



Invisible Z estimation in the monojet final state

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CMS, CERN Summer Student Program 2014





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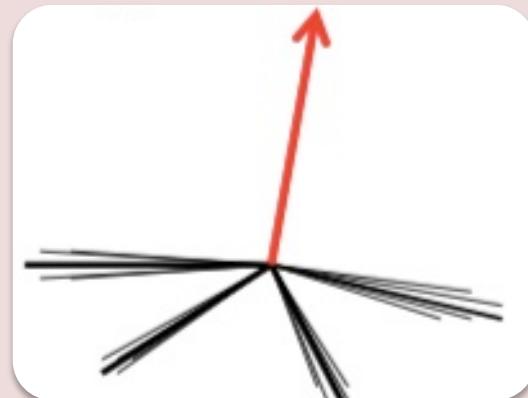
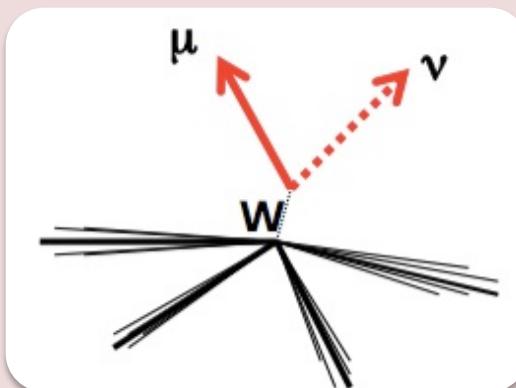
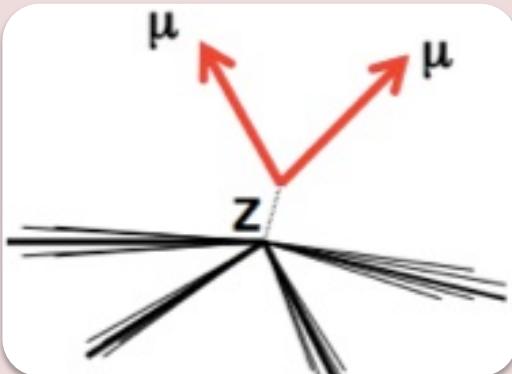
Objective

To estimate number of irreducible invisible Z background contribution in the monojet final state and obtain the exclusion limit in the study

Motivation

- Large missing transverse energy (MET) + jets has a great potential for discovery of new physics
- Invisible Z is one of dominant backgrounds remaining after applying monojet selection cuts.
- $Z(\mu\mu)$ sample is used to predict $Z(vv)$
 - High uncertainty because of small statistic

Candidates for invisible Z estimation



$Z(\mu\mu) + \text{jets}$

- Clean
- Easy to select
- Small statistic

$W(\mu\nu) + \text{jets}$

- Not so clean
- Larger statistic

$\gamma + \text{jets}$

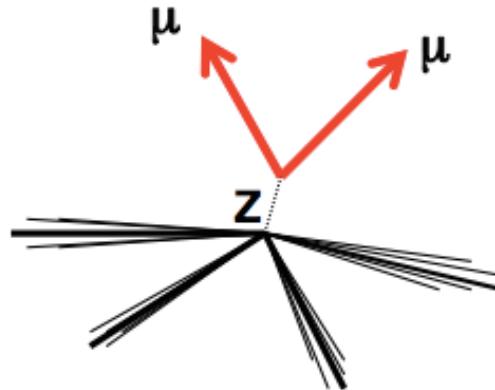
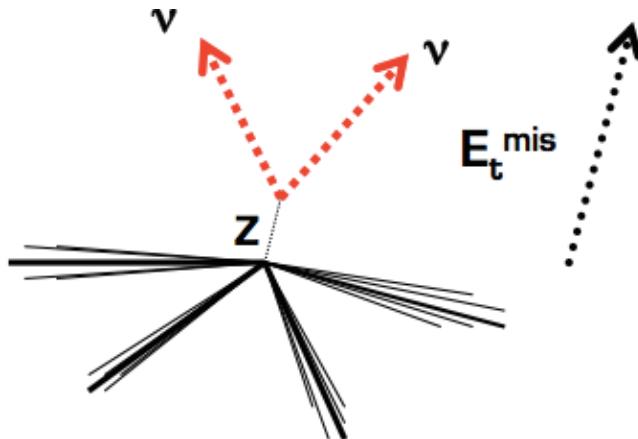
- Clean for high E_T
- Large statistic
- (Not so clean for < 100 GeV)

Outline

- Theoretical Background
- Event Selection
- Invisible Z prediction
- Result
- Summary
- Future Plans

Z+jets background estimation

$Z(\nu\nu)$ prediction using



The process of $Z(\nu\nu)$ is mimicked by removing muons in the $Z(\mu\mu)$ event to missing transverse energy.

Z+jets background estimation

Use $Z(\mu\mu)$ as control sample to estimate $Z(\nu\nu)$

$$N(Z \rightarrow \nu\nu) = \frac{(N_{Obs} - N_{Bgd})}{A \cdot \varepsilon} R \left(\frac{Z \rightarrow \nu\nu}{Z \rightarrow \mu\mu} \right)$$

N_{obs} = Number of observed $Z(\mu\mu)$ event

N_{bgd} = Number of estimated $Z(\mu\mu)$ background

A = Acceptance

ε = Selection efficiency

Acceptance and Efficiency

All generated events = X



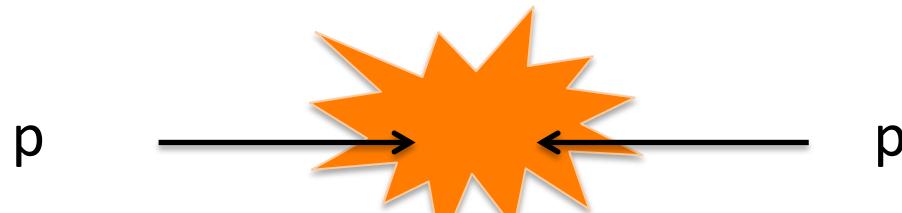
Y events



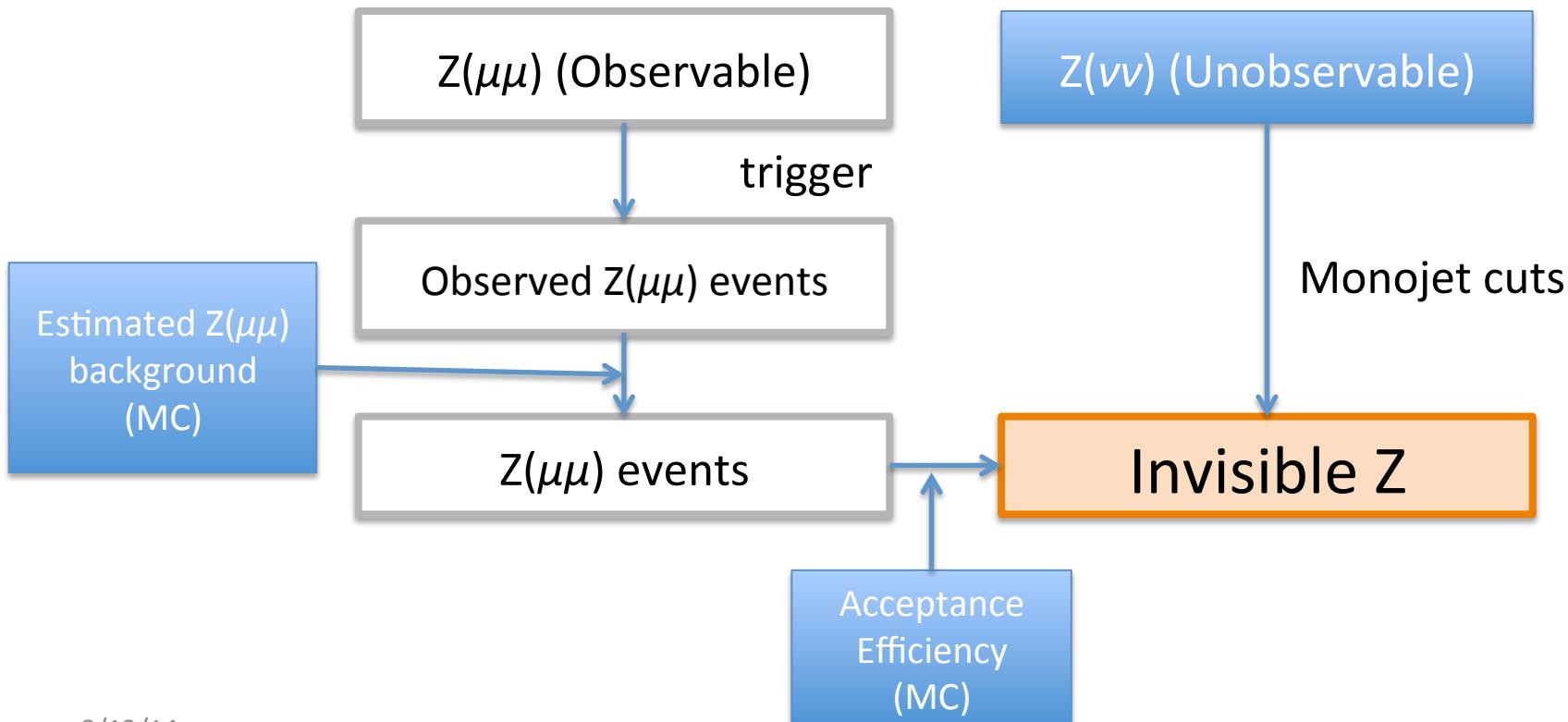
Z events

$$\text{Acceptance} = Y/X$$
$$\text{Efficiency} = Z/Y$$

Analysis workflow



p-p collision in the LHC



Event Selection

For $Z(\mu\mu)$

- Monojet selection
- $p_T > 20 \text{ GeV}$
- $|\eta| < 2.4$
- Transverse invariant mass between 60 – 120 GeV

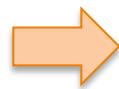
Require 2 muons and at least one of these has to be well identified and isolated.

$$M_T = \sqrt{2p_T^\mu E_T^{\text{miss}} (1 - \cos \Delta\phi)}$$

p_T^μ - transverse momentum of the muon
 $\Delta\phi$ - angle between the muon p_T and the E_T^{miss} vectors.

$Z(\nu\nu)$ Prediction

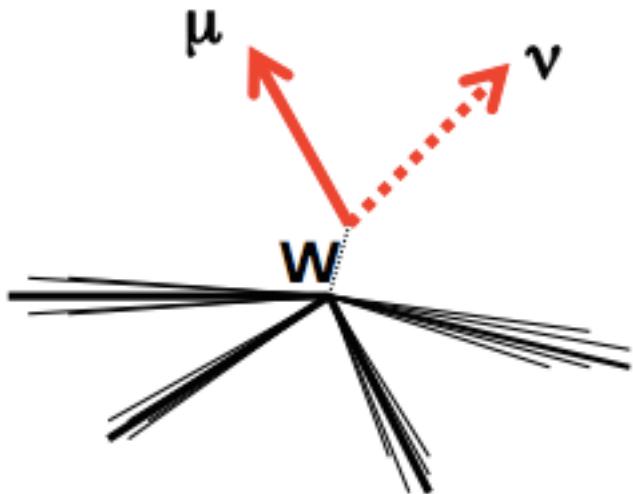
MET	N_{obs}	N_{bkgd}	Acceptance	Efficiency	$Z(\nu\nu)$
$> 250 \text{ GeV}$	3695	288	0.89 ± 0.02	0.74 ± 0.02	30700 ± 1898
$> 300 \text{ GeV}$	1538	128	0.91 ± 0.02	0.75 ± 0.02	12199 ± 831
$> 350 \text{ GeV}$	685	70	0.93 ± 0.02	0.76 ± 0.02	5174 ± 423
$> 400 \text{ GeV}$	348	31	0.94 ± 0.02	0.76 ± 0.02	2630 ± 238
$> 450 \text{ GeV}$	183	15	0.94 ± 0.02	0.75 ± 0.03	1420 ± 152
$> 500 \text{ GeV}$	96	7.7	0.94 ± 0.02	0.76 ± 0.03	732 ± 98
$> 550 \text{ GeV}$	47	5.1	0.95 ± 0.02	0.75 ± 0.04	349 ± 66



High uncertainty at high energy

$Z(\nu\nu)$ Prediction

Using W+jets



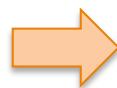
$$N(Z \rightarrow \nu\nu) = \frac{(N_{\text{Obs}} - N_{\text{Bgd}})}{A \cdot \varepsilon} R \left(\frac{Z \rightarrow \nu\nu}{W \rightarrow \mu\nu} \right)$$

Event selection

- Monojet selection
- $p_T > 20 \text{ GeV}$
- $|\eta| < 2.4$
- Well identified and isolated transverse invariant mass between 50 – 100 GeV

$Z(\nu\nu)$ Prediction

MET	N_{obs}	N_{bkgd}	Acceptance	Efficiency	$Z(\nu\nu)$
> 250 GeV	17191	1816	0.86 ± 0.02	0.39 ± 0.02	21874 ± 1090
> 300 GeV	6955	712	0.89 ± 0.02	0.39 ± 0.02	8719 ± 444
> 350 GeV	3104	319	0.90 ± 0.02	0.38 ± 0.02	3860 ± 210
> 400 GeV	1484	148	0.91 ± 0.02	0.38 ± 0.02	1851 ± 109
> 450 GeV	780	76	0.92 ± 0.02	0.37 ± 0.03	990 ± 66
> 500 GeV	402	44	0.93 ± 0.02	0.37 ± 0.03	521 ± 43
> 550 GeV	228	31	0.94 ± 0.02	0.38 ± 0.04	287 ± 31



Lower uncertainty in $Z(\nu\nu)$ prediction

Exclusion limit

MET (GeV)	> 250	> 300	> 350	> 400	> 450	> 500	> 550
Z($\nu\nu$) + jets	26287 \pm 1094	10459 \pm 471	4517 \pm 236	2241 \pm 131	1205 \pm 83	626 \pm 53	318 \pm 36
W + jets	17177 \pm 1030	5908 \pm 365	2333 \pm 154	1012 \pm 72	495 \pm 39	247 \pm 24	119 \pm 14
$t\bar{t}$	446 \pm 223	167 \pm 84	69 \pm 35	31 \pm 16	15 \pm 7.7	6.6 \pm 2.3	2.8 \pm 1.4
Z(ll) + jets	134 \pm 67	43 \pm 21	17 \pm 8.7	8.4 \pm 4.2	4.7 \pm 2.3	2.3 \pm 1.2	1.0 \pm 0.5
Single t	155 \pm 77	53 \pm 26	18 \pm 9.1	6.1 \pm 3.1	0.9 \pm 0.4	0.0 \pm 0.0	0.0 \pm 0.0
QCD Multijets	462 \pm 231	101 \pm 51	32 \pm 16	5.5 \pm 2.7	2.3 \pm 1.2	1.2 \pm 0.6	0.6 \pm 0.3
Diboson	3421 \pm 1710	1396 \pm 698	641 \pm 320	305 \pm 153	165 \pm 83	88 \pm 44	55 \pm 28
Total SM	48081 \pm 2302	12600 \pm 924	7628 \pm 429	3609 \pm 214	1889 \pm 124	971 \pm 73	497 \pm 48
Data	52157	19783	8324	3825	1828	934	519
Exp limit	4225	1707	787	368	199	122	90
Obs limit	7627	3103	1360	564	170	102	107

Summary

- One of useful methods to estimate the $Z(\nu\nu)$ is the use of $W(\mu\nu)+\text{jets}$ data.
- Taking the advantage of the similar kinematics of $Z(\nu\nu)$ and $W(\mu\nu)$
- Using $W(\mu\nu)$ control samples, the uncertainty was reduced.

Future Plans

- Expand range of event selection in order to reduce statistic uncertainty
- Invisible Z Estimation by using γ +jets.

References

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- CERN Summer Student Program 2014



Thank you for your kind attention.

Cr. CERN Summer Student Team



Monojet

- One jet plus nothing else
- Nothing is not well-identified in a real experiment.
 - The region that identification process is not reliable
- Monojet final state
 - parton + undetectable particles

Monojet Selection

- Exactly one quality-jet
- Without ...
 - second-quality jet above thresholds
 - Isolated electrons and muons above thresholds
 - No bad-jets above threshold