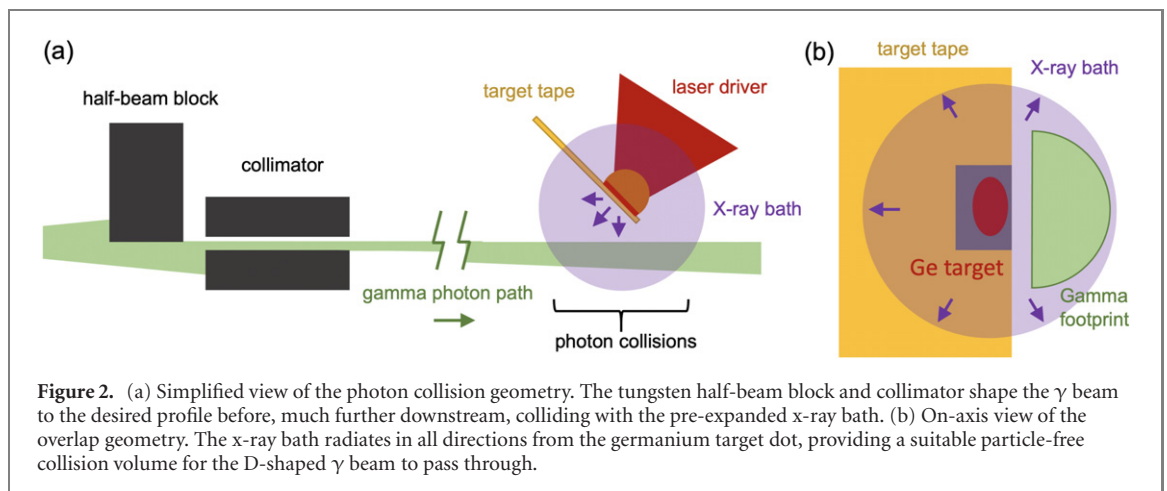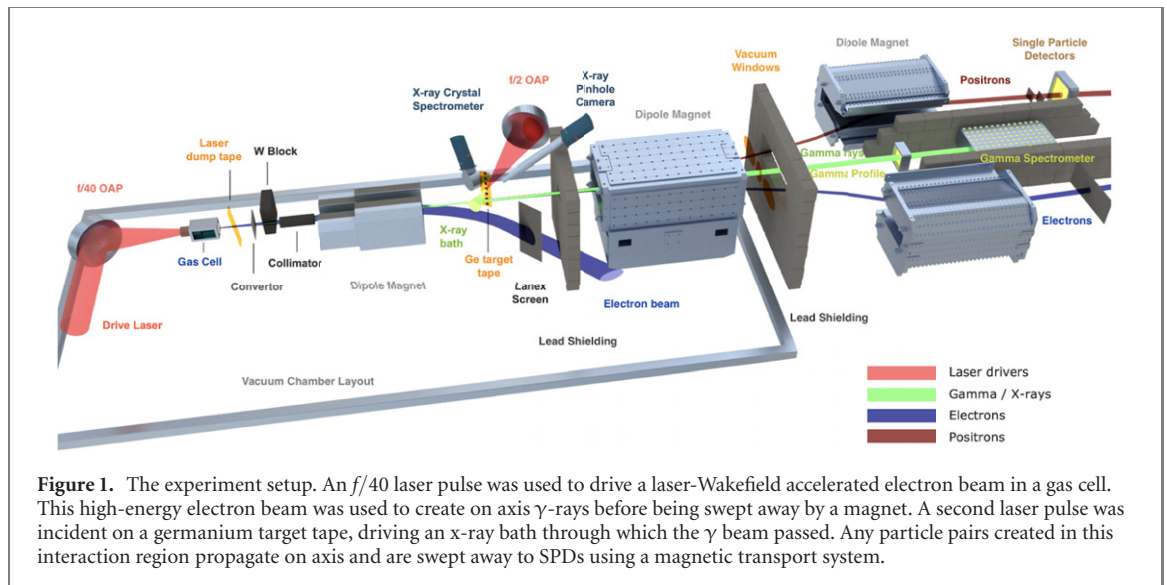
Although the predicted pair creation numbers per shot are lower than that of Pike *et al* (around unity, as discussed later), many shots can be taken per day to build a statistical significance. It is also worth noting that the platform we present here can also be modified to study the nonlinear BW process by colliding the γ ray photons directly with the intense optical laser field that would otherwise drive the x-ray bath.

In this paper we detail the design and commissioning of a laser plasma platform for photon–photon physics. Our characterisation and modelling of the background levels and expected signal level indicate that a measurement of the linear BW process is feasible within a reasonable time frame (1000s of laser shots), but show that this is very sensitive to the performance of the photon sources. Further developments of the platform are discussed, which should bring the number of required shots to less than a few hundred.

2. Experimental platform overview

The platform was developed at the Gemini laser system for strong field QED experiments [44] at the Central Laser Facility, UK and adapted to the specific requirements of this experiment. It is an optically synchronised two beam system, where each beam contains up to 15 J in $\approx$ 45 fs (FWHM), with a shot rate of once every 20 s. The synchronisation of the two beams allows for a temporal overlap and jitter of less than 10 fs [45]. One of the beams was used in a long focal length *f*/40 configuration to drive a laser-Wakefield accelerator. The high energy electron bream produced in this interaction is subsequently used to generate bremsstrahlung MeV γ photons. The second beam was stretched in duration to 40 ps (FWHM), and used to drive the laser–plasma generated keV x-ray field (using an *f*/2 focussing optic). These two sources are made to overlap spatially within the collision interaction point and a magnetic transport system relays any created electron–positron pairs to a set of single particle detectors (SPDs). An overview of the experimental layout can be seen in figure 1. Data used for this study is available at [46].

The long focal length arm of the laser, which drives the laser-Wakefield accelerator, is focussed onto a gas cell. The cell was 17.5 mm long and filled with helium and 2% nitrogen as a dopant at an electron density of $(2.6 \pm 0.4) \times 10^{18}$ cm^{-3}. The typical measured laser spot size was 44 μm $\times$ 53 μm FWHM at a pulse duration of 45 ± 5 fs and 5.5 ± 0.6 J on target, which corresponds to a normalised vector potential $a_0 = 1.1 \pm 0.2$. As the laser pulse propagates through the gas it strips electrons from the ions, forming a plasma. The ponderomotive force of the laser pulse pushes these electrons away from the intense pulse, while the ions remain stationary, forming a positively charged cavity that trails in the wake of the laser pulse [47]. The electrons can then be trapped and accelerated in the cavity to relativistic energies, travelling in the direction of the laser pulse. After driving the laser-Wakefield accelerator, the remaining drive laser exits the gas cell and is removed from the beamline using a refreshable tape drive.

Figure 1. The experiment setup. An *f*/40 laser pulse was used to drive a laser-Wakefield accelerated electron beam in a gas cell. This high-energy electron beam was used to create on axis γ-rays before being swept away by a magnet. A second laser pulse was incident on a germanium target tape, driving an x-ray bath through which the γ beam passed. Any particle pairs created in this interaction region propagate on axis and are swept away to SPDs using a magnetic transport system.

Figure 2. (a) Simplified view of the photon collision geometry. The tungsten half-beam block and collimator shape the γ beam to the desired profile before, much further downstream, colliding with the pre-expanded x-ray bath. (b) On-axis view of the overlap geometry. The x-ray bath radiates in all directions from the germanium target dot, providing a suitable particle-free collision volume for the D-shaped γ beam to pass through.

The generated high-energy electron beam continues along the beamline where it passes through a 1 mm thick bismuth *'converter'* target. While passing through the high-*Z* atoms of the converter, the electrons generate bremsstrahlung radiation, emitting γ photons with over 100 MeV energy along the propagation direction. The bremsstrahlung converter material and thickness was optimised so as to produce a bright γ beam with low divergence, i.e. without having many multiple scattering interactions which increase the background noise on. A limited number of materials and thicknesses were tested during commissioning to optimise the signal-to-noise ratio (signal measured on the γ-profile detector, noise level measured on the SPDs). The γ beam is collimated to ∼8 mrads divergence using a 10 cm long tungsten cylinder with a 2 mm aperture. One half of the γ beam is also blocked by a 5 cm thick tungsten shield. The remaining footprint of the γ beam is semi-circular in shape. This D shape is designed to allow the tape target of the x-ray bath source (further downstream) to be placed in close proximity to the centre of the γ beam, but hidden from direct photon collisions at the interaction point. See figure 2 for a depiction of the overlap geometry.

After the γ beam has been tailored to the desired shape, a $\int B\,\mathrm{d}x = 0.4$ Tm dipole magnet is used to sweep away the initial electron beam as well as any secondary particles generated from collisions with the tungsten shield elements. This combination of γ beam tailoring and particle removal provides a clean environment for photon–photon collisions.

The second laser beam is incident upon a plastic tape containing regularly spaced germanium targets, which are placed 1 mm from the γ beam footprint. The laser pulse rapidly heats and ionises the metal, which emits M–L shell transition radiation in the 1–2 keV region; see section 4 for more details. This emission radiates in all directions, but we are concerned with the x-ray plume that penetrates through the plastic and out of the rear of the target towards the interaction region. The two drive laser pulses should be timed so that the γ beam passes through the x-ray field at its highest photon density. The x-ray driver is also short enough in duration that no ions from the germanium plasma plume have penetrated through to the
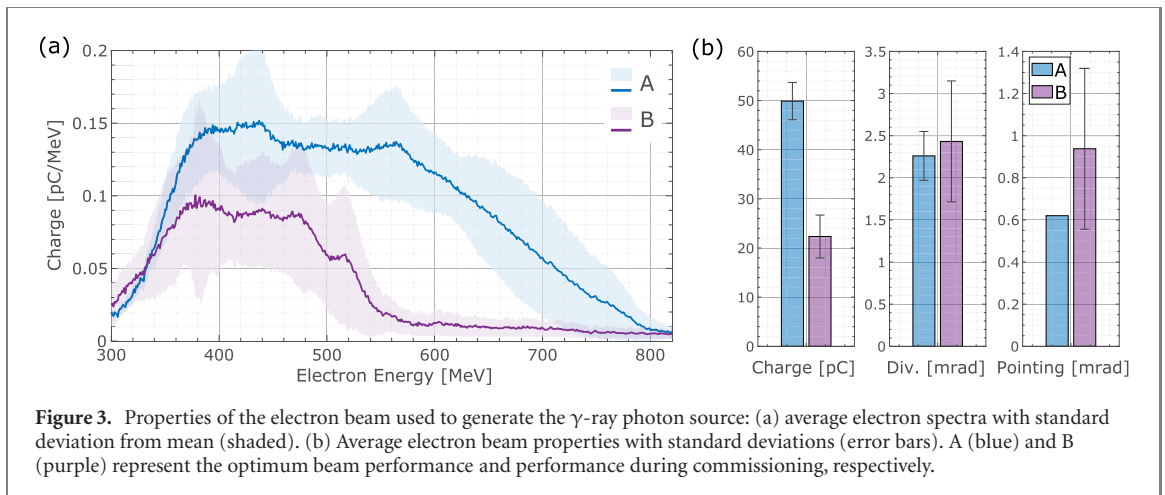
Figure 3. Properties of the electron beam used to generate the γ-ray photon source: (a) average electron spectra with standard deviation from mean (shaded). (b) Average electron beam properties with standard deviations (error bars). A (blue) and B (purple) represent the optimum beam performance and performance during commissioning, respectively.

collision volume by the time the γ beam is incident. The 5 μm thick Kapton tape substrate helps to prevent plasma from the front (laser-irradiated) side of the germanium target penetrating through to the rear (where photon collisions occur). For ions to propagate from the plasma to the interaction volume (1 mm away) within 100 ps of the laser driver, they would need to move at $\sim 10^7$ m s^{-1}, or $\sim 3\%$ the speed of light. However calculations of the laser–plasma interaction indicate a peak electron temperature of ~ 2.5 keV, near the peak of the pulse, with an associated ion sound speed ($c_s \sim 3.5$–6.3×10^5 m s^{-1}) that is a fraction of that required to reach the interaction by the collision time. Additional hydrodynamic simulations support this, showing that the Kapton substrate of the targets has not been burnt through by the time the γ beam collides with the interaction volume. Because the γ-rays are much higher energy than the x-rays, any electrons and positrons produced in the photon–photon collisions travel along the initial γ-ray beam direction. A second dipole magnet $\int B \, dx = 0.35$ Tm with a large aperture sweeps the generated positrons off axis. It is at this point the positrons exit the vacuum chamber through a 125 μm Kapton window. They pass through a third dipole magnet with opposite polarity to the second, redirecting the particles along a focussed beam path. This particle transport system is discussed in detail in section 5.

SPDs are placed on this beam path to measure the particle flux. Electrons generated within the interaction region are swept by the second dipole in the opposite direction, and a fourth magnet (identical to the third one) directs the electrons to a beam dump. The profile and spectrum of the on-axis γ-ray beam are diagnosed using scintillator-based detectors at the end of the beam line. The γ-ray diagnostics and SPDs are discussed in more detail in section 3 and section 6 respectively. Lead shielding is used extensively throughout the setup to mitigate the influence of unwanted high energy particle noise on the diagnostics, details of the shielding and the effect on background levels are discussed in section 8.

3. γ-ray photon source

During the commissioning of the platform the laser energy used to drive the accelerator was limited to 5 J, which is significantly below the maximum that can be achieved in Gemini (12–15 J on target). With 5 J laser energy the best electron beam performance achieved (data set 'A') reached on average a maximum energy of 710 ± 50 MeV and contained a charge of 50 ± 7 pC. The data set contained 10 shots and the average spectra of these beams is presented in figure 3(a), labelled dataset A (blue). The measurement had a low-energy cut off at ~ 350 MeV due to the yoke of the spectrometer magnet. The beams had a divergence of 2.3 ± 0.3 mrad and a shot-to-shot pointing standard deviation of 0.6 mrad. Additionally, during the commissioning phase of the platform (discussed later), the LWFA electron beam performed in a reduced capacity, with lower maximum energies achieved and less than half the total charge (data set 'B'). There were 79 shots in total and the average spectra for these runs are provided in figure 3, labelled dataset B (purple).

The γ profile diagnostic consists of an array of 20×20 caesium-iodide (CsI) crystals, each of dimensions 2 mm $\times$ 2 mm and 20 mm thick, positioned normal to the γ beam and 3.39 m downstream of the γ source. The scintillation of the CsI crystals upon exposure to high energy photons is observed with an optical camera. An example γ beam profile from this diagnostic can be seen in figure 4(a), the D shaped profile due to the combination of the collimator and tungsten block shielding the x-ray target are clearly apparent. Another larger array of 33×47 crystals, each 5 mm $\times$ 5 mm $\times$ 50 mm, is positioned along the beam axis behind the profile diagnostic and measures the decay of signal in the propagation direction to
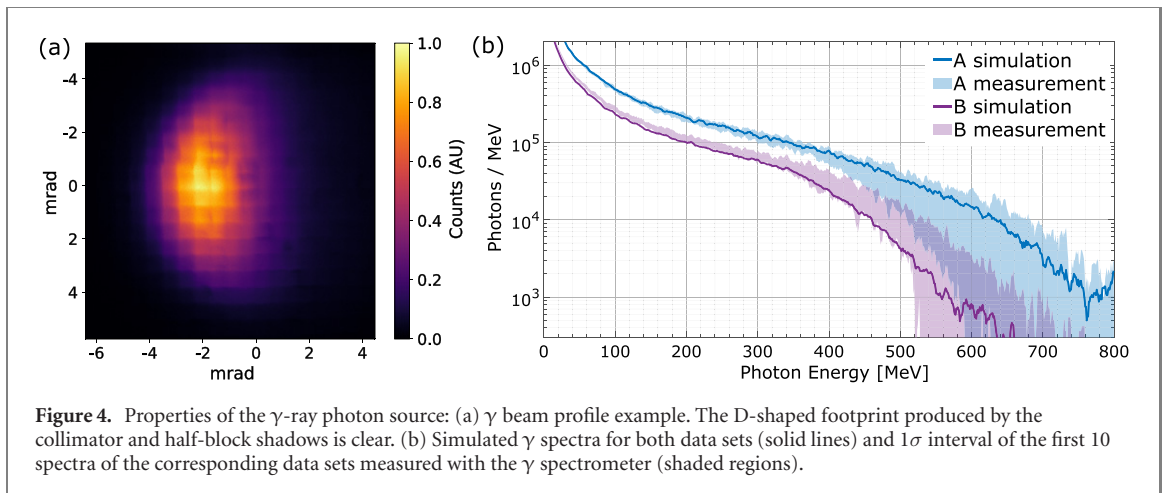
Figure 4. Properties of the γ-ray photon source: (a) γ beam profile example. The D-shaped footprint produced by the collimator and half-block shadows is clear. (b) Simulated γ spectra for both data sets (solid lines) and 1σ interval of the first 10 spectra of the corresponding data sets measured with the γ spectrometer (shaded regions).

deduce the spectrum [48]. The γ spectrum is inferred using a forward model to calculate the energy deposited in each crystal, finding the model parameters which most closely reproduce the observed signal using a least squares approach. In this case our forward model is based on Geant4 simulations of the signal produced due to the γ rays produced by a wide range of experimentally observed electron spectra. Both the profile and spectrum scintillator arrays are observed by cooled 14 bit EMCCD cameras (Andor iXon), with single photon sensitivity to the scintillation light.

The converter target scatters the incident electron beam such that the properties of the γ beam and the electrons cannot be measured simultaneously. Therefore the electron beam was characterised without the converter in place, and this was subsequently used to simulate the expected average γ spectrum. The γ spectra that were measured with the larger crystal array immediately after the characterisation of the electron beam (figure 4(b), shaded regions) agreed well with the simulated spectra (figure 4(b), solid lines).

Geant4 simulations using dataset A indicate the emission of $(3.5 \pm 0.5) \times 10^8$ photons or $(7 \pm 1) \times 10^6$ photons/pC, distributed in an exponentially falling spectrum with a high-energy cutoff at the maximum electron energy (see figure 4(b)). $(7 \pm 3) \times 10^7$ of these photons, or about 20% of the total photon yield, have an energy above 100 MeV. By matching the average electron beam and the average yield on the γ profiler, we obtain a calibration factor of $(5.0 \pm 0.4) \times 10^6$ γ-counts/pC. This calibration allows us to estimate the electron beam charge on shots where the convertor target is in place (and hence no electron spectrometer measurement is possible). The variations in the γ yield indicate stronger fluctuations in the photon numbers from shot to shot than expected from the electron charge. This is believed to be due to pointing fluctuations and the limited size of the γ detector occasionally missing outer portions of the beam. As described previously, for photon–photon collisions, the γ beam is collimated and one half is blocked to prevent interactions with the x-ray target (see figures 2 and 4(a)). This reduces the footprint of the beam and the effective yield of photons by a factor 3.6 ± 1.2.

4. X-ray photon source

The second arm of the laser system was used to generate the dense keV x-ray bath at the interaction point. The drive laser was focussed onto a refreshable Kapton tape target system using an *f*/2 off-axis parabola. Prior to reflecting from the final focussing optic, the 150 mm diameter beam passed through a distributive phase plate [49]. This allowed the laser energy to be deposited in a large quasi-flat-top focal spot 210 μm × 90 μm (major and minor diameters at $1/e^2$). Approximately 72% of the laser energy is contained in this spot, with the remaining light being diffracted away into higher orders. The laser pulse was also stretched to 40 ps duration, resulting in an intensity of 2×10^{15} W cm^{-2} on target. The targets themselves were ~100 nm coating of germanium metal on a 5 ± 2 μm etched square window (~1 mm^2) of the Kapton tape (which is otherwise nominally 25 μm thick) [50]. As the laser pulse is incident on the germanium coating, a hot dense plasma is created on the front surface. In this highly-ionised plasma, M–L shell transitions in the germanium efficiently produce photons in the range of 1−2 keV [51]. These x-rays are emitted into a 4π sphere, such that a few millimetres from the target a dense x-ray bath exists. A similar technique has previously been used to rapidly heat targets to extreme conditions [52]. The Kapton target substrate thickness was optimised to increase the transmission of these keV photons, while suppressing sub-keV photons and preventing any material contamination to the interaction region.
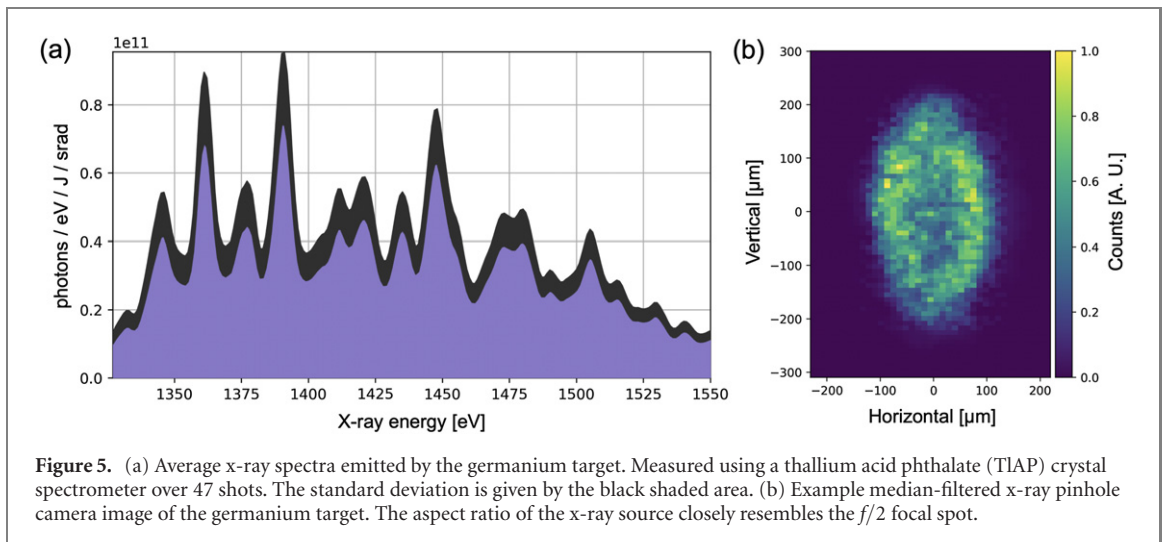
Figure 5. (a) Average x-ray spectra emitted by the germanium target. Measured using a thallium acid phthalate (TlAP) crystal spectrometer over 47 shots. The standard deviation is given by the black shaded area. (b) Example median-filtered x-ray pinhole camera image of the germanium target. The aspect ratio of the x-ray source closely resembles the $f/2$ focal spot.
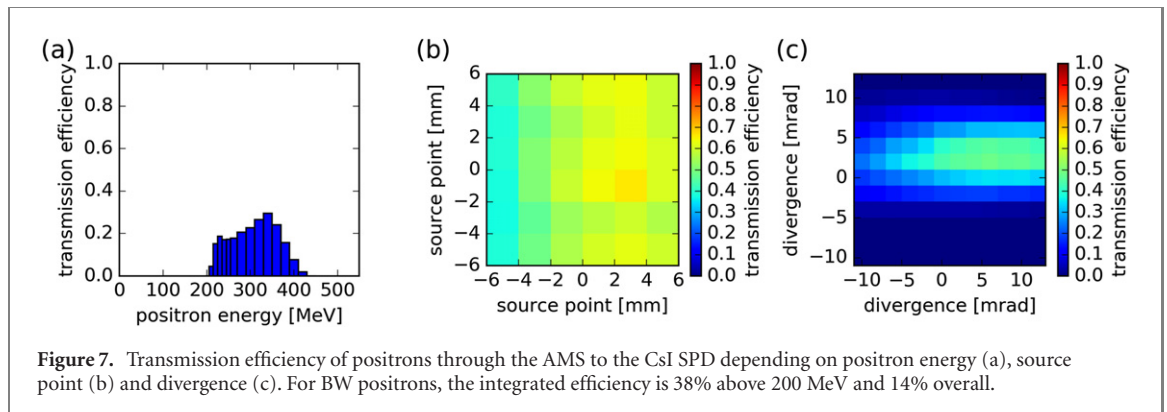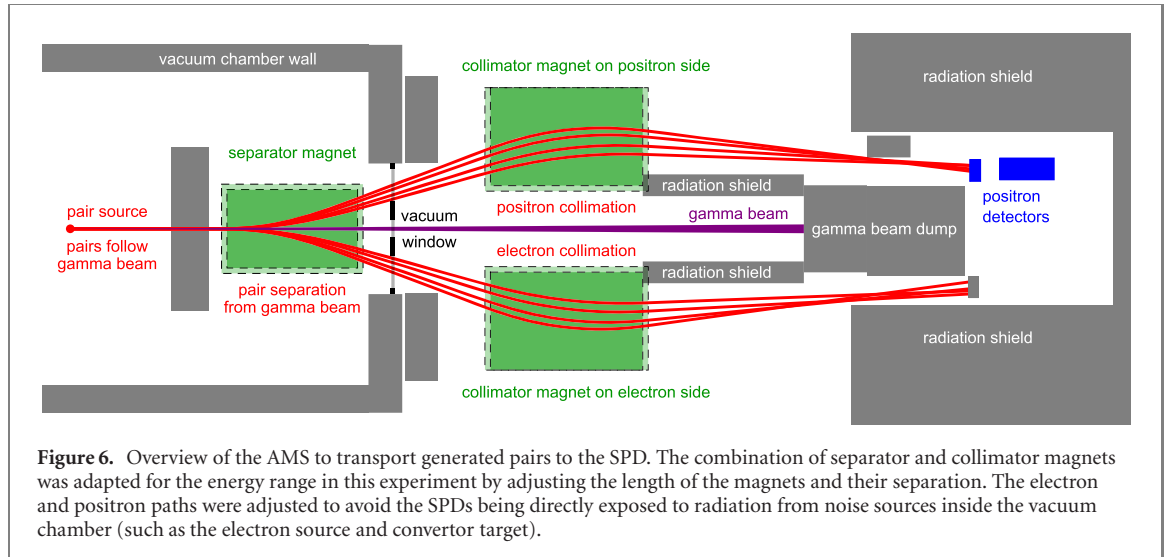
A TlAP crystal spectrometer viewed the front of the tape targets at angle of $\approx 55°$ (to normal) to diagnose the x-ray yields. Bragg reflections from the flat TlAP crystal onto an in-vacuum x-ray CCD (Andor DX-420-BN) allowed the x-ray spectrum to be measured between 1.3–1.5 keV with a spectral resolution of 4 eV. The average measured spectra from 47 shots can be seen in figure 5(a). The absorption edge from an aluminium filter at 1.56 keV provided a spectral marker allowing the dispersion across the CCD to be calculated. A variety of other filters were also used, allowing the signal to be corrected for background contributions from direct CCD irradiation, higher order Bragg reflections and crystal-substrate fluorescence.

In addition to the crystal spectrometer, a two-channel pinhole imaging system was employed. By imaging the x-ray producing plasma, the source size and target alignment were monitored on-shot. An example image from this diagnostic can be seen in figure 5(b). While the entire emission collected by the pinhole camera passed through a 25 μm Be filter (acting as a 600 eV high-pass filter), both channels of the imaging system had different spectral filters. This allowed for a broad x-ray yield quantification, which correlated well with the total energy measured by the crystal spectrometer.
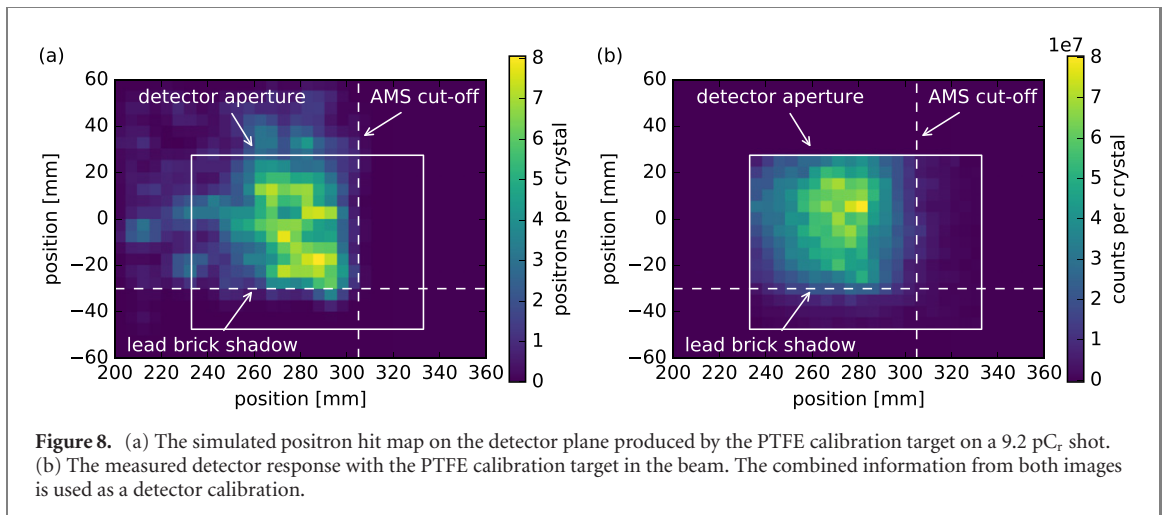
It was found that the total conversion efficiency from laser energy to 1.3–1.5 keV x-rays was $(2.4 \pm 0.3)\%$. This corresponds to $(3.7 \pm 0.4) \times 10^{10}$ photons eV^{-1} J^{-1} srad^{-1} emitted normal to the front surface of the germanium target. Taking into account the absorption in the kapton layer on the rear side of the target, at the interaction region (1 mm from the tape) this corresponds to a photon density of $(1.4 \pm 0.5) \times 10^{12}$ mm^{-3}. These photon numbers assume the Kapton tape remains cold throughout the interaction, and are supported by simulations of the x-ray source directionality and duration (97 ± 11 ps, FWHM). The average transmission of the Kapton tape is 47% over the range 1.3–1.5 keV. The flux stability from shot-to-shot was found to have a standard deviation of 28%. A tape drive was used to place a fresh target at the focal plane of the x-ray drive laser. The tape targets were able to be aligned at a repetition rate of 1 shot per minute. Note however that it should be possible to increase this rate with improvements to the target alignment system.

5. Positron transport system

In order to measure the produced pairs, a system composed of permanent magnets and single particle sensitive detectors has been developed [44] and specifically adapted for the energy range in this experiment. The overall geometry is shown in figure 6. The generated pairs initially follow the γ-ray beam path, but after passing through an aperture in a lead wall, they were separated from the beam using a dipole magnet. This magnet was 60 cm long with a field strength of 0.6 T ($\int B\,dx = 0.36$ Tm) and is referred to as the separator magnet. To avoid scattering of γ-rays on the poles, the magnet had a gap of 10 cm. The pairs were dispersed and exited the vacuum chamber through a 125 μm thick Kapton window. The pairs were then detected by two different types of SPDs, which have a limited spatial extent and are presented in section 6. To increase the collection efficiency of the generated pairs, two additional dipole magnets were used to focus the electrons and positrons to individual detectors. These magnets are referred to as the collimator magnets. They had an aperture of 9 cm, were 70 cm long and provided a magnetic field strength of 0.5 T ($\int B\,dx = 0.35$ Tm). The combination of separator and collimator magnets is called the analyser magnet system (AMS). The efficiency of this transport system is limited by various apertures along the beam path,

Figure 6. Overview of the AMS to transport generated pairs to the SPD. The combination of separator and collimator magnets was adapted for the energy range in this experiment by adjusting the length of the magnets and their separation. The electron and positron paths were adjusted to avoid the SPDs being directly exposed to radiation from noise sources inside the vacuum chamber (such as the electron source and convertor target).

Figure 7. Transmission efficiency of positrons through the AMS to the CsI SPD depending on positron energy (a), source point (b) and divergence (c). For BW positrons, the integrated efficiency is 38% above 200 MeV and 14% overall.

especially the vacuum exit window frame which has three openings, one for each particle beam and a central one for γ-rays, see figures 1 and 6. The AMS transport efficiency was evaluated using simulations performed with Radia [53, 54] for Wolfram Mathematica. The simulations model the magnetic fields and calculate the trajectory of positrons through the system for a range of possible source positions and initial momenta. The collection of all trajectories represents a five-dimensional function (the source point is fixed along the beam axis) that maps any positron to the detector plane depending on its initial phase space coordinate at the time of pair creation. Combined with a check if individual trajectories are blocked by apertures, this function, which is called the Radia filter, determines if and where a positron hits the detector. This is used to evaluate the transmission efficiency and the spatial SPD signal distribution for any incident positron population. This filter was used to map positrons produced by a 1 mm thick plastic PTFE calibration target to the SPD plane as a detector calibration, which is described in greater detail in section 6. Since it is a five-dimensional filter, the transport efficiency is hard to visualise due to correlations between all five initial coordinates and also depends on the phase space distribution of the incident positrons. A projection of the full filter to the energy axis as well as source point and divergence planes is shown on figure 7, which exhibits a narrow spectral acceptance between 200 and 420 MeV as well as a narrow angular acceptance between -1 and $+7$ mrad relative to γ beam axis. The source point does not affect the transmission significantly over the considered range from -6 to 6 mm relative to γ beam axis. The overall transport efficiency for positrons from the PTFE calibration target above 200 MeV is 28% and the transport efficiency for produced BW positrons above 200 MeV is 38%. For all BW positrons in the full spectral range, the transmission is 14%. The accuracy of calculated positron impact locations is approximately 1 cm at the detector plane, limited by the positioning accuracy of magnets, detectors and apertures inside the laboratory.

Figure 8. (a) The simulated positron hit map on the detector plane produced by the PTFE calibration target on a 9.2 pC$_r$ shot. (b) The measured detector response with the PTFE calibration target in the beam. The combined information from both images is used as a detector calibration.
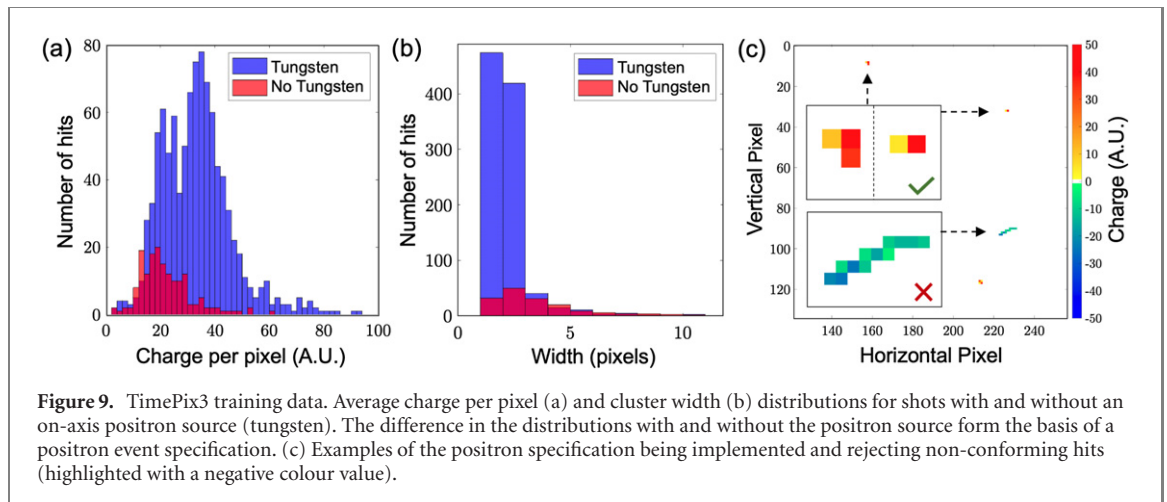
6. Single particle detectors

6.1. Caesium iodide single particle detector

A scintillation based SPD system was employed. It was composed of a thallium doped caesium iodide (CsI) crystal array coupled to a 4Picos ICCD from Stanford Computer Optics. The crystals were coated with titanium dioxide to avoid optical crosstalk between neighbouring crystals. The surface facing towards the camera was polished. The 4Picos camera achieve single photon sensitivity through a multi-channel-plate with photoelectron multiplication factors up to 10^6. This combination of CsI crystals and 4Picos camera gives the SPD single particle sensitivity.

The CsI SPD was calibrated with a PTFE calibration target in the beam close to the BW pair source point to produce a known positron signal on the SPD. As the γ beam passes through the PTFE, electron–positron pairs are generated in a narrow on-axis cone via the Bethe–Heitler process. The PTFE positron population was calculated using the measured LWFA beam properties as an input to a Geant4 simulation of the bremsstrahlung process, γ-ray beam shaping and Bethe–Heitler pair production. The measured electron beam performance (and subsequent Geant4 input) for these calibration shots are given in figure 3 dataset B. The obtained positron population was filtered and projected to the detector plane using the AMS phase space filter calculated with Radia as described in section 5. The simulated spatial distribution of positron impact locations on the SPD plane is shown on figure 8 together with the measured signal. The square box indicates the detector aperture and the dashed lines show positron cut-off's, one due to lead shielding that intentionally shadows the lower portion of the SPD and another one due to the AMS field distribution. The regions where no positrons impact the SPD are used to measure background radiation other than positrons transported through the AMS, which is described in section 7.

The simulated positron yield integrated over the detector aperture is (200 ± 7) positrons$_{ams}$, where the error is the standard error of the mean over three identical Geant4 simulations. The subscript ams indicates that we are referring to positrons transmitted through the AMS with the associated properties specified in figure 7. The LWFA beam spectral shape is assumed to be constant for all calibration shots, which allows us to assume a linear dependency between positron yield and LWFA charge. This normalisation to LWFA beam charge assuming the reference spectral shape is denoted with pC$_r$. The Geant4 simulations were performed with a total charge of 9.2 pC$_r$, so the positron yield produced by the PTFE target is obtained as (21.7 ± 0.8) positrons$_{ams}$/pC$_r$.

The mean measured signal integrated over the detector area is $(8.44 \pm 0.69) \times 10^8$ counts$_n$, where the subscript n denotes counts that are corrected for camera and imaging system effects. The PTFE data was taken with an average charge of (19.3 ± 1.1) pC$_r$, so the total measured PTFE positron yield is obtained as $(4.77 \pm 0.17) \times 10^7$ counts$_n$/pC$_r$ with errors being standard errors of the mean. This is then corrected for background radiation by taking shots without the PTFE target in the beam, which was measured to be $(0.109 \pm 0.004) \times 10^7$ counts$_n$/pC$_r$, and gives the BG corrected yield of $(4.66 \pm 0.17) \times 10^7$ counts$_n$/pC$_r$. This measured signal is now directly correlated to the simulated number of incident positrons to obtain the SPD calibration of $(2.15 \pm 0.03) \times 10^6$ counts$_n$/positron$_{ams}$.

Figure 9. TimePix3 training data. Average charge per pixel (a) and cluster width (b) distributions for shots with and without an on-axis positron source (tungsten). The difference in the distributions with and without the positron source form the basis of a positron event specification. (c) Examples of the positron specification being implemented and rejecting non-conforming hits (highlighted with a negative colour value).

6.2. Timepix3 single particle detectors

A second SPD system relied on two Timepix3-based hybrid detectors [55] arranged in a stacked geometry, with one 40 mm behind the other. The rear was offset laterally by 4 mm, closer to the beam axis. Both detectors were placed before the CsI SPD. Timepix3 are CMOS based hybrid pixel detectors with an active area of 14×14 mm^2 (256×256 pixels, each 55 μm $\times$ 55 μm). Timepix3 assemblies with silicon sensors were chosen with the front detector having a 300 μm thick sensor and the rear 150 μm thick. Any minimum ionising particles (MIPs) passing through the silicon deposit $\sim$2900 eV per micron, registering a trail of charge across the pixels. A MIP is considered to be a charged particle that has kinetic energy twice its rest mass, i.e. any positron particles over 1 MeV. As charged particles pass through the detector, each pixel records the *time-over-threshold* at a rate of up to 640 MHz. If fully calibrated the temporal resolution can be as low as $\sim$1 ns, but in our case we integrate the signal over one second during the shot period. A significant advantage of these detectors is that all charged particles of interest will register a signal, where as due to the silicon nature of the detector, they are relatively insensitive to high energy photons, i.e. they will be insensitive to background γ photon noise. For example, the probability of γ-ray absorption in 150 μm silicon is $\sim$8 $\times$ 10^{-4}–2 $\times$ 10^{-3} [56]. As an additional step to confirm that we are witnessing primarily charged particles rather than photons, the stacked Timepix3 setup was simulated in Geant4 with either an incoming particle beam, or an incoming photon beam. In the experiment calibration shots (discussed further below), an approximately equal number of strikes was measured on the front and rear detectors, but with 0.4 times the total energy absorbed in the rear (the sensor of half the thickness). This aligns with the Geant4 simulations for charged particle strikes, as opposed to photons which estimated approximately 2.5 times the energy should be deposited in the rear detector due to secondaries created from photon absorption in the first detector.

We deem each contiguously connected group of pixel charge registered on the detectors (a 'cluster') as a single charged particle or photon hit. In order to reduce unwanted noise, we can filter these clusters by their characteristic traits, to select only those that are consistent with positrons that have been transported through the AMS (as opposed to stray background particles). We do so by training an event identification criteria on a set of calibration data with an artificially created positron signal at the collision point (in a similar manner to the PTFE calibration of the CsI SPDs). For this, a 1 mm tungsten cube is placed on the γ beam axis, within a few cm of the collision point. The generated particle pairs are similar in energy to those created by our photon–photon collision (100's MeV). We collated the cluster statistics on 28 shots with the tungsten cube in place, and 28 null shots immediately before this, with the tungsten cube removed. The latter shots act as a comparable null reference for the background present on the Timepix3. The shots with the artificial positron source show that the clusters have a higher average charge per pixel, as well as narrower distributions for the width and height, see figures 9(a) and (b). These criteria can be applied to all clusters to identify AMS positron events and remove background hits. To build a complete AMS positron specification, we take the distribution for each metric (e.g. cluster width) individually and subtract the null shot series distribution. The limits for that metric are then taken as the 95% interval of the remaining distribution. The final AMS positron specification for our setup is given in table 1. It was observed on subsequent data and background shots that the cluster filtering was rejecting up to 60% of the total hits on the Timepix3 detectors. See figure 9(c) for an example of this AMS positron identification.

The AMS positron shot series was also used to infer the particle trajectory between the two Timepix3 detectors. With a known stream of incident particles passing through both detectors, it is possible in theory

Table 1. Timepix3 on-axis positron specifications. Applying
these criteria to each TimePix3 charge cluster allows the
identification of positrons travelling through the AMS, and rejects
events that are background noise.

Detector	Charge sum	Num. pixels	Height	Width
Front 300 μm	$40 \leqslant X \leqslant 145$	$\leqslant 5$	$\leqslant 2$	$\leqslant 3$
Rear 150 μm	$10 \leqslant X \leqslant 85$	$\leqslant 3$	$\leqslant 2$	$\leqslant 2$

to track the particles as they travel ballistically through the front detector and onto the rear. To do so, on
each shot a 2D map of positron hits was generated for each Timepix3 detector. A raster of horizontal and
vertical positions of one detector relative to the other was then calculated, and a correlation value given to
the overlapping region. Accounting for the fact the rear detector was offset laterally by 4 mm closer to the
beam axis, it was found that particles had a range of trajectories of $4°-12°$ between the two detectors. This
is supported by the Radia tracking simulations for the AMS. Unfortunately this variation in angular spread
(stemming from the variation in particle energy through the AMS) meant it was not possible to track
particles from the front Timepix3 through to the rear. It was possible however to treat the area of the second
TimePix3 outside of these trajectories as an additional detection area, without risking double counting
positrons passing through the first TimePix3. This increased the total detection area to 308 mm^2 (a 57%
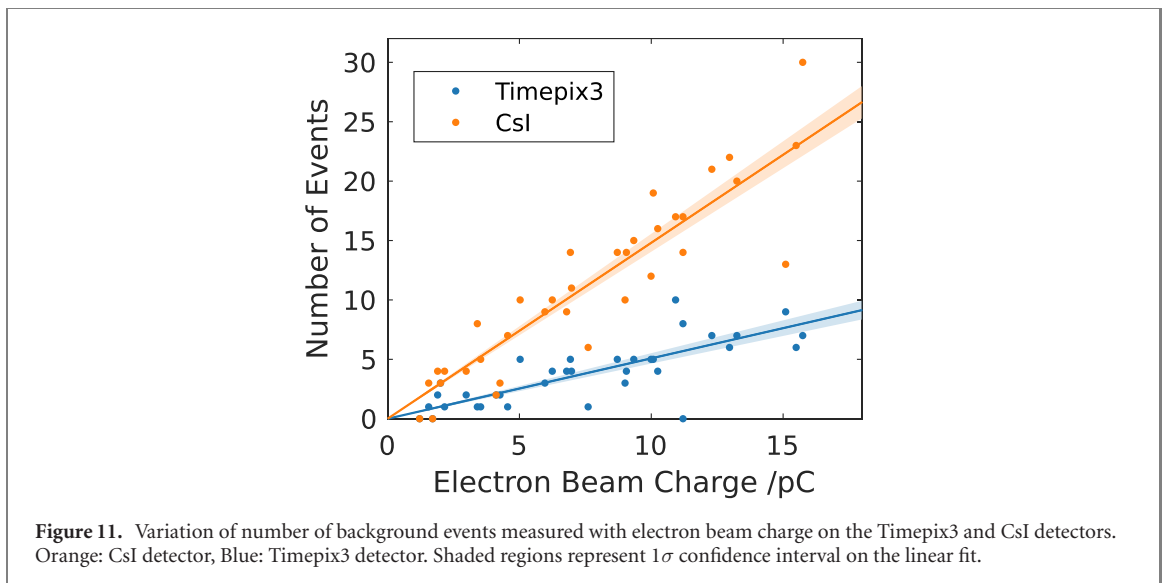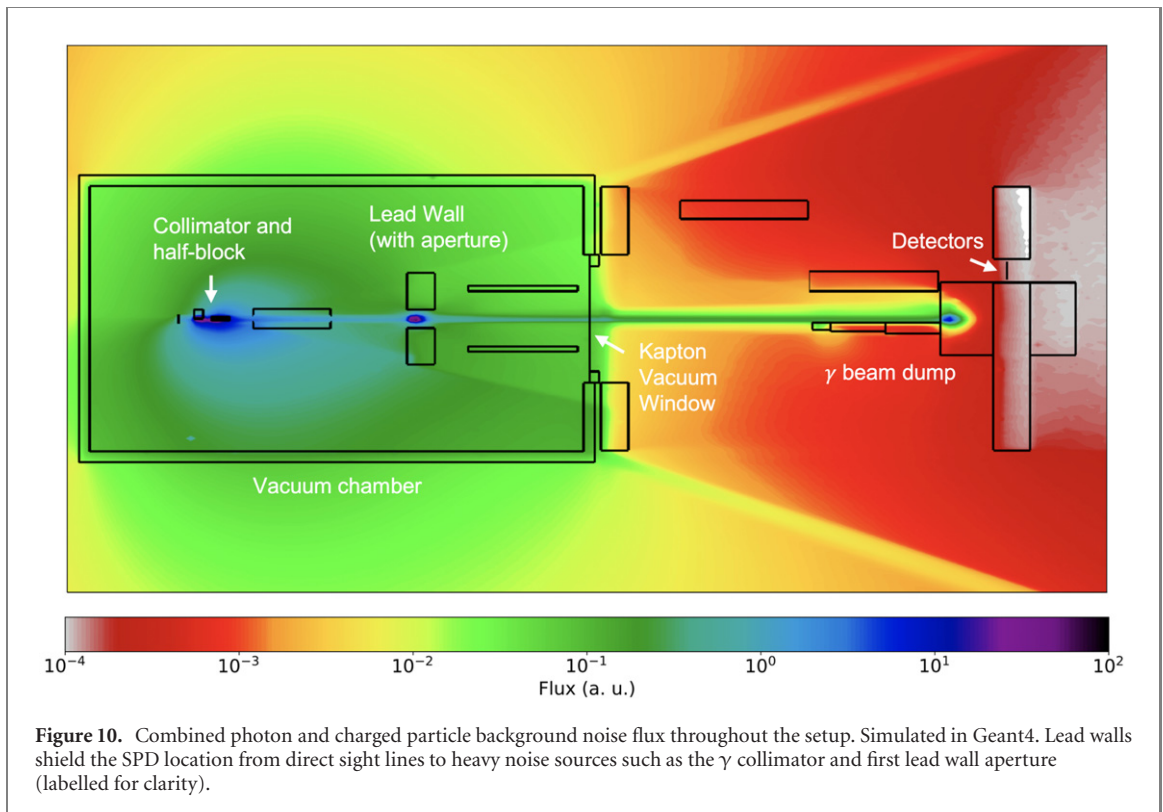increase).

7. Commissioning

7.1. Characterisation of background sources

The radiation background was characterised using the CsI SPD and found to be composed of two distinct
particle types, high energy positrons transported through the AMS and low energy diffuse radiation from
the γ-ray beam dump. The discrimination between these two was made by analysing the yield on different
spatial regions on the SPD shown on figure 8, where AMS positrons can only impact part of the SPD
aperture with the remaining area being shielded. The measured background in the shielded region below
the lead brick shadow and right of the AMS cut-off was equivalent to $BG_{diffuse} = 1.3 \pm 0.1 \times$
10^{-3} positrons$_{ams}$/pC$_r$/crystal. This is assumed to originate from diffuse low energy radiation leaking
through the γ-ray beam dump due to its homogeneous distribution over the detector aperture. The average
background level in the region open for positrons was equivalent to $BG_{open} = 2.5 \pm 0.1 \times$
10^{-3} positrons$_{ams}$/pC$_r$/crystal, almost twice as large compared to the shielded region. The difference,
$(BG_{open} - BG_{diffuse}) = BG_{AMS} = 1.2 \pm 0.1 \times 10^{-3}$ positrons$_{ams}$/pC$_r$/crystal is therefore assumed to
originate from high energy positrons that propagate through the AMS. For context, this corresponds to 1
background positron for every 4 pC$_r$ of reference LWFA beam charge integrated over the full 208 crystal
large CsI SPD region that is open for AMS positrons.

Sources of background AMS positrons are Bethe–Heitler pairs produced by the γ-ray beam in collisions
with residual gas inside the vacuum chamber induced by the LWFA target or with apertures close to the
beam line, for example the vacuum exit window or lead shielding in between the BW interaction point and
the separator magnet shown on figure 1. Geant4 simulations were used to understand these noise sources
and help design an appropriate shielding configuration. Figure 10 depicts the combined photon and
charged particle flux estimated throughout the final shielded setup. The main goal of any shielding is to
minimise the flux through the SPD location. The Geant4 simulations identify that the main sources of
background positrons which are transported through the AMS are the γ beam collimator and the aperture
in the lead wall before separator magnet. The aperture size was chosen to minimise the background
generated by particles (both γ rays and electrons from the LWFA) striking the edge of the aperture while
maintaining a low background caused by particles striking the pole pieces of the separator magnet. The
primary source of photon background at the detector was identified as the γ beam profile and γ spectrum
detectors and the beam dump.

7.2. Dependence of background noise on γ beam properties

To assess the capability of the photon–photon platform to measure the BW process, it is important to
quantify the background level of the SPDs on data shots. This measurement was made by taking a series of
shots where the γ beam was operating and the x-ray foil was at the collision position, but the second laser
beam, which generates the x-rays, was not fired. Since all the background sources are generated by the
interaction of the γ beam with components of the experiment, including the x-ray generation foil, these

Figure 10. Combined photon and charged particle background noise flux throughout the setup. Simulated in Geant4. Lead walls shield the SPD location from direct sight lines to heavy noise sources such as the γ collimator and first lead wall aperture (labelled for clarity).

Figure 11. Variation of number of background events measured with electron beam charge on the Timepix3 and CsI detectors. Orange: CsI detector, Blue: Timepix3 detector. Shaded regions represent 1σ confidence interval on the linear fit.

shots represent a direct measurement of the background that will be present on shots where the two photon sources are colliding. We call these type of shots *null shots*.

In a series of null shots it was found that the number of background events detected by both SPDs is linearly proportional to the effective charge in the electron beam, as measured on the calibrated γ profile diagnostic (see figure 11). A linear regression, assuming that the number of background events follows Poisson statistics, with the expected number of background events $\langle N_{bg} \rangle$ varying linearly with the electron charge, C, i.e. $\langle N_{bg} \rangle = \sigma_{bg} C$ shows that the signal on the CsI detector corresponds to 11.2 ± 0.6 positron events on a shot with the average charge. The number of events counted on the Timepix3 detector corresponds to 3.8 ± 0.3 events on an average shot. The CsI detector area (5200 mm^2) is much larger than the Timepix3 area (308 mm^2), so we might expect the CsI detector signal to be 17 times higher than the Timepix3 detector. However the spatial non-uniformity of the background on the CsI detector is expected to reduce this to a ratio of 9. The observed ratio is in fact close to 3.

The two detectors have very different sensitivity to the different background sources. The CsI detector analysis removes the spatially diffuse background caused by low energy γ events generated in the γ beam
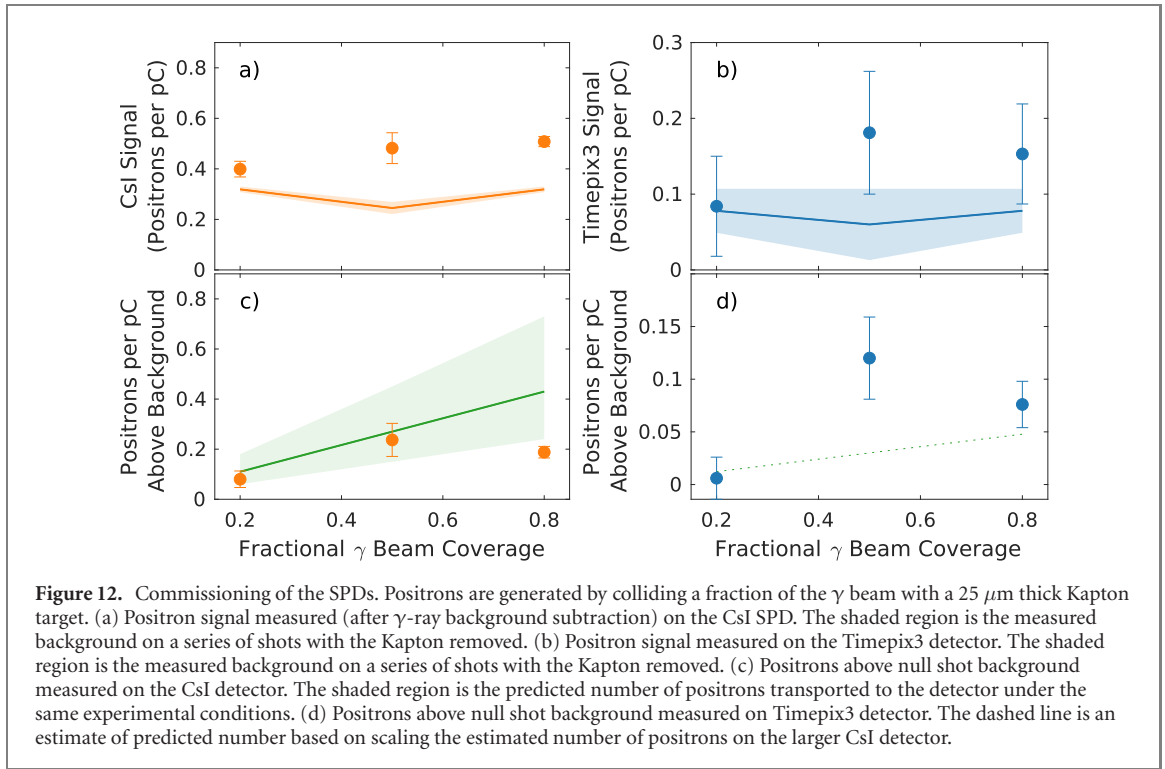
Figure 12. Commissioning of the SPDs. Positrons are generated by colliding a fraction of the γ beam with a 25 μm thick Kapton target. (a) Positron signal measured (after γ-ray background subtraction) on the CsI SPD. The shaded region is the measured background on a series of shots with the Kapton removed. (b) Positron signal measured on the Timepix3 detector. The shaded region is the measured background on a series of shots with the Kapton removed. (c) Positrons above null shot background measured on the CsI detector. The shaded region is the predicted number of positrons transported to the detector under the same experimental conditions. (d) Positrons above null shot background measured on Timepix3 detector. The dashed line is an estimate of predicted number based on scaling the estimated number of positrons on the larger CsI detector.

dump. This background subtraction is possible on the CsI detector because part of the detector is shielded from positrons transported by the AMS. However, this is not possible for the smaller Timepix3 detector. This difference in the background subtraction method on the detectors could be responsible for the difference in the observed rate on null shots. However, the thin silicon Timepix3 detector is quite insensitive to γ rays (at most 10^{-3} absorption), and the track shape detection is able to discount low energy events and events incident at an angle that is inconsistent with positrons transported through the AMS. This suggests that both detectors should be fairly insensitive to the diffuse γ background caused by the beam dump.

7.3. Commissioning with low yield positron source

To test the detection capability of the SPDs, a series of commissioning shots were performed. These involved inserting a 25 μm thick Kapton tape target into the γ beam at the photon–photon collision location (while not firing the x-ray drive laser). The tape target was 1 cm wide, and spooled in the vertical plane, presenting ≈ 20 cm of tape length. This meant it covered the entire vertical length of the γ footprint, but some fraction of the footprint in the horizontal plane. By translating the target so that it was only partially intercepting the γ beam profile we could vary the number of positrons produced and transported to the detectors. Only a limited number of shots ($\lesssim 10$) were taken at each position. The γ background was subtracted from the CsI data using the method describe in section 6. The background on both detectors owing to particles generated away from the collision point (e.g. in the lead aperture) being transported to the detectors was characterised by completely removing the Kapton target from the γ beam (null shots). The measured positron signal is shown in figure 12 for both detectors. After subtraction of the positron background, the CsI detector shows agreement within experimental uncertainty with the predicted number of positron reaching the detector calculated using Geant4, confirming the calibration procedure in this low signal regime. The larger area CsI detector was capable of measuring a signal above the positron background when only 20% of the γ beam was intercepted by the Kapton tape, corresponding to signal level of approximately 0.1 positrons per picocoulomb reaching the detector. The smaller area Timepix3 detector was able to discern a signal above background when 50% of the γ beam was intercepted, which is also a signal level of approximately 0.1 positrons per picocoulomb reaching the detector. Even with the limited amount of data taken in these commissioning shots, it is clear that the SPDs are capable of measuring small numbers of positrons per picocoulomb. However the detection limit of approximately 0.1 positrons per pC above background is very dependent on the number of shots taken to characterise both the signal and the background. In a data run with many shots, we would be able to achieve a significantly higher sensitivity.

8. Modelling

The key factor which determines whether a BW detection experiment is successful is the signal-to-noise ratio. In order to model this parameter and interpret it correctly, accurate predictions involving large scale numerical simulations are required. During the design and optimisation phase of this experiment, the Monte-Carlo platform Geant4 was used to perform background noise calculations. It is desirable to also perform signal calculations within this same framework, enabling faster design and optimisation of the experimental setup. However, the BW process is not included within any of the standard Geant4 physics packages. Therefore, a new photon-physics package was developed for this purpose which includes the BW process. Here a brief overview of the package is provided, however, a more in depth review can be found in reference [57].

8.1. Geant4 Breit–Wheeler module

Geant4 is a package for simulating the passage of particles through matter. To include the BW process within Geant4, it must be treated in a similar way. This has been achieved by handling the two interacting photon sources differently where one is treated as a dynamic beam of photons, and the other is treated as a static field fully defined by a spectral photon density per unit volume per unit solid angle, $n(\omega, \theta, \phi)$. Individual photons from the dynamic source are sampled and tracked through the static field. Calculations are then performed to determine whether the BW process has occurred, and if so, the properties of the emitted electron–positron pair are found.

The interaction between the dynamic particle beam and the static photon field occurs at a constant average rate making it a Poisson process. Therefore, the probability that a dynamic particle travels a length, *x*, is given by the following exponential distribution

$$P(x) = \lambda_d^{-1} \, e^{-x/\lambda_d}, \tag{1}$$

where λ_d is the mean free path of the dynamic particle. As this is an interaction between massless particles, the mean free path is given by [58]

$$\frac{1}{\lambda_d} = \int_0^{2\pi} d\phi \int_0^{\pi} d\theta \int_0^{\infty} d\omega \, \sigma(s) \, n(\omega, \theta, \phi) \, (1 - \cos\theta), \tag{2}$$

where the BW cross-section is

$$\sigma(s) = \frac{\pi r_e^2 (1 - \beta^2)}{2} \left[(3 - \beta^4) \log \frac{1+\beta}{1-\beta} - 2\beta(2 - \beta^2) \right], \tag{3}$$

r_e is the classical electron radius, $\beta = \sqrt{1 - 4m_e^2 c^4 / s}$, and $s = 2(1 - \cos\theta_s)E_1 E_2$ for two photons of energy E_1 and E_2 interacting at angle θ_s. For each dynamic photon that passes into the static field, *x* is sampled from equation (1). If this is longer than the static field, it will propagate through unaffected, if it is shorter, then it will propagate *x* before interacting.

If an interaction occurs, the dynamic photon is removed from the simulation and an electron–positron pair is added. However, before adding the new particles their momentum must be obtained, which can be defined through each particle's energy $E_\pm$ and polar, $\theta_\pm$ and azimuthal, $\phi_\pm$ scattering angles, where the $\pm$ indicates a positron or electron. It is easiest to first obtain these properties in the centre-of-mass frame. Here, both the electron and positron receive an energy of $E_\pm = \sqrt{s}/2$. The positron polar scattering angle is then obtained by sampling from the differential cross-section for the BW process
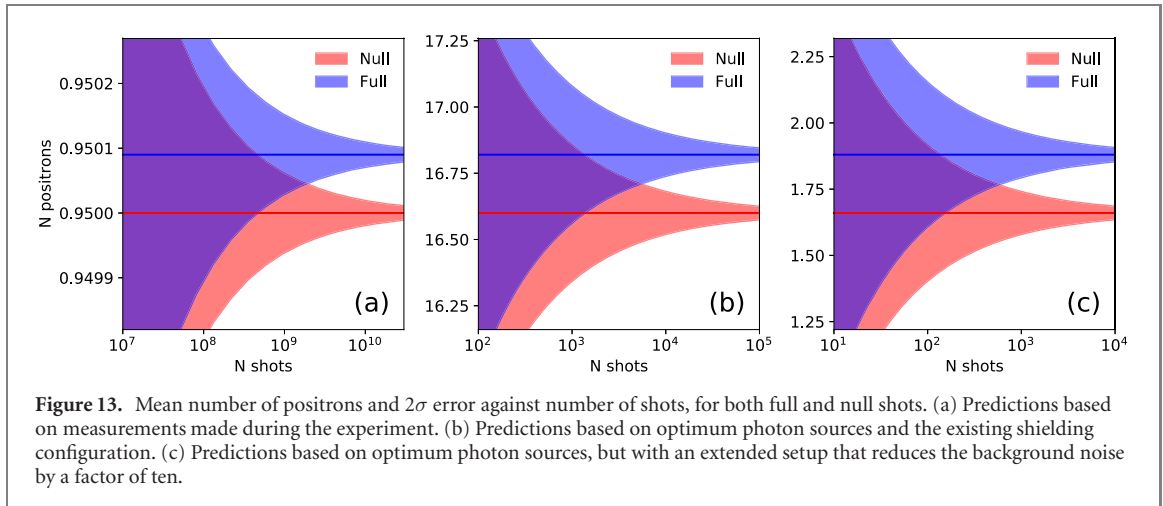
$$\frac{d\sigma}{d\Omega} = \frac{r_e^2 \beta m_e^2 c^4}{s} \left[\frac{1 + 2\beta \sin^2\theta - \beta^4 - \beta^4 \sin^4\theta}{(1 - \beta^2 \cos^2\theta)^2} \right], \tag{4}$$

whereas the positron azimuthal scattering angle is sampled uniformly from 0 to 2π. Due to conservation of momentum, this also defines the scattering angles of the electron, which are given by $\theta_- = \pi - \theta_+$ and $\phi_- = \phi_+ - \pi$. After applying a Lorentz boost, these properties are obtained in the laboratory frame, allowing the electron–positron pair to be added to the simulation.

This method does not account for either spatial gradients or a temporal evolution of the static photon field. However, if the full field is constructed by defining multiple sub fields, each with a different $n(\omega, \theta, \phi)$, then spatial gradients can be accounted for. It is challenging to account for a temporal evolution due to the static nature of the Geant4 computational domain. Therefore, this module will only provide accurate predictions if the time for the dynamic photons to traverse the static field is much shorter than the characteristic timescale of the static field. Due to the large difference in the duration of the sources, the

Table 2. Predicted number of background particles and signal Breit–Wheeler pairs reaching the TimePix3 detectors, per pC of primary electrons in the LWFA beam. The 'experiment' row shows predictions using the experimental measurements whereas the 'optimum' row shows predictions using the best obtained Gemini parameters.

	BW pairs/pC	BG leptons/pC	BG photons/pC
Experiment	$(1.87 \pm 0.69) \times 10^{-6}$	$(1.97 \pm 0.25) \times 10^{-2}$	1.37 ± 0.06
Optimal	$(1.37 \pm 0.05) \times 10^{-3}$	$(10.36 \pm 1.44) \times 10^{-2}$	4.06 ± 0.25

Figure 13. Mean number of positrons and 2σ error against number of shots, for both full and null shots. (a) Predictions based on measurements made during the experiment. (b) Predictions based on optimum photon sources and the existing shielding configuration. (c) Predictions based on optimum photon sources, but with an extended setup that reduces the background noise by a factor of ten.

x-ray source is designated as the static field and the γ-ray beam is treated dynamically. Ray tracing simulations of γ-rays propagating through a time varying x-ray field were performed to test the validity of the static x-ray field assumption. These show that the error introduced is between 5% and 15%.

8.2. Predictions

By developing a BW module for Geant4, start-to-end simulations of this experiment can be performed within a single framework, allowing for efficient calculations of the signal-to-noise ratio. Such calculations have been carried out using the electron beam and x-ray field measured during the experiment and discussed in sections 3 and 4 respectively. The results of these calculations are summarised in table 2, which show the predicted number of BW signal positrons, background leptons, and background photons which pass through the Timepix3 detectors. The background has been separated into leptons and photons as the former produces the same signature as the BW signal positrons on the Timepix3 detectors whereas the latter produces a different signature. Therefore, photon hits can be easily removed as a noise source, whereas lepton hits cannot. Note that the quoted numbers refer to the number of particles through the TimePix3 detectors, not the larger CsI SPD. For estimates of the number of particles through the CsI these numbers need to be scaled by the increased area, however, the BW signal is non-uniform over this increased area. From the Radia simulations described in section 5, it is estimated that the CsI would cover approximately nine times more BW signal than the TimePix3.

To make an observation of the BW process, a statistically significant difference between the mean number of detected positrons on full shots (both laser drivers with γ and x-rays colliding) and null shots (no collisions viable as no x-rays present) is required. Using the values in table 2, a coarse estimate of the number of shots needed to observe this difference can be made. Assuming an electron beam charge of 50 pC (corresponding to $\sim 3 \times 10^8$ primary particles), the mean number of events detected on full and null shots are $\mu_F = 0.95 + (0.9 \times 10^{-4})$ and $\mu_N = 0.95$ respectively. The number of events detected on a shot is a Poisson variable, so after N shots the error on $\mu_{N/F}$ is $\sigma_{N/F} = \sqrt{\mu_{N/F}/N}$. If a 2σ significance observation is desired, then it is required that $\mu_F - 2\sigma_F > \mu_N + 2\sigma_N$. Figure 13(a) shows this condition occurs after $\sim 10^9$ shots. This is far more shots than can be achieved over the duration of a Gemini experiment.

The signal-to-noise ratio obtained using the measured experimental parameters is far lower than what had been predicted from coarse calculations carried out prior to the experiment. This is due to both the electron beam and x-ray field behaving sub-optimally during the experiment. For example, figure 14 shows electron beam spectra obtained during a previous Gemini experiment (see reference [59]). These spectra extend to far higher energies than figure 3, with higher beam charges, in excess of 100 pC, also obtained. On

Figure 14. High energy electron beams obtained during a previous Gemini experiment. Taken from reference [59].
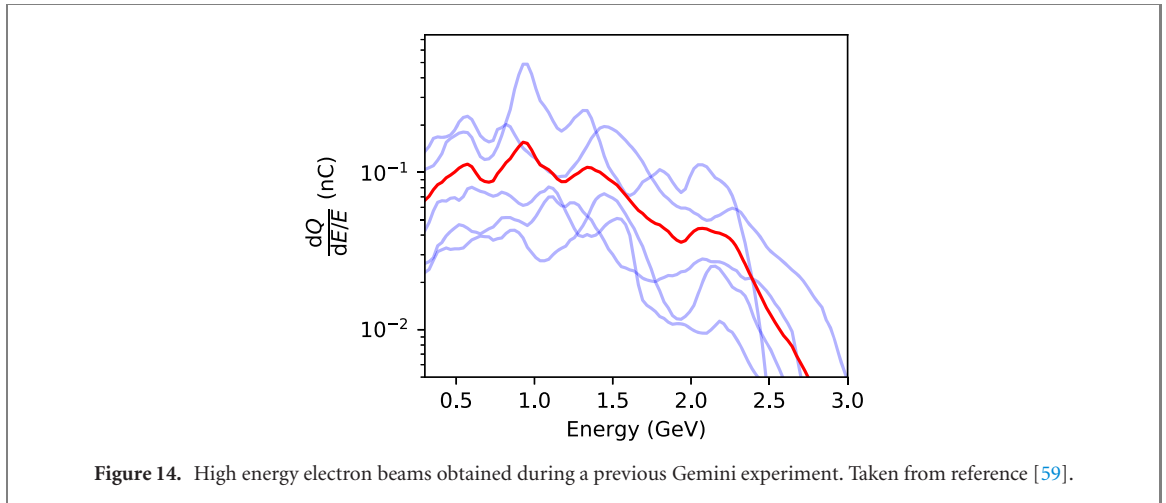
top of this, the energy in the laser pulse used to drive the x-ray field was also lower than expected. In previous experiments 15 J on target has been achieved, a factor of $\sim$1.5 times higher than this experiment. Including these factors leads to a significant increase in the BW yield as demonstrated in table 2. Using the same approach as before, the number of shots required to observe the BW process with optimum conditions can be estimated. Assuming a beam charge of 160 pC (corresponding to 10^9 primary electrons), figure 13(b) shows $\sim$5000 shots are required to make a 2σ observation with optimum experimental parameters.

9. Conclusions and future directions

A laser–plasma platform for investigating photon–photon physics has been presented. A high brilliance γ source with associated diagnostics for yield and spectrum was commissioned. It was found that $(7 \pm 3) \times 10^7$ γ photons/shot can be produced with energies over 100 MeV. In tandem, a high repetition rate tape-based keV x-ray source was also commissioned. Approximately $(1.4 \pm 0.5) \times 10^{12}$ x-ray photons/mm^3 in the energy range 1.3–1.5 keV were produced in the proximity of the photon collision interaction region. A magnetic transport system was implemented to successfully relay an estimated 14% of any pair particles produced in the interaction region to two SPD systems (38% of particles over 200 MeV). A CsI SPD was tested and was able to detect a signal level of 0.08 ± 0.03 positrons$_{ams}$/pC$_r$ over a background level equivalent to 0.32 ± 0.01 positrons$_{ams}$/pC$_r$. This agreed, within uncertainty, with the simulated signal level. Additionally a Timepix3 SPD system was also used, recording a signal level of (0.08 ± 0.07) positrons$_{ams}$/pC$_r$ over a background of (0.07 ± 0.04) positrons$_{ams}$/pC$_r$. These detection limits were achieved over only $\lesssim 10$ shots with the sensitivity expected to increase with the number of shots.

We estimated that with the achieved performance of the platform, over 10^9 shots would be needed to observe the linear BW process with 2σ significance. However, it is clear that the electron beam performance is crucial to the capabilities of the setup. Given optimal performance of the laser facility in question, electron beam energies of over 1.5 GeV are possible. Simulations indicate that the increased BW signal level associated with this increase in electron energy reduces the number of required shots to approximately 5000. With the repetition rate of the laser system used, this is achievable in under 28 h of operation, which would be an achievable feat. Note it has been shown recently that the stability and performance of LWFA electrons beams can be improved using machine learning techniques such as Bayesian optimisation [60–62].

In addition, we have identified possible improvements to the experiment setup to significantly increase the efficiency of the platform. Firstly, the Timepix3 SPD setup in this instance was limited to covering approximately 3 cm^2 of the signal area (just over one sensor). Increasing this detector area with the implementation of more sensors will reduce the number of required shots as more of the pair signal is collected. Under the assumption that detection rates follow Poisson statistics and that the signal and background are proportional to the detector area, the number of shots required to reach a desired significance level is proportional to the number of shots. For example doubling the number of sensors will halve the amount of data required. It is estimated that five sensors could cover the majority of the pair signal at the detector plane and reduce the number of required shots for 2σ significance to $\sim$1000; less than 6 h of data taking. With more investment and development it should also be possible to increase the efficiency of the SPDs further. This might include (but is not limited to) a fully functioning telescope system

[63] for the Timepix3 detectors to enable particle tracking, or tilting the chips to increase the capability of filtering by angular trajectory [64]. Fully calibrated Timepix3 sensors (with close to nanosecond time resolution [65]) could also employ time-of-flight distinction between signal and noise sources. An upgraded Timepix4 sensor will also soon be available which not only has a larger detection area ($\sim$7 cm^2 area) but can be interconnected to tile large areas [66].

Secondly, by making the setup more compact and bringing the interaction point closer to the γ source, a greater fraction of the γ rays can overlap with the regions of highest x ray density, increasing the pair yield. Alternatively, an active plasma lens [67] could be used to focus the electron beam, creating a narrower γ-ray beam at the collision point, allowing a greater fraction of the γ rays to interact with the regions of highest x-ray density.

Finally it is clear that reducing the detector background noise can be as important as increasing the number of pairs created or the platform efficiency. The number of shots required to make a significant measurement scales $\sim$noise/signal^2. Undertaking a more detailed design study or shielding construction to reduce the detector noise could have a major impact on the length of time that needs to be dedicated to data taking. As an example, figure 13(c) depicts the number of shots required to make a 2σ significance measurement with the optimum performance presented previously, but with a $\times 10$ reduction in background level. The number of shots required would now only be approximately 500.

It should be noted that future opportunities will be provided by increased laser capabilities coming online in the near future at other facilities across the globe [68–73]. For example, with 1 PW it is possible to produce electron beams of over 8 GeV energy [9]. As described previously the achievable electron energy is critical to the production rate of the pairs. The planned facilities are also in general closer to a repetition rate of 1 Hz, and will have increased laboratory space. The magnetic transport system and SPDs described for this platform were designed to fit within the size constraints of the Gemini target area which is less than 8 m long. With more available space it would be possible to separate the electron and γ beam dumps from the SPDs, and to introduce more shielding, reducing the background level further, which as mentioned above plays a major role in reducing the number of shots required for a measurement. More space would remove limitations in the spectral selection and focussing capabilities of the magnetic transport system.

These developments of the platform lead to exciting opportunities for investigating other photon–photon physics. This includes testing the nonlinear BW process by replacing the x-ray bath source with the direct high-intensity laser pulse. This 'clean' measurement of the nonlinear BW process without presence of an electron beam would be complementary to pair production measurements in electron–laser collisions where the trident process is expected to contribute in addition to nonlinear BW.

Improvements to the collimation of the γ ray beam and the use of higher energy x-rays to increase angular scattering will also enable the platform to explore elastic photon–photon scattering. The cross section of elastic photon–photon scattering is a factor of $\alpha^2 \approx (1/137)^2$ smaller than that of the BW process, making direct detection a significant challenge. This challenge is increased as photons which scatter from BW pairs produced in the collision would be difficult to distinguish from elastically scattered photons, unless the platform operates below the threshold for BW production. Nevertheless, this platform will be well suited to improving the experimental bounds on elastic photon photon scattering at $\sqrt{s} \approx m_e c^2$ which are currently many orders of magnitude above the predictions of QED [74, 75], and so may inform the ongoing search for physics beyond the standard model.

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Data availability statement

The data that support the findings of this study are openly available at the following URL/DOI: https://dx.doi.org/10.5281/zenodo.5027591.

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