

# The LHC machine - present and future

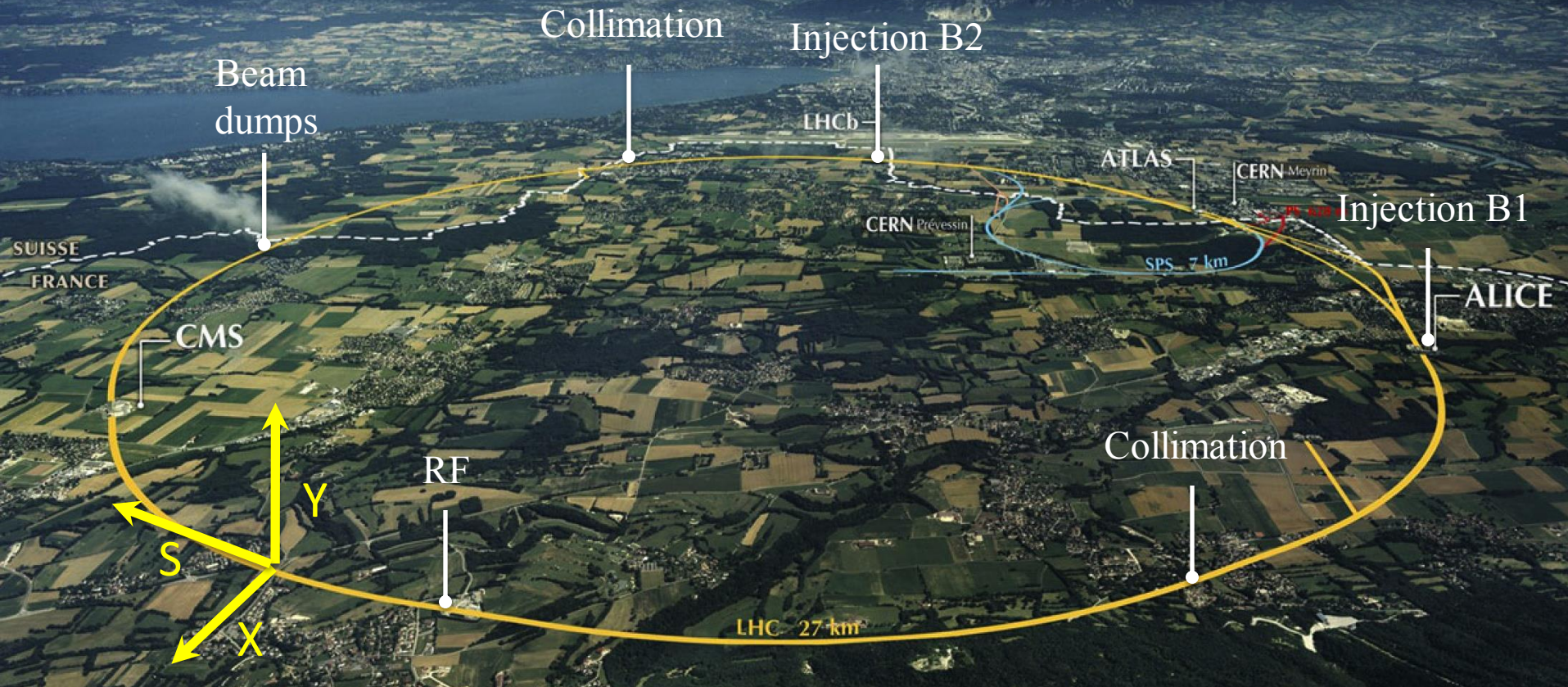


- Overview of the current machine, performance and limitations
- Upgrades towards ultimate luminosity
- Possibilities and challenges for higher energy

Mike Lamont

with acknowledgements to the people whose material I've used

# LHC: big, cold, high energy

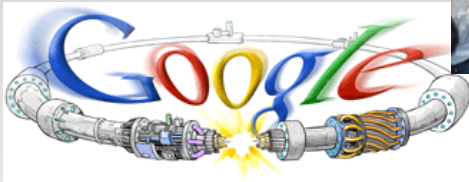
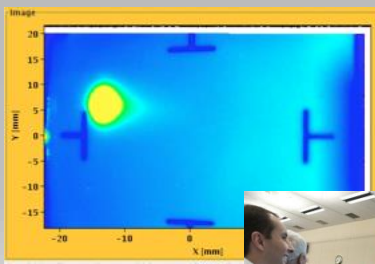


1720 Power converters  
> 9000 magnetic elements  
7568 Quench detection systems  
1088 Beam position monitors  
~4000 Beam loss monitors

150 tonnes Helium, ~90 tonnes at 1.9 K  
140 MJ stored beam energy in 2012  
450 MJ magnetic energy per sector at 4 TeV

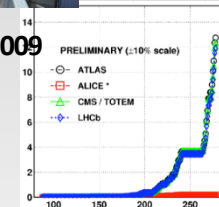


**August 2008**  
First injection test



**September 10, 2008**  
First beams around

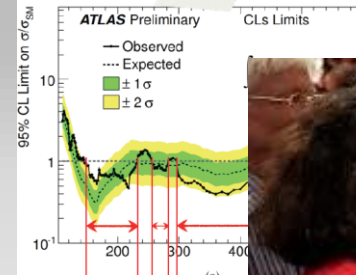
**November 29, 2009**  
Beam back



**October 14 2010**  
1e32  
248 bunches

**June 28 2011**  
1380 bunches

**1380**



**August 2011**  
2.3e33,  
1380 b



**4 July, 2012**

**6 June, 2012**  
6.8e33

2008

2009

2010

2011

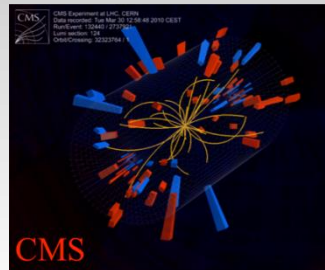
2012

**September 19, 2008**  
**Disaster**

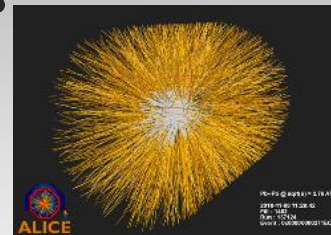
Accidental release of 600 MJ stored in one sector of LHC dipole magnets



**March 30, 2010**  
First collisions at 3.5 TeV



**November 2010**  
Ions



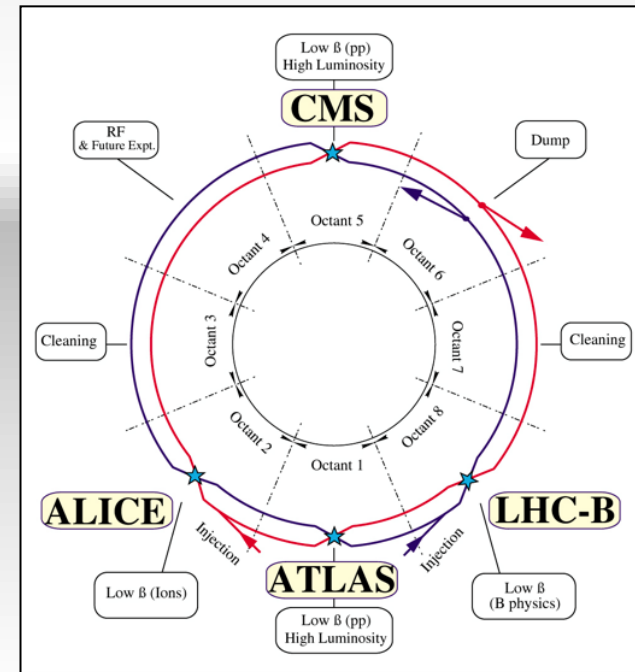
**18 June, 2012**  
6.6 fb<sup>-1</sup>  
to ATLAS & CMS

# LHC Timeline

- High energy
- High bunch intensity
- Many bunches
- Small beam size

# LHC – PRINCIPLES & REALITY

- LEP tunnel defines the bending radius
- Superconducting magnet technology
- 1.9 K cryogenics
- Bending radius & achievable field strength -> 7 TeV/c
- Two beam pipes - 2 in 1 magnet design
- Separated function - strong focusing
- Luminosity insertions



# Superconductivity

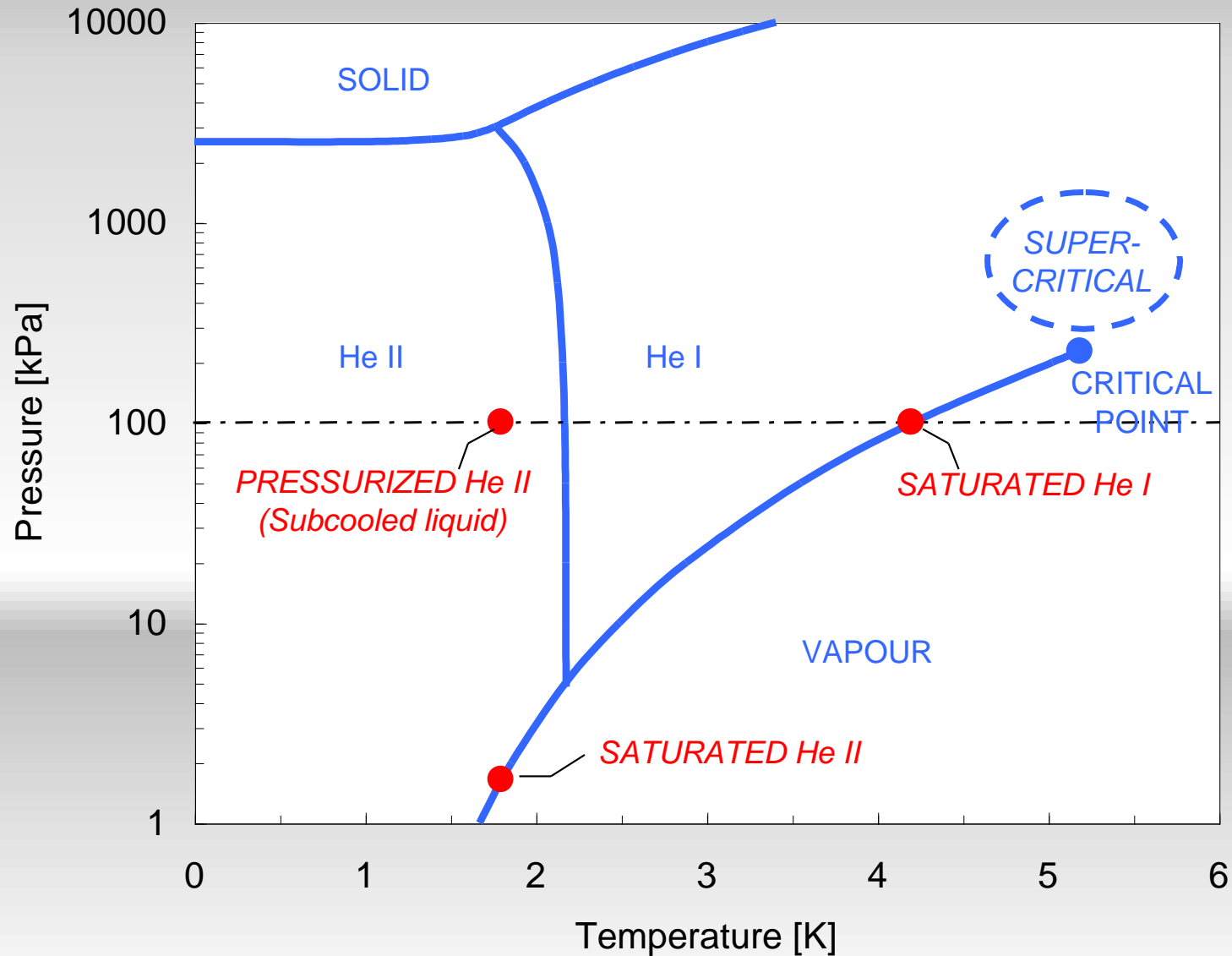
- To produce the high magnetic fields we need very high currents...
- Make use of the remarkable properties of He II
- Superfluid helium:
  - **Very high thermal conductivity** (3000 time high grade copper)
  - **Very low coefficient of viscosity**... can penetrate tiny cracks, deep inside the magnet coils to absorb any generated heat.
  - **Very high heat capacity**...stablizes small transient temperature fluctuations

How many Bose-Einstein condensates are there in the LHC?

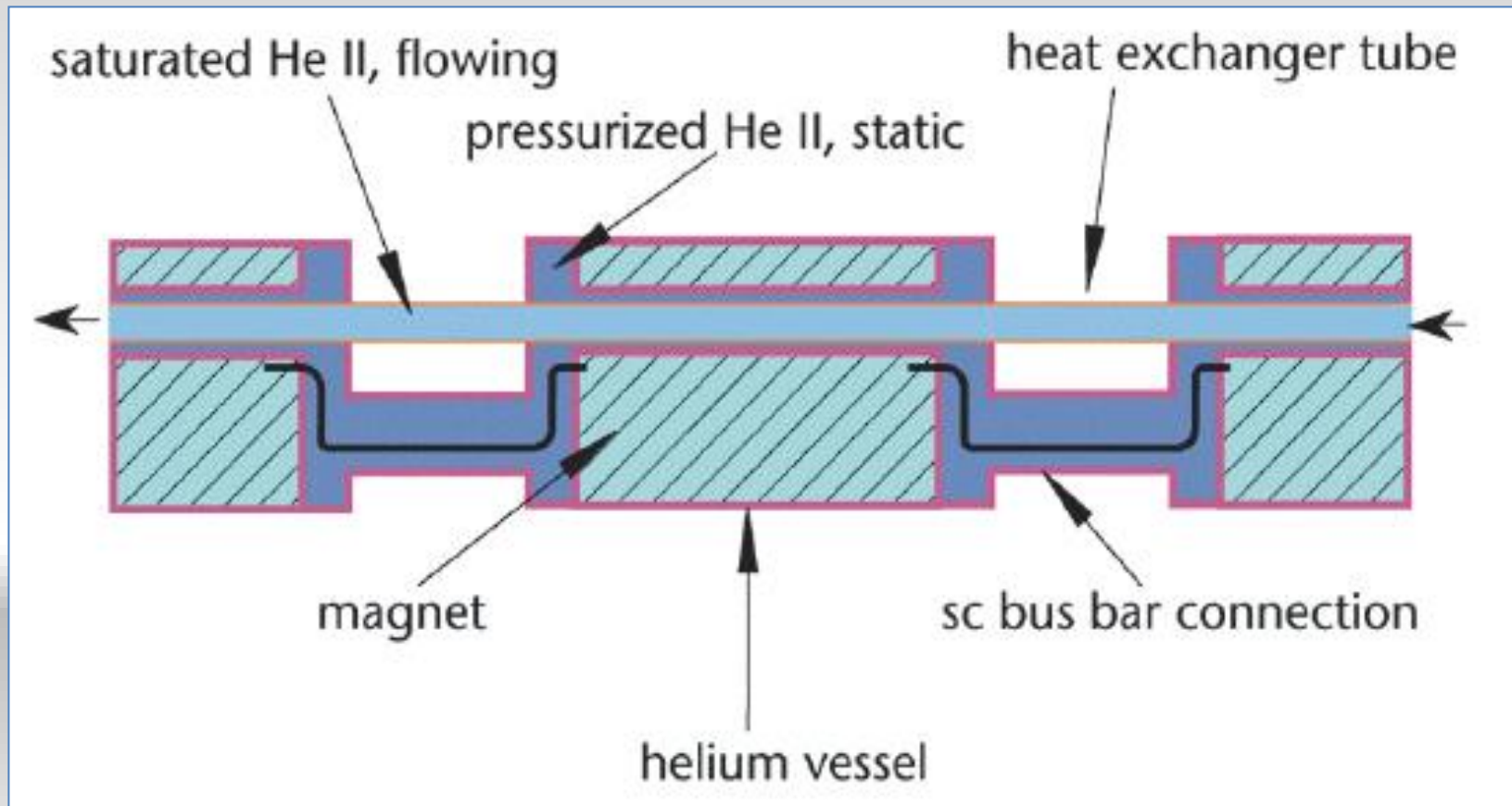
BBC FOUR



# Phase diagram of Helium



# Cooling magnets with superfluid helium

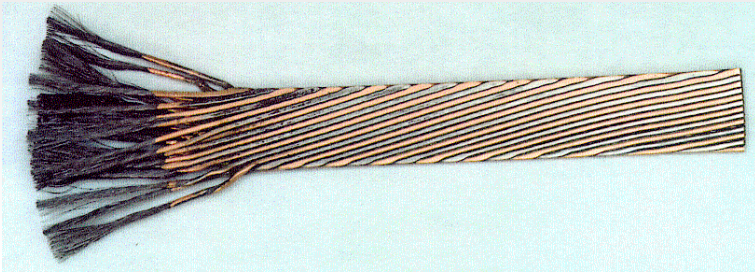




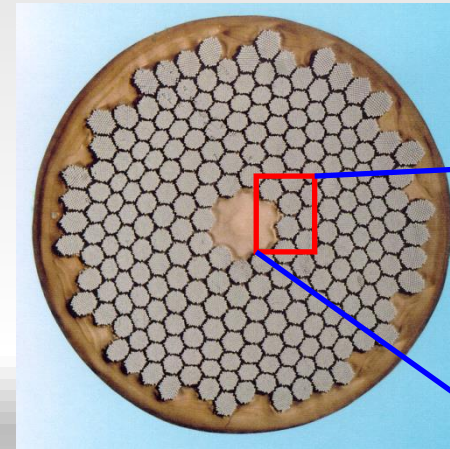
# The superconductor

## Niobium-titanium Rutherford cable

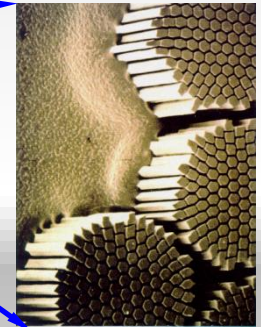
Cable



Strand

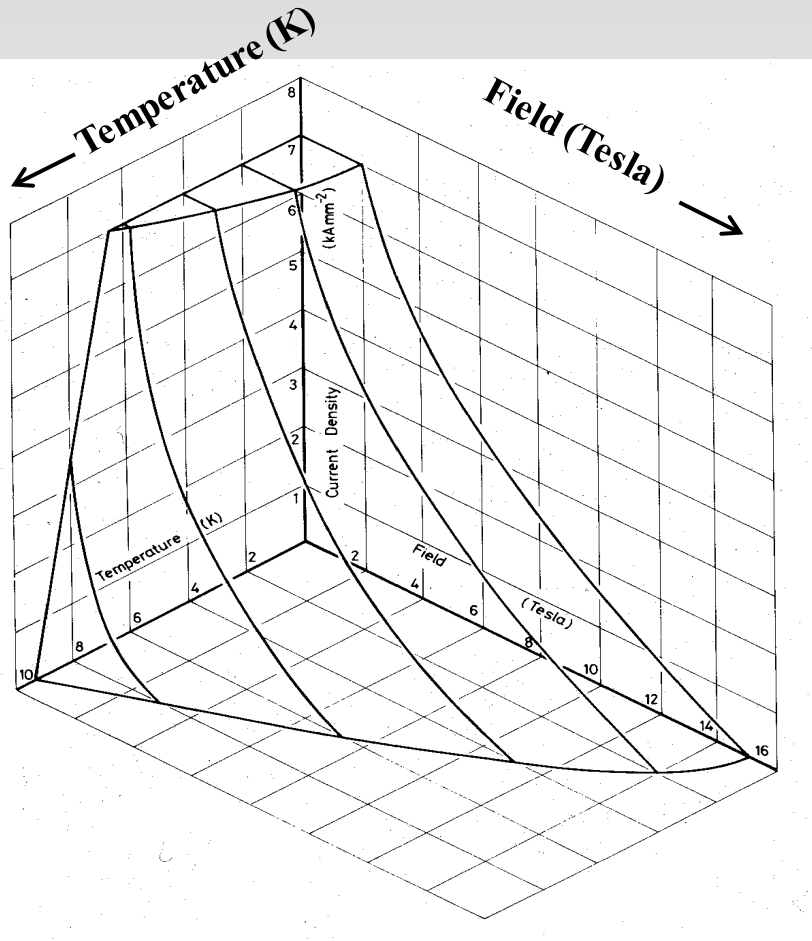


Filament



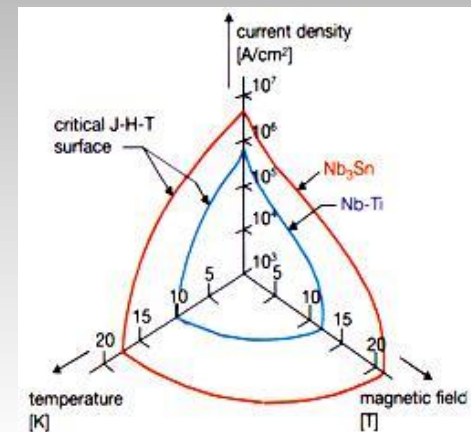
Used 1200 tonnes/7600 km of cable

# Critical surface of niobium-titanium

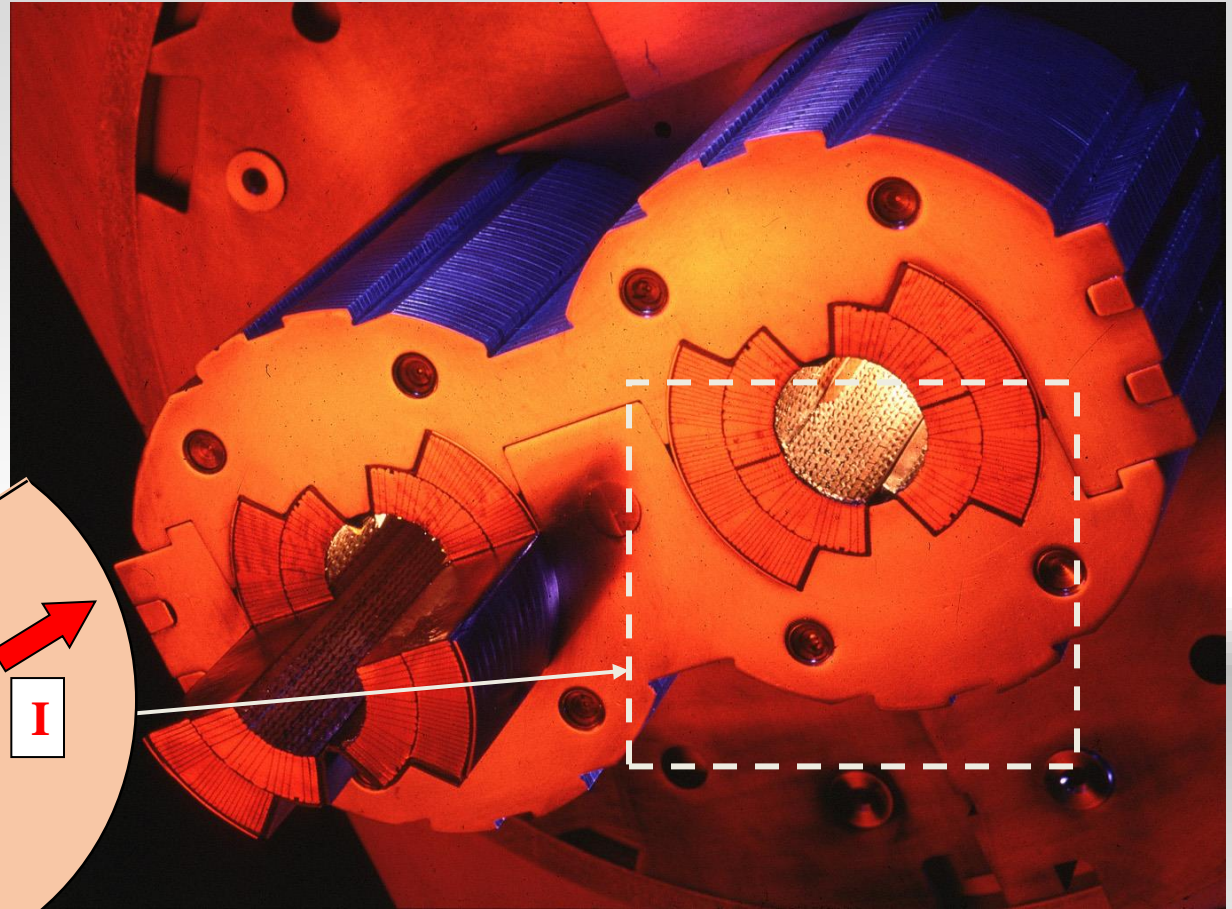
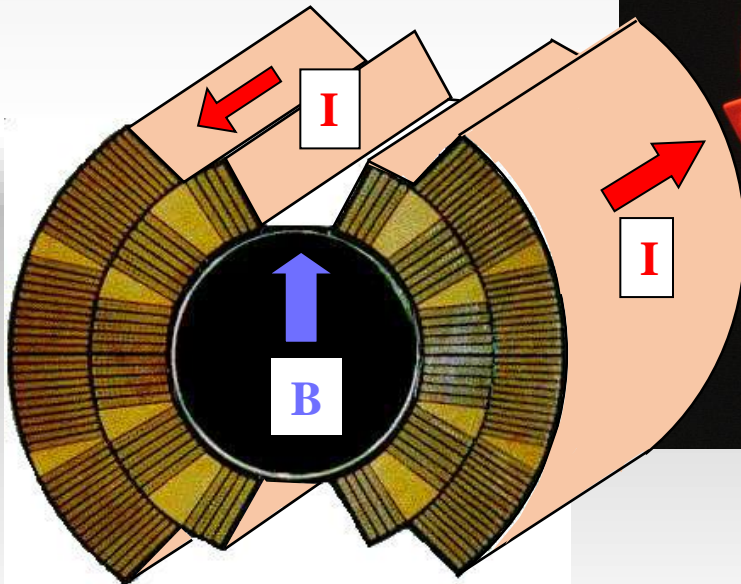
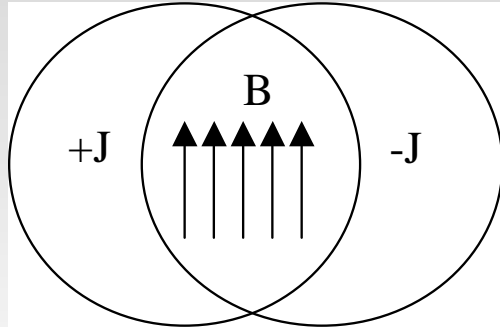


↑  
Current density (kA.mm<sup>-2</sup>)

- Niobium-titanium **NbTi** is the standard 'work horse' of the superconducting magnet business
- Picture shows the **critical surface**, which is the boundary between superconductivity and normal resistivity
- Superconductivity prevails everywhere below the surface, resistance everywhere above it

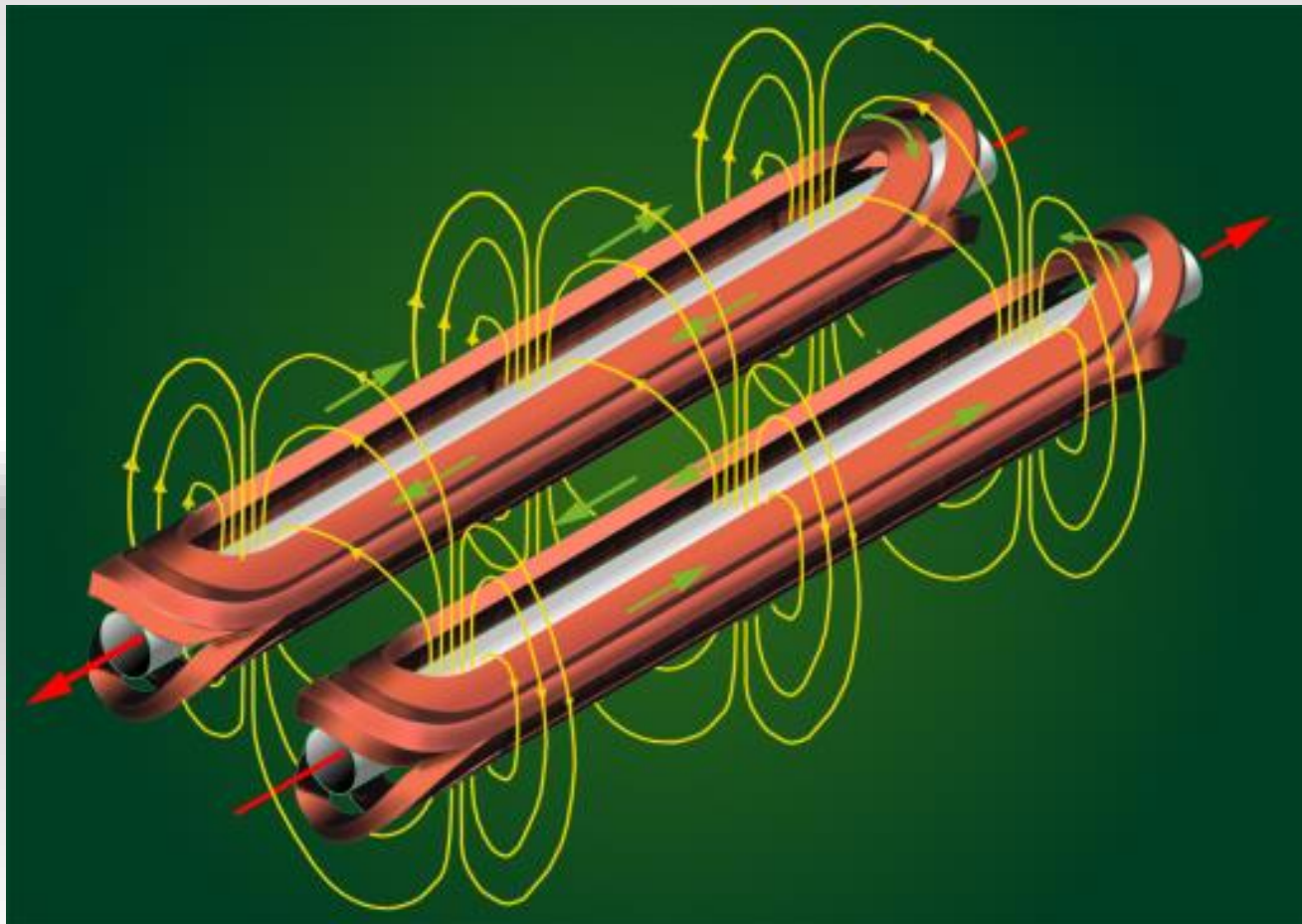


# Main components - dipole magnets





- Number of dipoles 1232
- Dipole field at 450 GeV 0.535 T
- Dipole field at 7 TeV 8.33 T
- Bending radius 2803.95 m
- Main Dipole Length 14.3 m



Horizontal force component per quadrant (nominal field) 1.7 MN/m

Force tends to “open” the magnet, hence the Austenitic steel collars



June 1994  
first full scale prototype dipole

June 2007 First sector cold

ECFA-CERN workshop



April 2008  
Last dipole down



1994 project approved by council (1-in-2)

SSC cancelled

Main contracts signed

198 90 91 92 93 94 95 96 97 98 99 00 01 02 03 04 05 06 07 08 09 10

First set of twin 1 m prototypes  
Over 9 T



2002 String 2



November 2006  
1232 delivered



September 19, 2008

# Energy stored in the magnets: quench

If not fast and safe ...

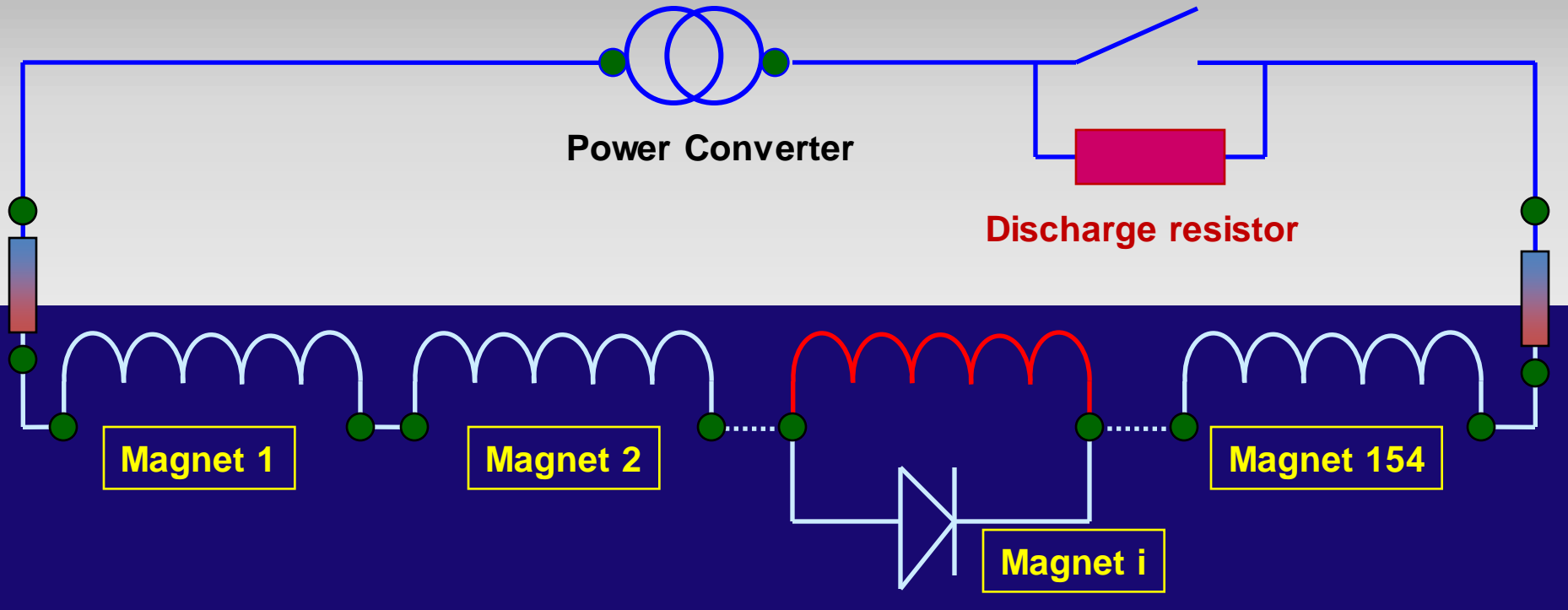
Quench in a magnet



During magnet test campaign, the 7 MJ stored in one magnet were released into one spot of the coil (inter-turn short)

P. Pognat

# Quench - discharge of the energy

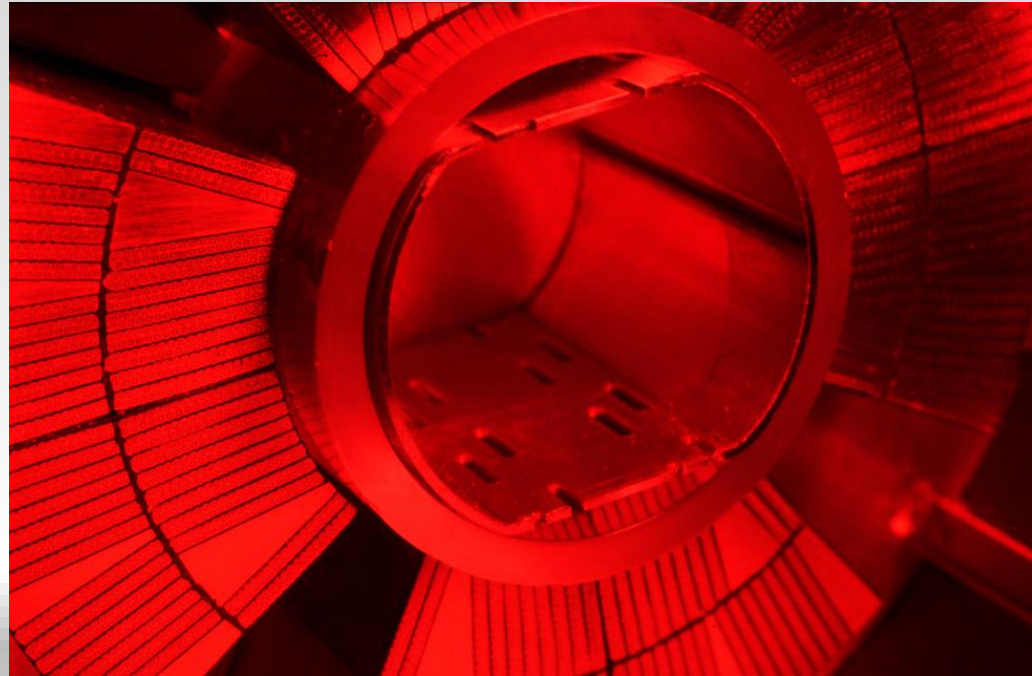
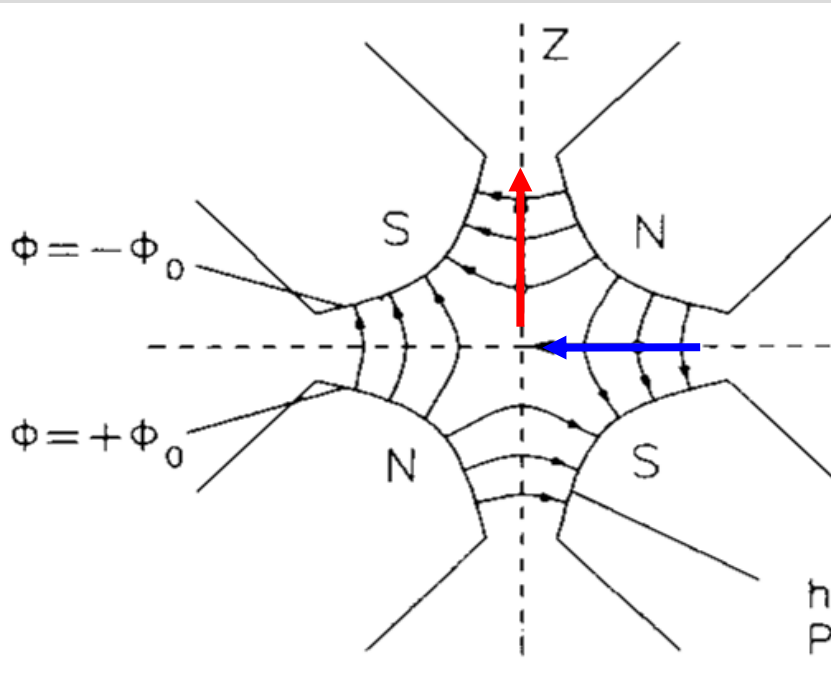


## Protection of the magnet after a quench:

- The quench is detected by measuring the voltage increase over coil.
- The energy is distributed in the magnet by force-quenching using quench heaters.
- The current in the quenched magnet decays in  $< 200$  ms.
- The current flows through the bypass diode (triggered by the voltage increase over the magnet).
- The current of all other magnets is discharged into the dump resistors.



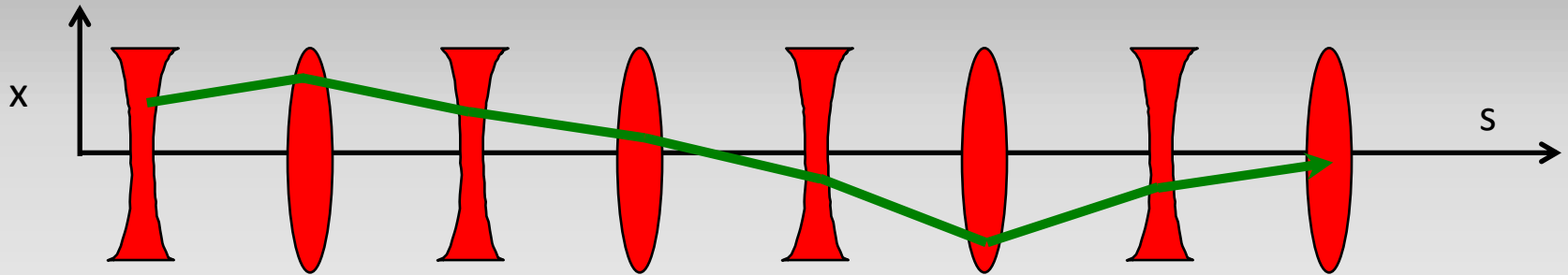
# OPTICS



- Recall: a quadrupole magnet will focus in plane and de-focus in the other.
- Convention: a “focusing” quadrupole focuses in the horizontal plane



# Alternate gradient focusing



The general linear magnet lattice can be parameterized by a 'varying spring constant',  $K=K(s)$

$$\frac{d^2 x}{ds^2} + K(s) x = 0 \quad (\text{and similarly for the vertical plane } y)$$

$K(s)$  describes the distribution of focusing strength along the lattice.

This is Hill's equation. The solution:

Amplitude term

$$x(s) = A \sqrt{\beta_x(s)} \cos(\phi(s) + \phi_0)$$

Oscillatory term

# Betatron function

$$x(s) = A\sqrt{\beta_x(s)} \cos(\phi(s) + \phi_0)$$

- A and  $\phi_0$  are constants, which depend on the initial conditions.
- $\beta(s)$  = the amplitude modulation due to the changing focusing strength.
- $\phi(s)$  = the phase advance, which also depends on focusing strength.

Stick the assumed solution into Hill's equation and turn the handle...

$$\phi' = \frac{d\phi}{ds} = \frac{1}{\beta}$$

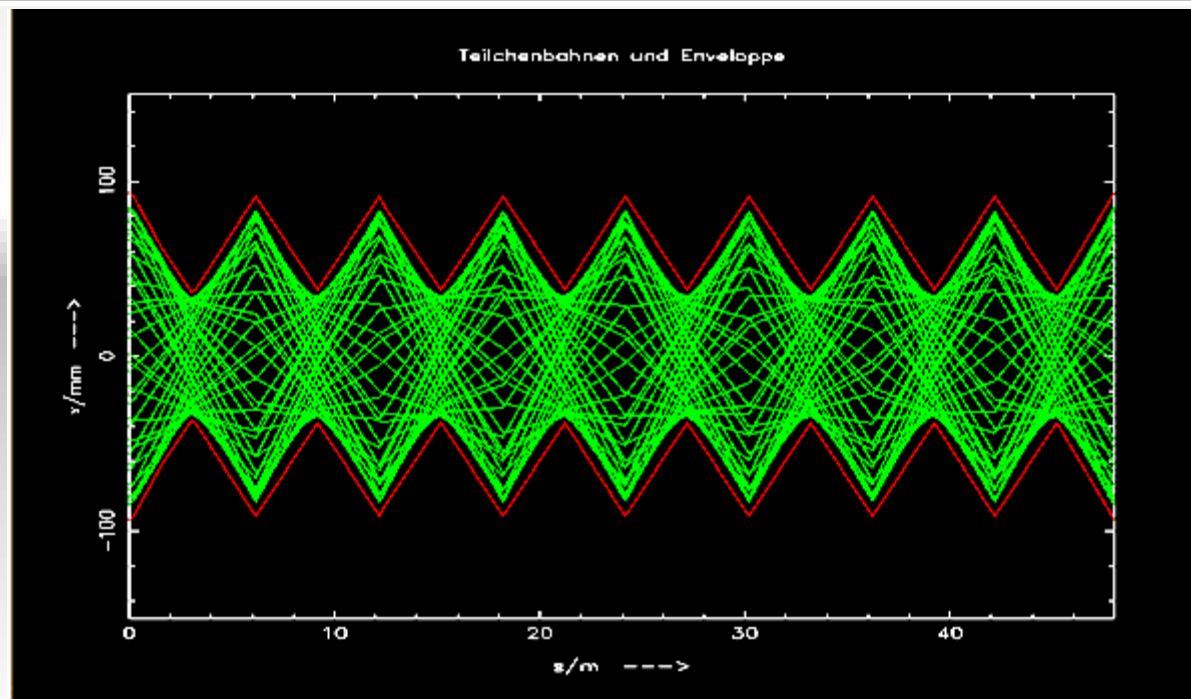
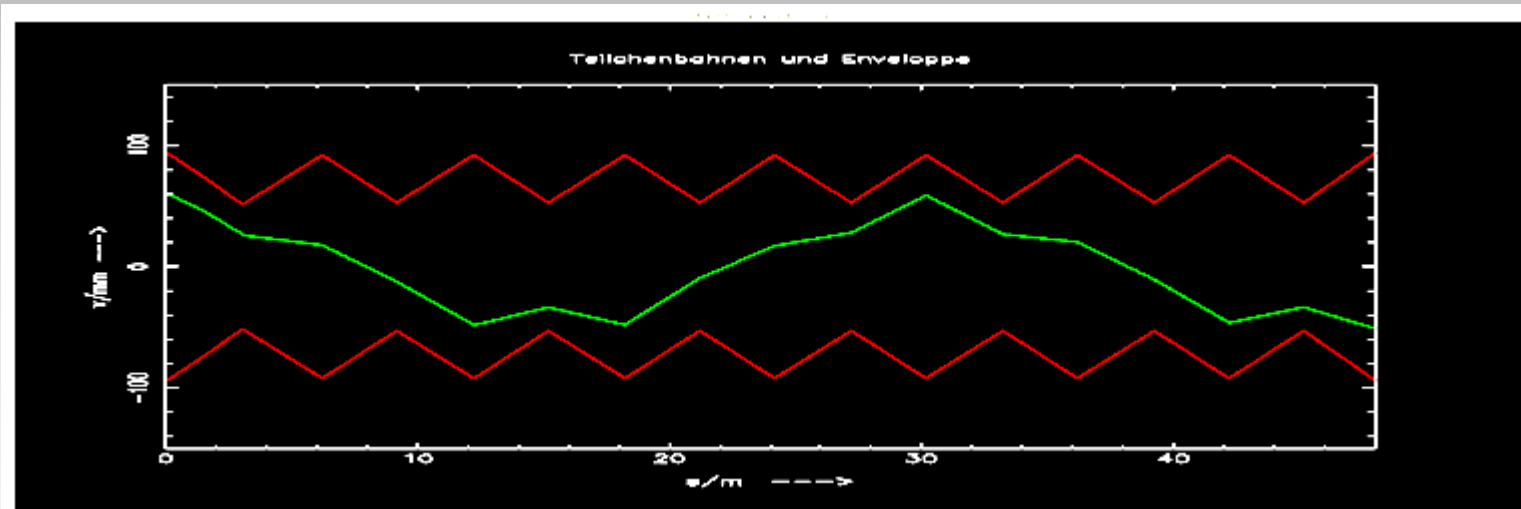
$$Df(s_1 \rightarrow s_2) = \int_{s_1}^{s_2} \frac{1}{b(s)}$$

- beta(s) maybe interpreted as the local wavelength of the oscillation (divided by 2 pi)

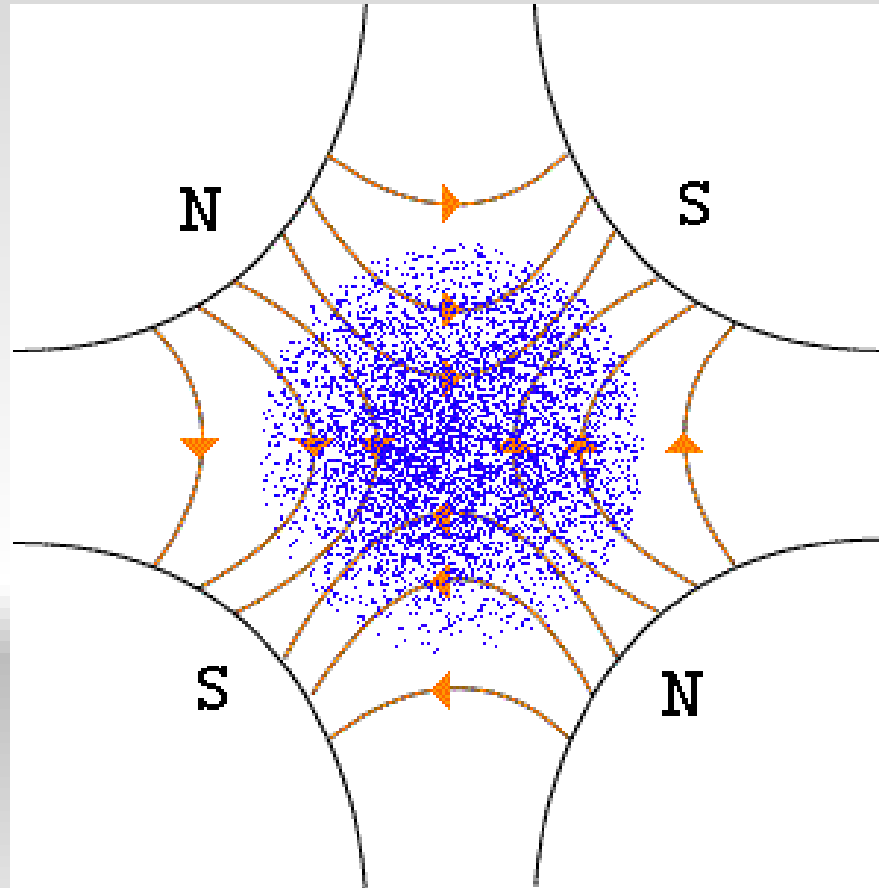
$$Q_x \approx \frac{1}{2\pi} \oint \frac{1}{\beta_x(s)}$$

Number of oscillations per turn is called the tune of the accelerator

# Betatron oscillations



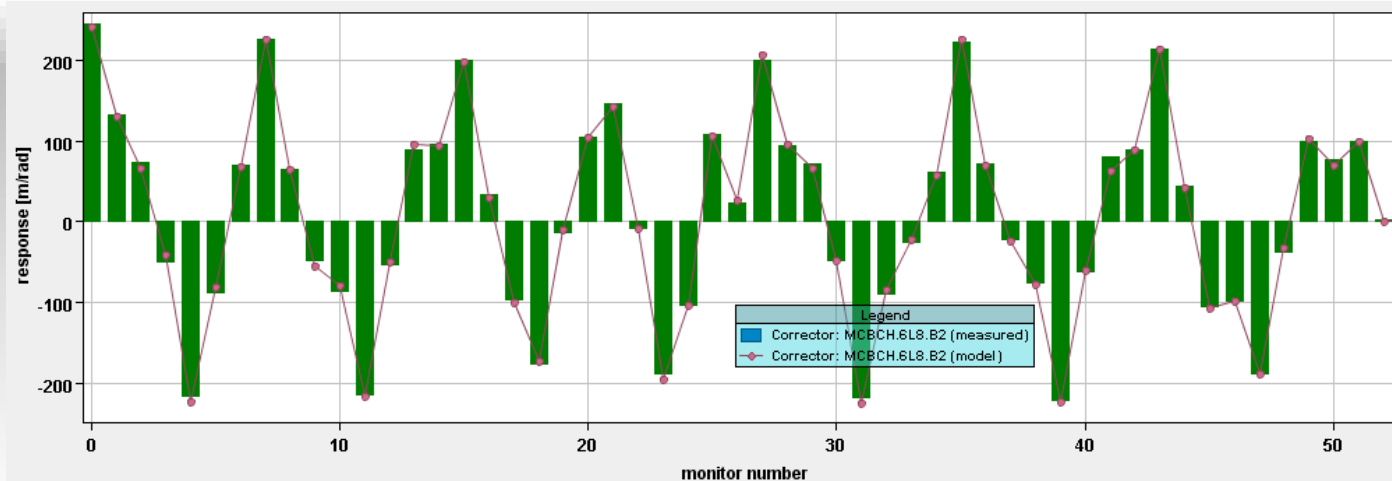
# Amplitude modulation



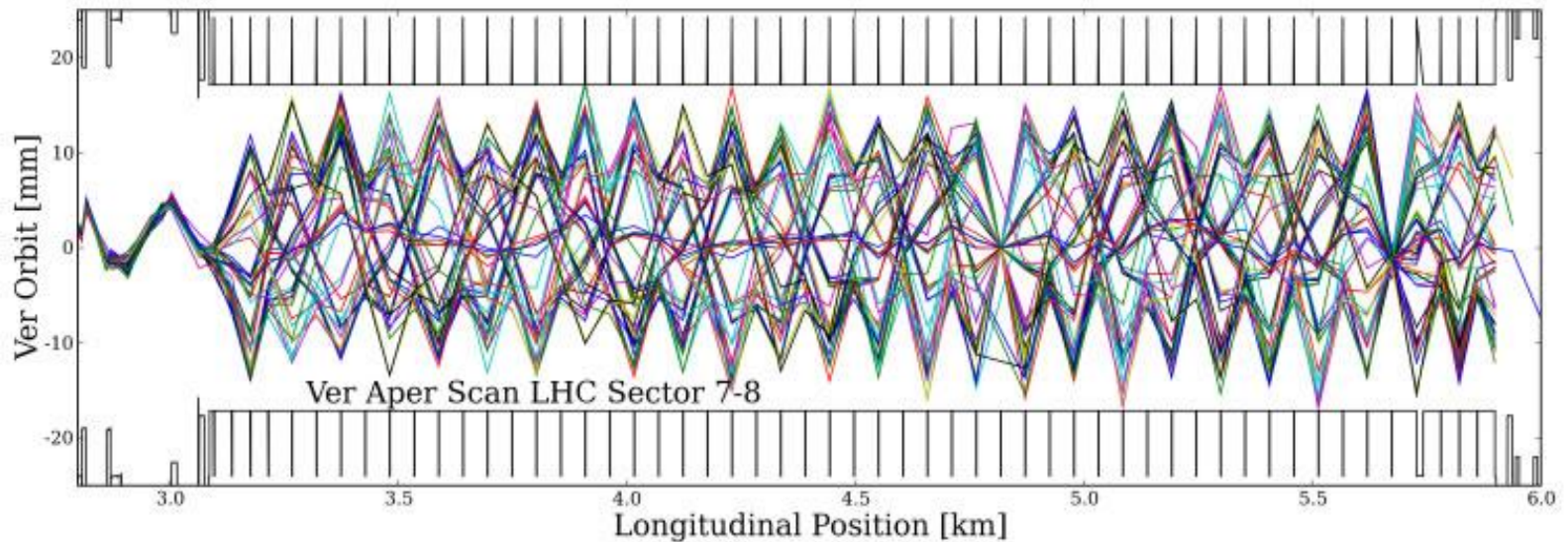
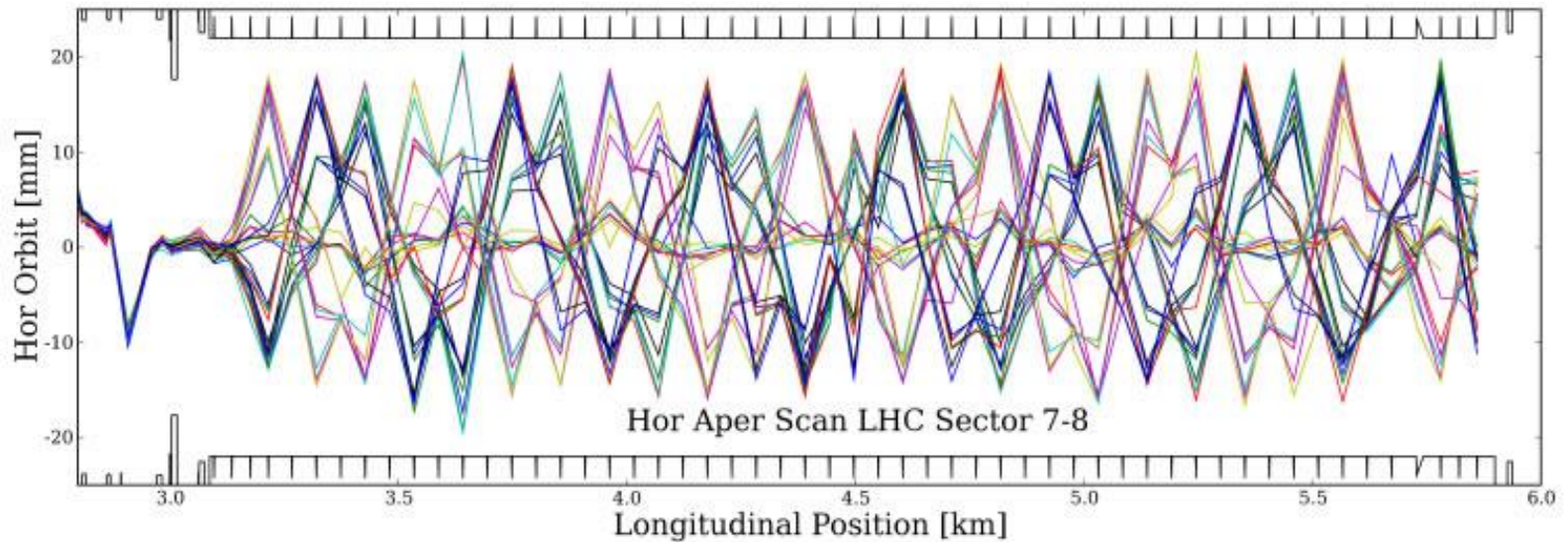


# Betatron oscillations

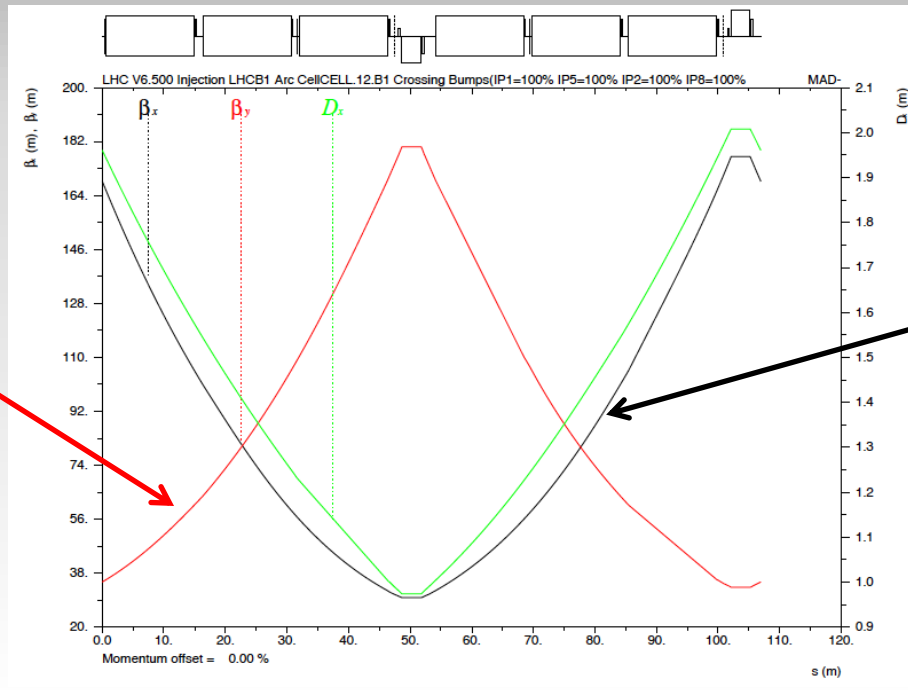
- Typical beam position monitor does not resolve individual particles.
- However, if the beam is deviated from the ideal trajectory it will oscillate due to the (de)-focusing fields of the accelerator.



# Injection test – 5 years ago



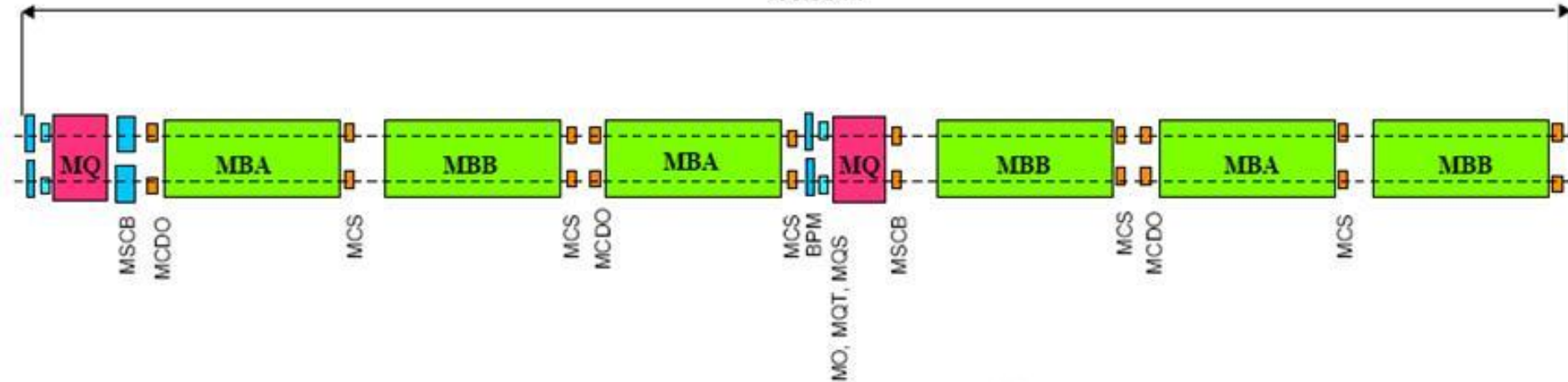
# LHC lattice in the arc



beta\_y

beta\_x

106.90 m



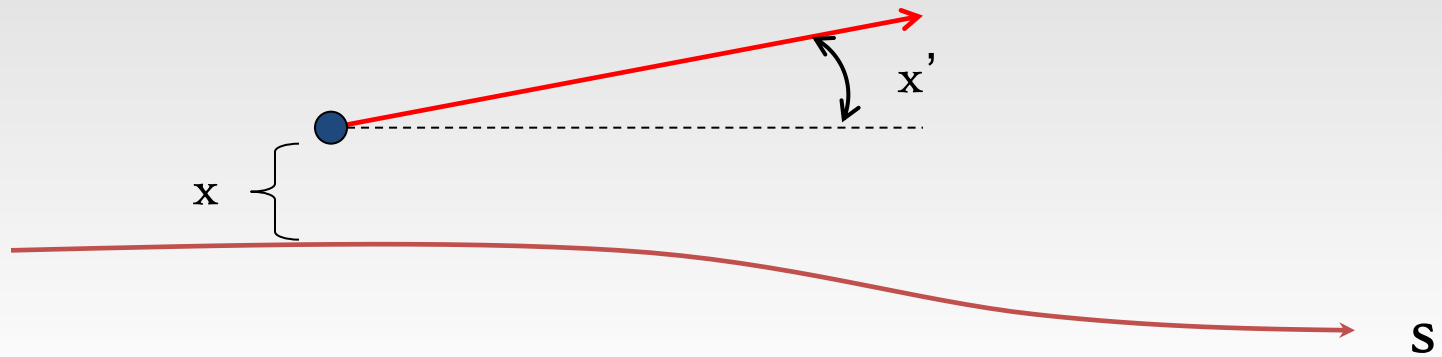


# For real



# TRANSVERSE ACCELERATOR COORDINATES

We are discussing the **paraxial** (small deflection angles => linear), **uncoupled** (x & y independent) **transverse motion** of **on-momentum** particles around the design trajectory/orbit for a magnet lattice:



(and similarly for y view)

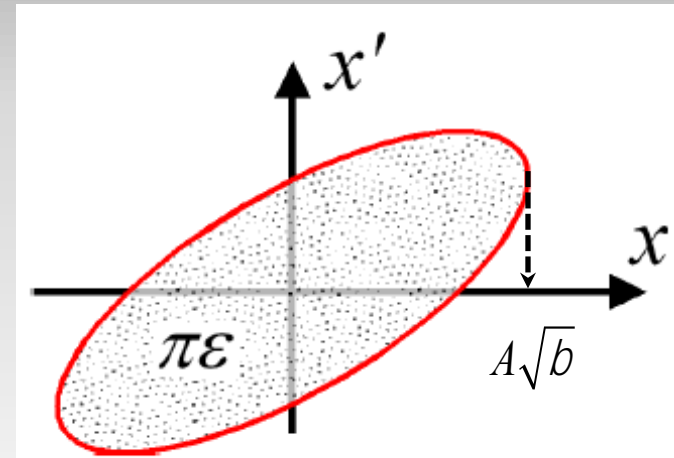
$$x' = \frac{dx}{ds}; \quad y' = \frac{dy}{ds}$$

$s$  = coordinate along design orbit

# Emittance

- **Linear** motion  $\rightarrow$  beam ellipse in 2-d-space
- Ellipse area in  $(x, x')$  plane:

$$A = \rho \times e$$

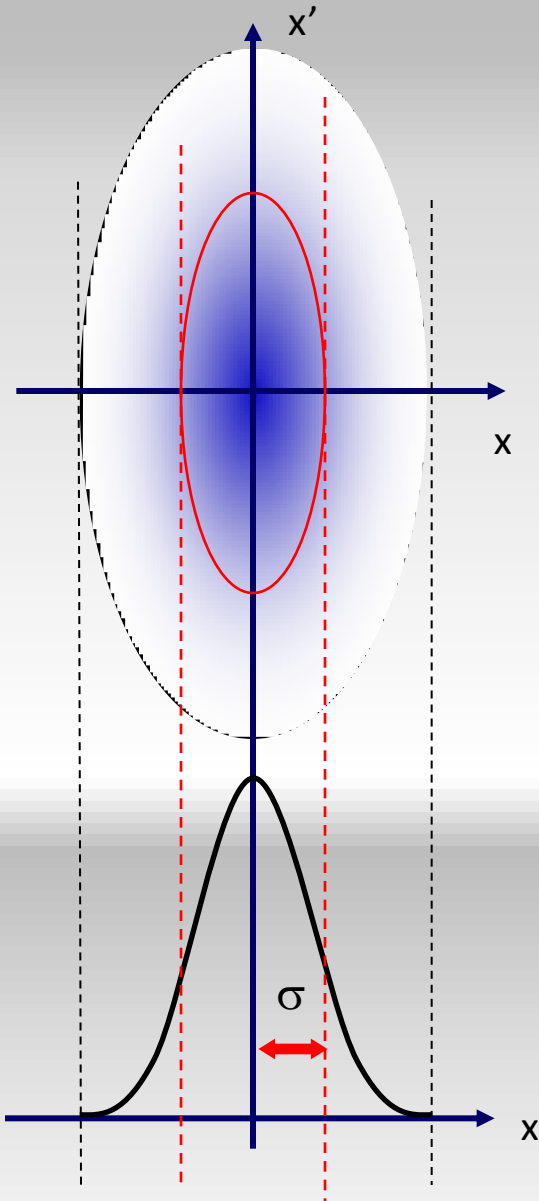


- Define emittance by a contour confining some fraction of particles depending on distribution (typically Gaussian)
- Units: mm.mrad (but you will see microns)
  - *Imagine sitting at one location and plotting  $x, x'$  of all particles*
  - *Or tracking one particle over many turns*



# Emittance

Area of phase space ellipse =  $\pi\epsilon$



$$\sigma = \sqrt{\epsilon\beta}$$

$$\epsilon = \frac{\sigma^2}{\beta}$$

Practical view:  $\sigma_x = \underbrace{\sqrt{\text{Betafunction } \beta}}_{\text{magnet structure}} \times \underbrace{\sqrt{\text{Emittance } \epsilon}}_{\text{particle ensemble}}$

Theoretical view:  $H = \frac{p_x^2}{2} + \frac{k(s)x^2}{2} \xrightarrow{\text{c.t.}} \tilde{H} = \frac{J}{\beta(s)}, \quad \epsilon = \langle J \rangle$

# Normalized emittance

Geometric emittance  $\{x, x'\}$  shrinks naturally as we go up in energy

$$Dp_s > 0, \quad Dp_x = 0, \quad e \sim \frac{1}{p_s}$$

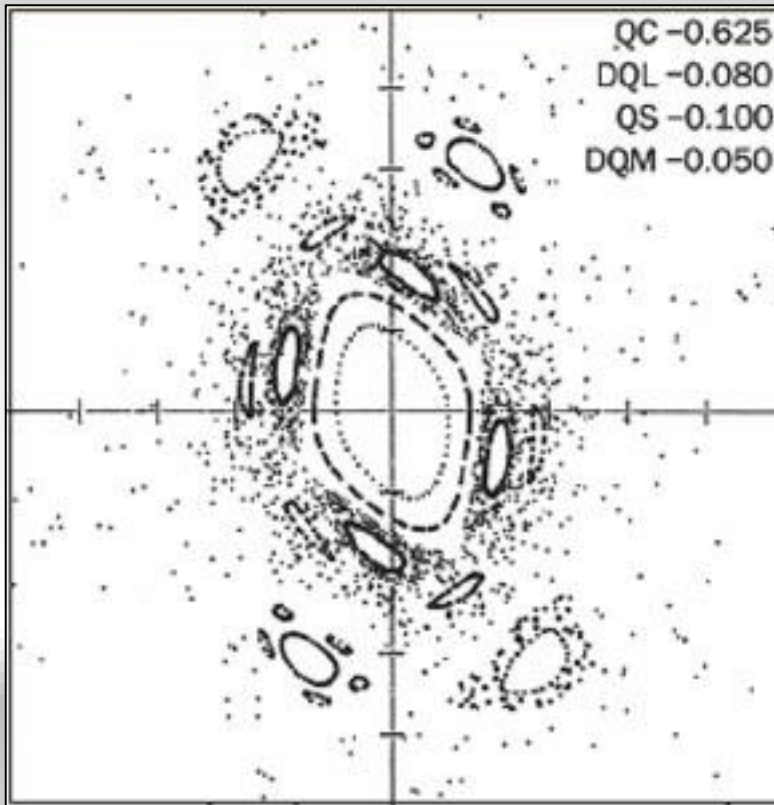
$$e_n = bge$$

**Important – it's energy independent and can be used across the accelerator complex and will show up later in a useful formulation of luminosity**

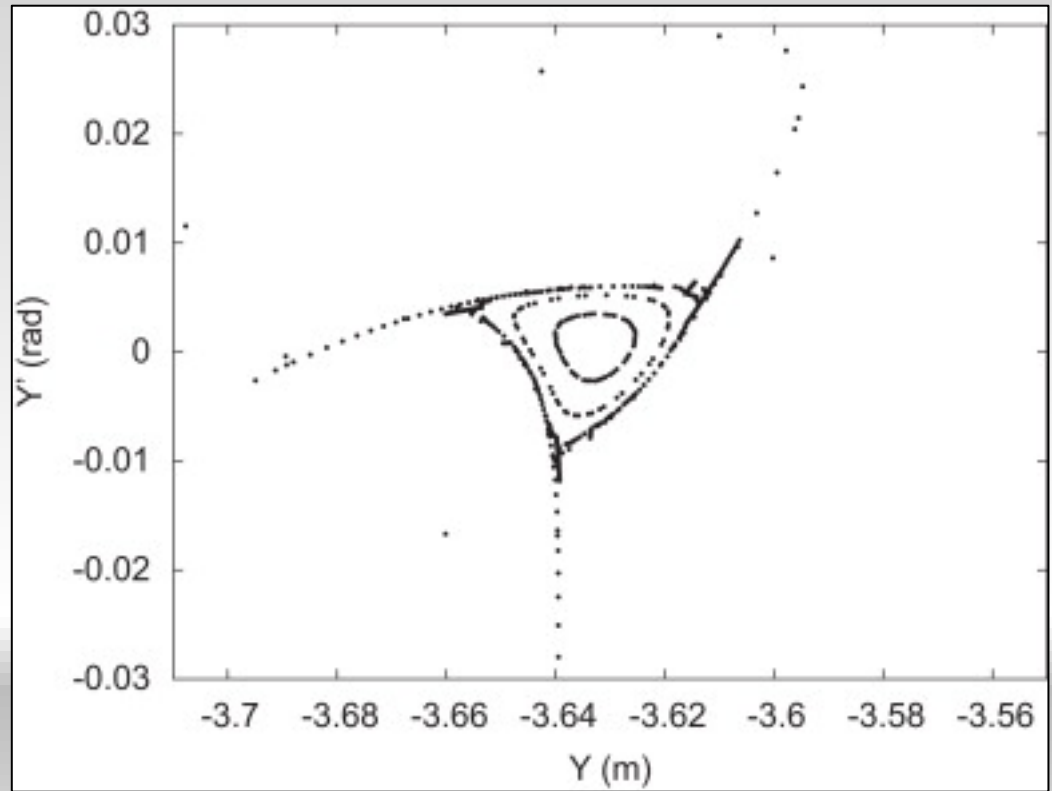
Normalized emittance	2.5 mm.mrad
Emittance 450 GeV	5.2 nm.rad
Emittance 7000 GeV	0.34 nm.rad

Beta [m]	Sigma 450 GeV micron	Sigma 7 TeV
180	967	246
30	395	100
4000	4566	1158
0.6	56	14

# Phase space



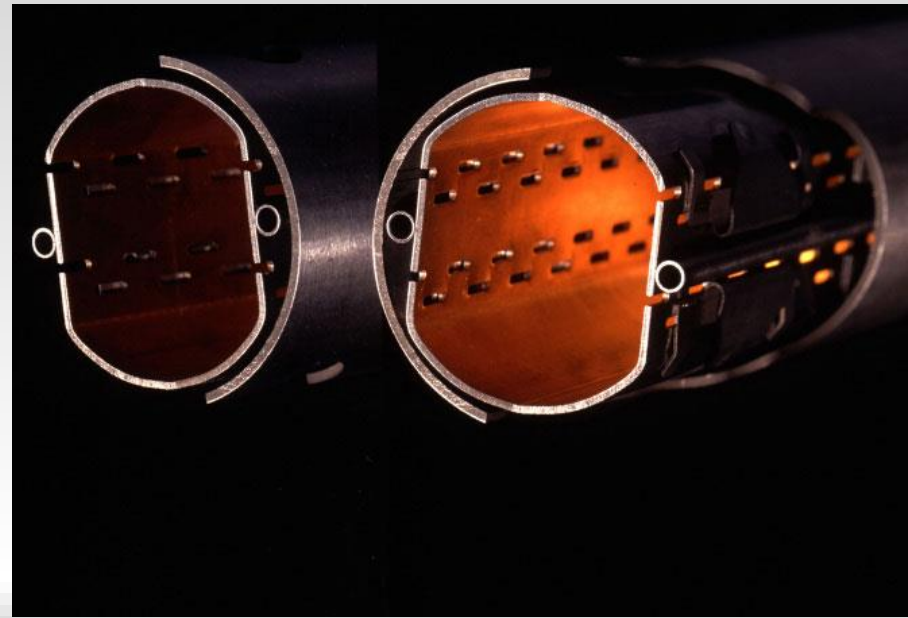
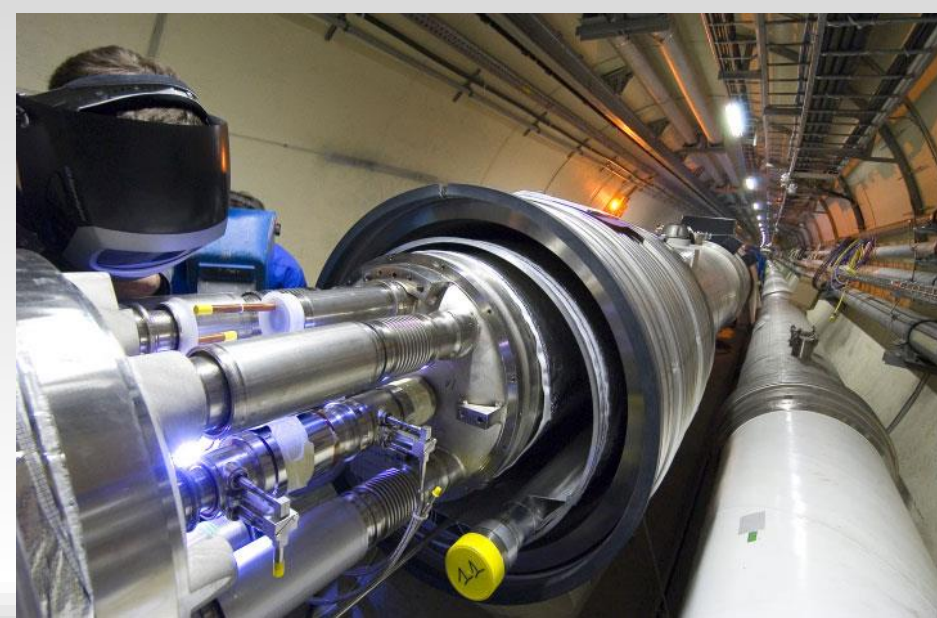
Simulation showing the chaotic effect of the beam-beam interaction



Phase space trajectories near the septum of a compact synchrotron during low extraction

# Vacuum

Beam vacuum  $\sim 10^{-10}$  mbar ( $\sim 3$  million molecules/cm<sup>3</sup>)



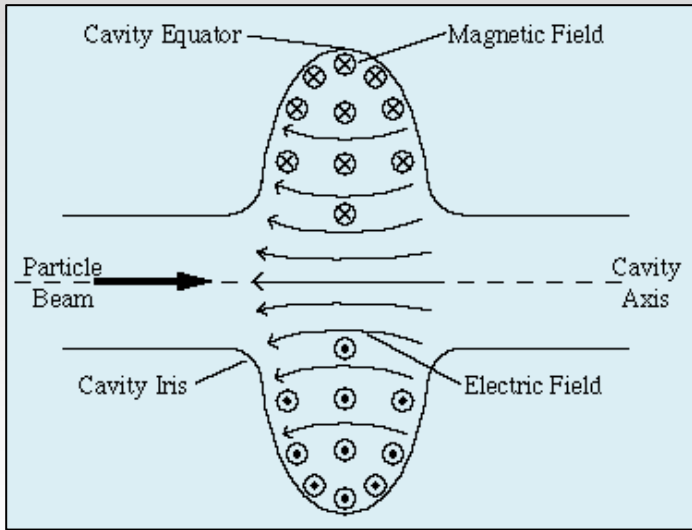
- **Cold**

- Pumping is insured by cold surfaces for all gases except helium. To avoid subsequent desorption, low initial pressures are required before cool-down, and this is ensured by turbo molecular pumps etc.

- **Warm**

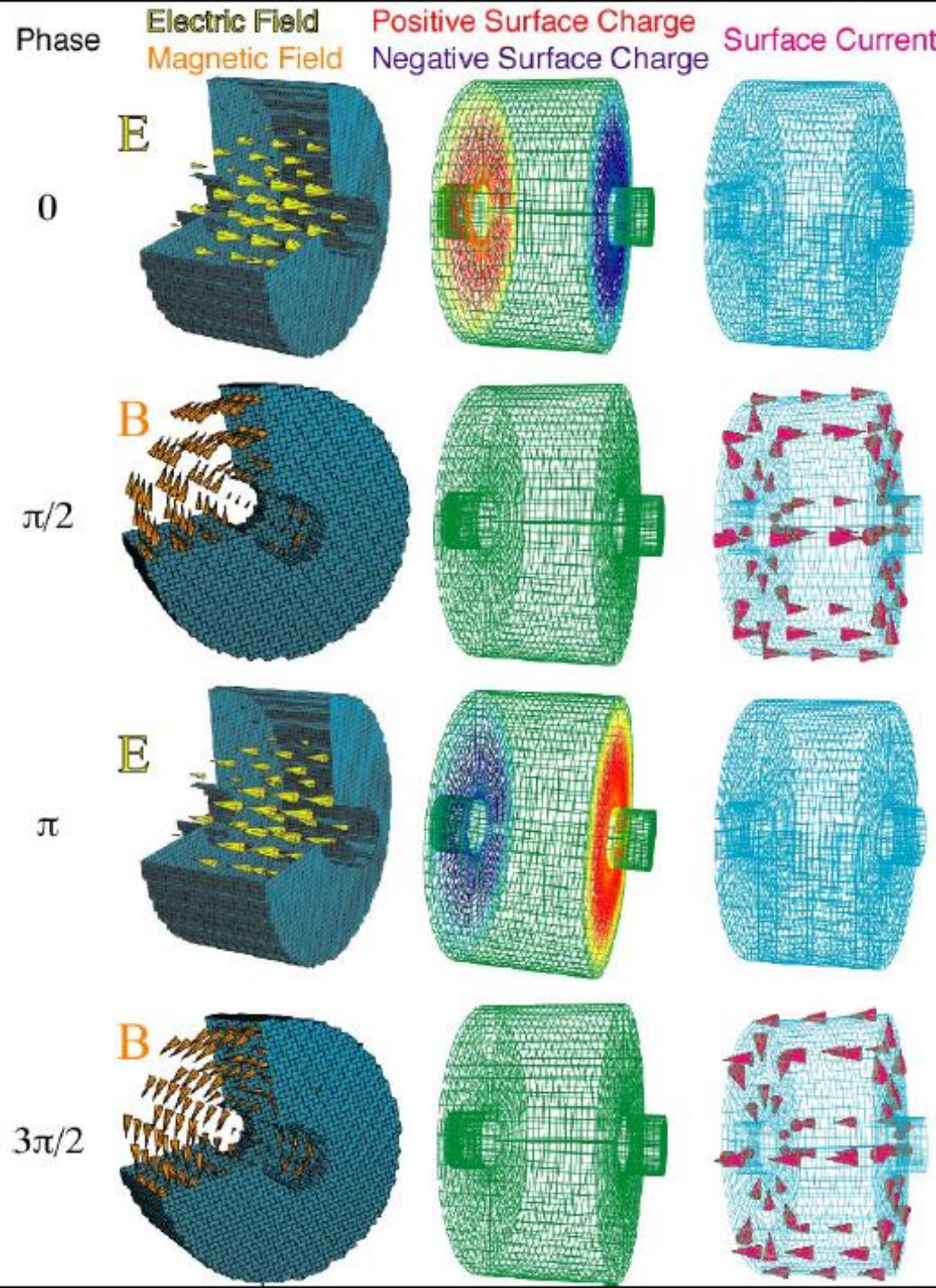
- NEG provides most of the pumping capacity, with additional ion pumps for the noble gases which are not pumped by the NEG



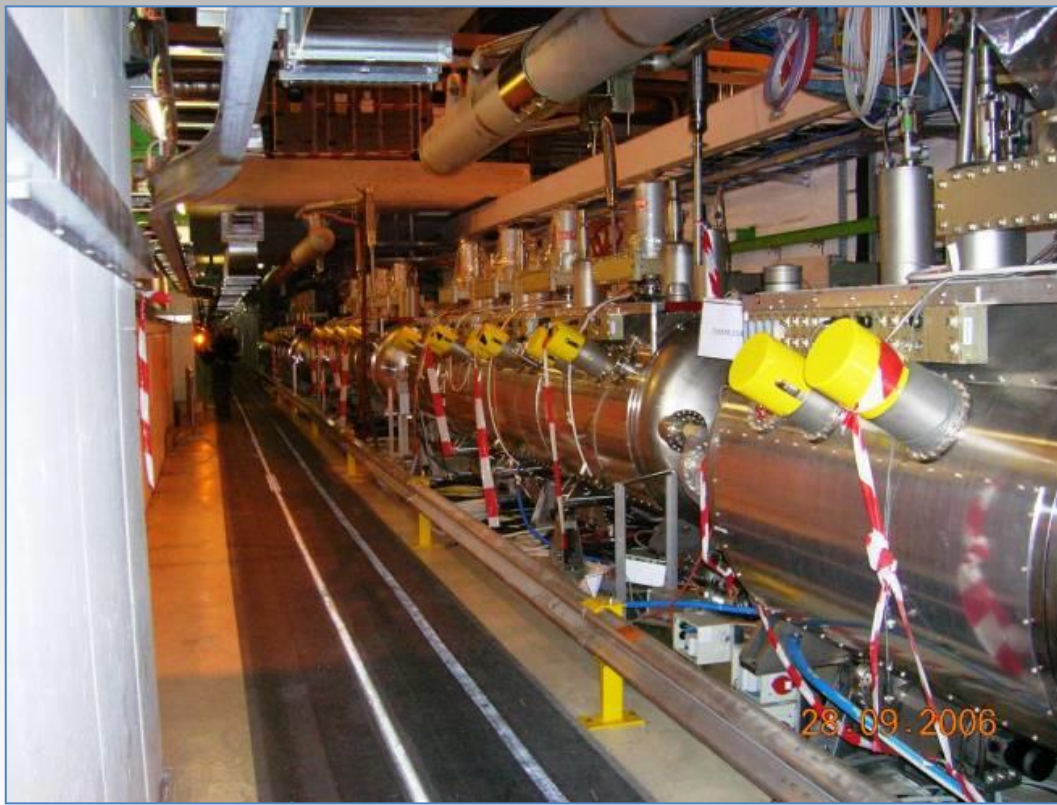


# RADIO FREQUENCY

Briefly!

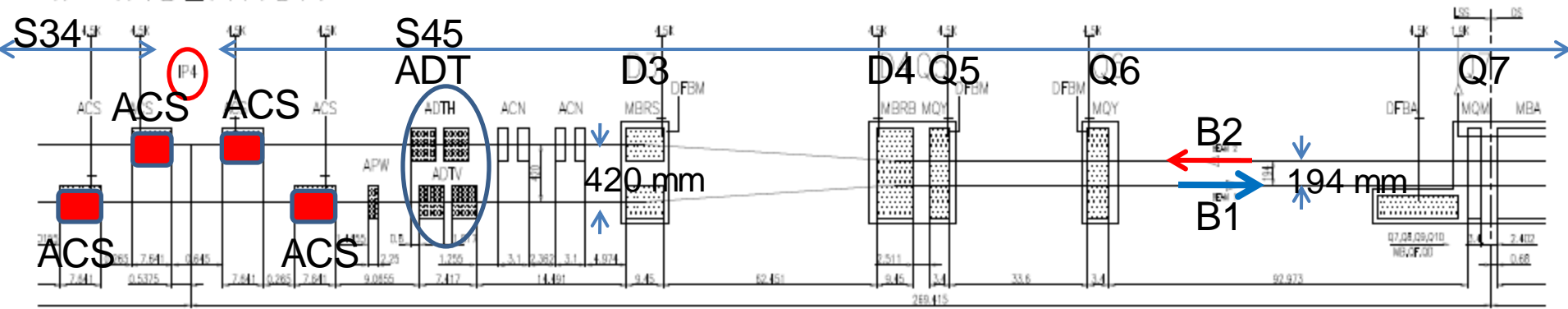






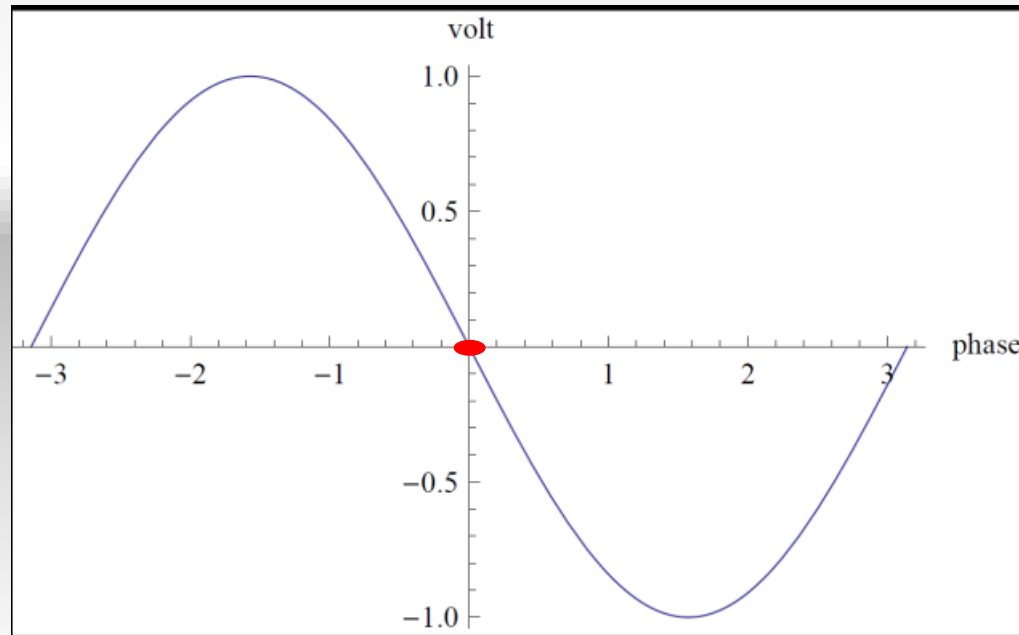
- 4xfour-cavity cryo module
- 400 MHz
- 16 MV/beam
- Nb on Cu cavities @4.5 K
- Beam pipe diam.=300 mm

RF INSERTION



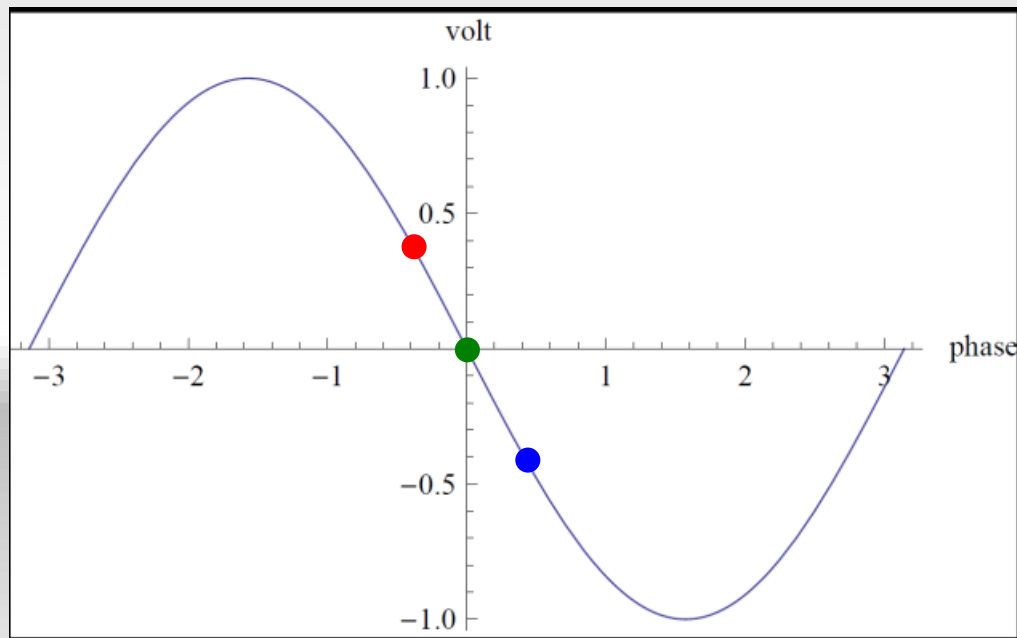
# RF basics

- We will use a RF cavities operating on resonance to produce an oscillating electric fields (400 MHz in the LHC)
- In the time a bunch takes to travel around the ring the RF performs 35640 cycles – the harmonic number (or  $f_{\text{rf}} = hf_{\text{rev}}$ )



# Now...

- A particle with lower momentum than  $p_0$  will go round faster and will arrive at the RF cavity earlier. It will get a higher energy kick and arrive relatively later the next turn.
- A particle with higher momentum than  $p_0$  will go round slower and will arrive at the RF cavity later. It will get a lower energy kick and arrive relatively earlier the next turn.



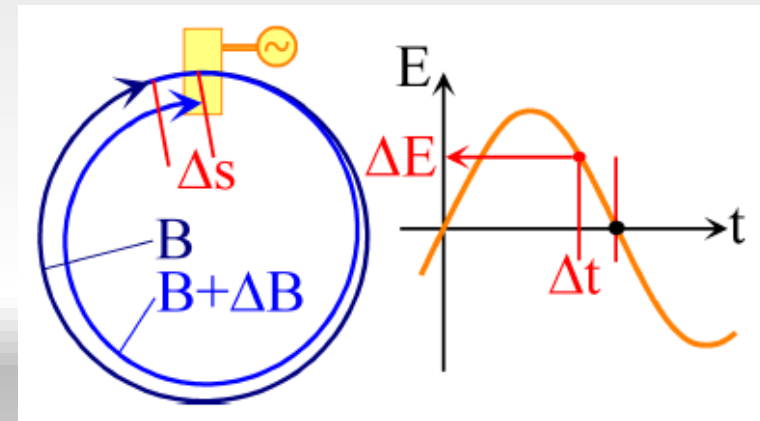
$$\phi = \phi_s$$

(In the LHC when  $dp/dt=0$ , the stable phase is 180 degrees. During the ramp it reaches 176.5 degrees)

$$\frac{mv^2}{R} = qvB \rightarrow p = qRB \rightarrow p(t) = qRB(t)$$

Momentum follows magnet field variation due to RF phase focussing:

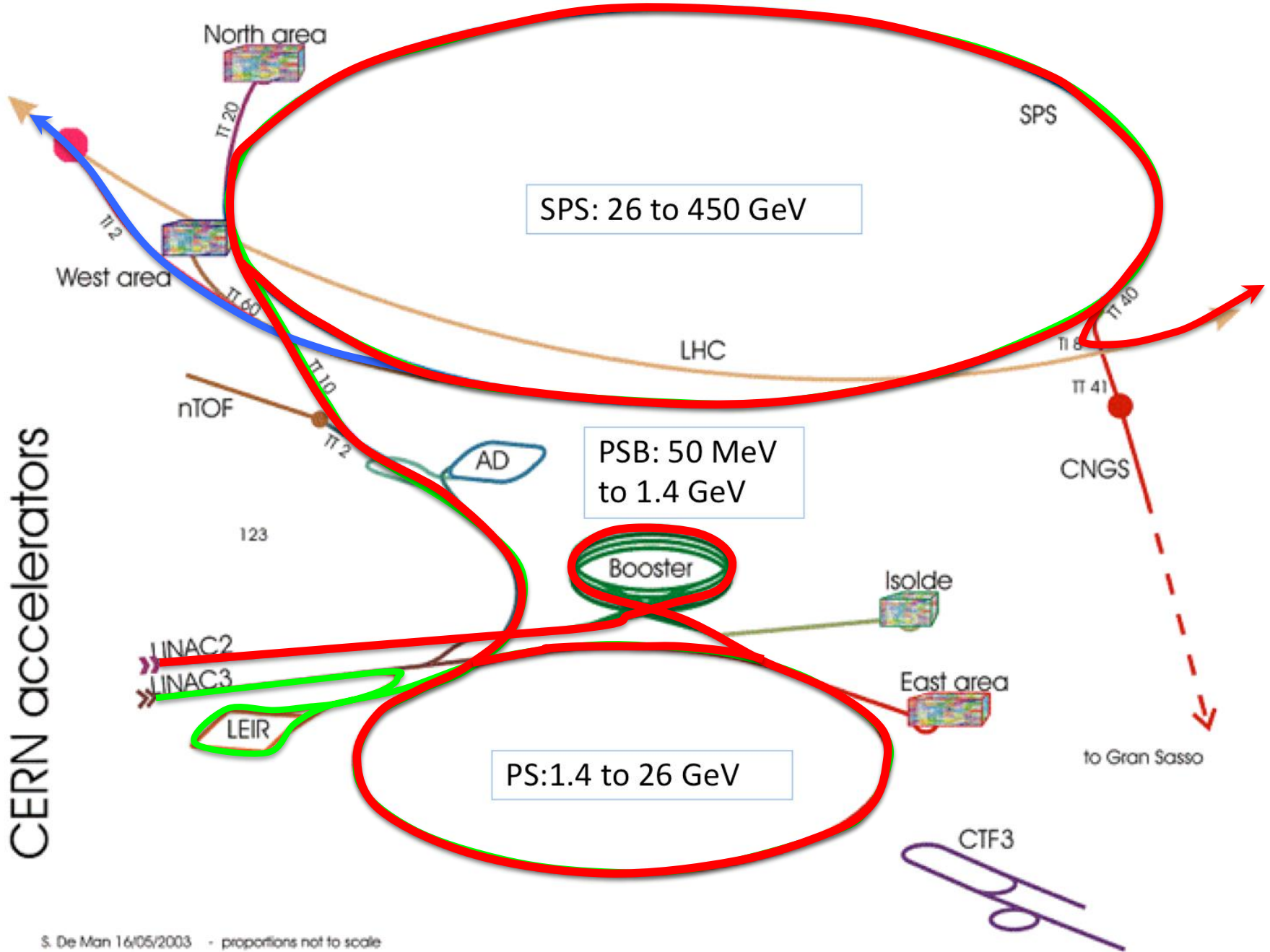
- inject beam into ring at  $B_0$  with momentum  $p_0 = qRB_0$
- increase B-field  $\rightarrow B + \Delta B$
- bending radius shrinks
- path becomes shorter by  $2\pi\Delta R$
- particles arrive earlier by  $\Delta t = (2\pi\Delta R)/\beta c$
- RF cavity:  $U(\Delta t) = U_0 \sin(\omega\Delta t + \phi) > 0$
- Acceleration by  $\Delta p = \beta q U(\Delta t)$



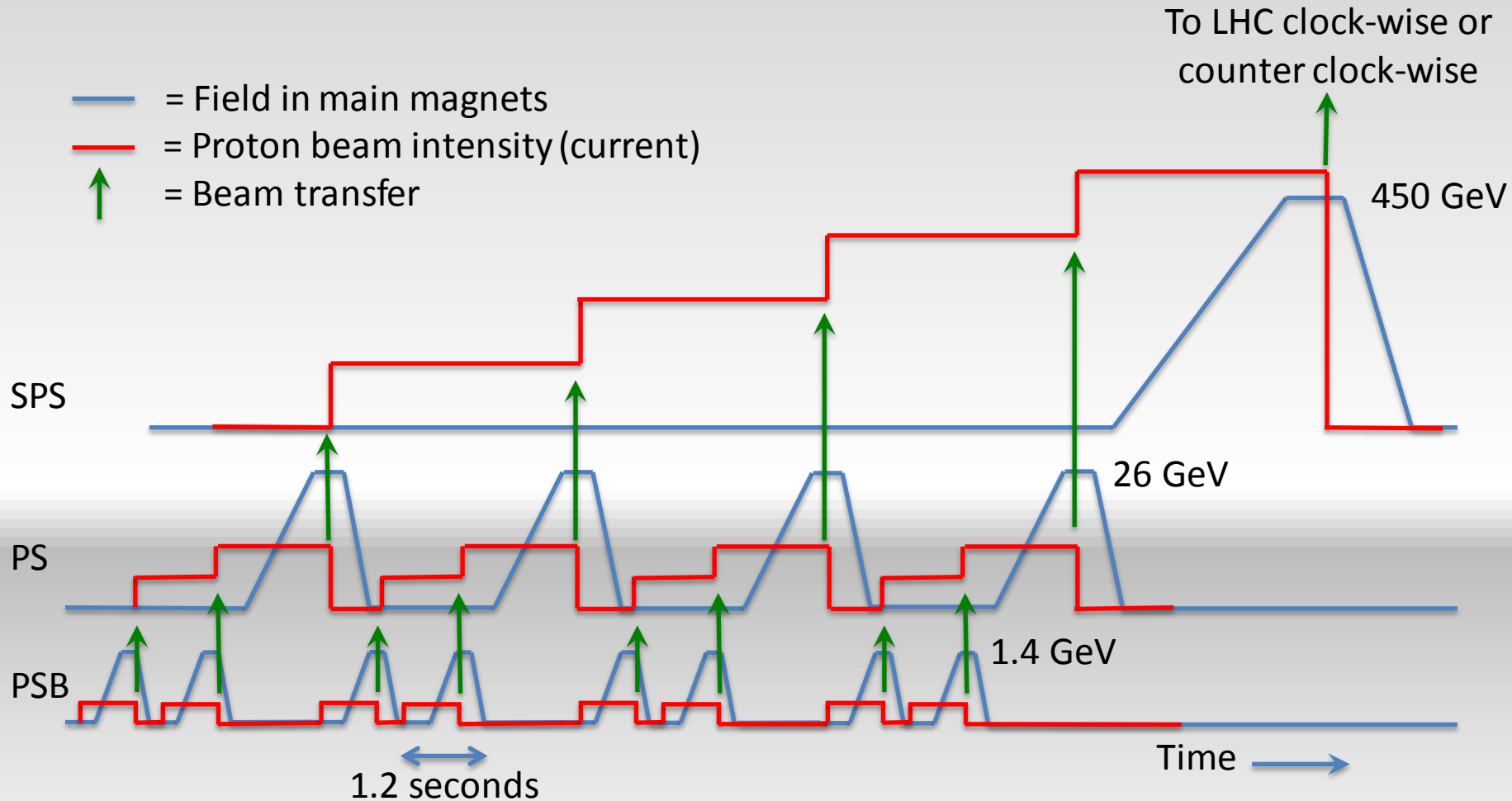
- $\Rightarrow$  self-synchronization of  $p(t)$  with  $B(t)$
- Constraints:  $\phi \approx \pi$  and  $2\pi R = n\beta\lambda_{rf}$



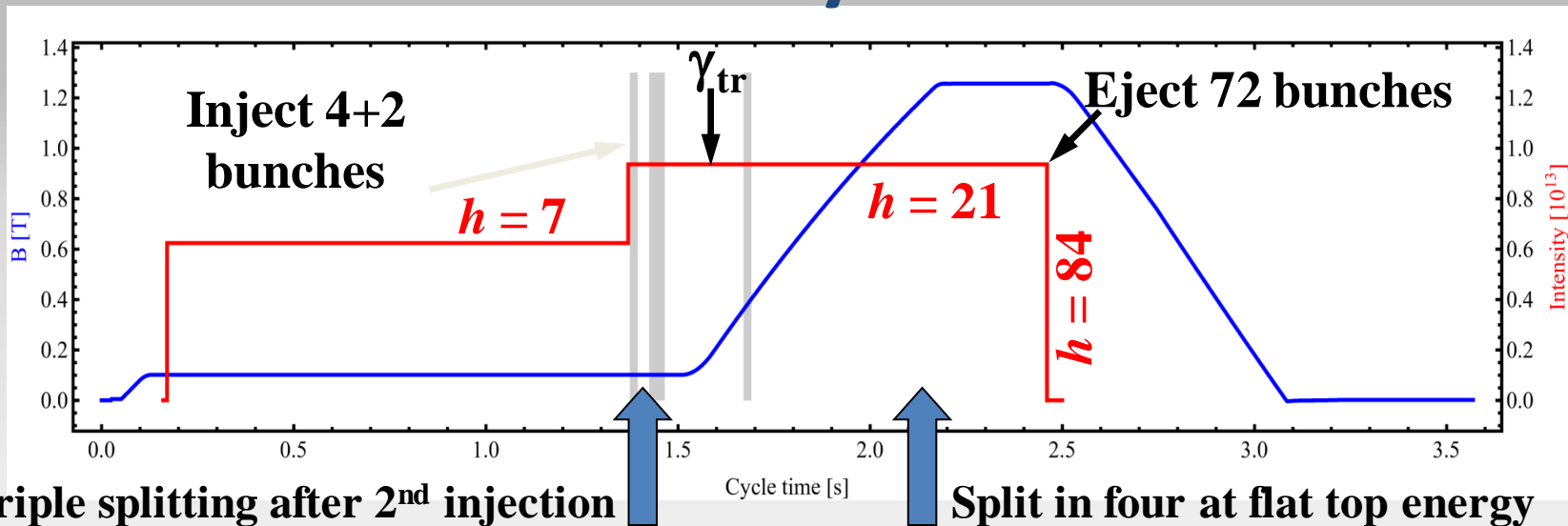
# CERN accelerators



# LHC Injector Cycling

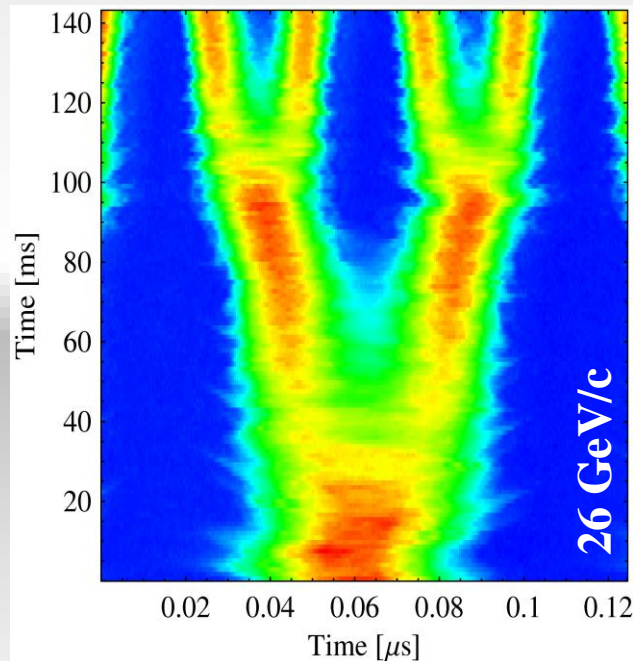
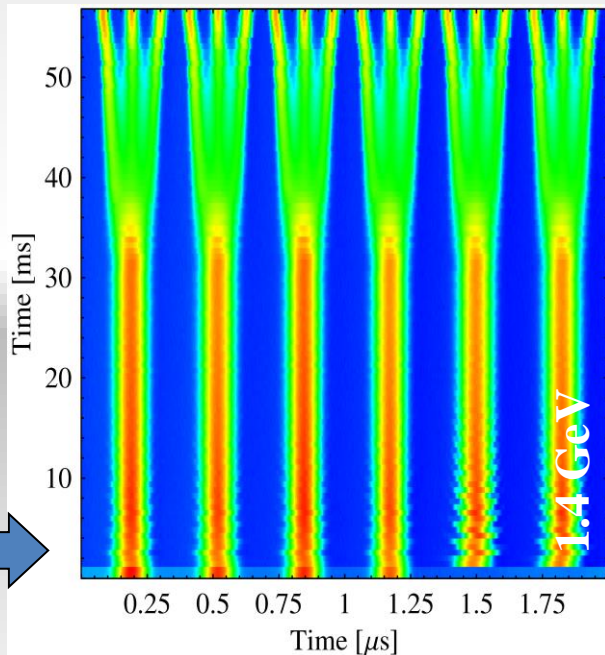
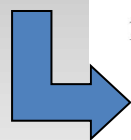


# The LHC 25 ns cycle in the PS



(sketched)

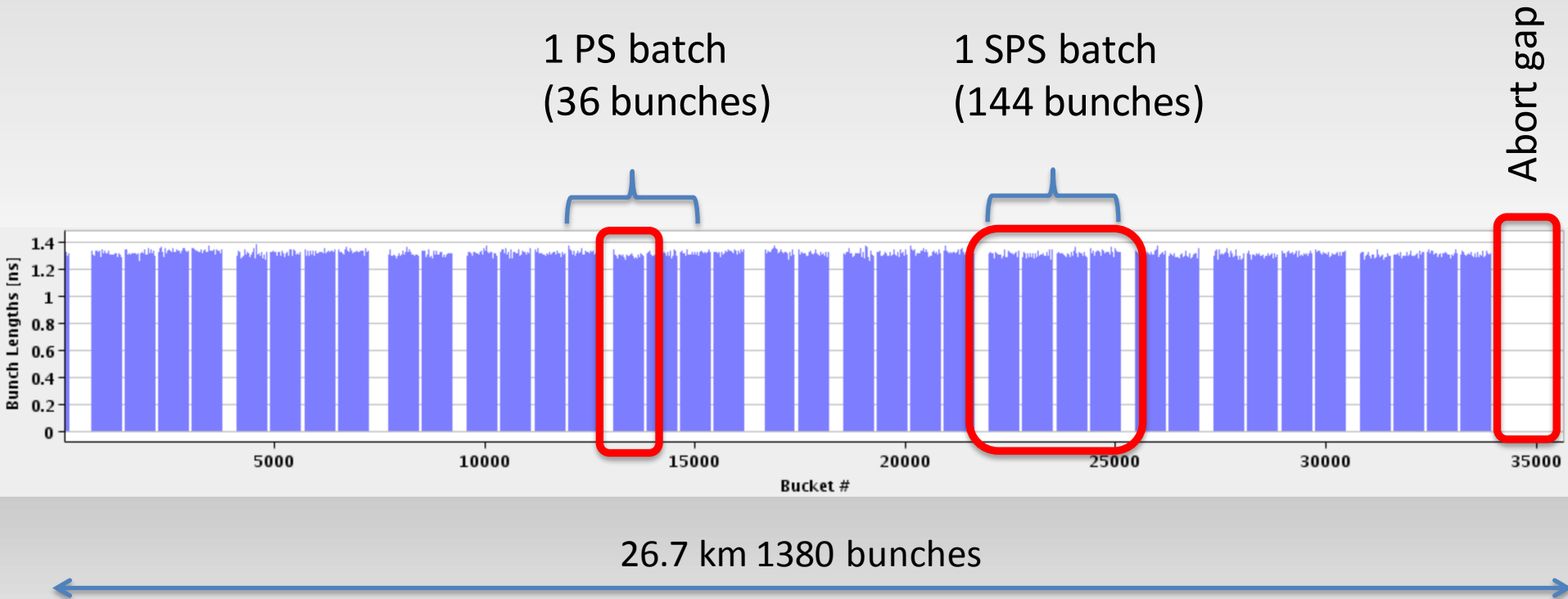
**2<sup>nd</sup> injection**



→ Each bunch from the Booster divided by 12 →  $6 \times 3 \times 2 \times 2 = 72$

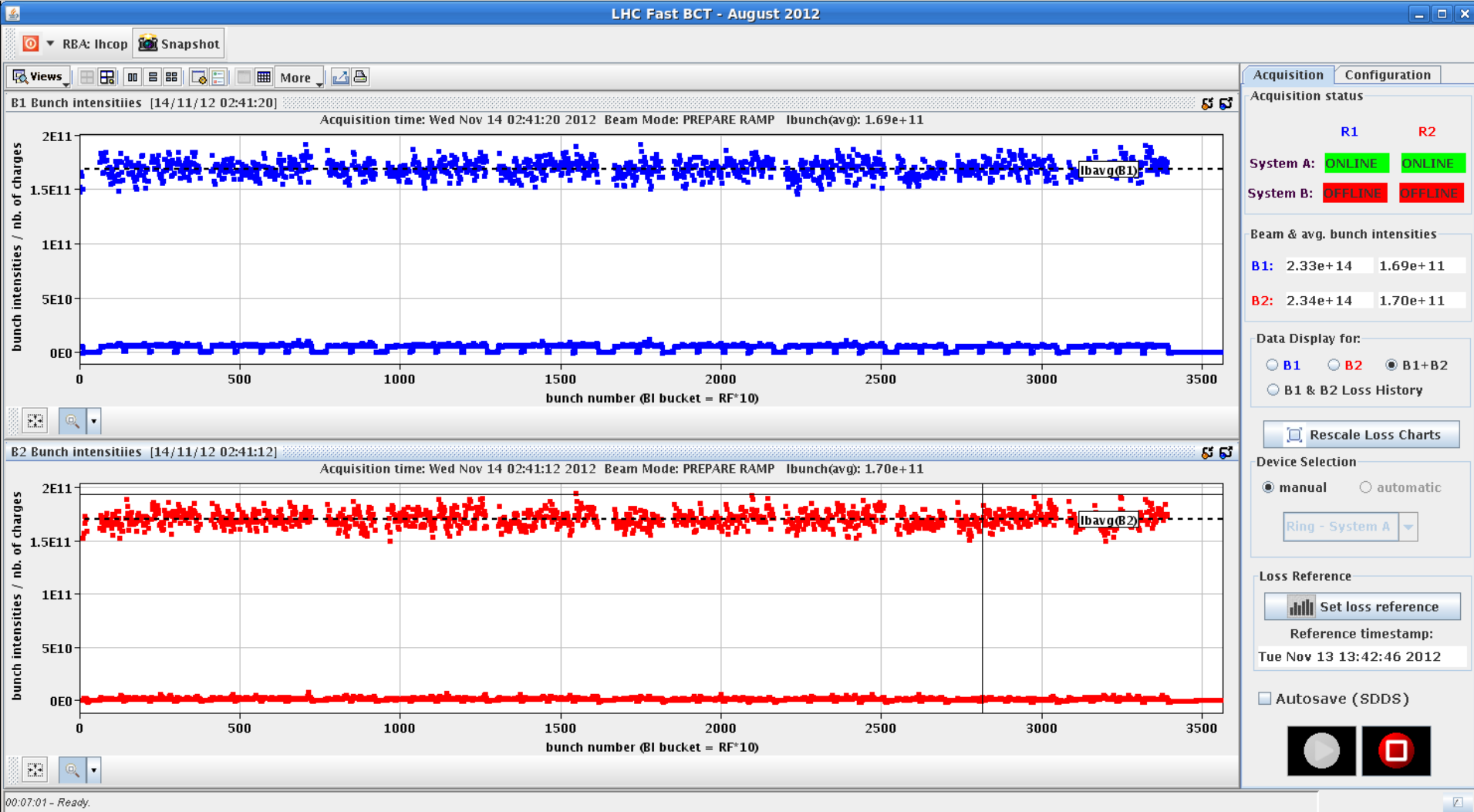
# LHC bunch structure - 2012

- 50 ns bunch spacing
- Maximum bunch intensity  $1.7 \times 10^{11}$  protons per bunch

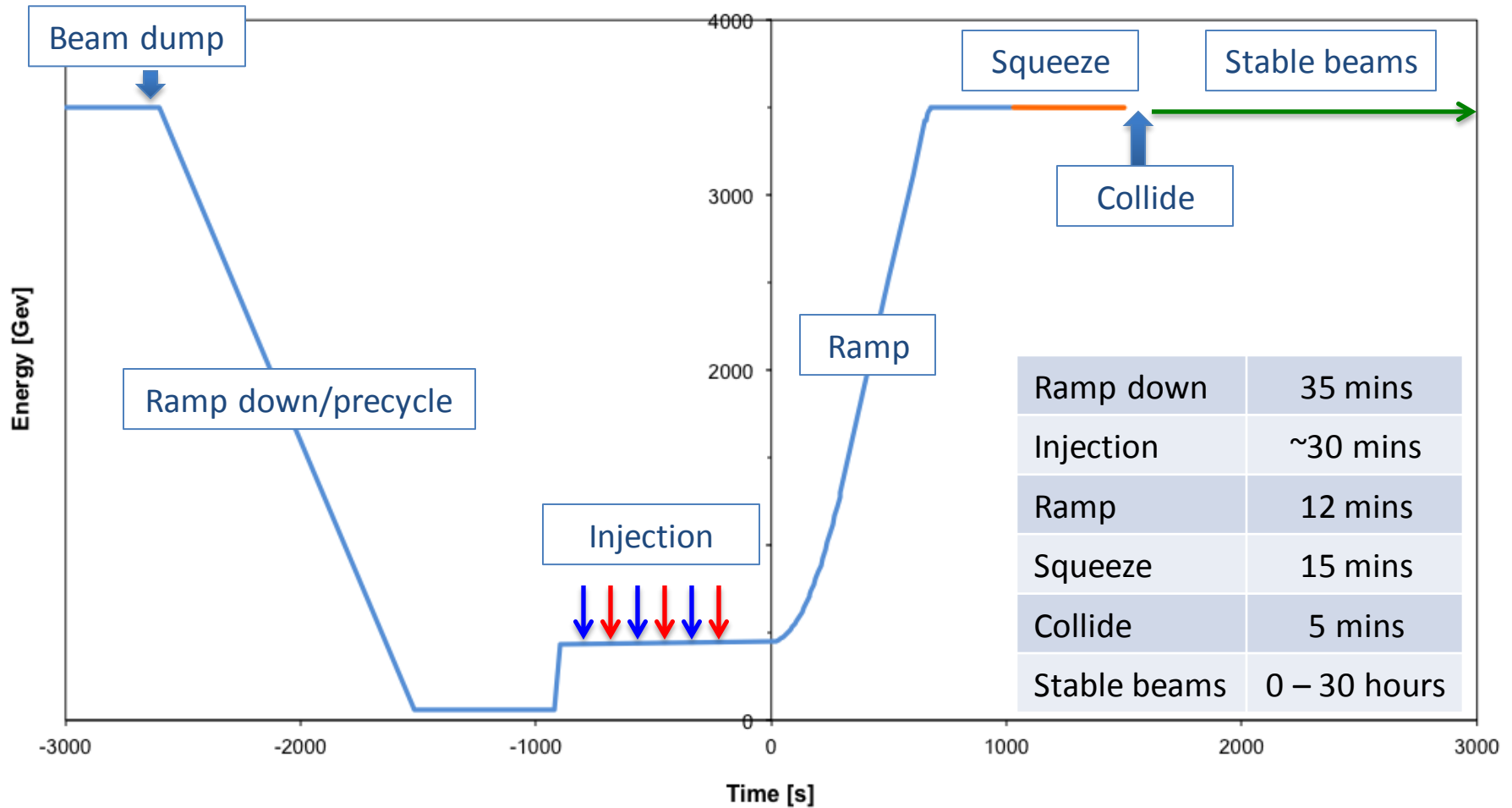




# From the control room



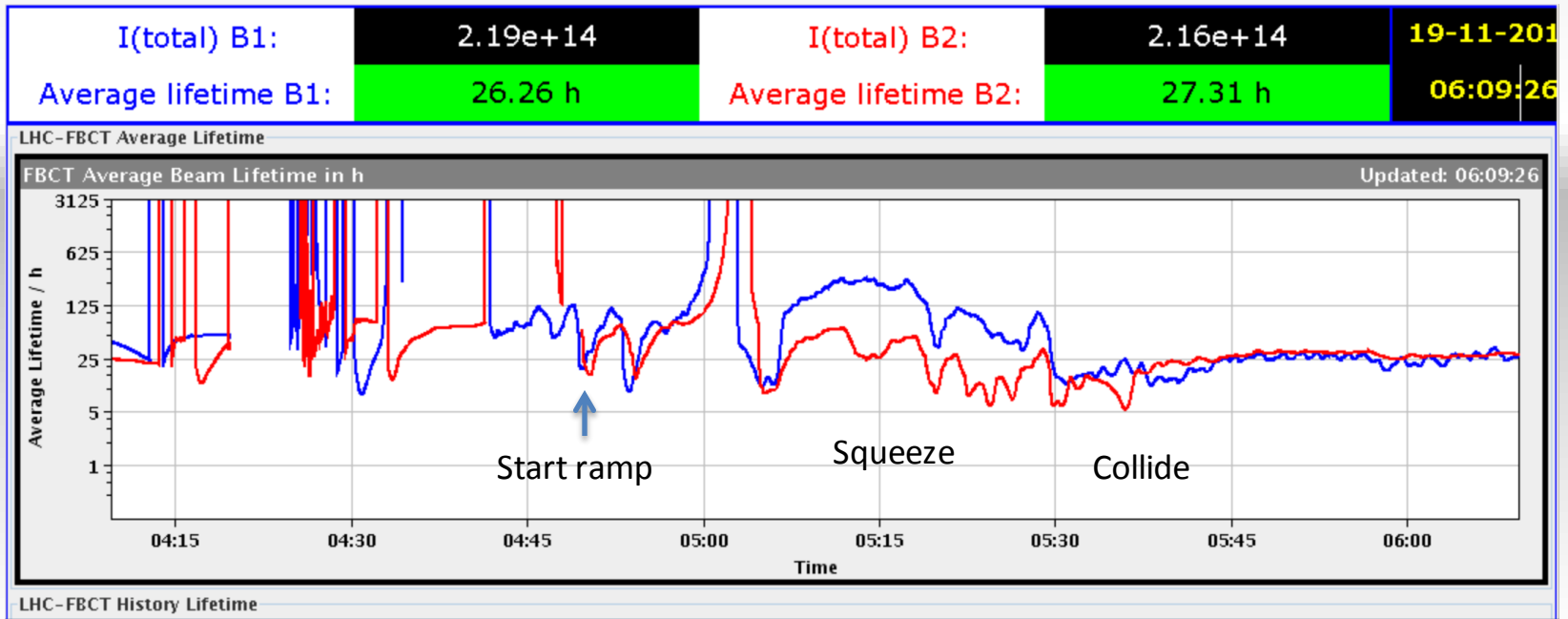
# Operational cycle



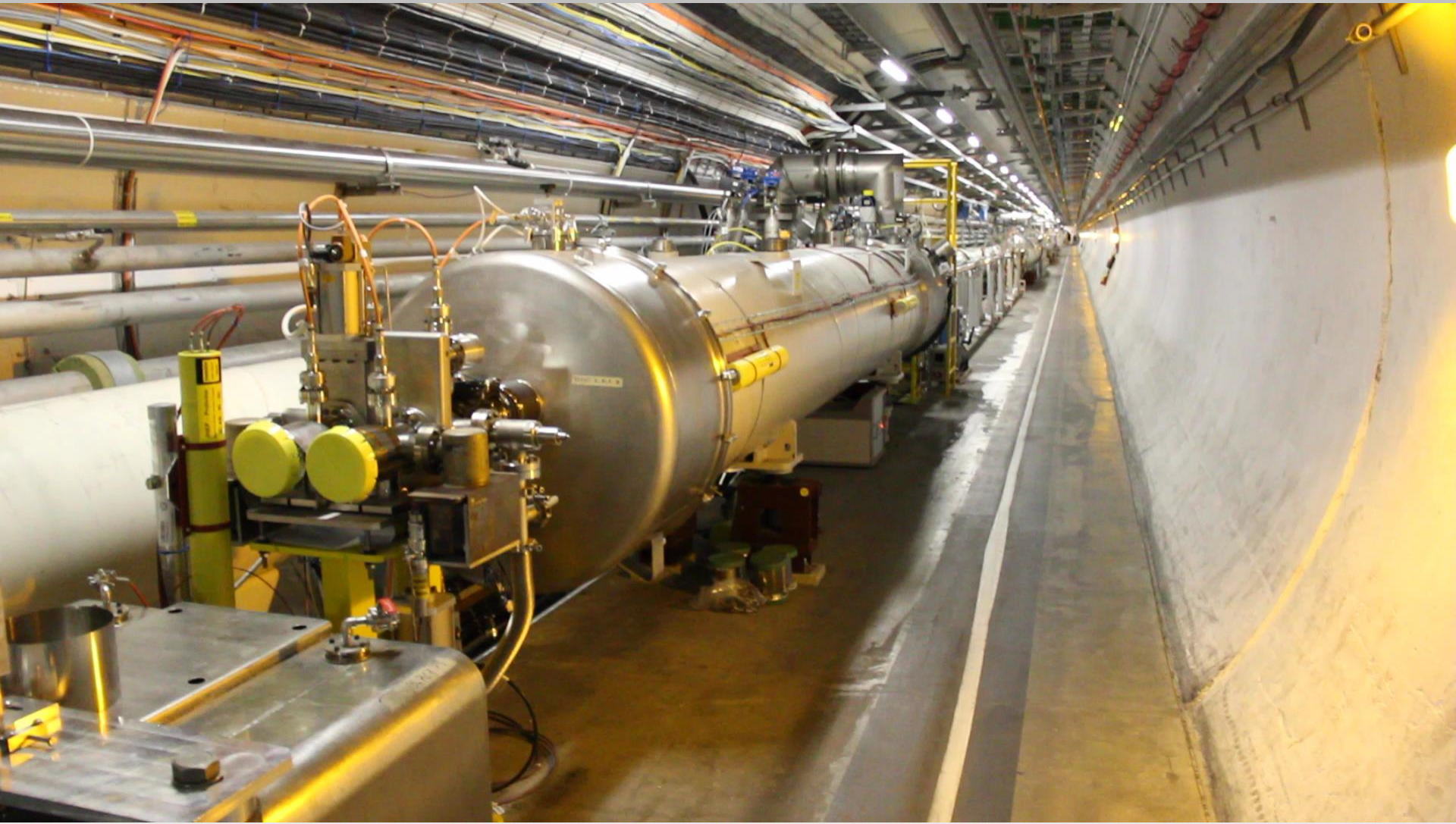
Turn around from stable beams to stable beams - 2 to 3 hours on a good day

# Beam

- Excellent single beam lifetime – good vacuum conditions
- Excellent field quality, good correction of non-linearities
- Low tune modulation, low power converter ripple, low RF noise

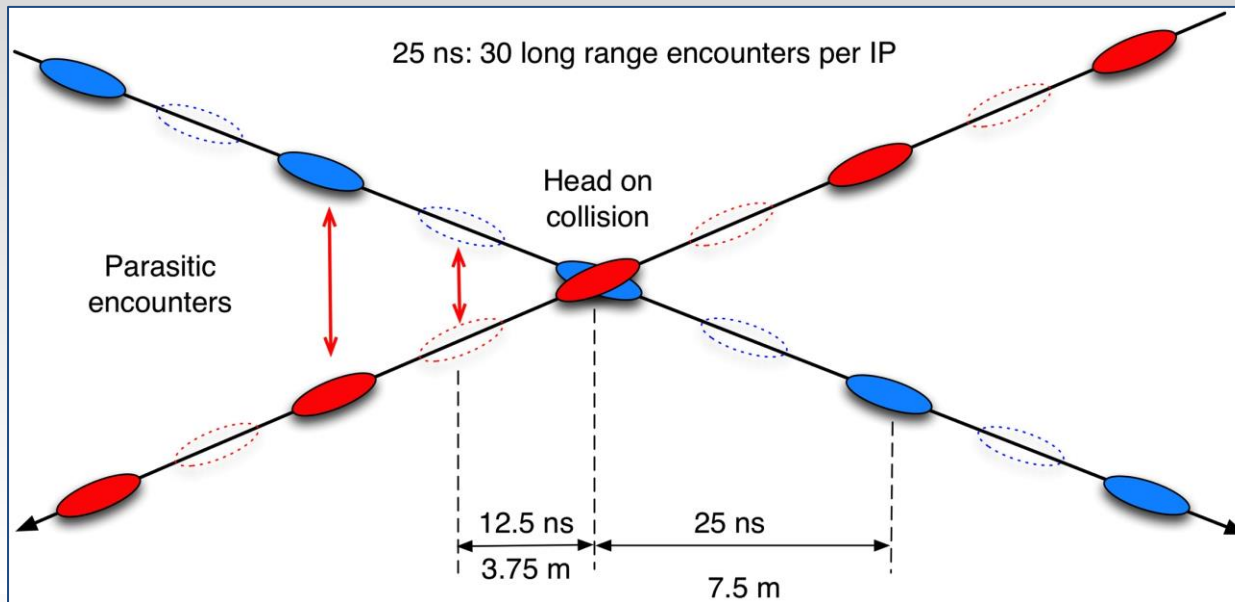


# In the long straight section



# Crossing angle

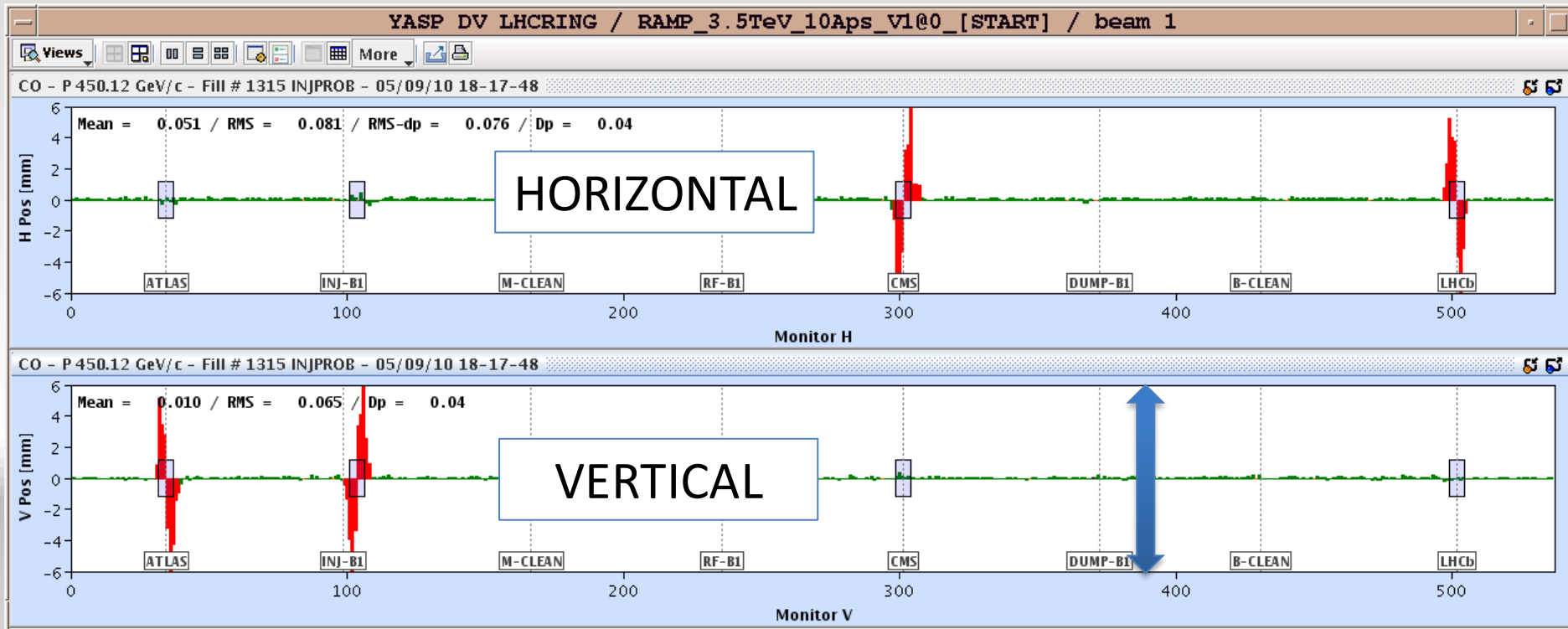
work with a crossing angle to avoid parasitic collisions.



- generates additional tune shift
- requires larger triplet magnet aperture
- breaks symmetry between x,y planes
- odd order resonances are excited
- couples longitudinal and transverse motion
- breaks the bunch symmetry
- lowers available luminosity

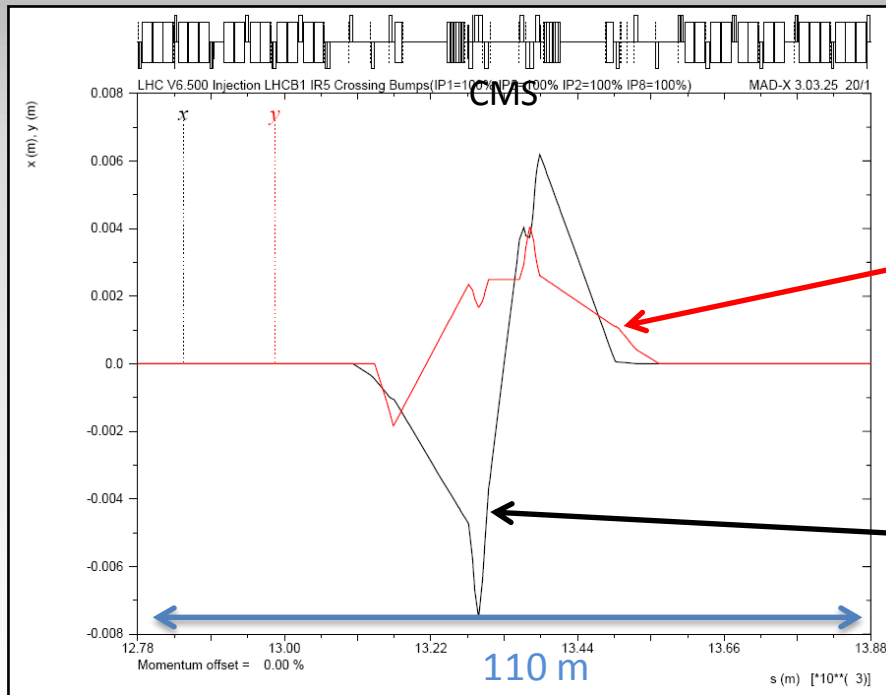


# In practice



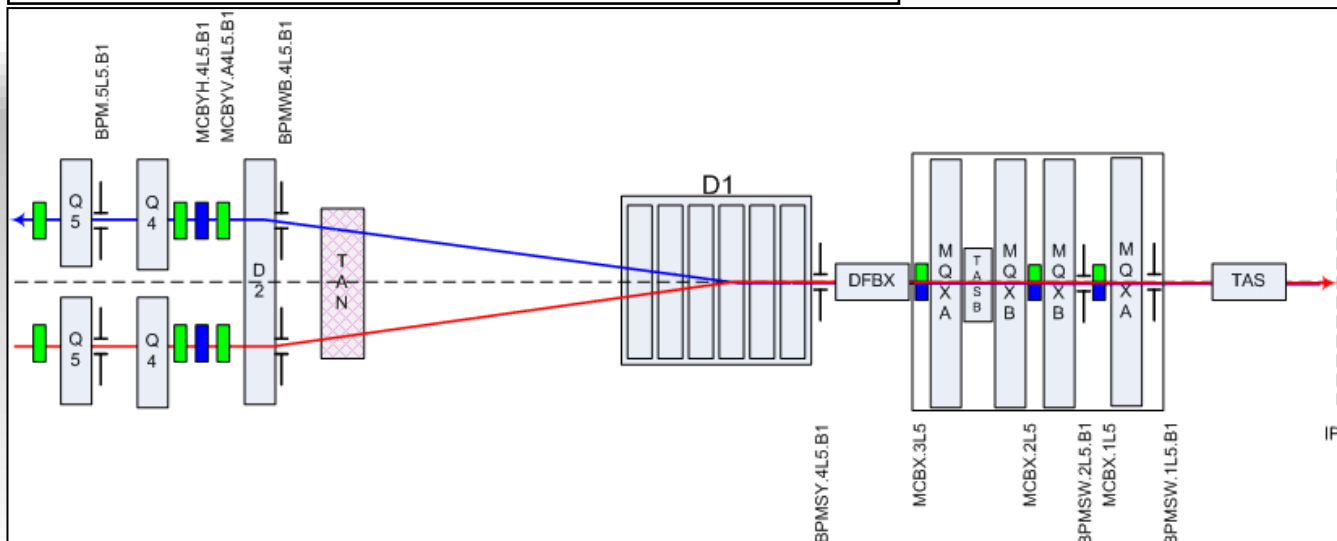
+/- 6 mm at 450 GeV

# Crossing and Separation Bumps

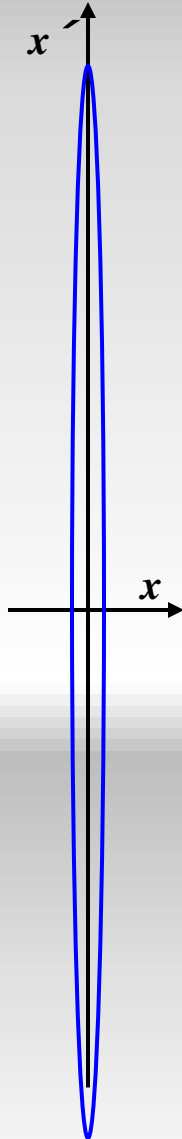
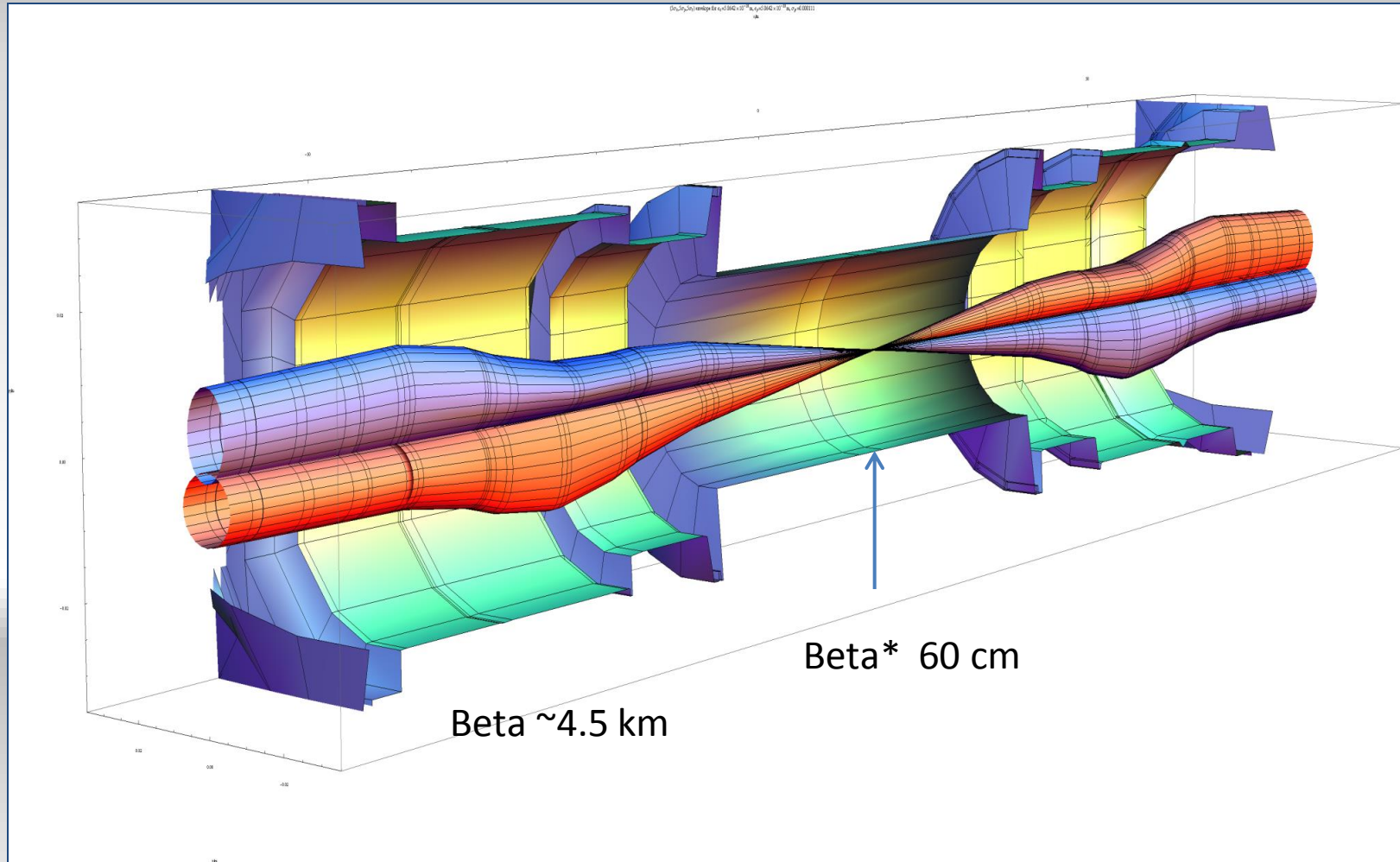


Collapse the **separation bump** after the squeeze to bring the beams into collision

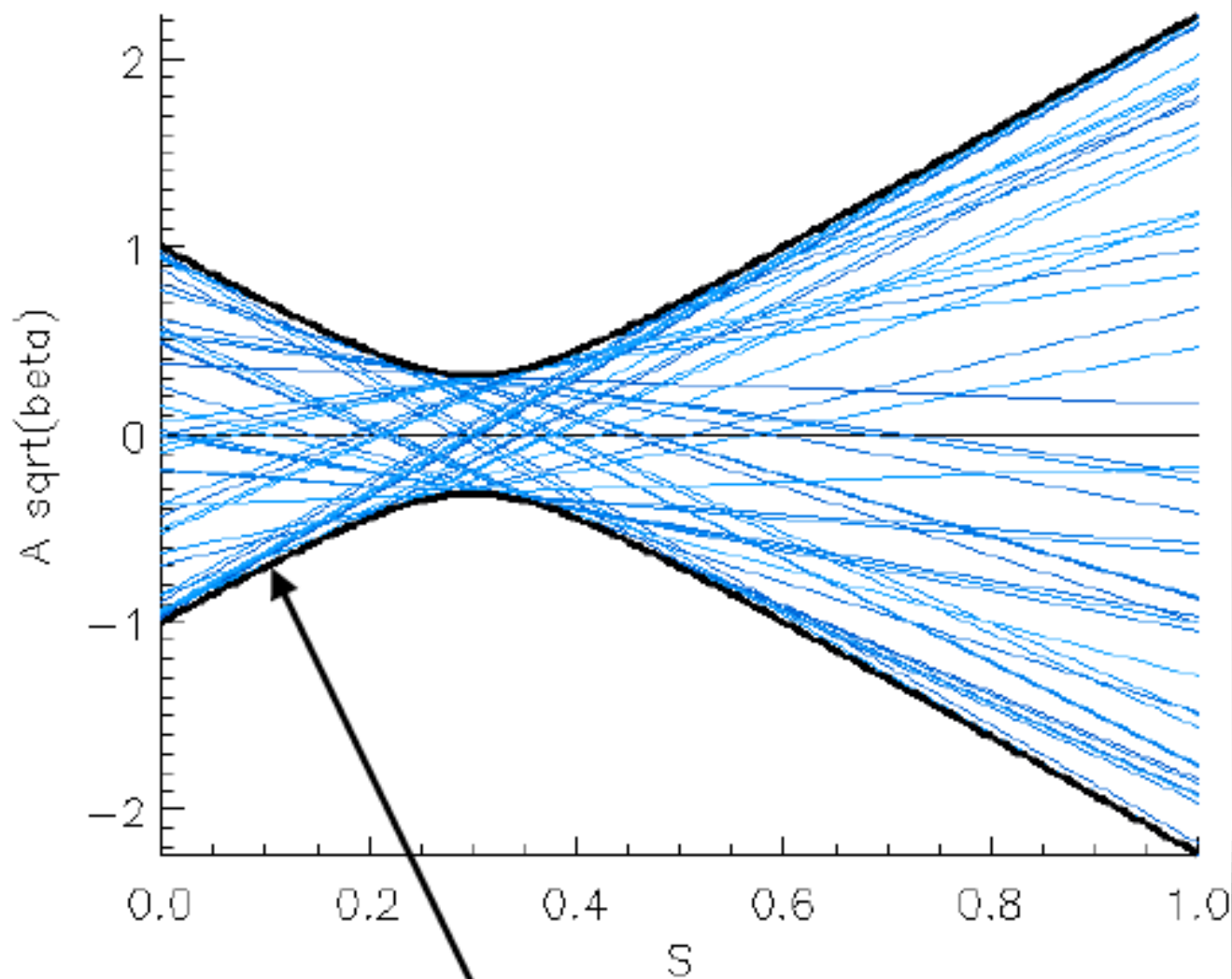
Crossing angle bump



# Squeeze in ATLAS



NB: round beams



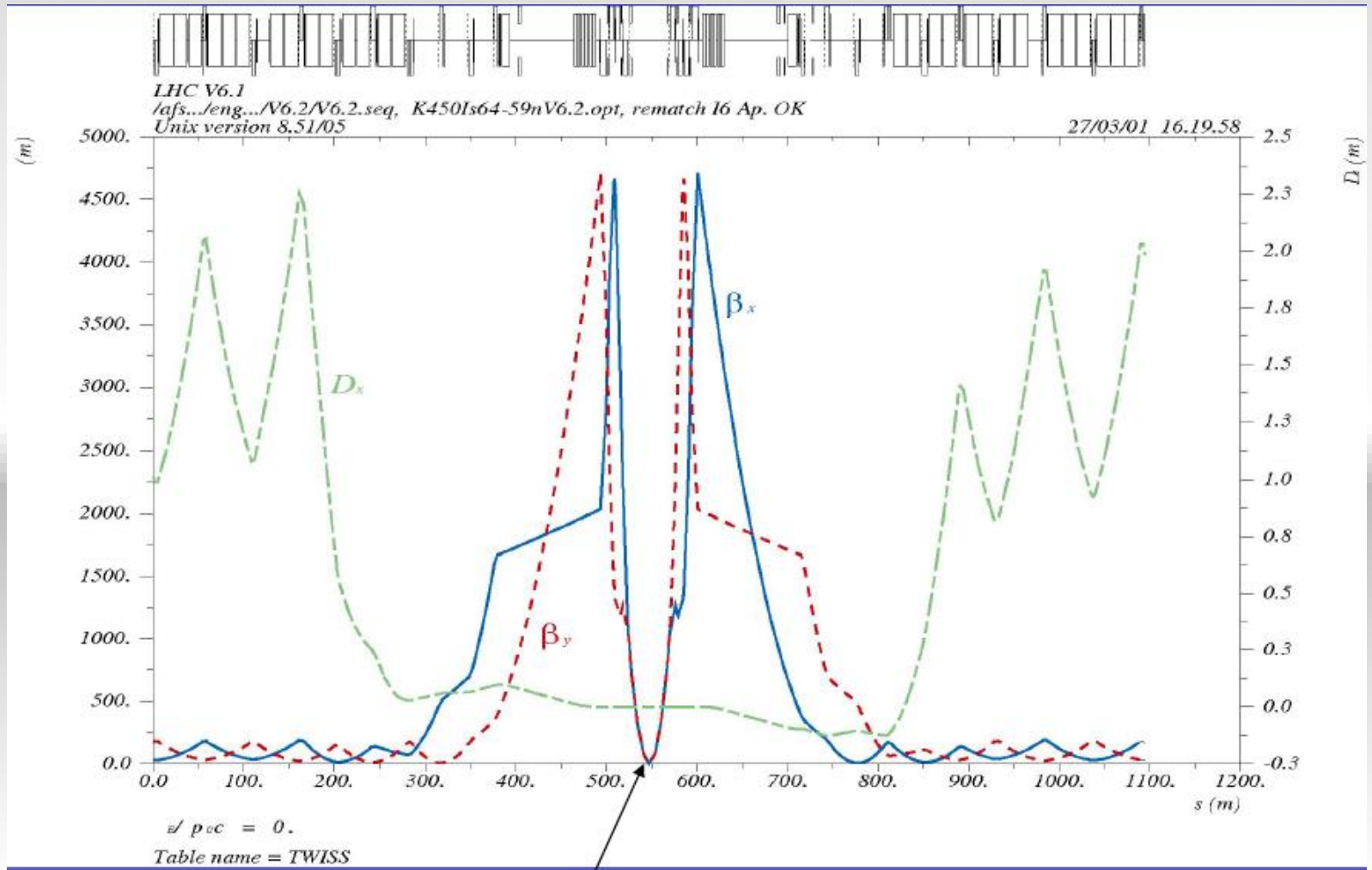
**beam envelope**

$$x_{\max}(s) = A\sqrt{\beta(s)}$$

# Beam - Squeeze

Small beam in the IP  $\rightarrow$  big beams in the inner triplets  $\rightarrow$  reduced aperture

Therefore inject & ramp with bigger beam sizes at IP.





# Triplets



# Squeeze in practice

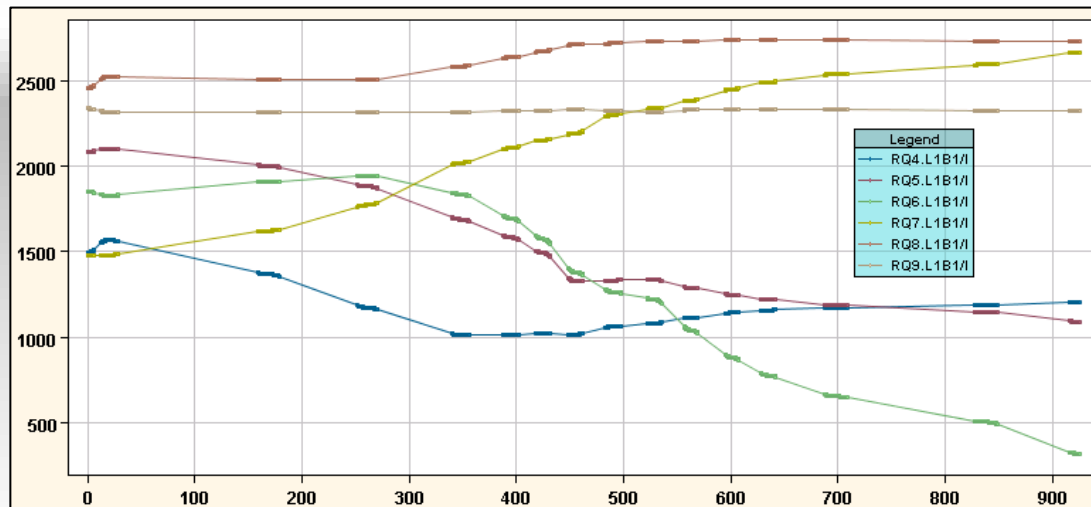
Matched optics

Time in seconds

Optic Name	Energy	Time
A1100C1100A1000L1000_INJ_2012	4000.0	0
A1100C1100A1000L1000_2012	4000.0	19
A900C900A900_0.00915L750_0.00932_2012	4000.0	169
A700C700A750_0.00897L600_0.00909_2012	4000.0	262
A400C400A600_0.00889L500_0.00900_2012	4000.0	348
A300C300A500_0.00889L375_0.00888_2012	4000.0	396
A250C250A450_0.00889L350_0.00882_2012	4000.0	425
A200C200A400_0.00889L325_0.00878_2012	4000.0	455
A160C160A350_0.00889L300_0.00875_2012	4000.0	491
A150C150A300_0.00889L300_0.00875_2012	4000.0	529
A120C120A300_0.00889L300_0.00875_2012	4000.0	563
A100C100A300_0.00889L300_0.00875_2012	4000.0	602
A90C90A300_0.00889L300_0.00875_2012	4000.0	634
A80C80A300_0.00889L300_0.00875_2012	4000.0	696
A70C70A300_0.00889L300_0.00875_2012	4000.0	840
A60C60A300_0.00889L300_0.00875_2012	4000.0	925

← Beta\* - 11 m ATLAS, CMS; 10 m in ALICE, LHCb

← Beta\* - 0.6 m ATLAS, CMS; 3 m in ALICE, LHCb



Current during the squeeze in a few quads at point 1

# Luminosity

$$L = F \frac{N_{b1} N_{b2} f_{rev} k_b}{2\pi \sqrt{(\sigma_{x1}^2 + \sigma_{x2}^2)(\sigma_{y1}^2 + \sigma_{y2}^2)}} \cdot \exp \left\{ -\frac{(\bar{x}_1 - \bar{x}_2)^2}{2(\sigma_{x1}^2 + \sigma_{x2}^2)} - \frac{(\bar{y}_1 - \bar{y}_2)^2}{2(\sigma_{y1}^2 + \sigma_{y2}^2)} \right\}$$

$$F = \frac{1}{\sqrt{1 + \frac{q_c s_z \ddot{\theta}^2}{e 2 S^* \dot{\theta}}}}$$

$N_1, N_2$  – number of particles per bunch  
 $k$  – number bunches per beam  
 $f$  – revolution frequency  
 $\sigma^*$  – beam size at IP  
 $\theta_c$  – crossing angle  
 $\sigma_z$  – bunch length

**Make some simplifying assumptions:**

- **beam 1 = beam 2**
- **round beams at interaction point**

# Luminosity

$$L = \frac{N^2 k_b f}{4\rho s_x^* s_y^*} F = \frac{N^2 k_b f g}{4\rho e_n b^*} F$$

**N** Number of particles per bunch

**k<sub>b</sub>** Number of bunches

**f** Revolution frequency

**σ\*** Beam size at interaction point

**F** Reduction factor due to crossing angle

**ε** Emittance

**ε<sub>n</sub>** Normalized emittance

**β\*** Beta function at IP

$$s^* = \sqrt{b^* e}$$

$$e_N = 2.5 \cdot 10^{-6} \text{ m.rad}$$

$$e = 3.35 \cdot 10^{-10} \text{ m.rad}$$

$$s^* = 11.6 \cdot 10^{-6} \text{ m}$$

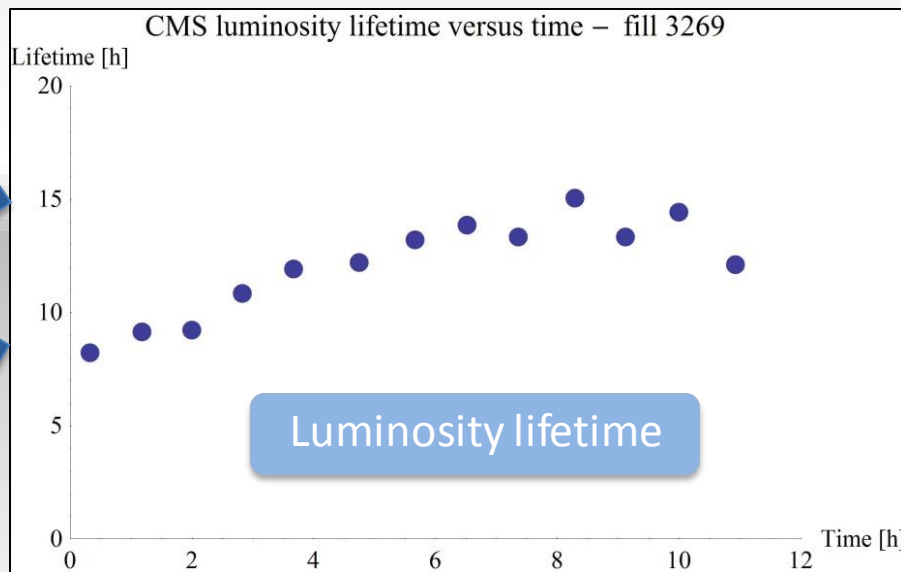
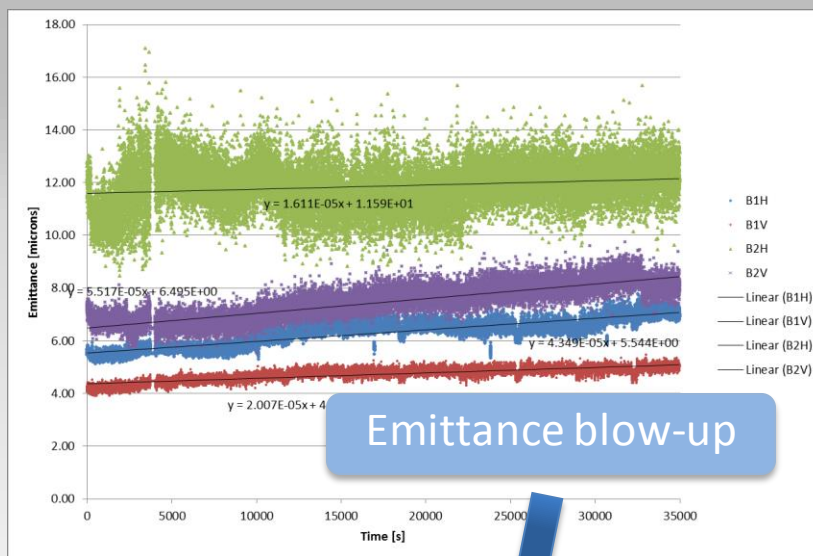
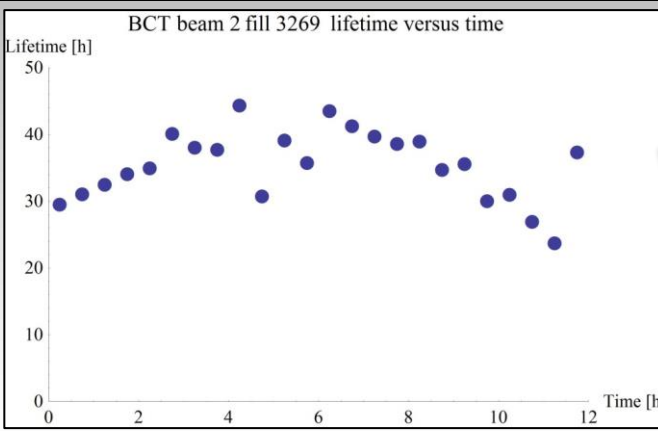
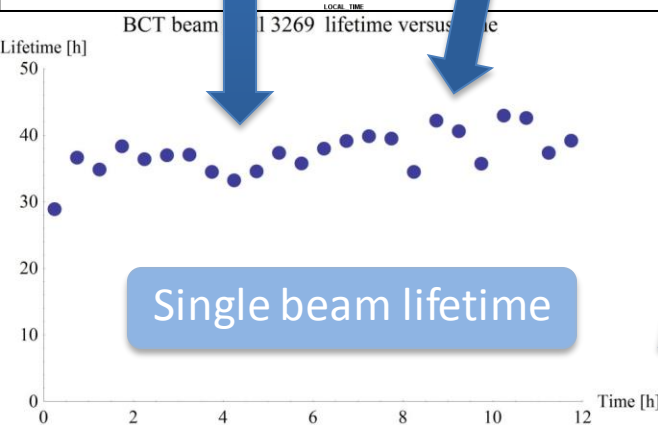
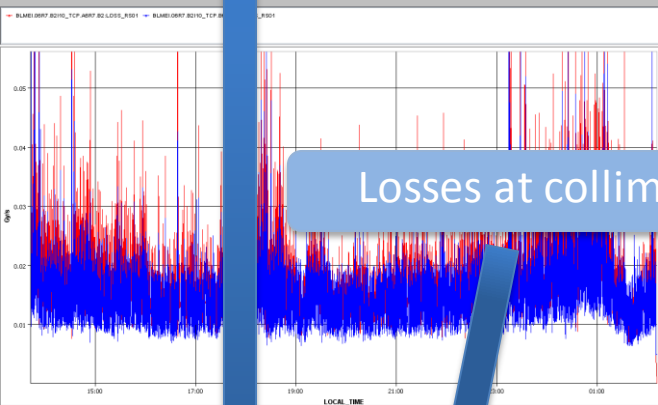
$$(p = 7 \text{ TeV}, b^* = 0.4 \text{ m})$$

# Reasonably comfortable life in Stable Beams

Luminosity burn

Beam-gas

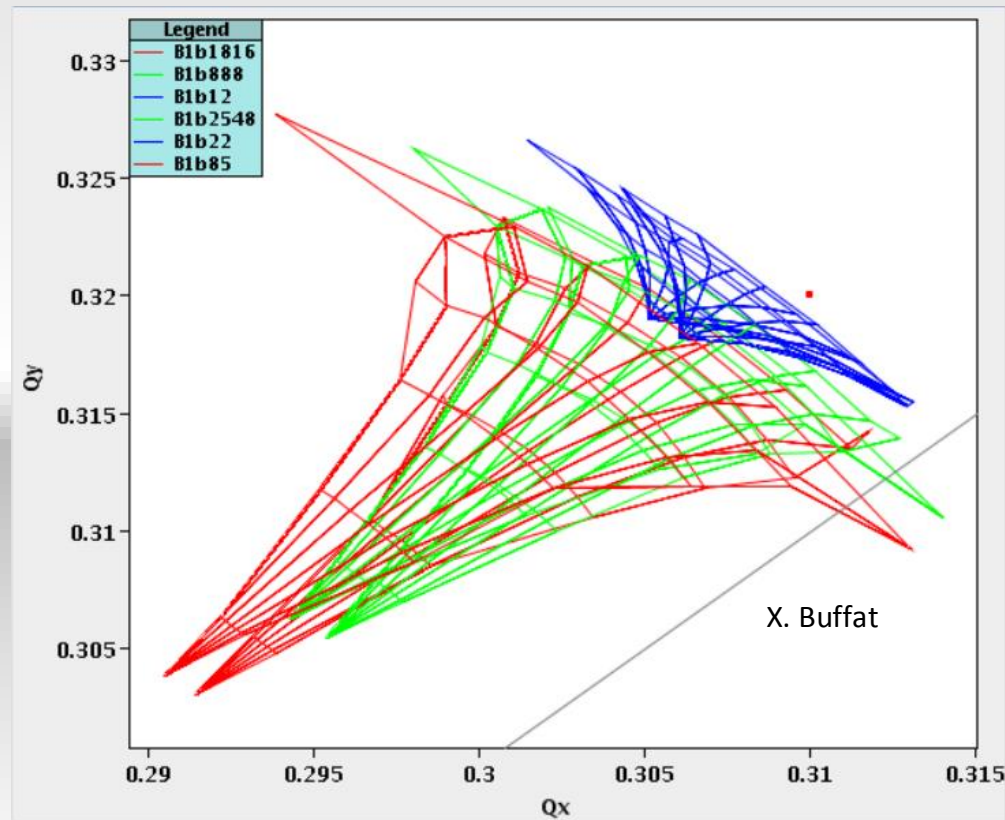
Losses at collimators





# Beam-beam

- Head-on beam-beam is not an operational limitation
- Linear head-on parameter in operation  $\sim 0.02$  (up to 0.034 in MD)
- Long range taken seriously
- Interesting interplay with the instabilities seen in 2012...





***A well-deserved toast to all who have built such a marvelous machine, and to all who operate it so superbly  
(first 7 TeV collisions on 30<sup>th</sup> March 2010)***



# First 7 TeV collisions – another view



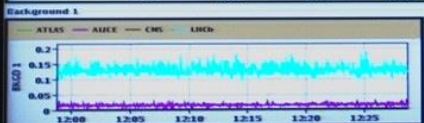
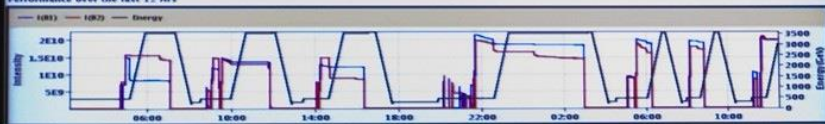
You lucky, lucky buggers!!!



30-Mar-2010 12:28:46 Fill #: 1005 Energy: 2982.0 GeV I(B1): 2.00e+10 I(B2): 2.01e+10

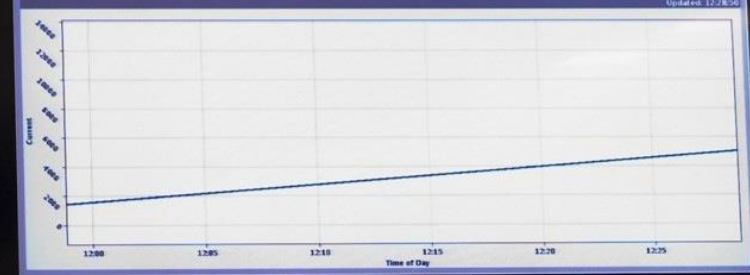
Experiment Status	ATLAS	ALICE	CMS	LHCb
Instantaneous Luminosity	1.746e-05	0.000e+00	1.621e-07	0.000e+00
BRAN Count Rate	2.987e+00	6.158e-02	2.166e-01	1.779e+00
BKGD 1	0.002	0.011	0.002	0.150
BKGD 2	0.000	0.000	0.000	0.373
BKGD 3	0.000	0.005	0.000	0.034

LHC# STANDBY Count(Chz): 0.000 LHCb VELO Position 200 Gap: 58.0 mm TOTEM: CALIBRATION



Monitoring set: S34-A34-RB 30 March, 2010, 12:28:49

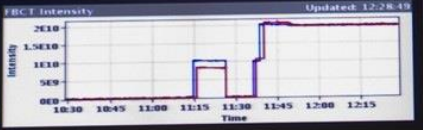
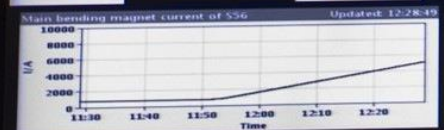
# RB.A34 [5023 A] [2.97 TeV]



LHC Page1 Fill: 1005 E: 2986 GeV 30-03-2010 12:28:49

## BEAM SETUP: RAMP

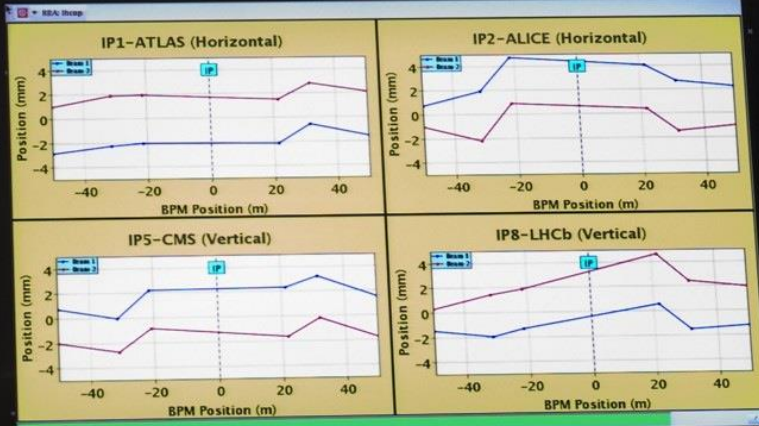
Energy: 2986 GeV I(B1): 1.82e+10 I(B2): 1.85e+10



Comments 30-03-2010 11:52:19 :  
Ramping again

BIS status and SMP flags	B1	B2
Link Status of Beam Permits	false	false
Global Beam Permit	true	true
Setup Beam	true	true
Beam Presence	true	true
Moveable Devices Allowed In	false	false
Stable Beams	false	false

LHC Operation in CCC : 77600, 70480 PM Status B1: ENABLED PM Status B2: ENABLED



12:28:49 Review

# **MACHINE PROTECTION**



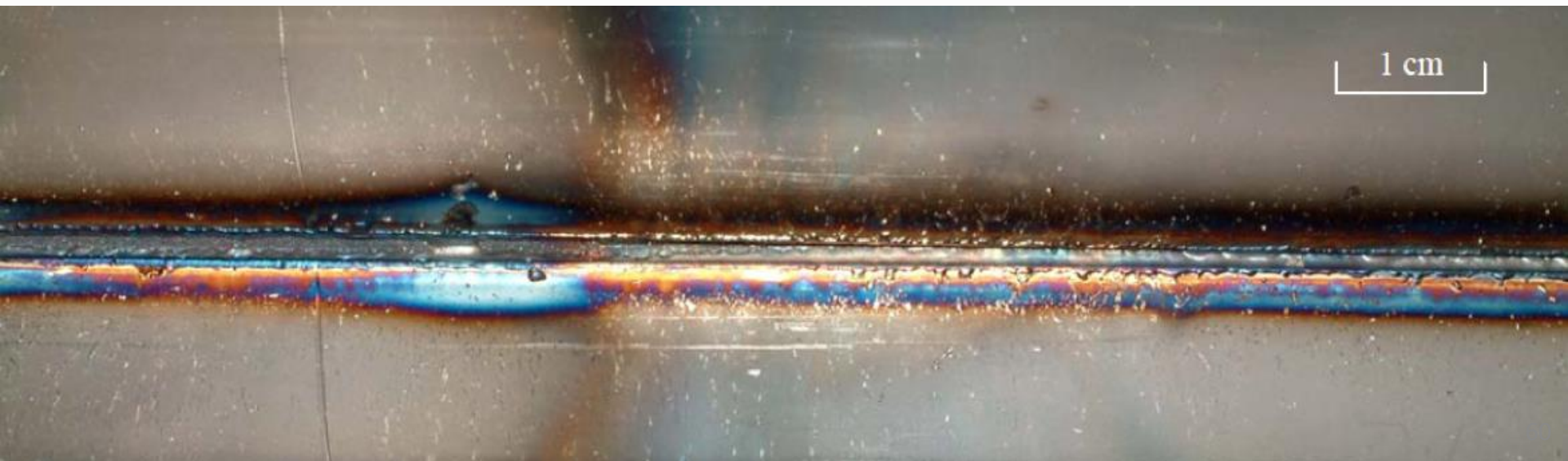
# Energy

4 TeV with 1380 bunches – 2012

- **~3.6 GJ** of energy stored in the main dipoles
- **140 MJ** stored in each beam ~21 kg of TNT.

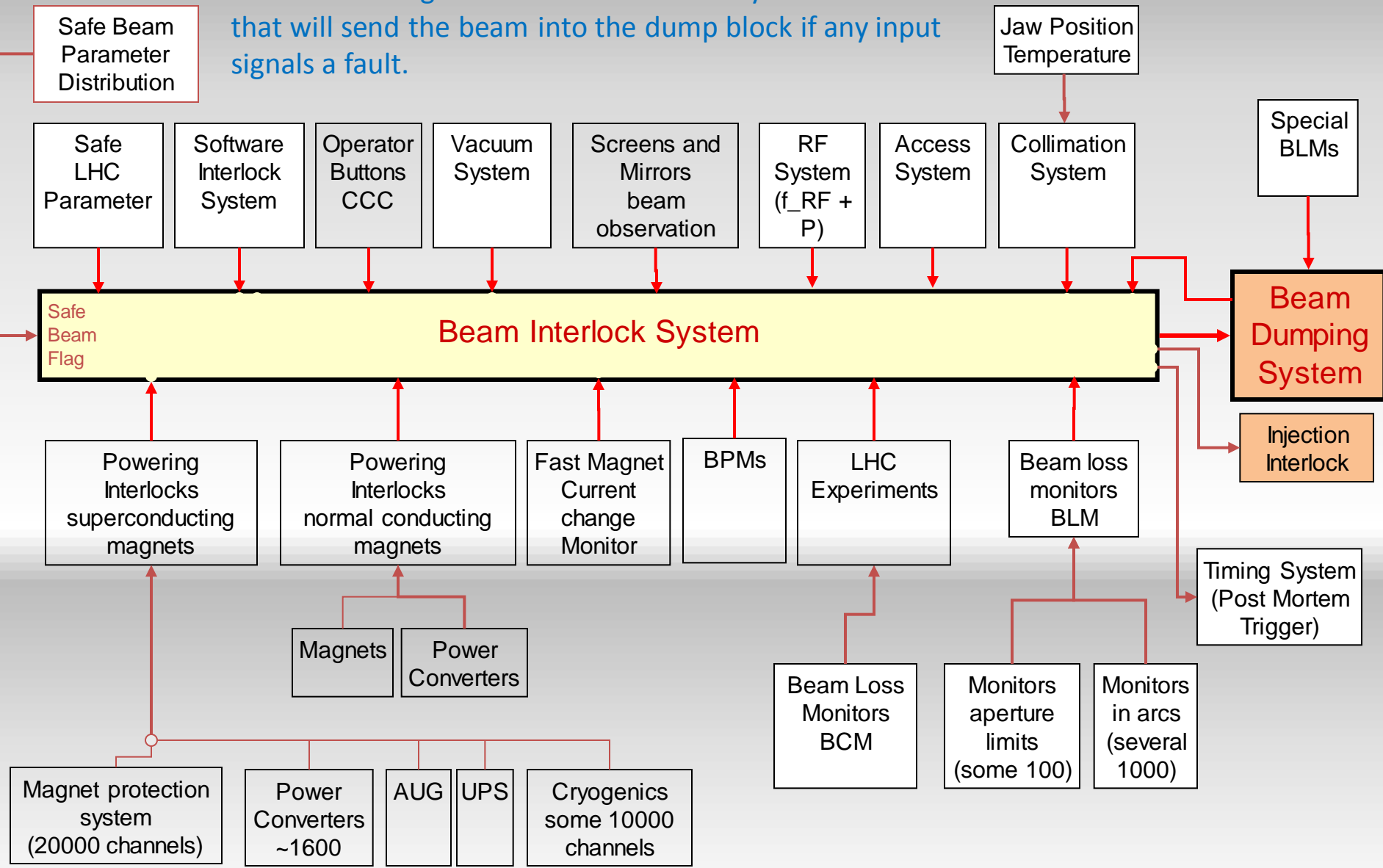
During an SPS extraction test in 2005...

The beam was a 450 GeV full LHC injection batch of  $3.4 \cdot 10^{13}$  p+ in 288 bunches **[2.5 MJ]**

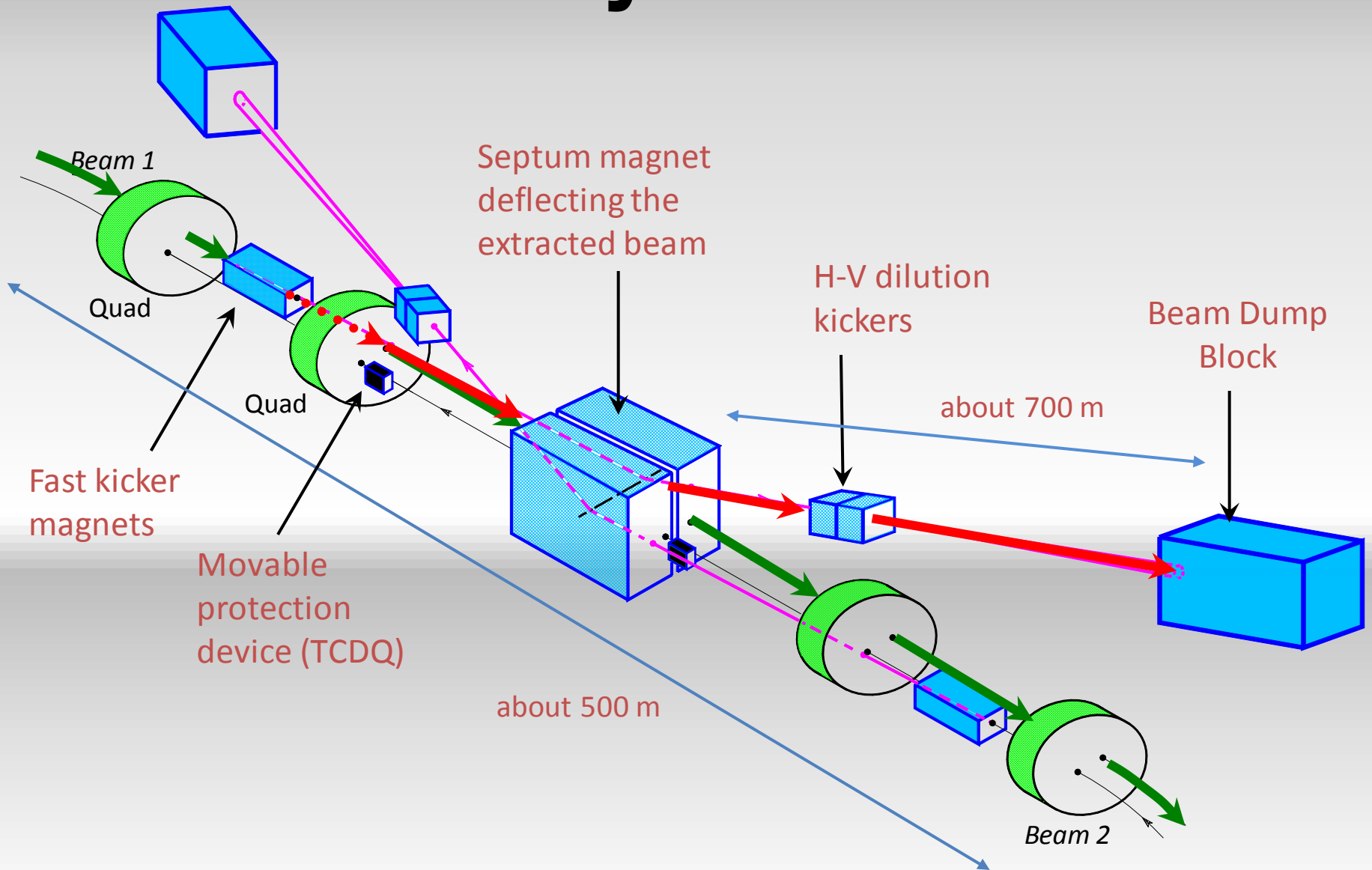


# Beam Interlock System

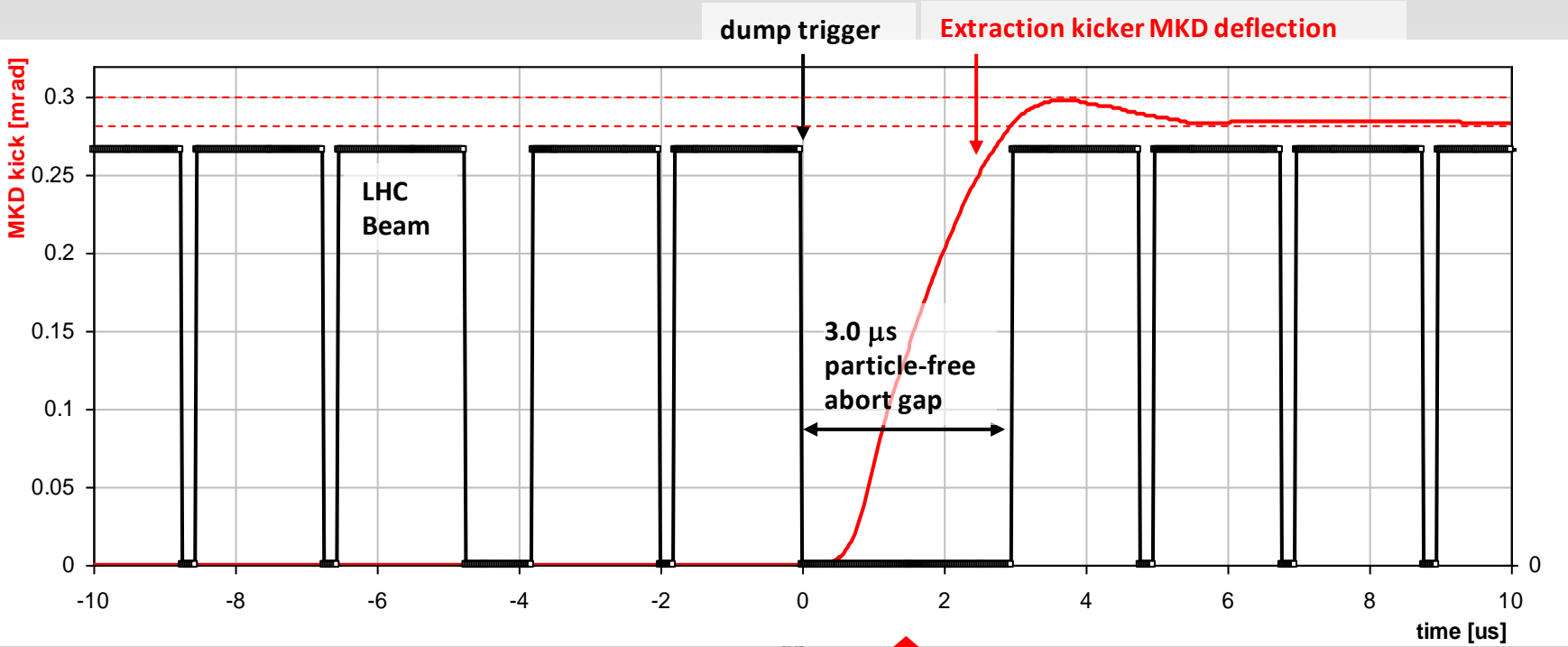
Over 10' 000 signals enter the interlock system of the LHC that will send the beam into the dump block if any input signals a fault.



# Layout of LHC beam dumping system



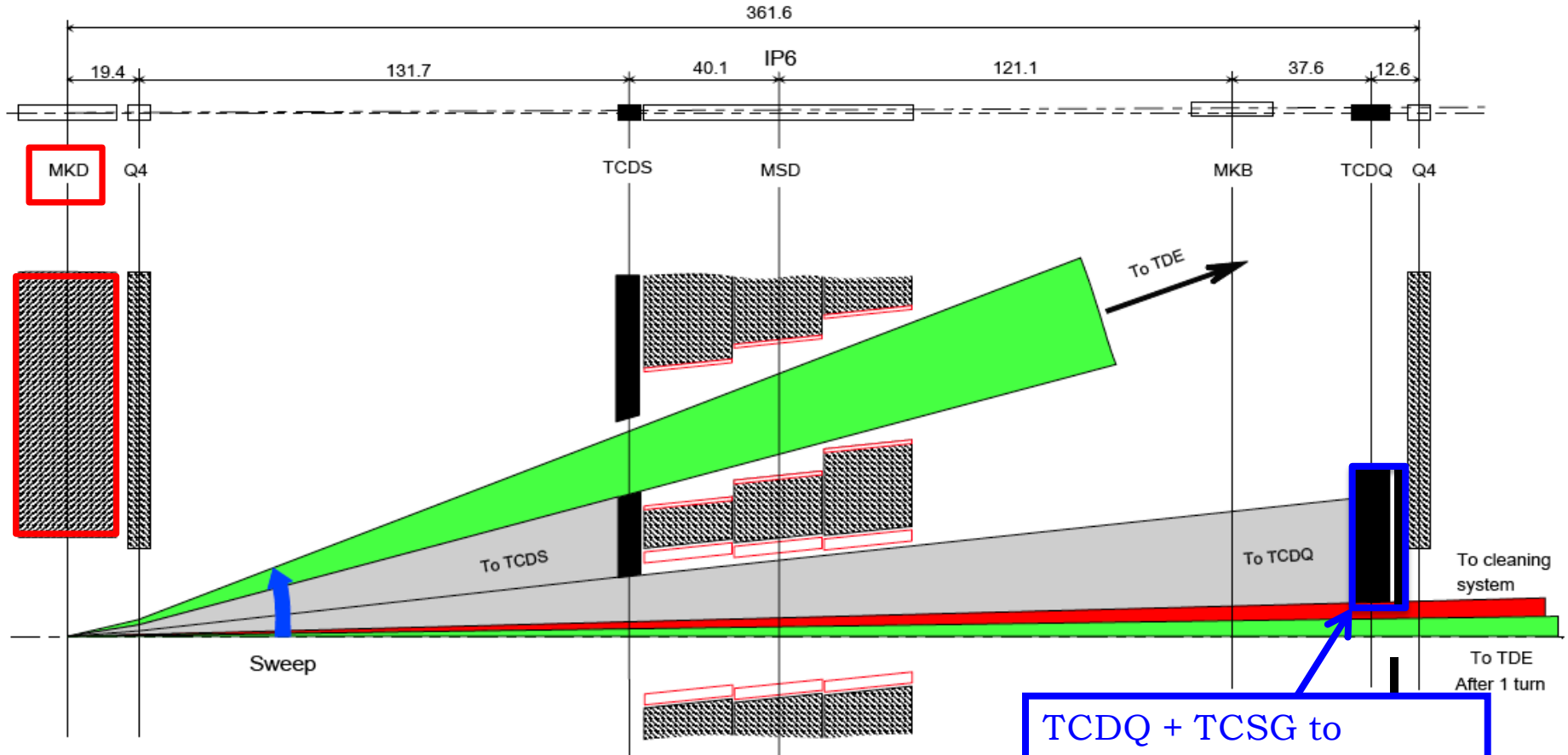
# Abort Gap



63



# Asynchronous Beam Dump



TCDQ = 6 m long CFC\* one-sided collimator  
 TCSG = 1 m long CFC\* two-sided collimator

\*CFC = carbon fibre compound

TCDQ + TCSG to  
 protect downstream  
 superconducting  
 magnets (Q4)

Estimated occurrence : at least once per year, 0 events up to now!



- **Two warm cleaning insertions:**

- ▶ IR3: momentum cleaning

- 1 Primary (H)
    - 4 Secondaries (H/S)
    - 4 Shower Abs. (H/V)

- ▶ IR7: betatron cleaning

- 3 Primaries (H/V/S)
    - 11 Secondaries (H/V/S)
    - 5 Shower Abs. (H/V)

- **Local cleaning at triplets**

- ▶ 8 tertiaries: 2 per IP per Beam

- **Physics debris absorption**

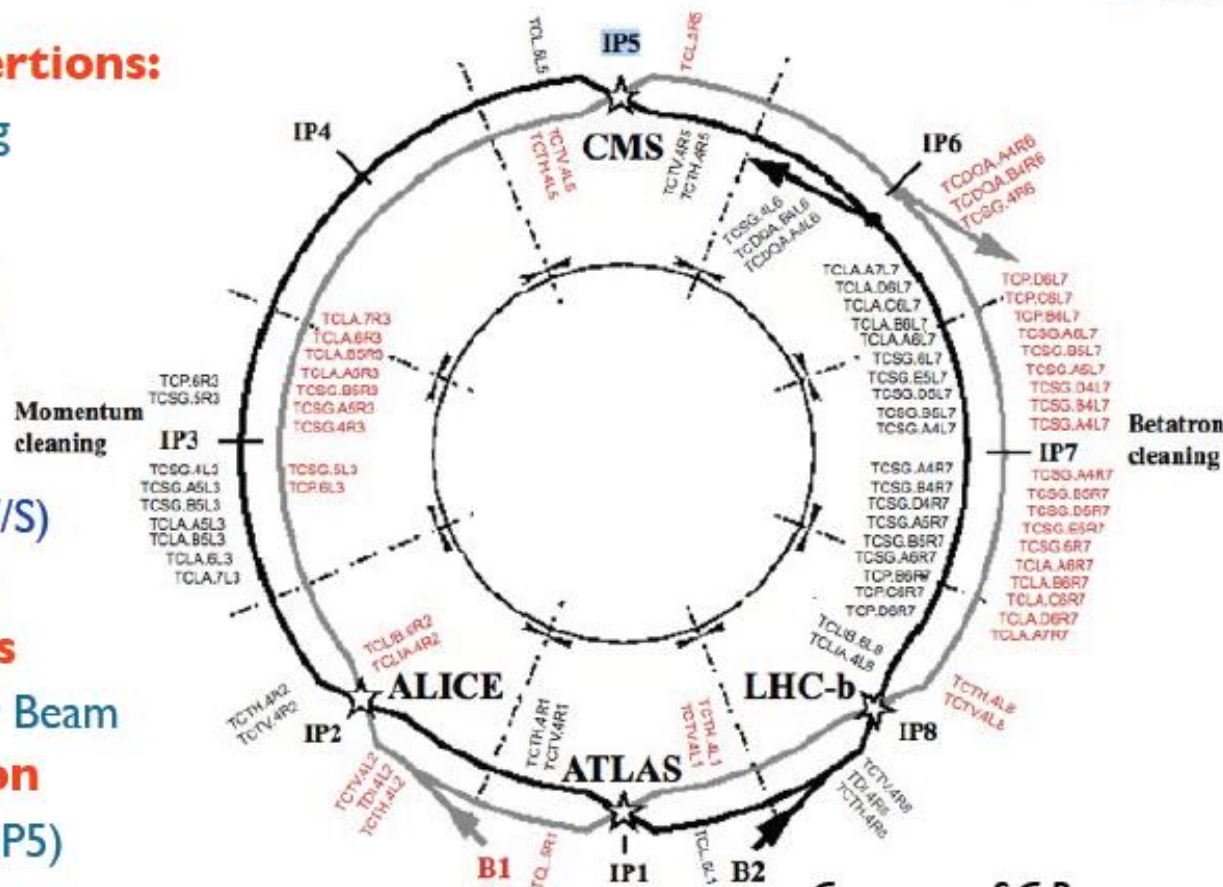
- ▶ 2 TCL (1 per beam IP1/IP5)

8 passive absorbers for warm magnets in IP3/IP7

Transfer lines (13 collimators)

Injection and dump protection (10 collimators)

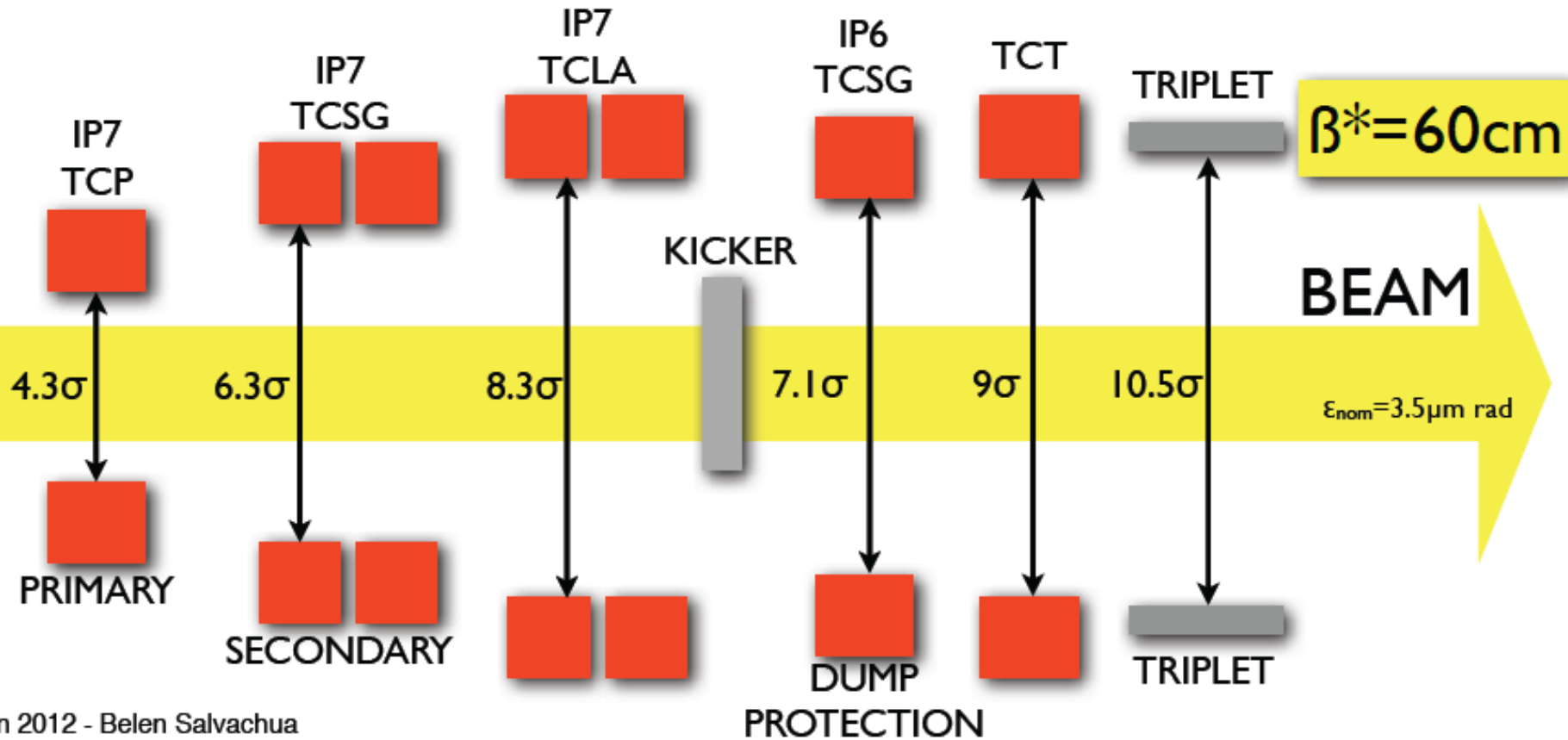
**Total of 108 collimators  
(100 movable)**



Courtesy of C.Bracco



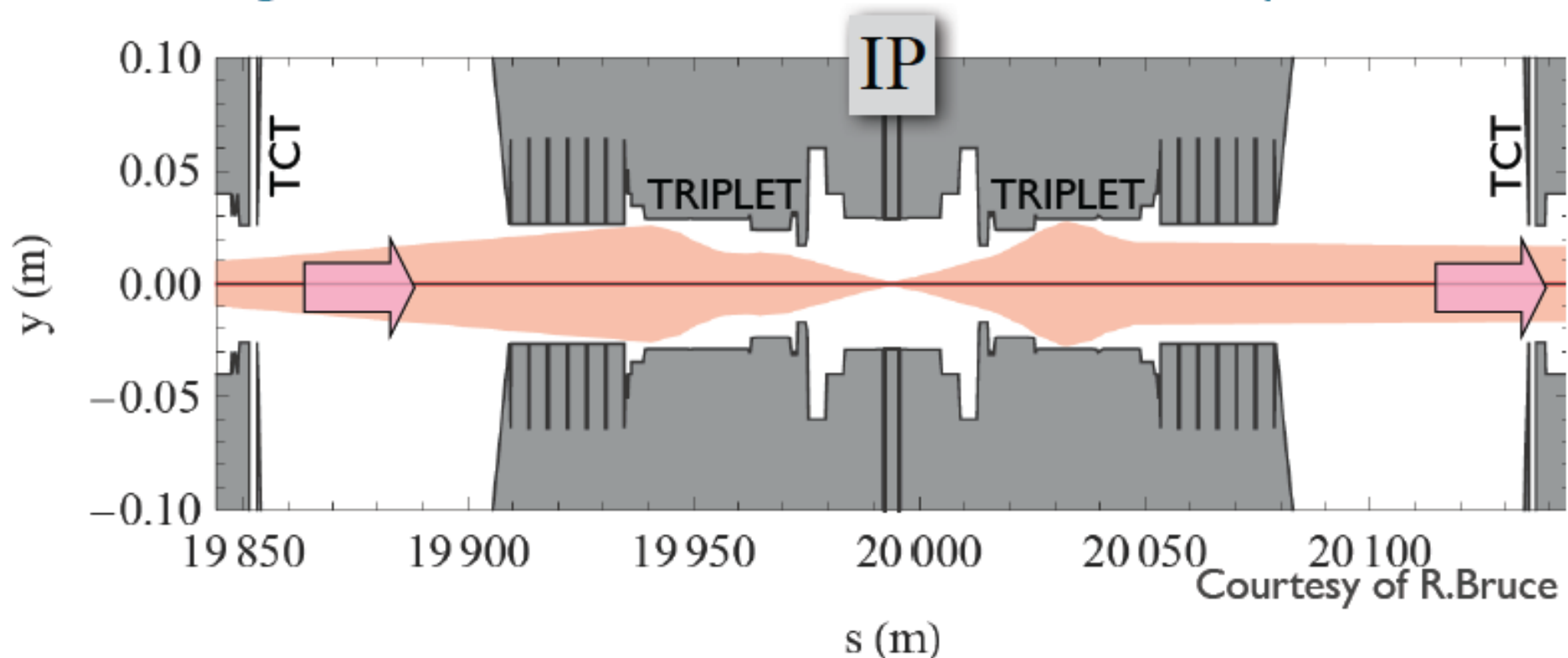
# Collimator hierarchy



Evian 2012 - Belen Salvachua

# Collimator hierarchy

- Normalized Triplet aperture decreases when reducing  $\beta^*$
- Triplet aperture MUST be protected by the tertiary collimators (TCTs)
- At the same time, TCTs must be shadowed by the dump protection
- Dump protection must be outside the primary and secondary collimators





# Collimator hierarchy

- The hierarchy must be respected at all times.
- The collimators and protection devices are positioned with respect to the closed orbit
- Therefore the closed orbit must be in tolerance at all times.
- This includes the ramp and squeeze.
  - Orbit feedback becomes mandatory
  - Interlocks on orbit position become mandatory



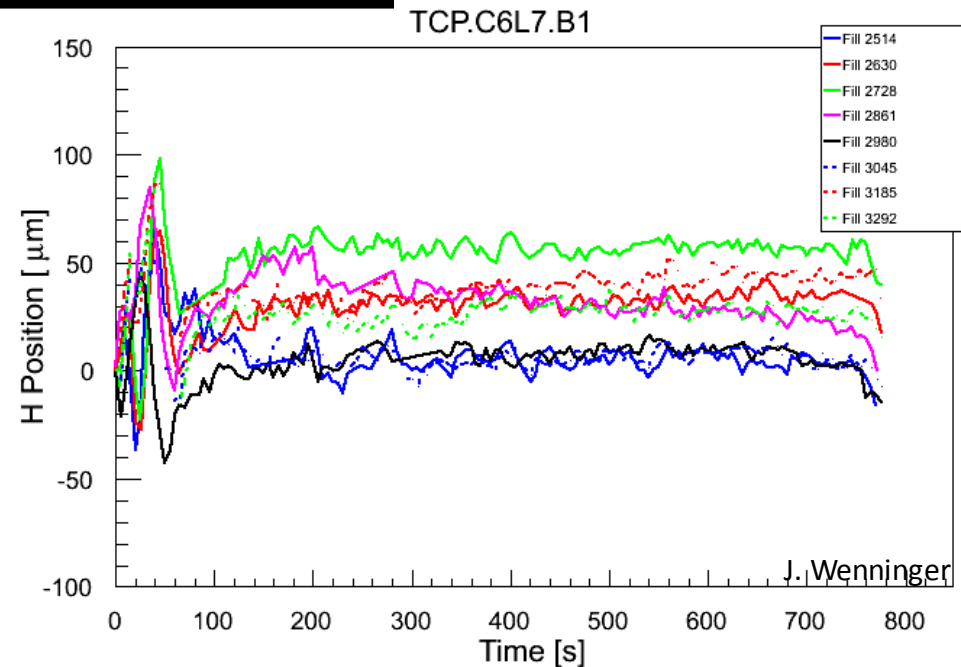
# Collimation/reproducibility

1.33	TCP.D6L7.B1	-0.84
1.33	TCP.C6L7.B1	-1.69
0.94	TCP.B6L7.B1	-1.61
1.85	TCSG.A6L7.B1	-2.01
1.92	TCSG.B5L7.B1	-2.66
2.1	TCSG.A5L7.B1	-2.58
1.42	TCSG.D4L7.B1	-1.55
2.98	TCSG.B4L7.B1	-1.29
2.93	TCSG.A4L7.B1	-1.27
2.8	TCSG.A4R7.B1	-1.4
2.78	TCSG.B5R7.B1	-2.02
2.22	TCSG.D5R7.B1	-2.66
2.48	TCSG.E5R7.B1	-2.39
3.08	TCSG.6R7.B1	-3.54
2	TCLA.A6R7.B1	-1.34
2.66	TCLA.B6R7.B1	-3.36
4.37	TCLA.C6R7.B1	-1.5
1.7	TCLA.D6R7.B1	-2.14
1.5	TCLA.A7R7.B1	-2.32

Orbit at primary

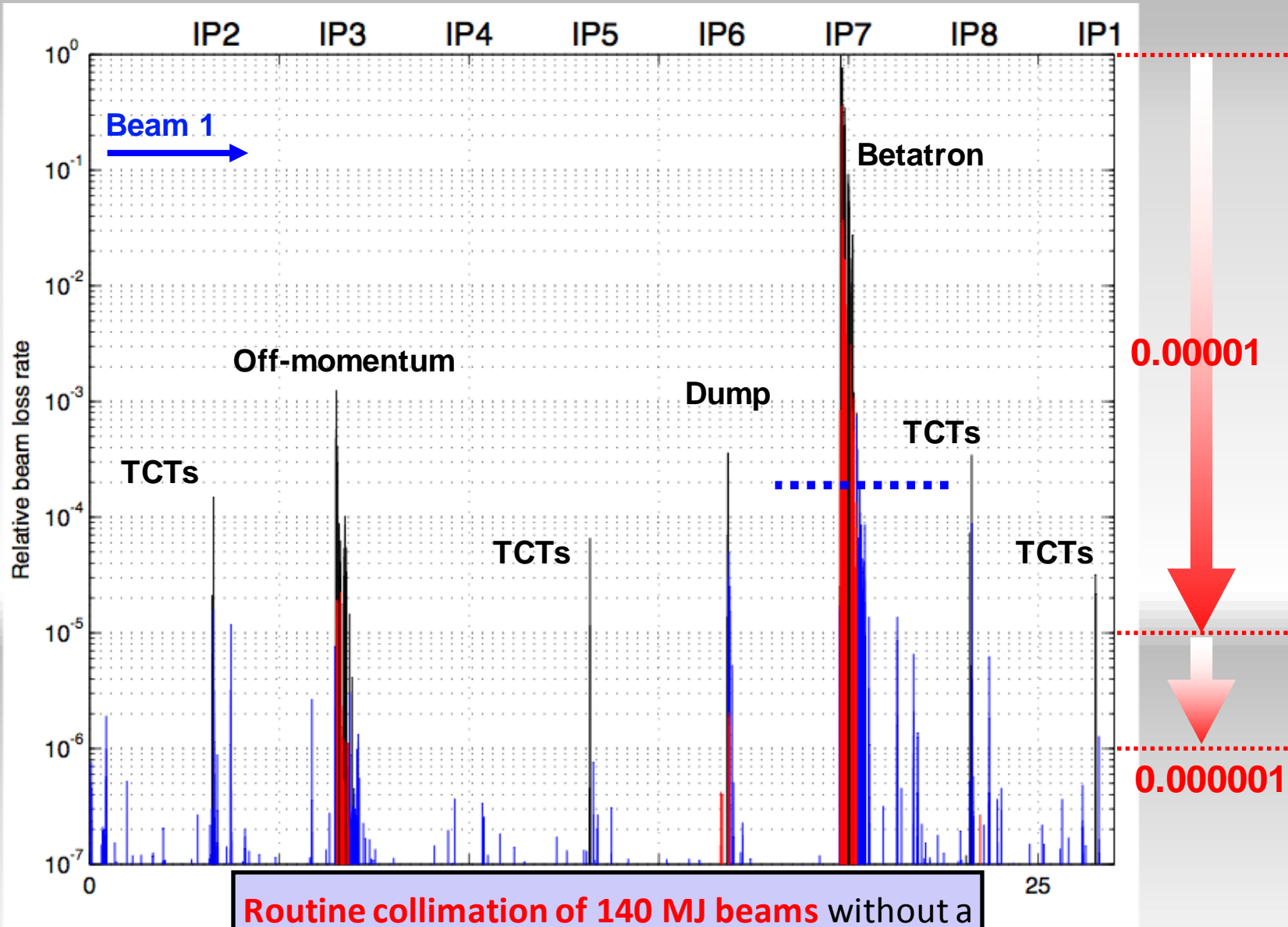
IR7 collimators – beam 1

2011-2012: only **ONE** full alignment in IR3/IR7



# Collimation

Generate higher loss rates: excite beam with transverse dampers



Routine collimation of 140 MJ beams without a single quench from stored beam