

# Determination of the Higgs properties with the ATLAS detector

---

**Tamara Vazquez Schroeder, on behalf of the ATLAS Collaboration**

*McGill University*

*E-mail:* [tamara.vazquez.schroeder@cern.ch](mailto:tamara.vazquez.schroeder@cern.ch)

Combined measurements of Higgs boson production cross sections and branching fractions are presented using the  $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ^* \rightarrow 4\ell$  decay channels based on  $36.1 \text{ fb}^{-1}$  of proton-proton collision data recorded at  $\sqrt{s} = 13 \text{ TeV}$  by the ATLAS experiment at the LHC. No significant deviations from the Standard Model expectations are observed.

*EPS-HEP 2017, European Physical Society conference on High Energy Physics  
5-12 July 2017  
Venice, Italy*



## 1. Motivation

The Higgs boson discovery in July 2012 by the ATLAS [1] and CMS [2] experiments was a major milestone in the Large Hadron Collider’s (LHC) physics program and a breakthrough in our knowledge of high-energy particle physics. Since then, a wide range of measurements of its properties have been performed using proton-proton collision data produced by the LHC at centre-of-mass energies of  $\sqrt{s} = 7$  and 8 TeV in Run 1 [3], and at  $\sqrt{s} = 13$  TeV in Run 2 using  $13.3 \text{ fb}^{-1}$  of data by the ATLAS experiment [4] and using  $35.9 \text{ fb}^{-1}$  of data by the CMS experiment [5], all of them consistent with the Standard Model (SM) predictions. The combination of Higgs-boson decay channels maximises the precision of these measurements, allows more model-independent measurements, and constrains the coupling to the Higgs boson with less assumptions.

The combined measurements of Higgs boson production cross sections and branching fractions presented here use the  $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ^* \rightarrow 4\ell$  decay channels and an integrated luminosity of  $36.1 \text{ fb}^{-1}$  of proton-proton collision data recorded at  $\sqrt{s} = 13$  TeV by the ATLAS experiment at the LHC during 2015 and 2016 [6]. The dataset consists of approximately three times the luminosity of the results in Ref. [4], allowing improved measurements of Higgs boson production and decay.

## 2. Input analyses

The input decay channels used in this combination are the “golden” Higgs discovery channels:  $H \rightarrow \gamma\gamma$  [7] and  $H \rightarrow ZZ^* \rightarrow 4\ell$  [8]. The analyses of the individual channels separate the measured events into exclusive kinematic and topological categories. Each of these categories targets a specific Higgs boson production mode: gluon fusion (ggF),  $gg \rightarrow H$ , weak vector-boson fusion (VBF),  $qq' \rightarrow qq'H$ , associated production with a W or Z boson,  $q\bar{q}' \rightarrow VH$ , and associated production with a pair of top quarks,  $q\bar{q}/gg \rightarrow t\bar{t}H$ .

The  $H \rightarrow \gamma\gamma$  analysis defines 31 exclusive event categories. The categories are structured hierarchically to prioritise the selection of signal events from processes with the smallest production cross sections,  $t\bar{t}H$  and  $tHX$ , followed by  $VH$  where categories are distinguished by the vector-boson decays, and finally VBF and ggF.

The  $H \rightarrow ZZ^* \rightarrow 4\ell$  analysis defines 9 event categories to distinguish the  $t\bar{t}H$ ,  $VH$ , VBF, and ggF production modes. The analysis reconstructs the intermediate Z bosons using their decays to electrons and muons, and requires the four-lepton invariant mass to be between 118 and 129 GeV.

## 3. Measurements and results

Combined cross section measurements are reported in various global fit models in order of increasing granularity of Higgs boson production modes and kinematics. First, the total inclusive production cross section ( $pp \rightarrow H + X$ ) and the ratio of the total signal yield with respect to the SM prediction ( $\mu$ ) are presented in Sections 3.1 and 3.2. Then, the inclusive cross sections of individual production processes and their ratios are reported in Sections 3.3 and 3.4. Finally, cross sections in kinematic regions of the production processes in the framework of simplified template cross sections (STXS) are detailed in Section 3.5. To characterise deviations from the SM expectations,

39 the results are translated into measurements of the multiplicative coefficients  $\kappa_i$  of Higgs-boson  
40 couplings in the SM, as shown in Section 3.6.

### 41 3.1 Total cross section

42 The total cross section is measured based on the inclusive event yields in each decay channel.  
43 The event yields are corrected for detector effects, the fiducial acceptance relative to the full phase  
44 space, and branching fractions. The corrections are derived using the SM predictions for the cross  
45 section ratios between the different production modes. The total  $pp \rightarrow H + X$  cross section at  
46 centre-of-mass energy of 13 TeV is  $57_{-5.9}^{+6.0}(stat.)_{-3.3}^{+4.0}(syst.)$  pb, in good agreement with the SM  
47 prediction at N3LO QCD of  $55.6_{-3.4}^{+2.4}$  pb [9]. The measurement is statistically limited.

### 48 3.2 Global signal strength

49 The global signal strength  $\mu$  is determined from a fit to a single parameter defined as the ratio  
50 of the total observed yield to its SM expectation,  $\mu = \frac{\sigma \times B}{(\sigma \times B)_{SM}}$ . This parameter is applied as a  
51 single scaling factor to all production processes and decay modes. The global signal strength is  
52 measured to be  $\mu = 1.09 \pm 0.12 = 1.09 \pm 0.09(stat.)_{-0.05}^{+0.06}(exp.)_{-0.05}^{+0.06}(th.)$ . The combined  $\mu$  value  
53 lies within the signal strengths of the individual measurements in each decay mode:  $\mu = 0.99 \pm 0.14$   
54 in  $H \rightarrow \gamma\gamma$  and  $\mu = 1.28_{-0.19}^{+0.21}$  in  $H \rightarrow ZZ^* \rightarrow 4\ell$ . The event categorisation reduces the statistical  
55 uncertainty relative to the total cross section measurement. The measurement is consistent with  
56 the SM prediction with a p-value of  $p_{SM} = 47\%$ . The dominant theoretical uncertainties originate  
57 from the QCD scale and PDF variations, whereas the dominant experimental uncertainties are due  
58 to luminosity and electron/photon energy resolution.

### 59 3.3 Production cross sections

60 A simultaneous fit is performed for the cross sections of ggF, VBF,  $VH$ , and  $t\bar{t}H$  for  $|y_H| <$   
61 2.5 and assuming SM branching fractions. The measurement of  $bbH$  is included in ggF, and the  
62 measurement of  $tH$  is included in  $t\bar{t}H$ . The process  $gg \rightarrow ZH$  is fully attributed to  $ZH$ .

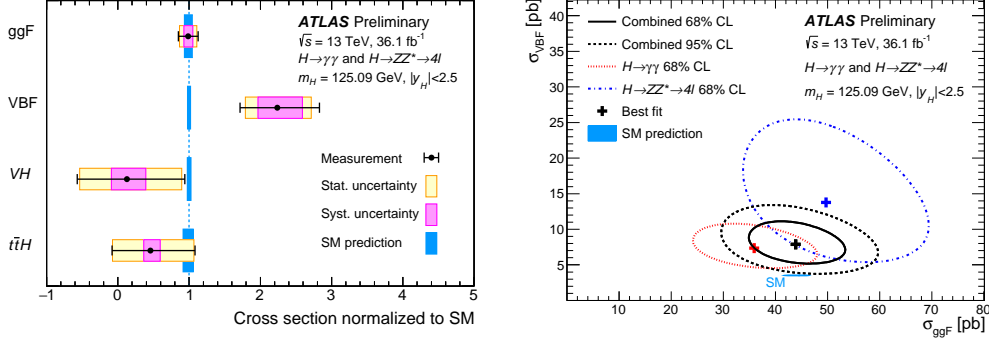
63 Figure 1 (left) shows the production cross section measurement results. The cross section for  
64 the VBF production is  $7.9_{-1.8}^{+2.1}$  and has the largest deviation from its SM expectation,  $3.52_{-0.07}^{+0.08}$  [9].  
65 The four-dimensional compatibility between the measurement and the SM prediction corresponds  
66 to a p-value of  $p_{SM} = 5\%$ .

67 Figure 1 (right) shows the measured likelihood contours in the VBF and ggF cross section  
68 plane from the individual fits in  $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ^* \rightarrow 4\ell$ , as well as for the combined fit.  
69 The cross sections for  $VH$  and  $t\bar{t}H$  are profiled with the data. The two-dimensional compatibility  
70 between the measurement and the SM prediction corresponds to a p-value of  $p_{SM} = 3\%$ .

### 71 3.4 Ratios of cross sections and branching fractions

72 A combined fit to data is performed with the production cross sections of VBF,  $VH$ , and  $t\bar{t}H$   
73 normalised to ggF, and with the branching ratio of  $H \rightarrow \gamma\gamma$  normalised to  $H \rightarrow ZZ^* \rightarrow 4\ell$ . The  
74 product of the cross section and the branching fraction can be expressed in terms of these ratios:

$$\sigma_i \times BR_f = \sigma_{ggF} \times BR_{H \rightarrow 4\ell} \times \frac{\sigma_i}{\sigma_{ggF}} \times \frac{BR_f}{BR_{H \rightarrow 4\ell}} \quad (3.1)$$



**Figure 1:** (Left) cross sections for ggF, VBF, VH, and  $t\bar{t}H$  normalised to the SM predictions and measured with the assumption of SM branching fractions. The black error bars and pink and yellow boxes show the total, systematic, and statistical uncertainties in the measurements, respectively. The blue bands indicate the theoretical uncertainties in the predictions [6]. (Right) measured likelihood contours in the  $\sigma_{VBF}$  versus  $\sigma_{ggF}$  plane in  $H \rightarrow \gamma\gamma$  (red) and  $H \rightarrow ZZ^* \rightarrow 4\ell$  (blue), as well as their combination (black) [6], together with the SM prediction (light blue) [9].

75 where  $i$  is a Higgs-boson production process,  $f$  is a final state, and  $BR_f$  is the branching  
 76 fraction for the Higgs boson to decay into  $f$ .

77 Table 1 shows the measurements of  $\sigma_{ggF} \times BR_{H \rightarrow 4\ell}$ ,  $\sigma_{VBF}/\sigma_{ggF}$ ,  $\sigma_{VH}/\sigma_{ggF}$ ,  $\sigma_{t\bar{t}H}/\sigma_{ggF}$ , and  
 78  $BR_{H \rightarrow \gamma\gamma}/BR_{H \rightarrow 4\ell}$ . The five-dimensional compatibility between the measurements and the SM pre-  
 79 dictions corresponds to a p-value of  $p_{SM} = 3\%$ .

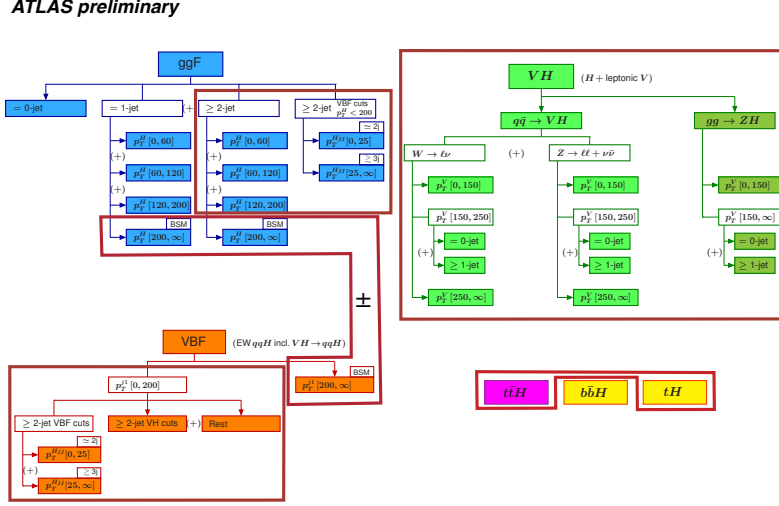
Quantity	Result	Uncertainty			SM prediction		
		Total	Stat.	Exp.		Th.	
$\sigma_{ggF} \cdot B_{4\ell}$	[fb]	6.6	$+1.2$ $-1.0$	$+1.1$ $-1.0$	$\pm 0.4$ $\pm 0.2$	$5.6^{+0.3}$ $-0.4$	
$B_{\gamma\gamma}/B_{4\ell}$		12.5	$+2.8$ $-2.3$	$+2.6$ $-2.2$	$+0.9$ $-0.7$	$\pm 0.2$	$18.1 \pm 0.2$
$\sigma_{VBF}/\sigma_{ggF}$	$[10^{-2}]$	21.5	$+8.5$ $-6.3$	$+7.3$ $-5.6$	$+2.8$ $-1.7$	$+3.6$ $-2.2$	$7.9^{+0.4}$ $-0.6$
$\sigma_{VH}/\sigma_{ggF}$	$[10^{-2}]$	0.2	$+4.5$ $-3.4$	$+4.2$ $-3.2$	$+1.2$ $-0.9$	$+0.9$ $-0.4$	$4.5^{+0.2}$ $-0.3$
$\sigma_{t\bar{t}H}/\sigma_{ggF}$	$[10^{-2}]$	0.7	$+1.0$ $-0.9$	$+1.0$ $-0.9$	$+0.2$ $-0.1$	$\pm 0.1$	$1.3 \pm 0.1$

**Table 1:** Best-fit values and uncertainties of  $\sigma_{ggF} \times BR_{H \rightarrow 4\ell}$  and the ratios of cross sections and branching fractions.

### 80 3.5 Simplified template cross sections

81 The simplified template cross section (STXS) framework defines a set of kinematic regions for  
 82 each production process and combines these with the ratios of branching fractions for the various  
 83 decay channels in order to probe the properties of the Higgs boson. The splitting of the produc-  
 84 tion modes is based on kinematic variables such as the jet multiplicity, the transverse momentum  
 85 of the Higgs boson, and the transverse momentum of the leading jet. The basic STXS scheme  
 86 has been agreed between the ATLAS experiment, the CMS experiment, and theorists within the  
 87 LHCXSWG [9] and is schematised in Figure 2 (“stage-1”). The cross sections are then measured  
 88 in these mutual exclusive regions of the phase space called “truth bins”. As a temporary solution to

89 increase the sensitivity to SM production and minimise correlations, some bins have been merged  
 90 for the current results. All regions require  $|y_H| < 2.5$ . This framework allows to maximise experi-  
 91 mental sensitivity while minimising theory dependence. Additionally, these cross sections can be  
 92 interpreted in various beyond-the-SM models.



**Figure 2:** The merged STXS stage-1 regions defined for the measurements [9]. All regions surrounded by red boxes are merged, except for the sum and difference indicated by the “±” sign connecting two merged  $gg \rightarrow H$  regions with one  $qq \rightarrow Hqq$  region. The  $bbH$  region is merged with the  $gg \rightarrow H$  bins [6].

93 Table 2 shows the fitted values of the ratio of branching ratios  $B_{H \rightarrow \gamma\gamma}/B_{H \rightarrow 4\ell}$  and of the sim-  
 94 plified template cross sections. The results show good overall agreement with the SM predictions.  
 95 The ten-dimensional compatibility between the measurement and the SM prediction corresponds  
 96 to a p-value of  $p_{SM} = 9\%$ .

Measurement region	Result	Uncertainty			SM prediction
		Total	Stat.	Syst.	
$B_{\gamma\gamma}/B_{4\ell}$	12.5	$^{+2.8}_{-2.3}$	$\left( \begin{smallmatrix} +2.6 \\ -2.2 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +0.8 \\ -0.6 \end{smallmatrix} \right)$	$18.1 \pm 0.2$
$gg \rightarrow H$ (0-jet)	29.7	$^{+7.3}_{-6.4}$	$\left( \begin{smallmatrix} +6.6 \\ -6.0 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +3.1 \\ -2.4 \end{smallmatrix} \right)$	27.6 ± 1.9 pb
$gg \rightarrow H$ (1-jet, $p_T^H < 60$ GeV)	4.4	$^{+4.8}_{-4.5}$	$\left( \begin{smallmatrix} +4.4 \\ -4.1 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +1.7 \\ -1.8 \end{smallmatrix} \right)$	6.6 ± 0.9 pb
$gg \rightarrow H$ (1-jet, $60 \leq p_T^H < 120$ GeV)	4.6	$^{+2.8}_{-2.4}$	$\left( \begin{smallmatrix} +2.7 \\ -2.4 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +0.7 \\ -0.5 \end{smallmatrix} \right)$	4.6 ± 0.7 pb
$gg \rightarrow H$ (1-jet, $120 \leq p_T^H < 200$ GeV)	1.6	$^{+1.1}_{-0.9}$	$\left( \begin{smallmatrix} +1.0 \\ -0.9 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +0.3 \\ -0.2 \end{smallmatrix} \right)$	0.75 ± 0.15 pb
$gg \rightarrow H$ ( $\geq 2$ -jet, $p_T^H < 200$ GeV or VBF-like)	10.6	$^{+4.7}_{-4.2}$	$\left( \begin{smallmatrix} +4.3 \\ -3.9 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +1.9 \\ -1.4 \end{smallmatrix} \right)$	4.8 ± 1.0 pb
$gg \rightarrow H$ ( $\geq 1$ -jet, $p_T^H \geq 200$ GeV) + $qq \rightarrow Hqq$ ( $p_T^H \geq 200$ GeV)	1.9	$^{+0.9}_{-0.7}$	$\left( \begin{smallmatrix} +0.8 \\ -0.7 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +0.3 \\ -0.2 \end{smallmatrix} \right)$	0.81 ± 0.16 pb
$qq \rightarrow Hqq$ ( $p_T^H < 200$ GeV)	9.8	$^{+4.3}_{-3.5}$	$\left( \begin{smallmatrix} +4.0 \\ -3.2 \end{smallmatrix} \right)$	$\left( \begin{smallmatrix} +1.5 \\ -1.4 \end{smallmatrix} \right)$	4.58 $^{+0.15}_{-0.18}$ pb
$gg/qq \rightarrow H\ell\ell/H\ell\nu$	0.2	$^{+0.9}_{-0.7}$	$\left( \begin{smallmatrix} +0.8 \\ -0.7 \end{smallmatrix} \right)$	$\pm 0.2$	0.63 $^{+0.03}_{-0.06}$ pb
$q\bar{q}/gg \rightarrow t\bar{t}H$	0.3	$^{+0.5}_{-0.4}$	$\left( \begin{smallmatrix} +0.5 \\ -0.4 \end{smallmatrix} \right)$	$\pm 0.1$	0.59 $^{+0.04}_{-0.05}$ pb

**Table 2:** Best-fit values and uncertainties of the ratio of branching ratios  $B_{H \rightarrow \gamma\gamma}/B_{H \rightarrow 4\ell}$  and of the simplified template cross sections. The SM predictions are shown for each region.

### 97 3.6 $\kappa$ framework

98 The  $\kappa$  framework expresses the product of cross sections and branching ratios ( $\sigma \times B$ ) in terms  
99 of coupling modifiers  $\kappa$  [10]:

$$\sigma_i \times B(H \rightarrow f) = \frac{\kappa_i^2 \kappa_f^2}{\kappa_H^2} \sigma_i^{SM} \times B^{SM}(H \rightarrow f) \quad (3.2)$$

100 where  $i$  and  $f$  are the initial and final states, respectively,  $\sigma_i^{SM}$  is the SM production cross  
101 section, and  $B^{SM}(H \rightarrow f)$  is the SM value of the branching ratio ( $H \rightarrow f$ ). In the absence of non-  
102 SM decays the coefficient to the Higgs boson width,  $\kappa_H$ , is a function of the other  $\kappa$  parameters;  
103 similarly, in the absence of non-SM loops in the effective  $Hgg$  and  $H\gamma\gamma$  couplings, the correspond-  
104 ing coefficients  $\kappa_g$  and  $\kappa_\gamma$  are functions of the other parameters. Three models are tested using  
105 progressively relaxed assumptions on the coupling relationships.

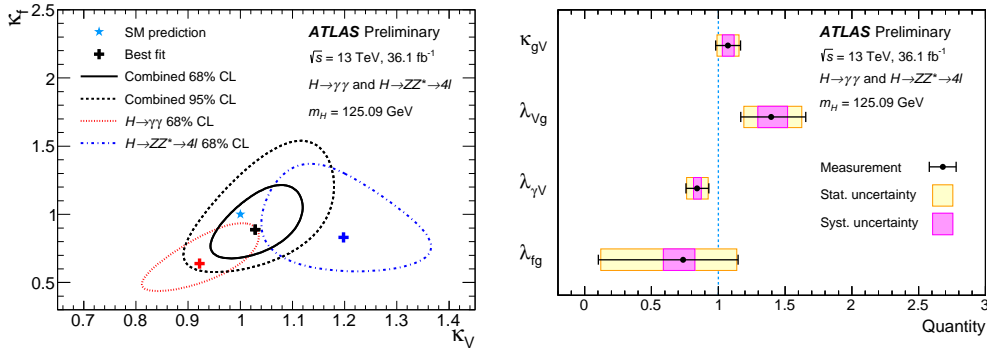
106 In the first model, a two-parameter fit of  $\kappa_f$  and  $\kappa_V$  is performed. The  $H \rightarrow ZZ^* \rightarrow 4\ell$  branch-  
107 ing fraction is proportional to  $\kappa_V^2$ , while the  $H \rightarrow \gamma\gamma$  branching fraction depends on  $\kappa_f^2$ ,  $\kappa_V^2$ , and  
108  $\kappa_V \kappa_f$  due to significant contributions from top-quark and W-boson loops, and their interference, in  
109 the decay. Both branching fractions are inversely dependent on these three  $\kappa$  combinations through  
110 the total width of the Higgs boson. The dominant production mechanisms of ggF and VBF de-  
111 pend on  $\kappa_f^2$  and  $\kappa_V^2$ , respectively. The branching fraction to new states is assumed to be zero. The  
112 fit results are summarised in Figure 3 (left) and show a positive correlation of 54% due in part  
113 to the destructive interference between the top-quark and W-boson loops in the  $H\gamma\gamma$  decay. The  
114 best-fit values and uncertainties are  $\kappa_V = 1.03 \pm 0.06$  and  $\kappa_f = 0.89_{-0.15}^{+0.20}$ . The two-dimensional  
115 compatibility with the SM prediction is  $p_{SM} = 52\%$ .

116 In the second model, the effective couplings  $\kappa_g$  and  $\kappa_\gamma$  capture all loop contributions to the  
117 Higgs-boson interaction with gluons and photons, respectively. New loop processes would appear  
118 in these modifiers rather than being absorbed by the  $\kappa_f$  and  $\kappa_V$  modifiers. In this model, production  
119 and decay modes other than ggF,  $H \rightarrow gg$  and  $H \rightarrow \gamma\gamma$  are fixed to their SM expectations. Similar  
120 to the first model, the branching fraction to new states is assumed to be zero. The two-parameter fit  
121 for  $\kappa_g$  and  $\kappa_\gamma$  shows a strong anti-correlation of -64% because the leading constraint comes from  
122  $H \rightarrow \gamma\gamma$  in the gluon fusion channel. The best-fit values and uncertainties are  $\kappa_g = 1.08_{-0.10}^{+0.11}$  and  
123  $\kappa_\gamma = 0.93_{-0.08}^{+0.09}$ . The two-dimensional compatibility with the SM prediction is  $p_{SM} = 68\%$ .

124 Finally, a generic model based on a set of four ratios is constructed to probe the loop vertices  
125 ( $\kappa_g$ ,  $\kappa_\gamma$ ), total width ( $\kappa_H$ ), and the fermion and vector couplings ( $\kappa_f$ ,  $\kappa_V$ ):  $\kappa_{gV} = \kappa_g \kappa_V / \kappa_H$ ,  $\lambda_{Vg} =$   
126  $\kappa_V / \kappa_g$ ,  $\lambda_{fg} = \kappa_f / \kappa_g$ , and  $\lambda_{\gamma V} = \kappa_\gamma / \kappa_V$ . The inclusion of  $\kappa_H$  in the parameterisation allows for  
127 non-SM decays of the Higgs boson. The results are summarised in Figure 3 (right). The four-  
128 dimensional compatibility with the SM prediction is  $p_{SM} = 15\%$ .

## 129 4. Conclusions and Outlook

130 Measurements of Higgs boson production cross sections and branching ratios have been pre-  
131 sented for the combination of  $H \rightarrow \gamma\gamma$  and  $H \rightarrow ZZ^* \rightarrow 4\ell$  decay channels based on  $36.1 \text{ fb}^{-1}$  of  
132 proton-proton collision data recorded during 2015 and 2016 at  $\sqrt{s} = 13 \text{ TeV}$  by the ATLAS exper-  
133 iment at the LHC. No significant deviations from the Standard Model expectations are observed.



**Figure 3:** (Left) contours at 68% and 95% CL in the  $(\kappa_f, \kappa_V)$  plane for the individual decay modes  $H \rightarrow ZZ^* \rightarrow 4\ell$  (blue) and  $H \rightarrow \gamma\gamma$  (red), and the combination (black). The SM prediction is shown as a blue star and lies within the 68% contour of the best fit combined values [6]. (Right) best-fit values and uncertainties of  $\kappa_{gV}$ ,  $\lambda_{Vg}$ ,  $\lambda_{fg}$ , and  $\lambda_{\gamma V}$ . All parameters are defined to be unity in the SM [6].

134 Some tensions at the level of  $2\sigma$  can be observed in the VBF production cross section compared to  
 135 the SM prediction, but all compatibility p-values are a few percents or higher. The first measure-  
 136 ments using a simplified version of “stage-1” STXS have been performed. The sensitivity to finer  
 137 splitting of kinematics is expected to increase with higher data statistics.

## 138 References

- 139 [1] ATLAS Collaboration, Phys. Lett. B 716 (2012) 1.  
 140 [2] CMS Collaboration, Phys. Lett. B 716 (2012) 30.  
 141 [3] ATLAS and CMS Collaborations, JHEP 1608 (2016) 045.  
 142 [4] ATLAS Collaboration, <https://cds.cern.ch/record/2206272>, ATLAS-CONF-2016-081 (2016).  
 143 [5] CMS Collaboration, <https://cds.cern.ch/record/2257530>, CMS-PAS-HIG-17-015 (2017).  
 144 [6] ATLAS Collaboration, <https://cds.cern.ch/record/2273854>, ATLAS-CONF-2017-047 (2017).  
 145 [7] ATLAS Collaboration, <https://cds.cern.ch/record/2273852>, ATLAS-CONF-2017-045 (2017).  
 146 [8] ATLAS Collaboration, <https://cds.cern.ch/record/2265796>, ATLAS-CONF-2017-032 (2017).  
 147 [9] D. de Florian et al., LHC Higgs Cross Section Working Group, arXiv: 1610.07922 [hep-ph] (2016).  
 148 [10] S. Heinemeyer et al., LHC Higgs Cross Section Working Group, arXiv: 1307.1347 [hep-ph] (2013).