



ADIABATIC DAMPING OF LARGE PHASE OSCILLATIONS

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The phase equation  $\frac{d}{dt} \left( \frac{E_s \dot{\varphi}}{-K} \right) + \frac{h \omega_0^2 e V}{2\pi} (\sin \varphi - \sin \varphi_0) = 0$

can be derived from the Hamiltonian  $H = \frac{1}{2} \left( \frac{-K}{E_s} \right)^2 P^2 + \frac{h \omega_0^2 e V}{2\pi} [u(\varphi) - u(\varphi_0)] = W$

where  $P = \frac{E_s \dot{\varphi}}{-K}$  ;  $K = \alpha - \left( \frac{M c^2}{E_s} \right)^2$  ;  $u(\varphi) = \cos \varphi + \varphi \sin \varphi_0$

Stable phase oscillations occur about the stable phase  $\varphi_0$

if the maximum phase  $\varphi_m < \pi - \varphi_0$

The action of a stable oscillation is  $J = \oint p d\varphi = \sqrt{\frac{h \omega_0^2 e V E_s \cos \varphi_0}{2\pi - K}} F(\varphi_m, \varphi_0)$

where  $F(\varphi_m, \varphi_0) = \int_{\cos \varphi_0}^{\cos \varphi_m} \sqrt{u(\varphi) - u(\varphi_m)} d\varphi$

The phase oscillation angular frequency  $\Omega = \frac{2\pi W}{J} = \omega_0 \sqrt{\frac{h e V \cos \varphi_0 (-K)\pi}{E_s}} \frac{G}{F}$

The maximum relative energy spread  $\frac{\Delta E}{E_s} = \beta_s \sqrt{\frac{e V \cos \varphi_0}{-K \pi h E_s}} \cdot \sqrt{G}$

The " " momentum spread  $\frac{\Delta P}{P_s} = \frac{1}{\beta_s^2} \frac{\Delta E}{E_s}$

The " " radial "  $\frac{\Delta r}{r_s} = \frac{\alpha}{\beta_s^2} \frac{\Delta E}{E_s}$

The functions  $F, \frac{G}{F}, \sqrt{G}$  are tabulated below, for  $\varphi_0 = 30^\circ$

J is invariant under adiabatic change of its parameters. Thus given certain conditions, say  $\varphi_m + E_s$  at injection, we can use the table to compute the change in  $\varphi_m$  and hence  $\frac{\Delta E}{E_s}$  as  $E_s$  rises.

As soon as  $\varphi_m - \varphi_0$  has damped down to less than  $30^\circ$  the formulae in Courant and Snyder's paper are valid.

NOTE.  $\frac{\Delta E}{E_s}$  gives the maximum tolerable spread in injection energy in terms of  $E_{inj}, \frac{eV}{h}, \varphi_0, \varphi_m$  for sudden application of the r.f. It is a more stringent requirement in most cases than that requiring the particles to stay within the doughnut. Symbols have the same meaning as in Courant's and Snyder's paper.

$$\varphi_0 = 30^\circ$$

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$\varphi_m - \varphi_0$	$\varphi_m - \varphi_0$	$F(\varphi_m, \varphi_0)$	$G(\varphi_m, \varphi_0)$	$\sqrt{G}(\varphi_m, \varphi_0)$	$\frac{G}{F}$
$0^\circ$	0 rad	.0000	.0000	.0000	.225
<del>10</del> $10^\circ$	.1745	.065	.0148	.1217	.225
$20^\circ$	.3490	.251	.0561	.2369	.224
$30^\circ$	.5235	.539	.1203	.3468	.224
$40^\circ$	.6980	.909	.2020	.4494	.224
$45^\circ$		1.090	.2477	.4977	.223
$50^\circ$	.8725	1.339	.2956	.5437	.221
$60^\circ$	1.047	1.807	.3954	.6288	.219
$75^\circ$	1.309	2.532	.5431	.7370	.214
$90^\circ$	1.571	3.258	.6704	.8188	.206
$105^\circ$	1.833	3.793	.7584	.8609	.200
$120^\circ$	2.094	4.049	.7908	.8893	.176