

63

JINR E 93-397
SW 9416



2-93-307

A. I. GOLDBERG

THE PROCESS OF POLARIZATION
ON THE POLARIZED HYDROGEN TARGET
AS AN ANALYZER OF THE ^3He POLARIZATION

Report at International Symposium «Dubna, Deuteron-93»,
14—18 September, 1993

1993

In connection with the interest in experiments with polarized beams of ${}^3\text{He}$ nuclei the problem of measuring the ${}^3\text{He}$ polarization appears. It seems that the reaction $p({}^3\text{He}, {}^4\text{He})\pi^+$ on the polarized hydrogen target is appropriate for this purpose.

1. It follows from the conservation of the angular momentum and the space parity that in the case of reactions of the type $\frac{1}{2} + \frac{1}{2} \rightarrow 0 + 0$, when two fermions with spins of $1/2$ transform into two spinless bosons, the transitions from fermion singlet state are forbidden on condition that the internal parity products of the initial fermions and the final bosons are opposite [1,2,3]. Therefore, in view of the fact that π^+ is the pseudoscalar meson, the process ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ is possible only when the total spin of the ${}^3\text{He}$ nucleus and the proton equals unity (the triplet state).

Let \vec{P}_1 and \vec{P}_2 be the independent polarization vectors of the ${}^3\text{He}$ nucleus and the proton, respectively. It is easy to show that the probabilities of detecting the $({}^3\text{He}, p)$ -system in the triplet states with the total spin projections $M = \pm 1, 0$ to the ${}^3\text{He}$ momentum are the following:

$$W_{\pm 1}^{(t)} = \frac{1}{4} (1 \pm \vec{P}_1 \vec{\ell}) (1 \pm \vec{P}_2 \vec{\ell}), \quad (1)$$

$$W_0^{(t)} = \frac{1}{4} [1 + \vec{P}_1 \vec{P}_2 - 2(\vec{P}_1 \vec{\ell})(\vec{P}_2 \vec{\ell})], \quad (2)$$

where $\vec{\ell}$ is the unity vector in the direction of the ${}^3\text{He}$ momentum.

The total probability of detecting the triplet states is equal to

$$W^{(t)} = \frac{1}{4} (3 + \vec{P}_1 \vec{P}_2), \quad (3)$$

and the probability of detecting the singlet state

$$W^{(s)} = \frac{1}{4} (1 - \vec{P}_1 \vec{P}_2). \quad (4)$$

2. Let us consider the case of the flight of the final ${}^4\text{He}$ nucleus, produced in the reaction ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$, in the forward direction. Then, in consequence of the conservation of the projection of the total angular momentum, this process is possible only at the zero value of the total spin projection of $({}^3\text{He}, p)$ -system to the reaction axis, and it is forbidden for projections equalling (+1) and (-1). As a result, in accordance with formula (2) the differential cross-section of the process ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ at the zero angle depends simply on polarizations of the ${}^3\text{He}$ beam and the proton target:

$$\frac{d\sigma}{d\Omega}(0) = \left(\frac{d\sigma}{d\Omega}(0)\right)_{\text{unpol}} \left[1 + \vec{P}_1 \vec{P}_2 - 2(\vec{P}_1 \vec{\ell})(\vec{P}_2 \vec{\ell}) \right], \quad (5)$$

where $(d\sigma/d\Omega)_{\text{unpol}}(0)$ is the differential cross-section of the same reaction in the case when both the beam and the target are unpolarized.

Under standard methods of the beam polarization the vector \vec{P}_1 is perpendicular to the momentum ($\vec{P}_1 \vec{\ell} = 0$). Then the relation (5) gives

$$\frac{d\sigma}{d\Omega}(\theta) = \left(\frac{d\sigma}{d\Omega}(\theta) \right)_{\text{unpol}} \left(1 + |\vec{P}_1| |\vec{P}_2| \cos \alpha \right), \quad (6)$$

where α is the angle between the polarization directions of the ${}^3\text{He}$ nucleus and the proton. The relative change of the cross-section with the ${}^3\text{He}$ polarization reversal (or the proton polarization one) equals

$$\eta = \left(\frac{d\sigma^{(+)}(\theta)}{d\Omega} - \frac{d\sigma^{(-)}(\theta)}{d\Omega} \right) / \left(\frac{d\sigma^{(+)}(\theta)}{d\Omega} + \frac{d\sigma^{(-)}(\theta)}{d\Omega} \right) = |\vec{P}_1| |\vec{P}_2| \cos \alpha. \quad (7)$$

Thus, it is possible to determine the degree of the ${}^3\text{He}$ beam polarization, knowing the asymmetry η , the degree of the proton target polarization and the angle between the ${}^3\text{He}$ and proton polarization vectors.

3. At non-zero angles of the flight of the ${}^4\text{He}$ nucleus, taking into account the parity conservation, the cross-section of the reaction ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$, integrated over azimuth angles (as well as the total cross-section of the reaction), has the following structure:

$$\sigma(\theta) = \frac{1}{4} \sigma_{t,0}(\theta) [1 + \vec{P}_1 \vec{P}_2 - 2 (\vec{P}_1 \vec{\ell})(\vec{P}_2 \vec{\ell})] + \frac{1}{2} \sigma_{t,1}(\theta) [1 + (\vec{P}_1 \vec{\ell})(\vec{P}_2 \vec{\ell})], \quad (8)$$

where $\sigma_{t,\mu}(\theta)$ is the cross-section corresponding to the triplet state of the $({}^3\text{He}, p)$ -system with the projection μ to the direction of the ${}^3\text{He}$ momentum. Here we have taken into account that the interference between states with different projections of the total spin disappears after the inte-

gration over azimuth angles, and, besides, in consequence of the parity conservation the equality

$$\overline{\sigma}_{t,+1}(\theta) = \overline{\sigma}_{t,-1}(\theta) \quad (9)$$

is satisfied. At very small angles the triplet cross-section corresponding to projections (± 1) tends to zero being proportional to the square of the angle ($\overline{\sigma}_{t,1} \sim \theta^2$).

The formula (8) can be rewritten in the form

$$\sigma(\theta) = \overline{\sigma}_1(\theta) + \overline{\sigma}_2(\theta)(\vec{P}_1 \vec{P}_2) + \overline{\sigma}_3(\theta)(\vec{P}_1 \vec{P}_2)(\vec{P}_2 \vec{P}_1), \quad (10)$$

where

$$\begin{aligned} \overline{\sigma}_1(\theta) &= \frac{1}{4} \overline{\sigma}_{t,0}(\theta) + \frac{1}{2} \overline{\sigma}_{t,1}(\theta), \quad \overline{\sigma}_2(\theta) = \frac{1}{4} \overline{\sigma}_{t,0}(\theta), \\ \overline{\sigma}_3(\theta) &= \frac{1}{2} (\overline{\sigma}_{t,1}(\theta) - \overline{\sigma}_{t,0}(\theta)). \end{aligned} \quad (11)$$

It is evident that values $\overline{\sigma}_1(\theta)$ and $\overline{\sigma}_2(\theta)$ are positive, and $\overline{\sigma}_1(\theta) \geq \overline{\sigma}_2(\theta)$. When the polarization vector of a ${}^3\text{He}$ nucleus is perpendicular to its momentum the asymmetry at the polarization reversal is expressed as

$$\eta = A |\vec{P}_1| |\vec{P}_2| \cos d, \quad A = \frac{\overline{\sigma}_2}{\overline{\sigma}_1} \leq 1. \quad (12)$$

In accordance with the relations (11) the following equality is valid:

$$\overline{\sigma}_1 = 3\overline{\sigma}_2 + \overline{\sigma}_3. \quad (13)$$

The same result takes place for the annihilation processes $\bar{p}p \rightarrow \pi^+\pi^-, K^+K^-, K^0\bar{K}^0$ [3].

4. The process ${}^3\text{He}(p,\pi^+){}^4\text{He}$ on the helium target, as well as the time-reversal reactions $\pi^+ + {}^4\text{He} \rightarrow p + {}^3\text{He}$, $\pi^- + {}^4\text{He} \rightarrow n + {}^3\text{H}$, were studied experimentally in a number of works (see, for example, [4-8]).

In accordance with experimental data, when both the beam and the target are unpolarized, at ${}^3\text{He}$ laboratory kinetic energies in the range of (1+2) GeV the small-angle differential cross-section of the process ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ is of $\sim 10 \mu\text{b}/\text{sr}$ in the center-of-mass system. The cross-section maximum of about $15 \mu\text{b}/\text{sr}$ at the ${}^3\text{He}$ laboratory kinetic energy $T \approx 1.25 \text{ GeV}$ and $\theta_{c.m.} = 0$ corresponds, apparently, to the formation of $\Delta(33)$ -resonance on a bound nucleon [5] (in the proton rest frame $d\sigma/d\Omega^{(max)}(0) \approx 250 \mu\text{b}/\text{sr}$).

Author would like to thank N.M.Piskunov, I.M.Sitnik and E.A.Strokovsky for their interest in this work and valuable discussions.

References

1. Bylenky S.M., Ryndin R.M. Phys.Lett., 1963, v.6, p.217.
2. Bylenky S.M., Ryndin R.M. Zh.Teor.Exp.Fiz., 1963, v.45, p. 1192.

3. Lyuboshitz V.L. Yad.Fiz., 1970, v.12, p.199
(Sov. J.Nucl.Phys., 1970, v.12, p.107).
4. Cabathuler K. et al. Nucl.Phys., 1972, v. B40, p. 32.
5. Tatischeff B. et al. Phys.Lett., 1976, v.B63, p. 158.
6. Höistad B. et al. Phys.Rev., 1984, v. C29, p.553.
7. Källne J. et al. Phys.Rev., 1981, v. C24, p. 1102.
8. Källne J. et al. Phys.Rev., 1983, v. C28, p.304.

Received by Publishing Department
on November 2, 1993.

SUBJECT CATEGORIES OF THE JINR PUBLICATIONS

Index	Subject
1.	High energy experimental physics
2.	High energy theoretical physics
3.	Low energy experimental physics
4.	Low energy theoretical physics
5.	Mathematics
6.	Nuclear spectroscopy and radiochemistry
7.	Heavy ion physics
8.	Cryogenics
9.	Accelerators
10.	Automatization of data processing
11.	Computing mathematics and technique
12.	Chemistry
13.	Experimental techniques and methods
14.	Solid state physics. Liquids
15.	Experimental physics of nuclear reactions at low energies
16.	Health physics. Shieldings
17.	Theory of condensed matter
18.	Applied researches
19.	Biophysics

Принимается подписка на препринты, сообщения Объединенного института ядерных исследований и «Краткие сообщения ОИЯИ».

Установлена следующая стоимость подписки на 12 месяцев на издания ОИЯИ, включая пересылку, по отдельным тематическим категориям:

Индекс	Тематика	Цена подписки на год
1.	Экспериментальная физика высоких энергий	915 р.
2.	Теоретическая физика высоких энергий	2470 р.
3.	Экспериментальная нейтронная физика	365 р.
4.	Теоретическая физика низких энергий	735 р.
5.	Математика	460 р.
6.	Ядерная спектроскопия и радиохимия	275 р.
7.	Физика тяжелых ионов	185 р.
8.	Криогеника	185 р.
9.	Ускорители	460 р.
10.	Автоматизация обработки экспериментальных данных	560 р.
11.	Вычислительная математика и техника	560 р.
12.	Химия	90 р.
13.	Техника физического эксперимента	720 р.
14.	Исследования твердых тел и жидкостей ядерными методами	460 р.
15.	Экспериментальная физика ядерных реакций при низких энергиях	460 р.
16.	Дозиметрия и физика защиты	90 р.
17.	Теория конденсированного состояния	365 р.
18.	Использование результатов и методов фундаментальных физических исследований в смежных областях науки и техники	90 р.
19.	Биофизика	185 р.
	«Краткие сообщения ОИЯИ» (6 выпусков)	560 р.

Подписка может быть оформлена с любого месяца года.

По всем вопросам оформления подписки следует обращаться в издательский отдел ОИЯИ по адресу: 141980, г.Дубна, Московской области

Любошиц В.Л.

E2-93-397

Процесс ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ на поляризованной водородной мишени как анализатор поляризации ${}^3\text{He}$

Исследуется зависимость эффективного сечения процесса ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ от векторов поляризации ядра ${}^3\text{He}$ и протона. Показано, что реакция $p({}^3\text{He}, {}^4\text{He})\pi^+$ на поляризованной водородной мишени может быть использована для измерения поляризации ${}^3\text{He}$.

Работа выполнена в Лаборатории высоких энергий ОИЯИ.

Препринт Объединенного института ядерных исследований. Дубна, 1993

Lyuboshitz V.L.

E2-93-397

The Process ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ on the Polarized Hydrogen Target as an Analyzer of the ${}^3\text{He}$ Polarization

The cross-section dependence of the process ${}^3\text{He} + p \rightarrow {}^4\text{He} + \pi^+$ on the ${}^3\text{He}$ and proton polarization vectors is investigated. It is shown that the reaction $p({}^3\text{He}, {}^4\text{He})\pi^+$ on the polarized hydrogen target can be used for measuring the ${}^3\text{He}$ polarization.

The investigation has been performed at the Laboratory of High Energies, JINR.

Preprint of the Joint Institute for Nuclear Research. Dubna, 1993

7 p.

Редактор Е.И.Хижняк. Макет Р.Д.Фоминой

Подписано в печать 24.12.93

Формат 60x90/16. Офсетная печать. Уч.-изд. листов 0,53

Тираж 470. Заказ 46881

Издательский отдел Объединенного института ядерных исследований
Дубна Московской области