



ROADMAP OF DARK MATTER MODELS FOR RUN3

CERN WORKSHOP, 13-17 MAY 2024

Cosmological perspectives and constraints on t-channel models

Cosmology of t-channel models

- The work presented here is based on several contributions
 - T-channel white paper (cosmo section effort), together with M. Becker, E. Copello, M. Garny, J. Harz, J. Heisig, A. Ibarra, S. Khalil, M. Kirtiman, Y. Koay, L. Lopez Honorez, T. Murphy, L. Panizzi, D. Sengupta and S. Tentori
 - CA, B. Fuks, J. Heisig, M. Kraemer, L. Mantani, L. Panizzi, Phys. Rev. D 108 (2023) [arXiv:2307.10367 [hep-ph]]
 - CA, B. Fuks, L. Mantani, H. Meis, L. Panizzi, J. Salko, Phys. Lett. B 813 (2021) [arXiv:2010.136038 [hep-ph]]
 - CA, B. Fuks, L. Mantani, Eur. Phys J. C 80 (2020) [arXiv:2001.05024 [hep-ph]]

Cosmology of t-channel models

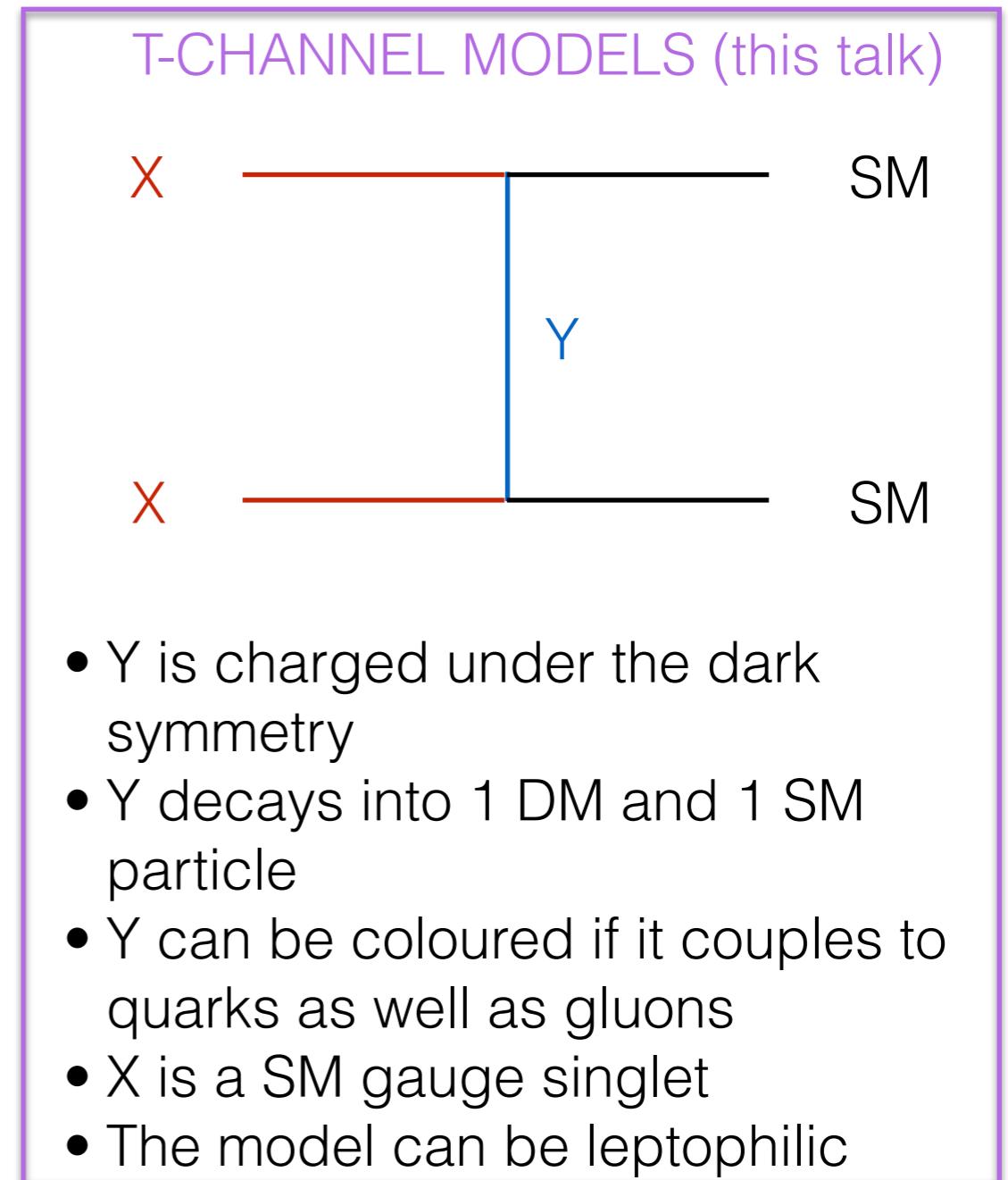
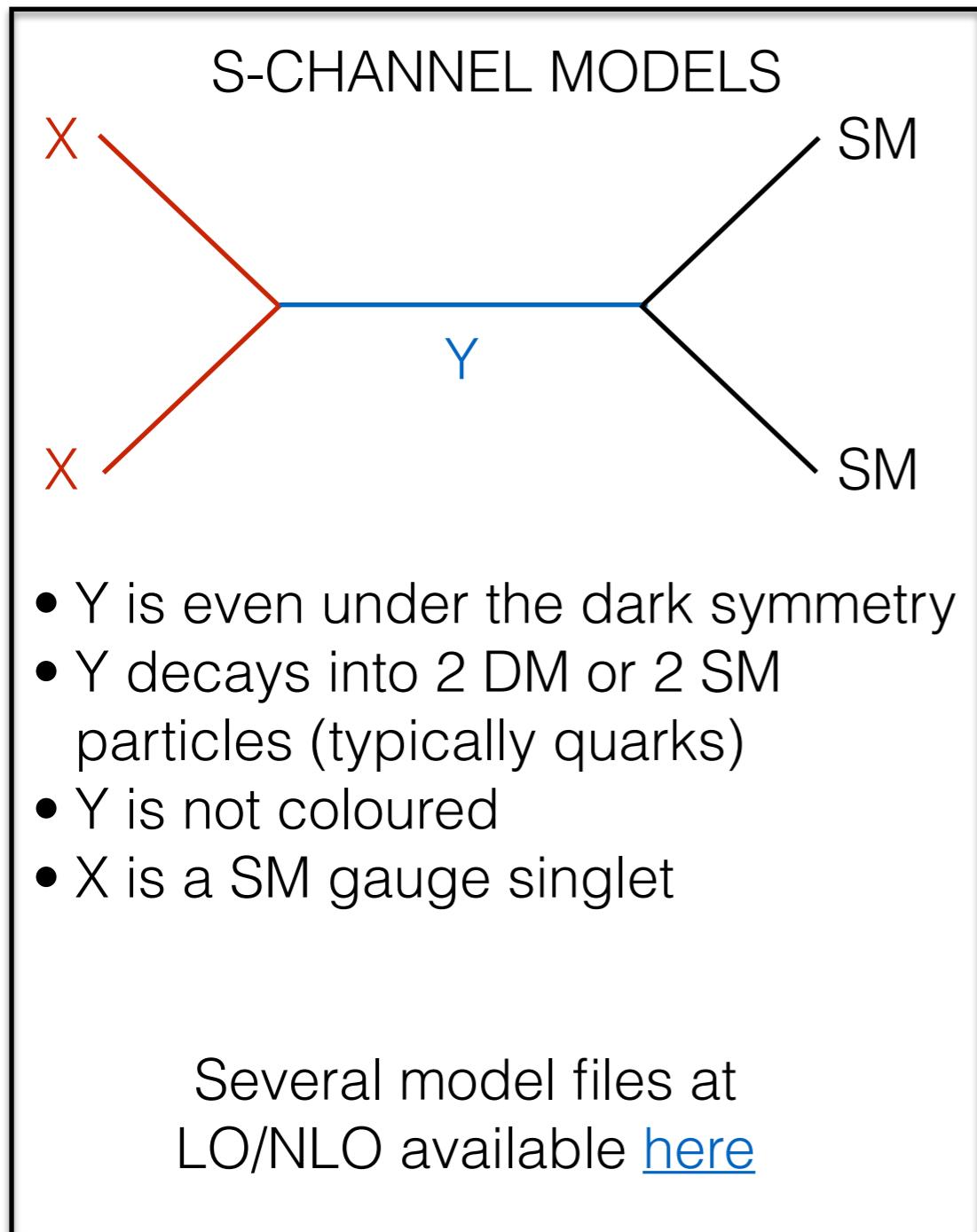
- Brief overview of the minimal simplified t-channel models
- Early universe physics and constraints on such models
- Today's probes from astroparticle and dark matter experiments
- Complementarity with collider searches for few selected benchmarks

HERE CONSIDERED ONLY COLORED MEDIATORS COUPLING TO QUARKS

UNDERLINE THE DIFFERENCES WHEN RELEVANT THAT DEPEND ON THE QUARK GENERATION

**HOWEVER T-CHANNEL CAN APPEAR UNDER MANY FORMS (SEE TALKS THIS MORNING) AND CAN BE
LEPTOPHILIC AS WELL**

Minimal t-channel model



Minimal t-channel model

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{kin}} + \mathcal{L}_F(\chi) + \mathcal{L}_F(\tilde{\chi}) + \mathcal{L}_S(S) + \mathcal{L}_S(\tilde{S}) + \mathcal{L}_V(V) + \mathcal{L}_V(\tilde{V})$$

VERY GENERIC MODEL WITH 6 DARK MATTER CANDIDATES AND 24 MEDIATORS OF DIFFERENT SPIN

Field	Spin	Repr.	Self-conj.
\tilde{S}	0	(1, 1, 0)	yes
S	0	(1, 1, 0)	no
$\tilde{\chi}$	1/2	(1, 1, 0)	yes
χ	1/2	(1, 1, 0)	no
\tilde{V}_μ	1	(1, 1, 0)	yes
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$$\mathcal{L}_F(X) = \left[\lambda_Q \bar{X} Q \varphi_Q^\dagger + \lambda_u \bar{X} u \varphi_u^\dagger + \lambda_d \bar{X} d \varphi_d^\dagger + \text{h.c.} \right]$$

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$$\mathcal{L}_V(X) = \left[\hat{\lambda}_Q \bar{\psi}_Q \not{X} Q + \hat{\lambda}_u \bar{\psi}_u \not{X} u + \hat{\lambda}_d \bar{\psi}_d \not{X} d + \text{h.c.} \right]$$

X

Y

$\varphi_Q = \begin{pmatrix} \varphi_Q^{(u)} \\ \varphi_Q^{(d)} \end{pmatrix}$	0	(3, 2, $\frac{1}{6}$)	no
φ_u	0	(3, 1, $\frac{2}{3}$)	no
φ_d	0	(3, 1, $-\frac{1}{3}$)	no

**UNIQUE IMPLEMENTATION FOR
COLLIDER AND DM STUDIES**

**MODEL FILES AND DOCUMENTATION ARE AVAILABLE HERE:
[HTTP://FEYNRULES.IRMP.UCL.AC.BE/WIKI/DMSIMPT](http://FEYNRULES.IRMP.UCL.AC.BE/WIKI/DMSIMPT)**

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New couplings
3x3 matrices in flavour space
real and flavour diagonal

**UNIQUE IMPLEMENTATION FOR
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Minimal t-channel model

Name	DM	Mediators	Parameters
S3M_uni	$\tilde{\chi}$	$\varphi_{Q_f}, \varphi_{u_f}, \varphi_{d_f}$	
S3D_uni	χ		
S3M_3rd	$\tilde{\chi}$	$\varphi_{Q_3}, \varphi_{u_3}, \varphi_{d_3}$	$M_\varphi, M_\chi, \lambda_\varphi$
S3D_3rd	χ		
S3M_uR	$\tilde{\chi}$	φ_{u_1}	
S3D_uR	χ		
F3S_uni	\tilde{S}	$\psi_{Q_f}, \psi_{u_f}, \psi_{d_f}$	
F3C_uni	S		
F3S_3rd	\tilde{S}	$\psi_{Q_3}, \psi_{u_3}, \psi_{d_3}$	$M_S, M_\psi, \hat{\lambda}_\psi$
F3C_3rd	S		
F3S_uR	\tilde{S}	ψ_{u_1}	
F3C_uR	S		
F3V_uni	\tilde{V}_μ	$\psi_{Q_f}, \psi_{u_f}, \psi_{d_f}$	
F3W_uni	V_μ		
F3V_3rd	\tilde{V}_μ	$\psi_{Q_3}, \psi_{u_3}, \psi_{d_3}$	$M_V, M_\psi, \hat{\lambda}_\psi$
F3W_3rd	V_μ		
F3V_uR	\tilde{V}_μ	ψ_{u_1}	
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THE GENERIC MODEL HAS SEVERAL RESTRICTIONS WHERE THE UNDESIRED FIELDS ARE DECOUPLED AND INTERACTIONS ARE SET TO ZERO

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Any restriction has
3 free model parameters:
DM and mediator masses + coupling
(M_χ, M_Y, λ)

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coupling only to quark up-right

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coupling only to quark up-right

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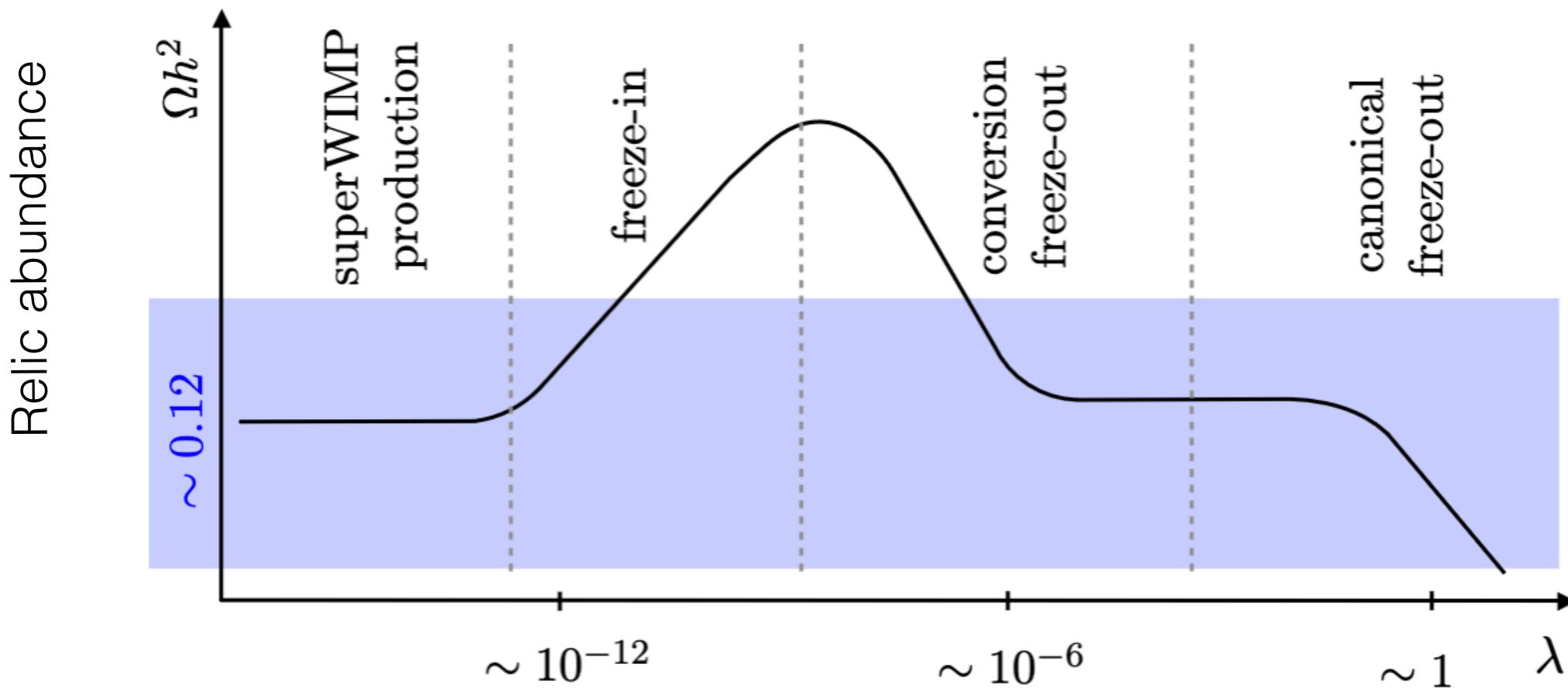
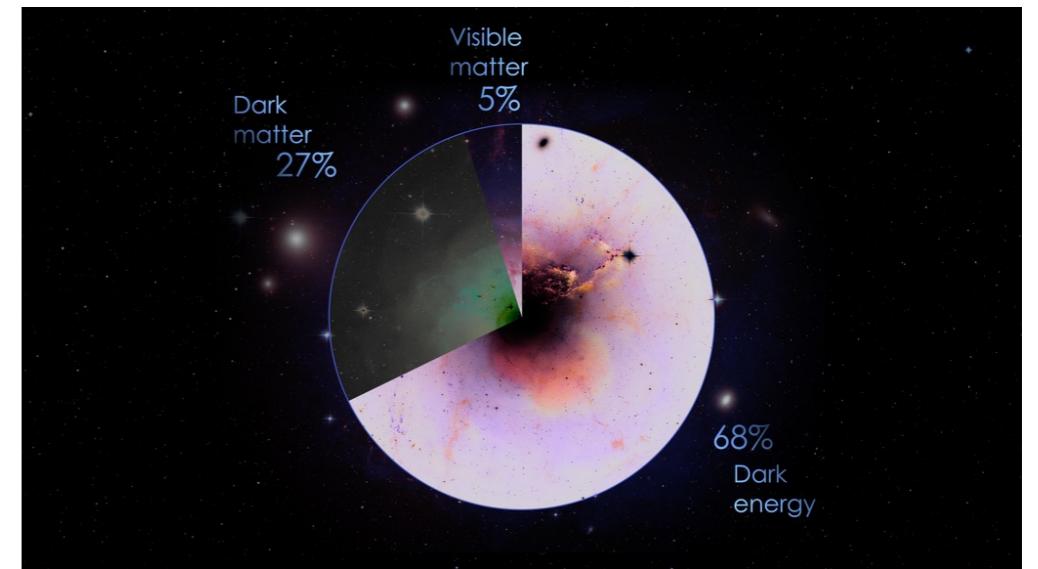
coupling to all quarks

coupling only to b and t quarks

coupling only to quark up-right

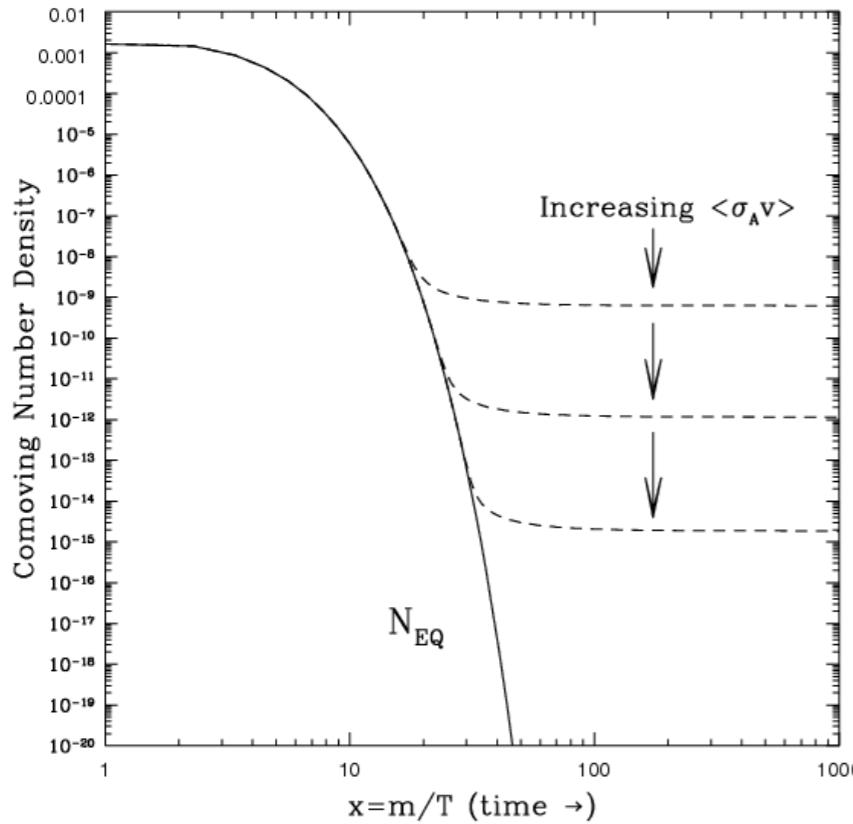
Early universe

**$\Omega h^2 = 0.12$ PLANCK SATELLITE
[ARXIV:1807.06209]**

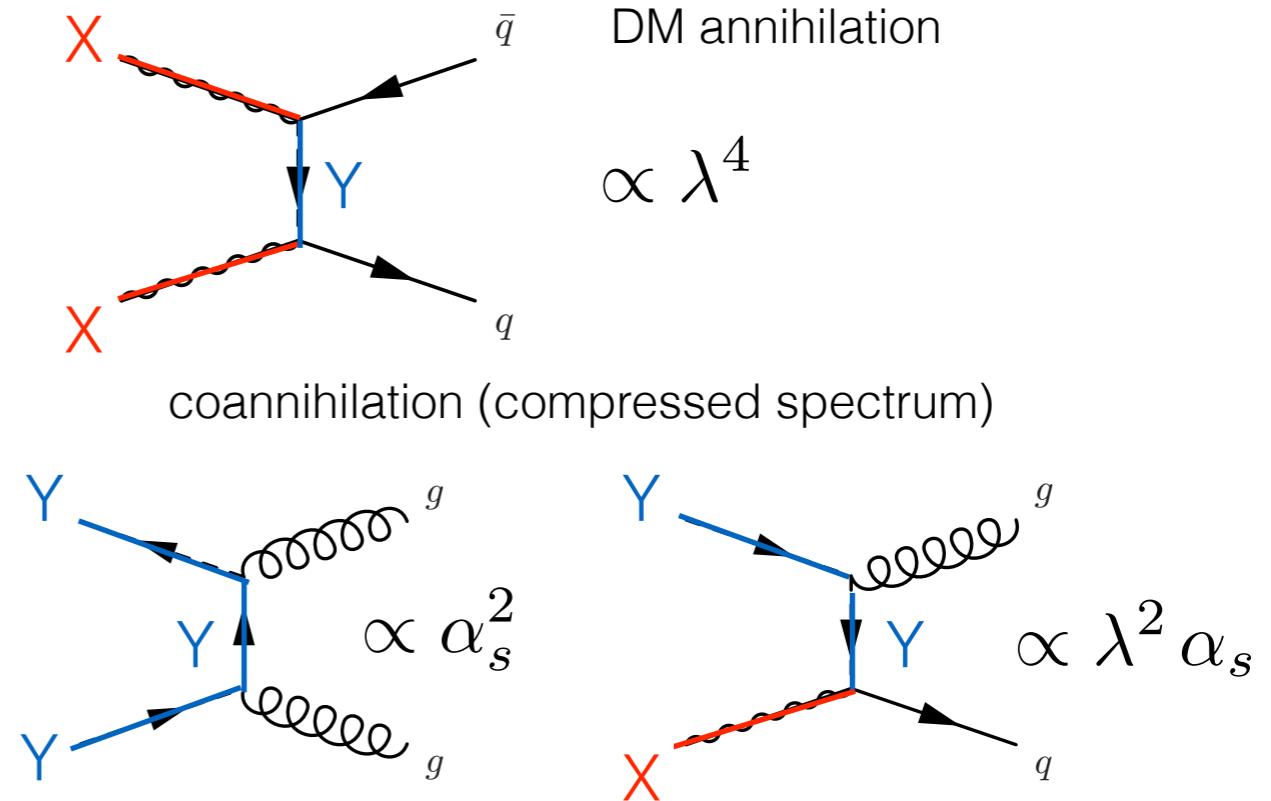


Early universe: freeze-out

TYPICALLY LO PROCESSES ARE DOMINANT

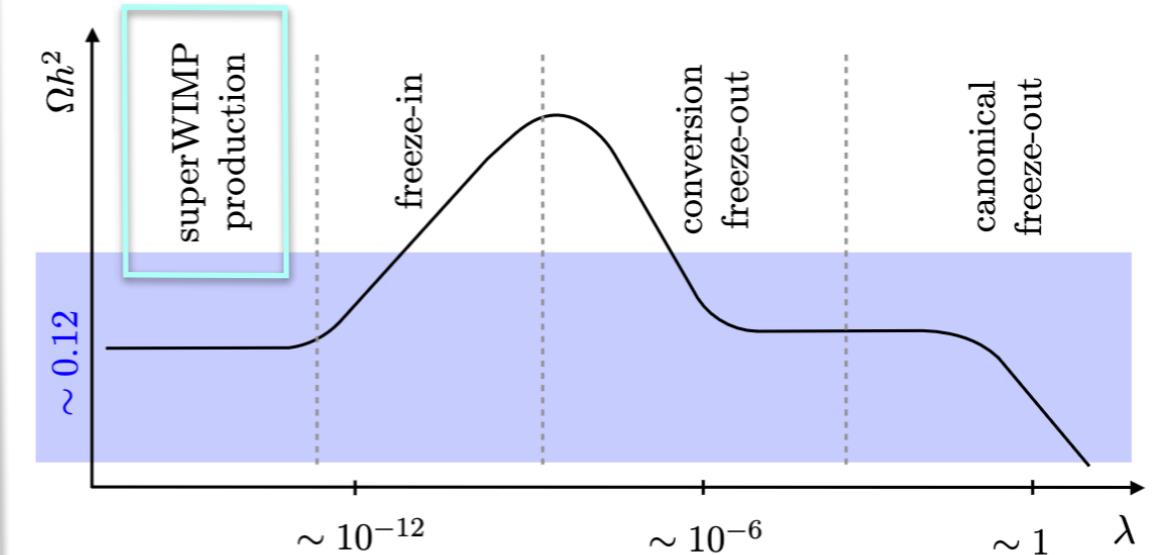
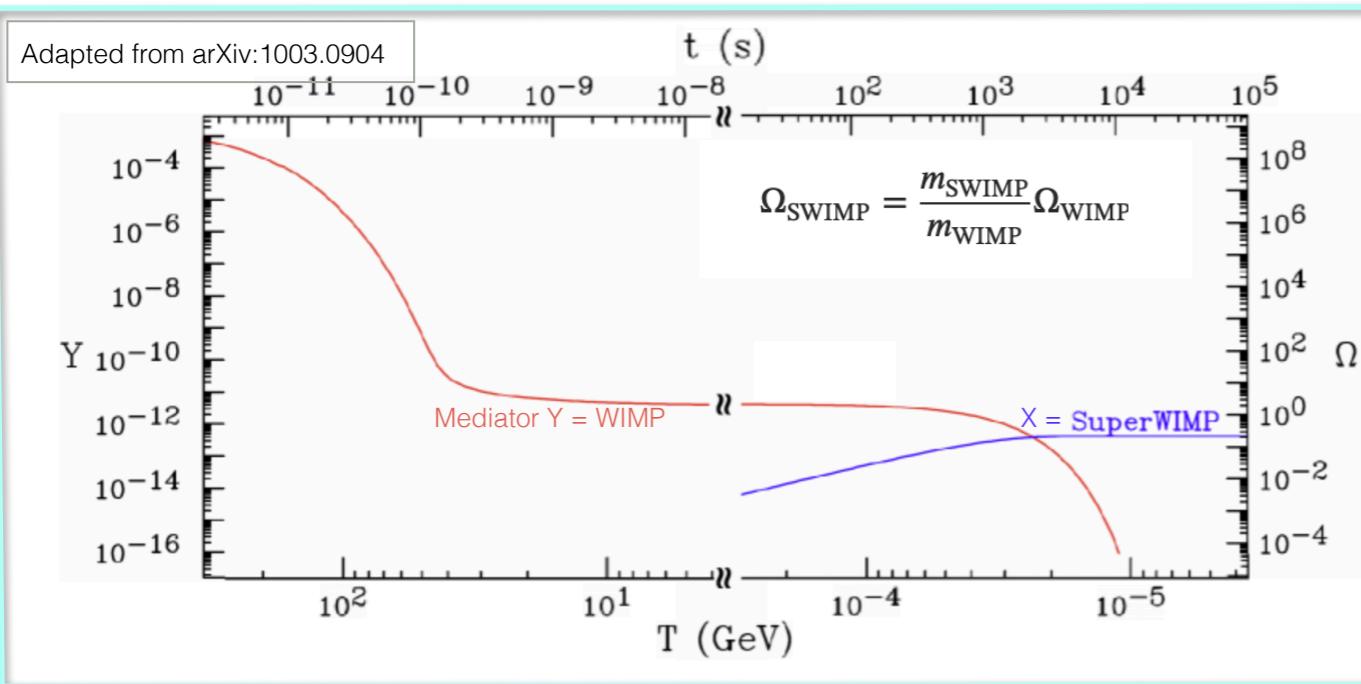


LARGE COUPLING = WIMP
DETECTABLE SIGNATURES IN
DIRECT, INDIRECT SEARCH
EXPERIMENT FOR DARK MATTER
AND AT COLLIDERS

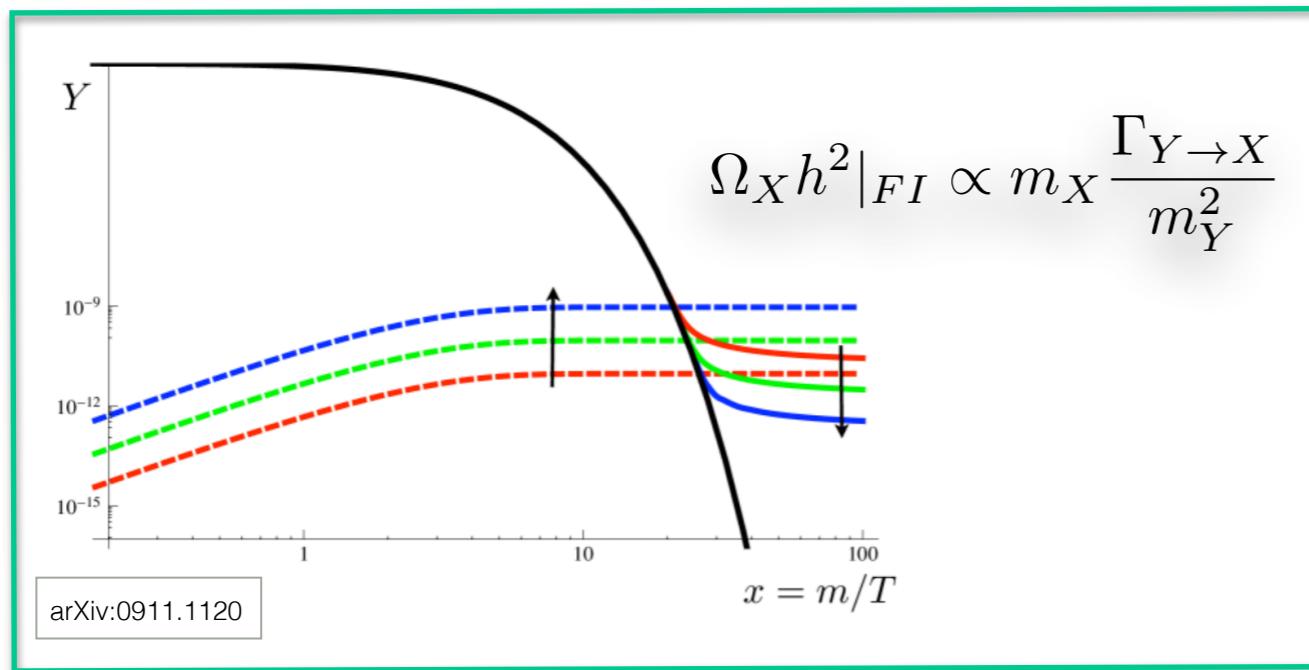
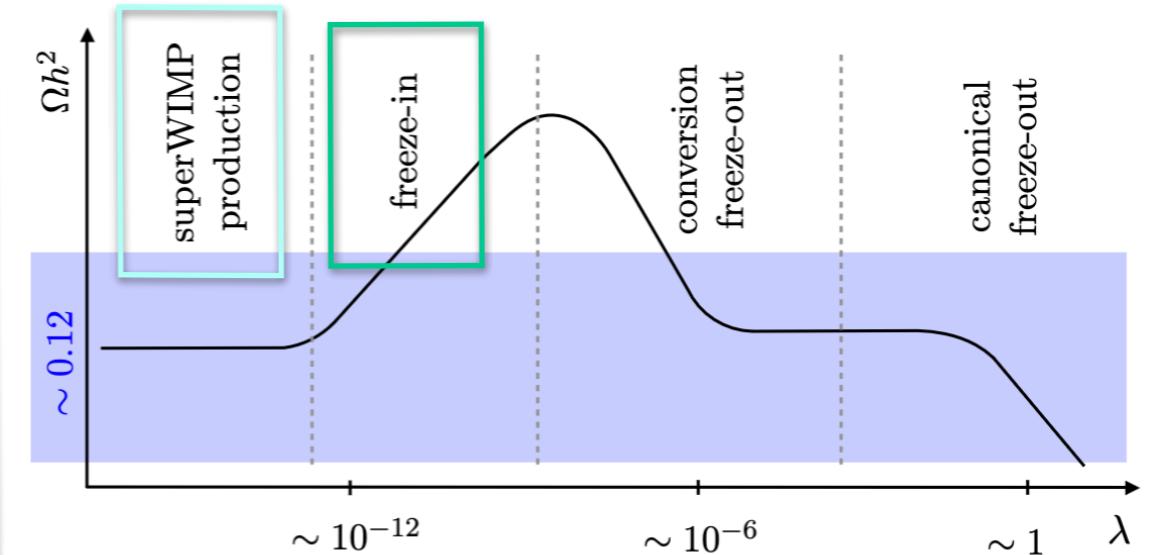
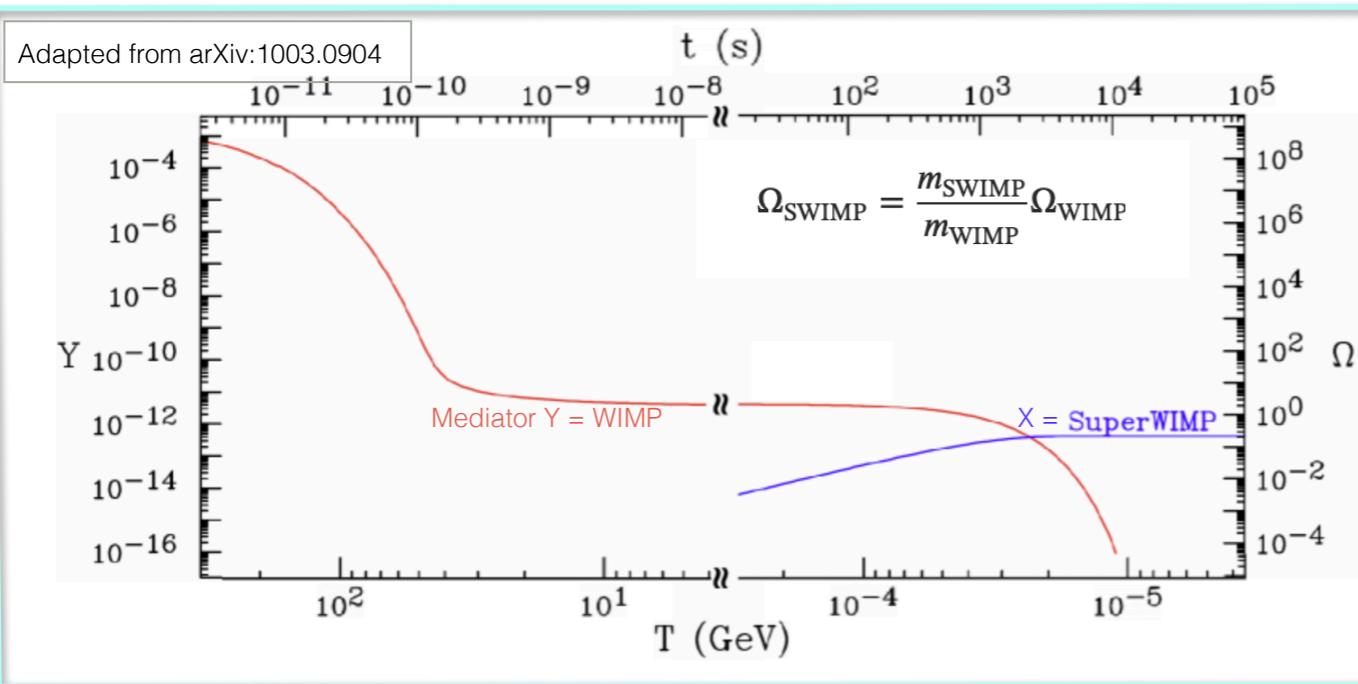


- Sommerfeld enhancement and bounds state formation can be relevant when mediators are light ($YY \rightarrow qq$, $YY \rightarrow gg$)
- For d-wave suppressed annihilation cross-sections NLO corrections are relevant
- Pheno is the basically the same for all quark flavours (threshold effects for heavy quarks)

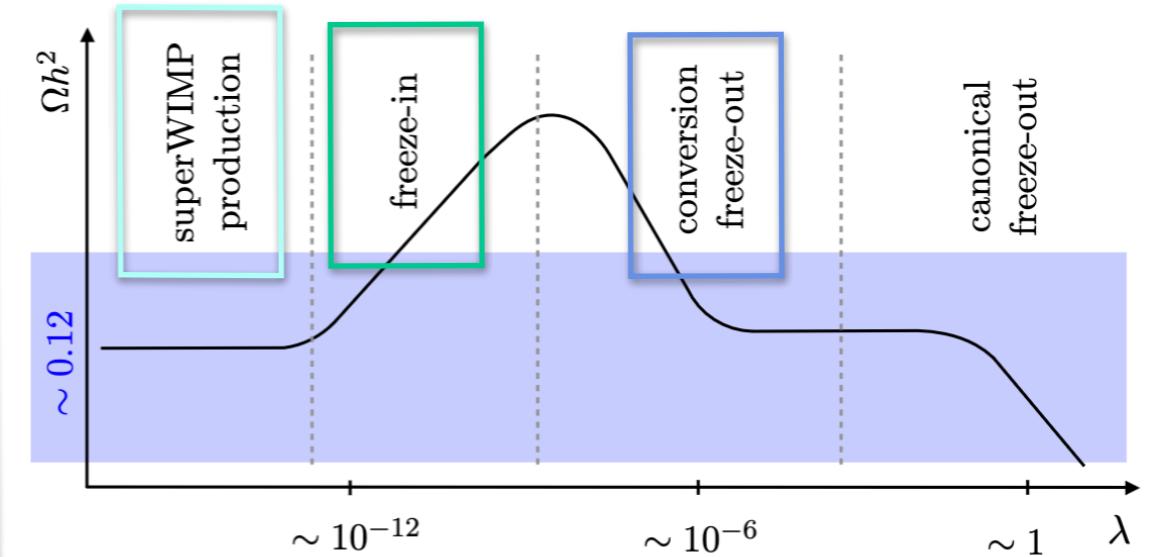
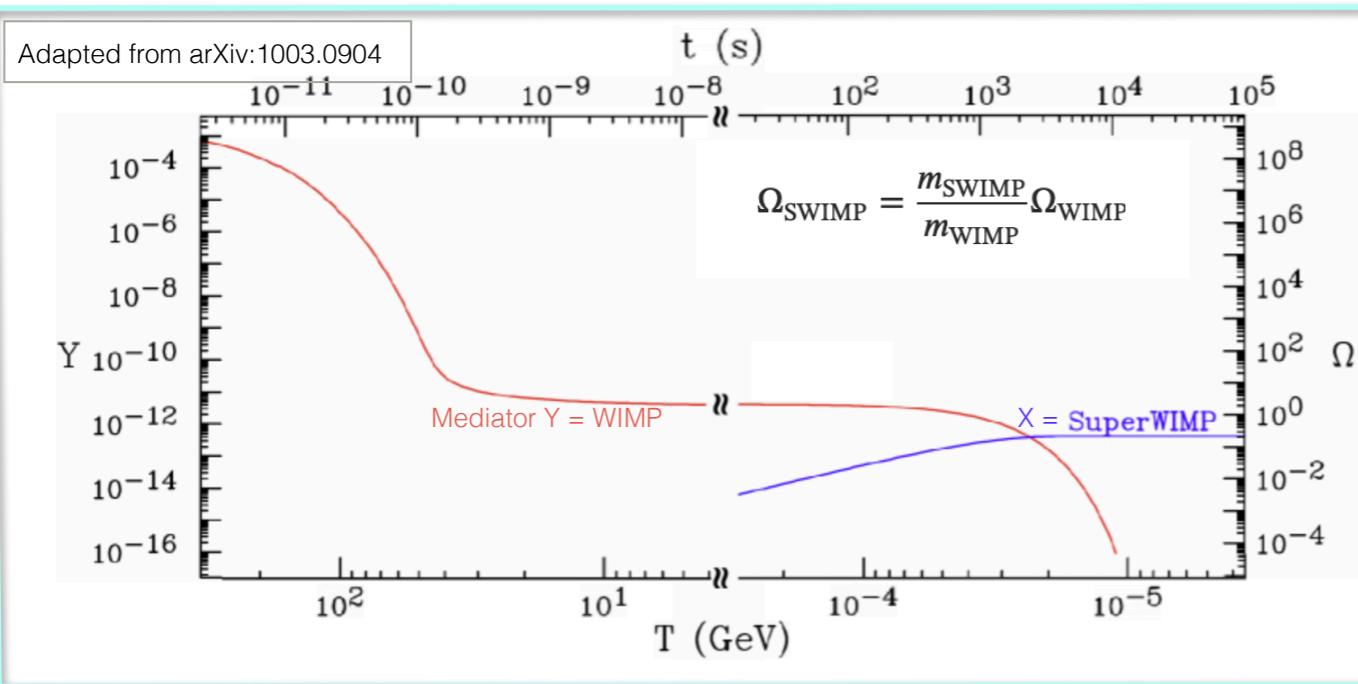
Early universe: small couplings (FIMPs)



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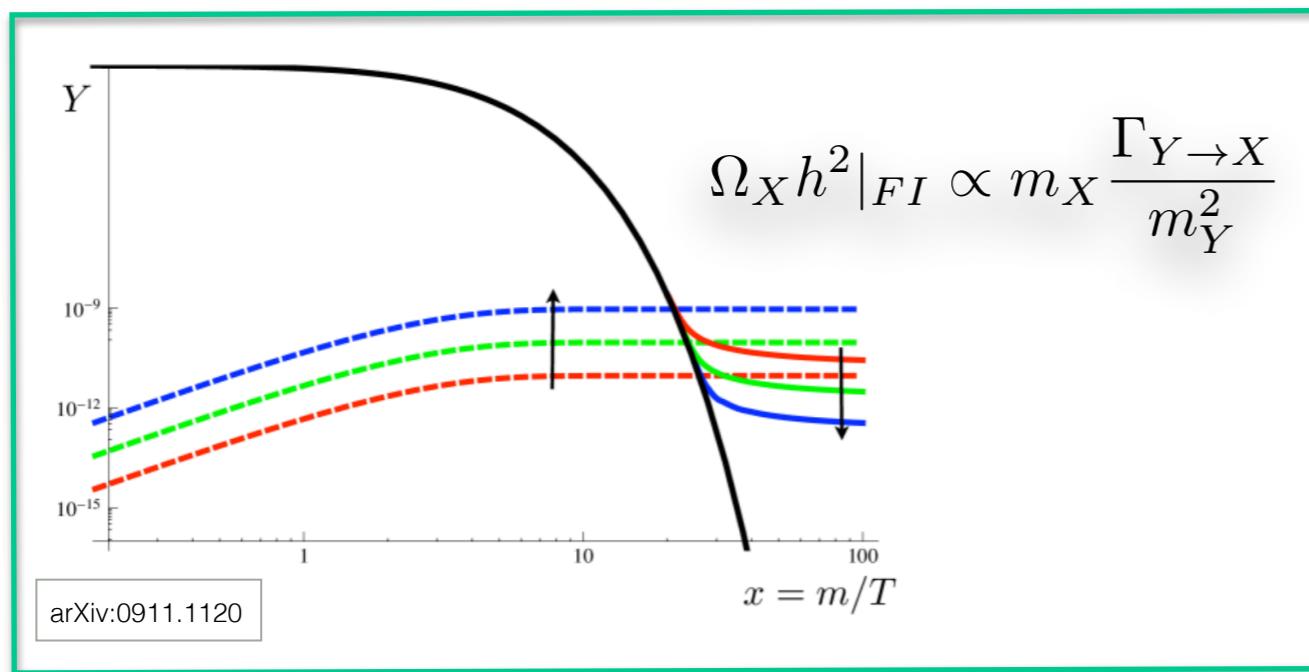


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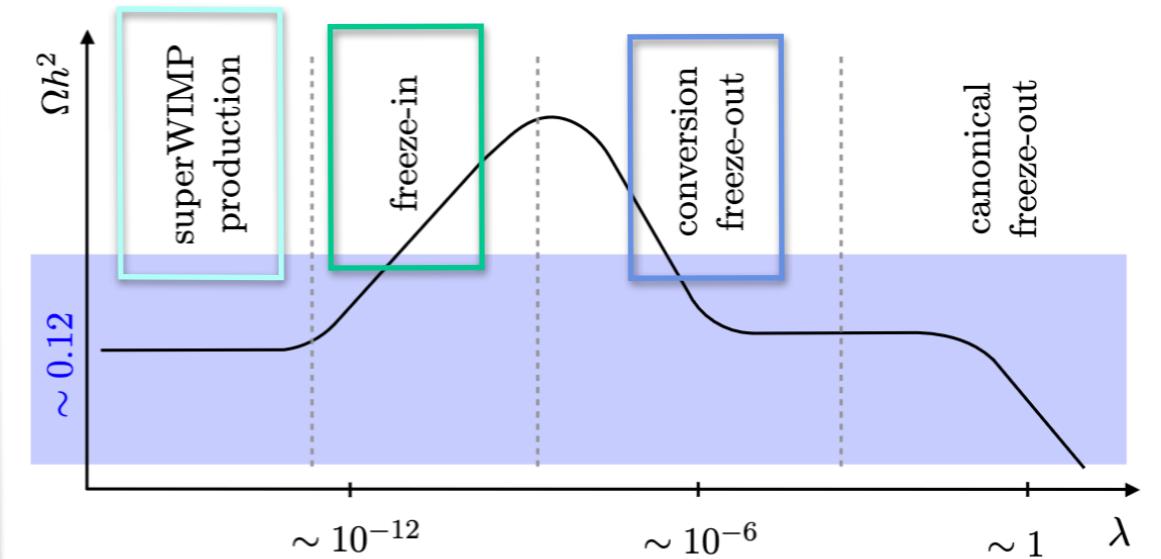
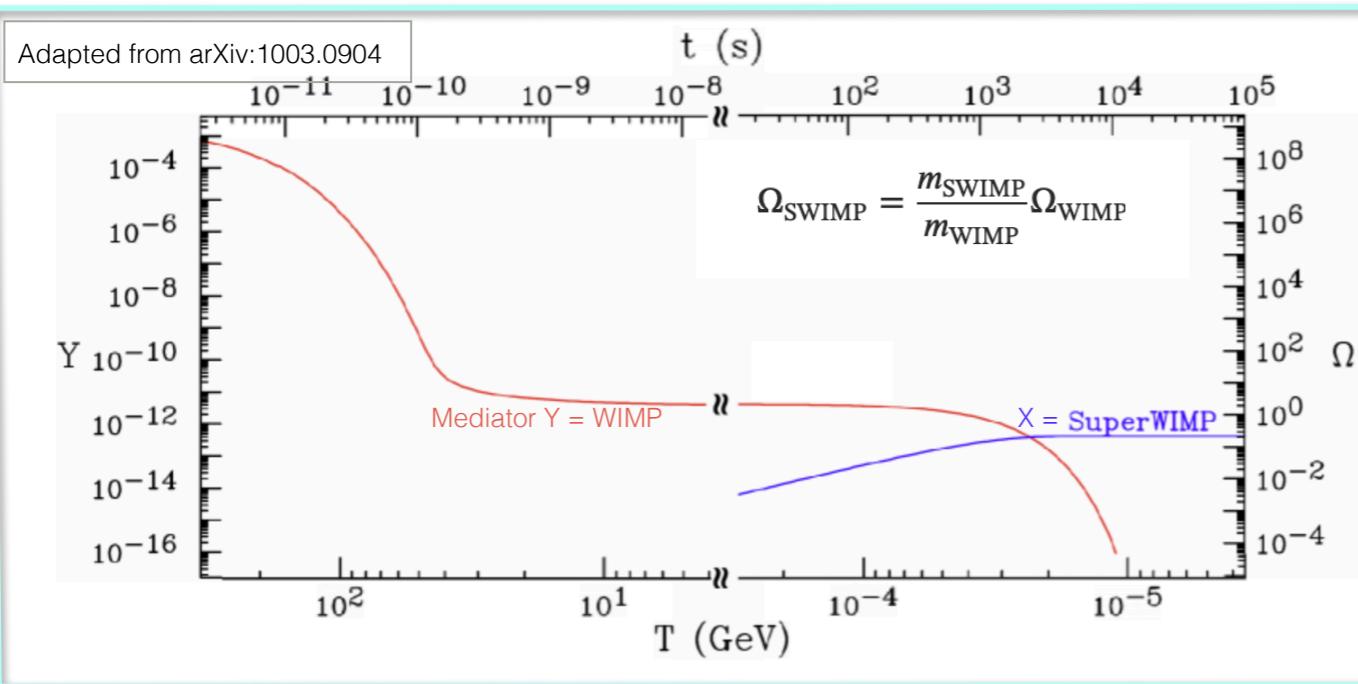


- Annihilation $XX, YX, YY \rightarrow AB$ becomes inefficient
- Conversions such as $Y \rightarrow XA, YA \rightarrow XB$ ($A, B = \text{SM}$) lead the freeze-out process

Much more in J. Heisig talk

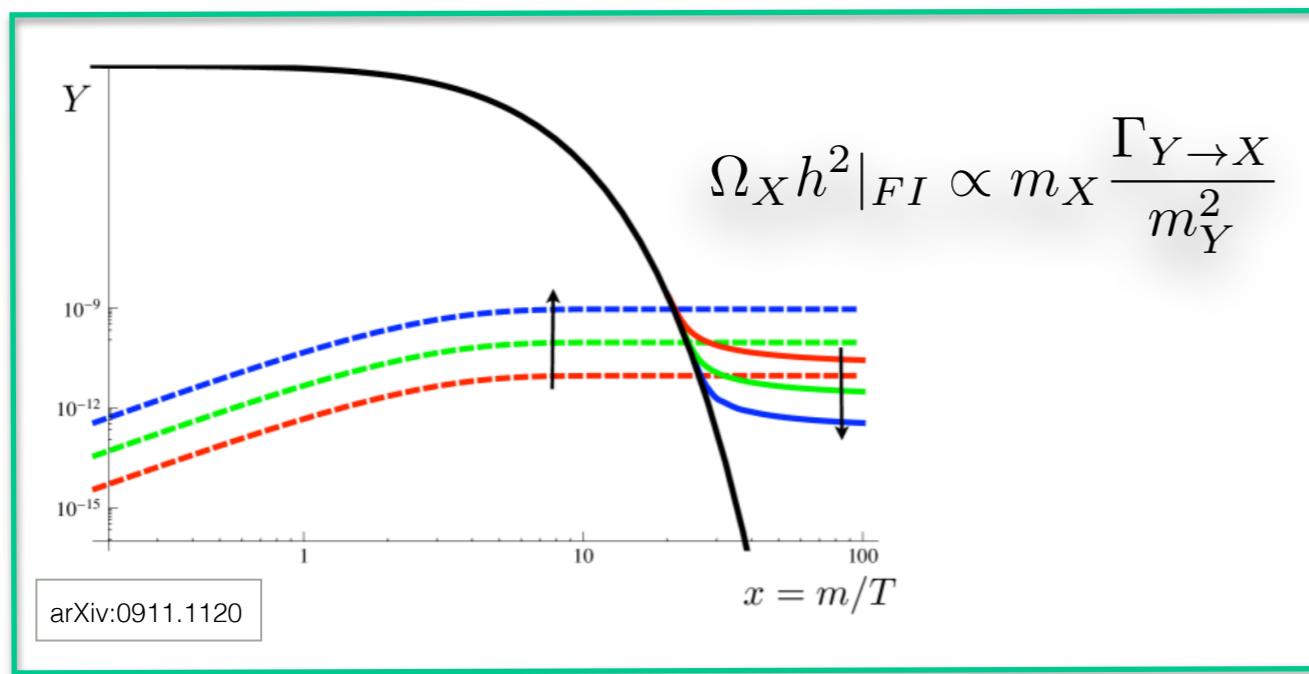


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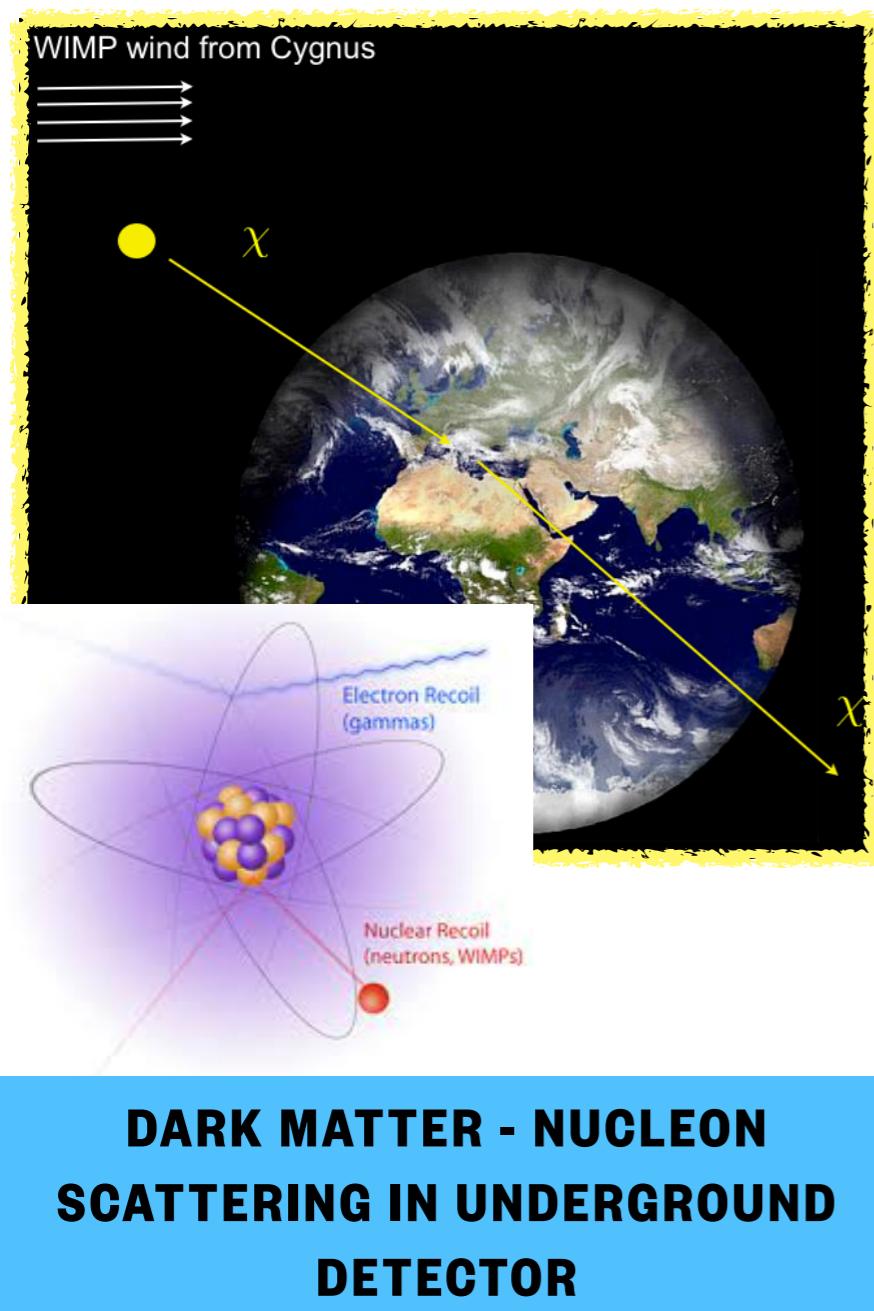
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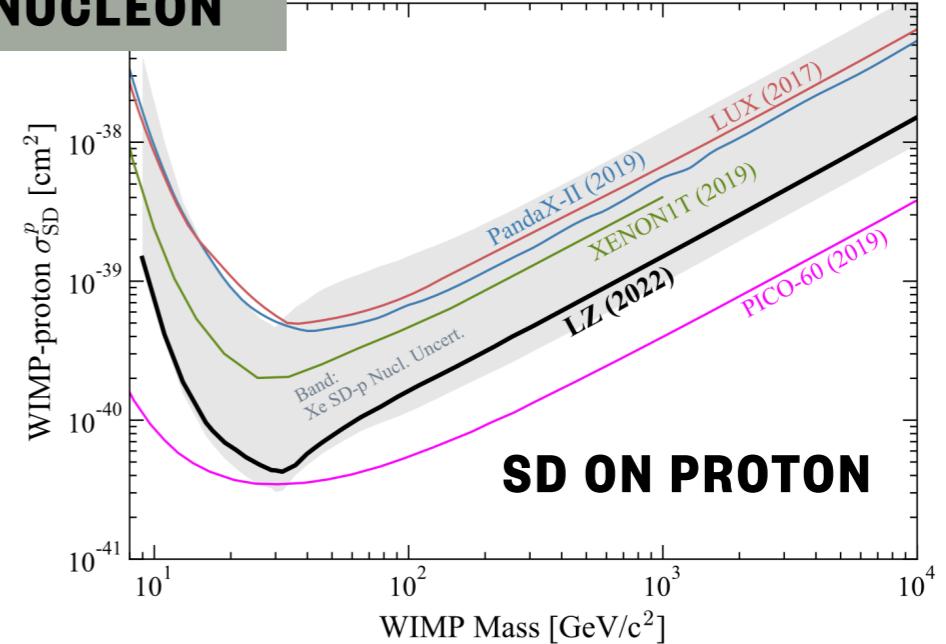
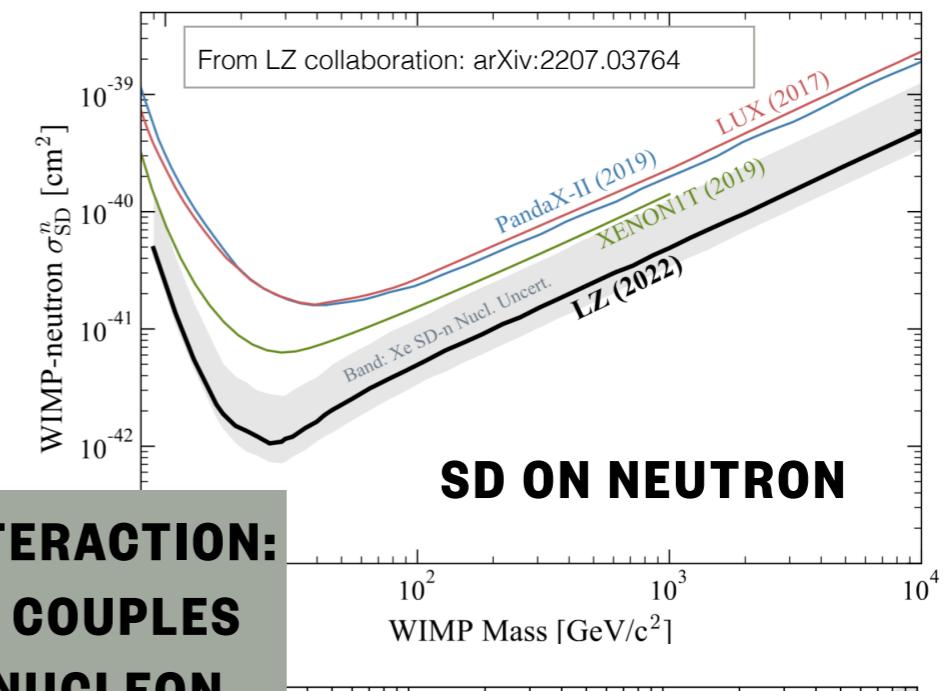
SMALL COUPLINGS

- CHALLENGING TO SEE IN DARK MATTER SEARCHES
- PROVIDE LLP SIGNALS AT LHC
- CAN HAVE VERY COMPRESSED SPECTRUM
- CONSTRAINTS FROM EARLY UNIVERSE (E.G. LYMAN- α , CMB, BBN)

Direct detection of WIMPs

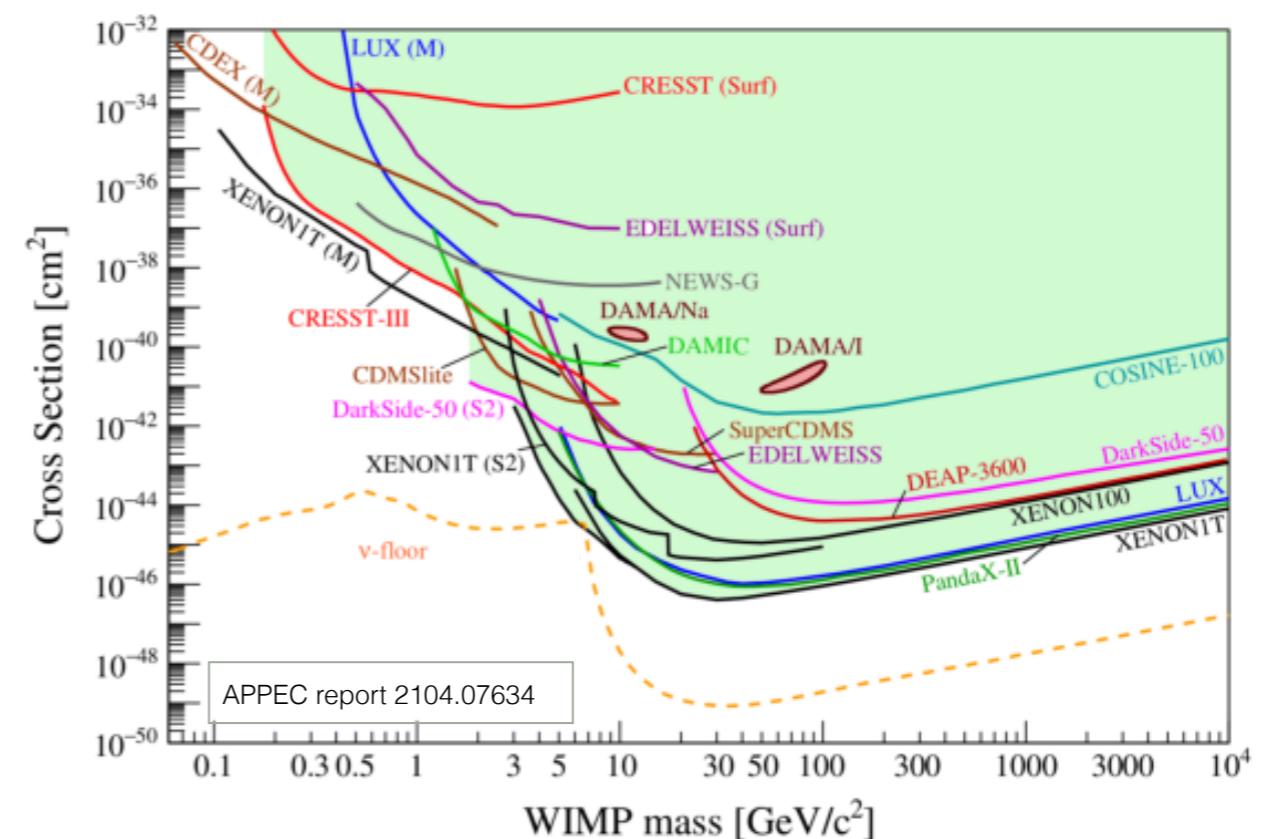
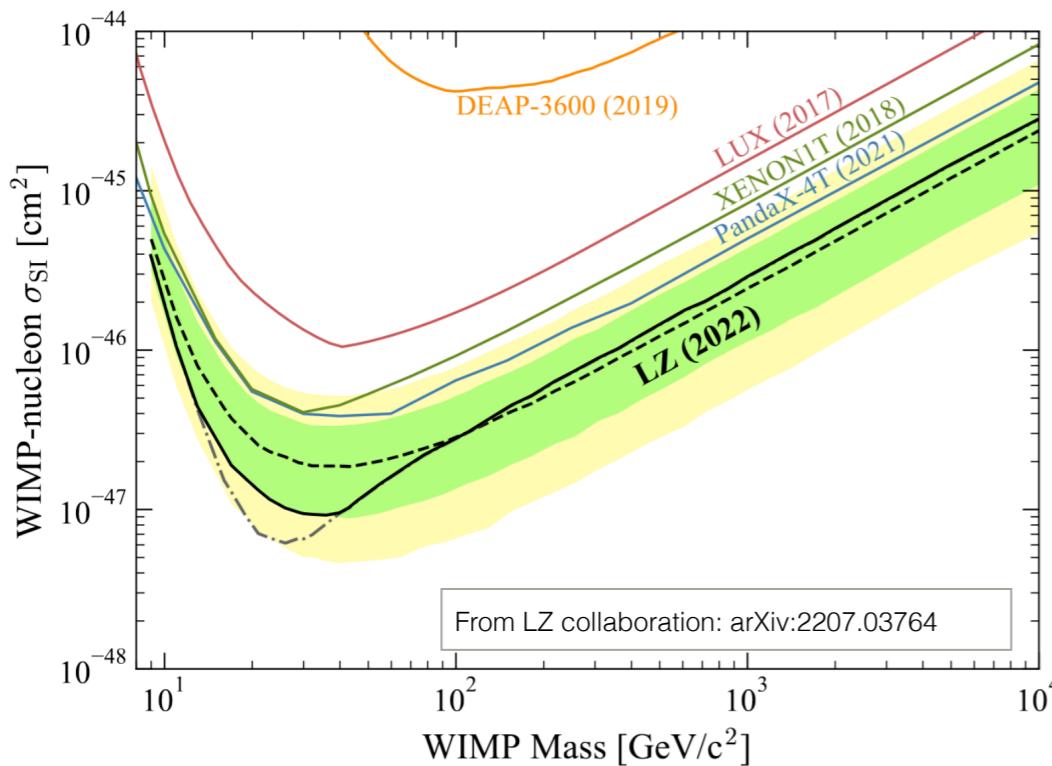


**SPIN-DEPENDENT INTERACTION:
DARK MATTER SPIN COUPLES
TO THE UNPAIRED NUCLEON**



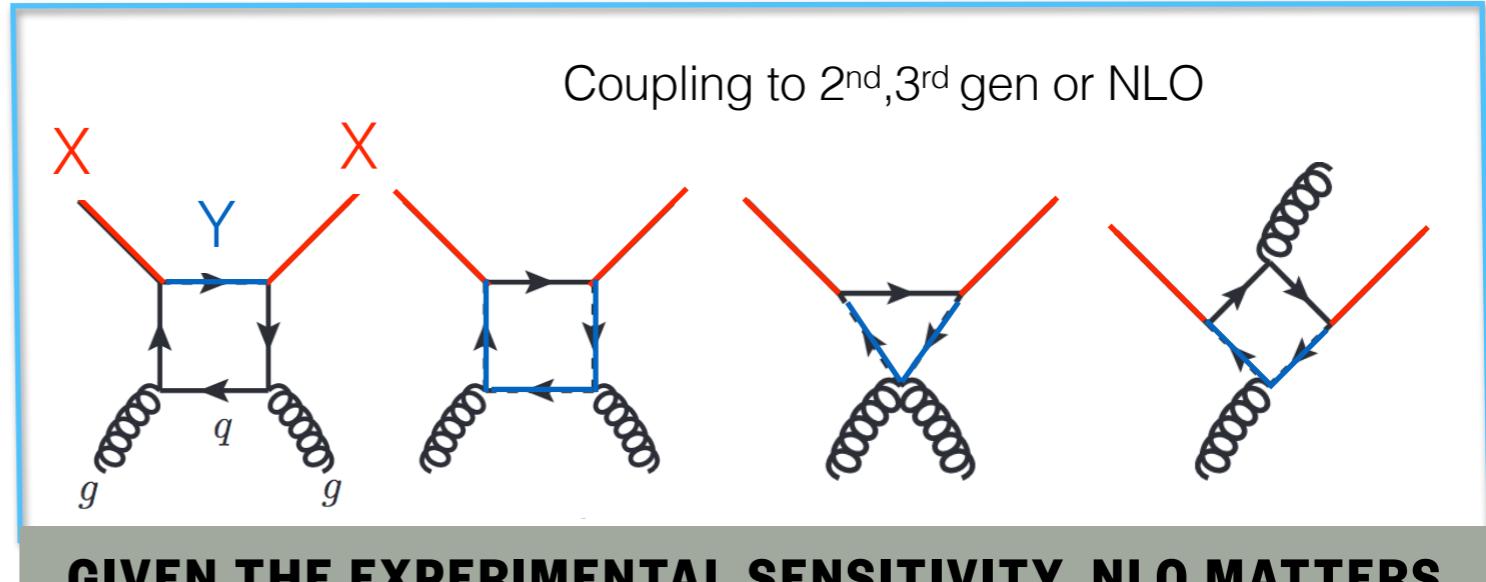
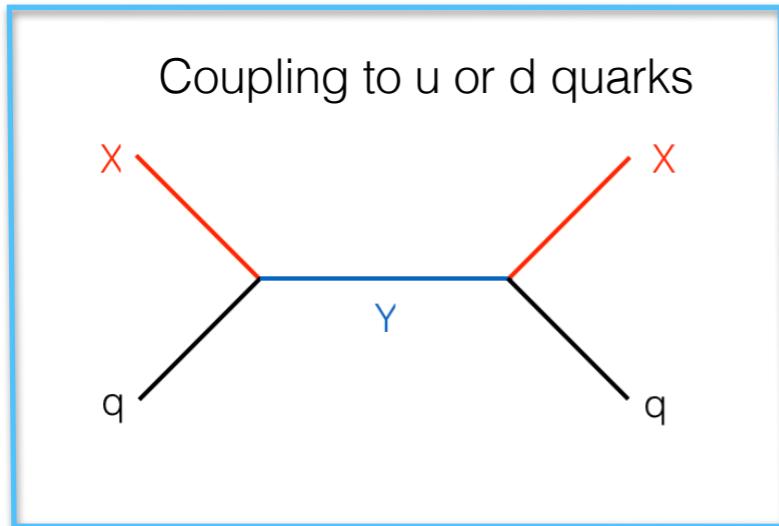
Direct detection of WIMPs

SPIN-INDEPENDENT INTERACTION: DARK MATTER COUPLES TO ALL NUCLEONS (SENSITIVE TO A^2)



DIRECT DETECTION SENSITIVE UP TO TENS OF TEV IN DARK MATTER MASS ($\propto m_X^{-2}$)

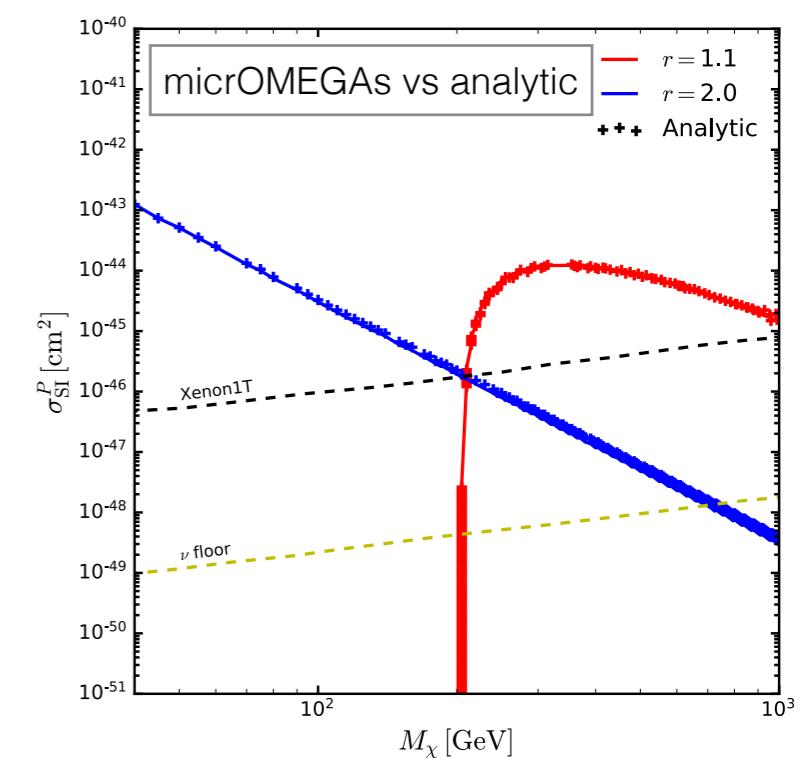
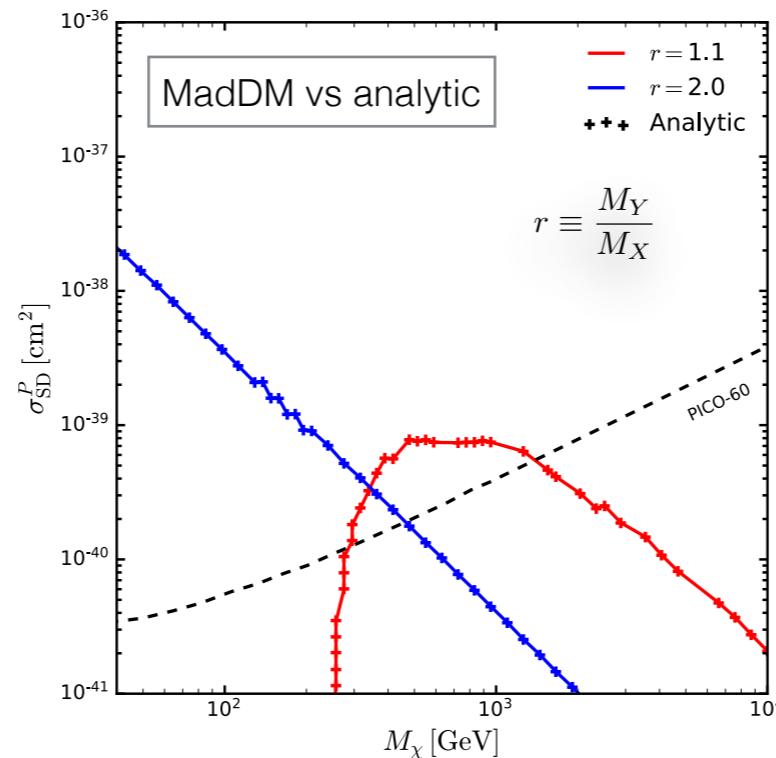
Direct detection of WIMPs



GIVEN THE EXPERIMENTAL SENSITIVITY, NLO MATTERS

FOR MAJORANA COUPLING TO UR (S3M_UR) DIRECT DETECTION IS:

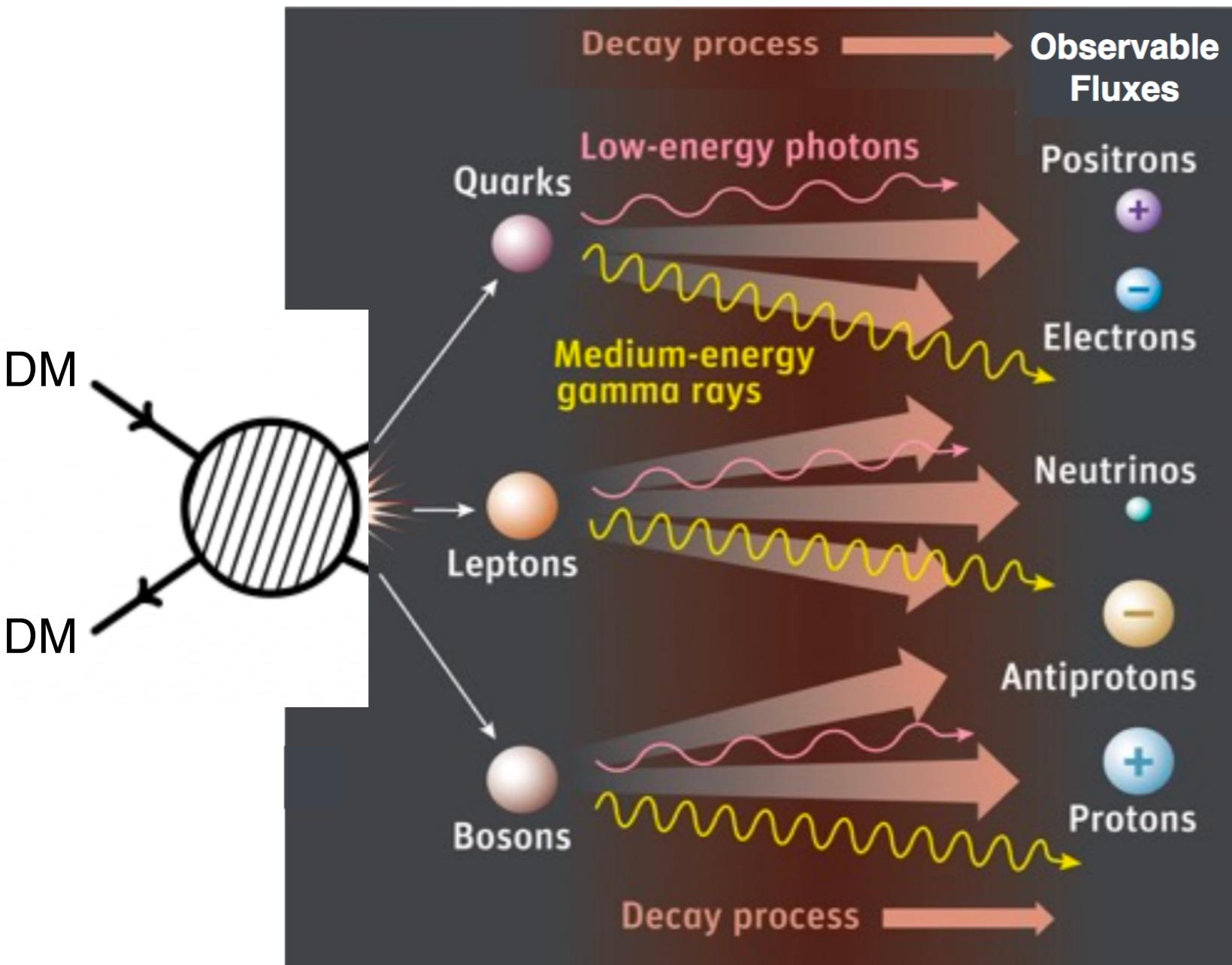
- LO FOR SPIN-DEPENDENT (MADDM TOOL)
 - NLO FOR SPIN-INDEPENDENT (MICROMEGAS TOOL)
 - ANALYTIC EXPRESSIONS [HISANO ET AL. (JHEP 2015)]



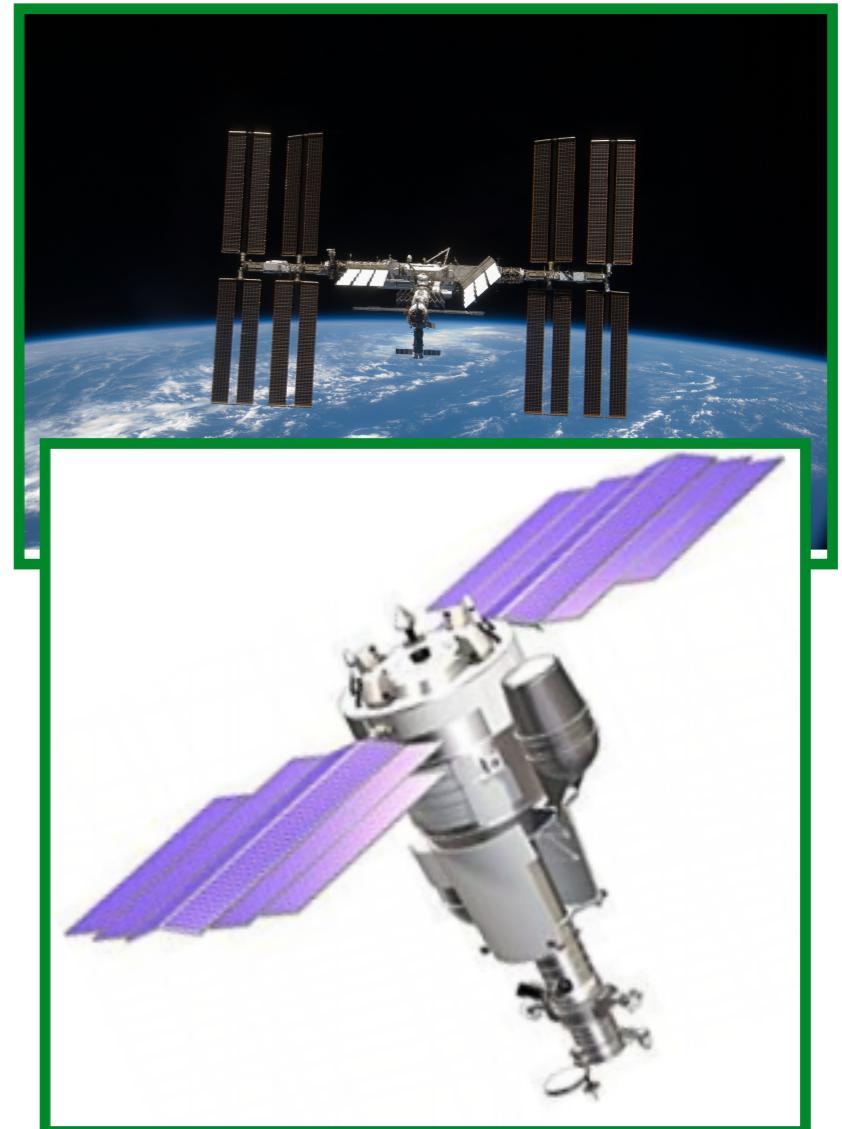
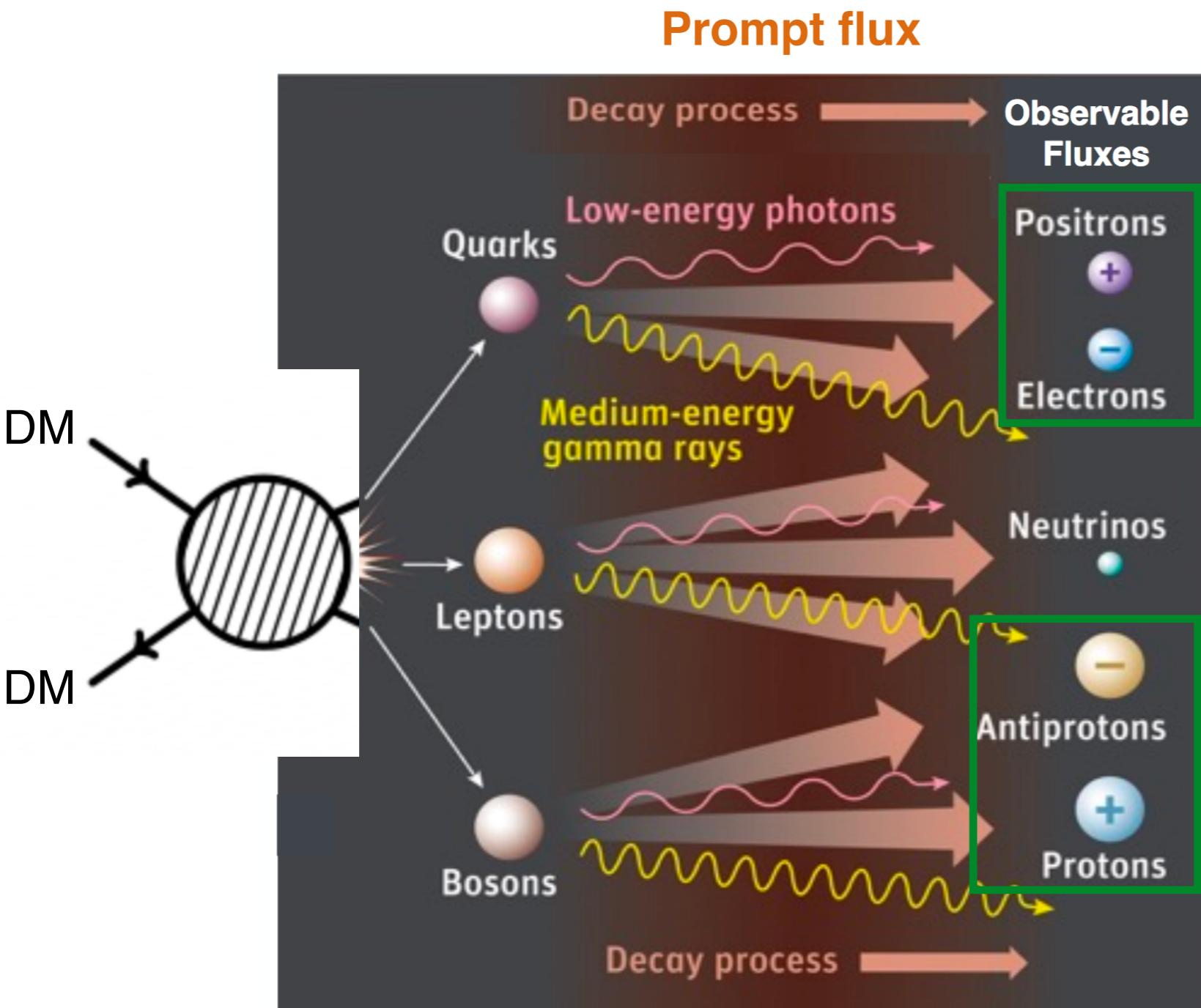
UFO AT LO AND NLO CAN BE USED DIRECTLY INTO DARK MATTER TOOLS

Indirect detection

Prompt flux

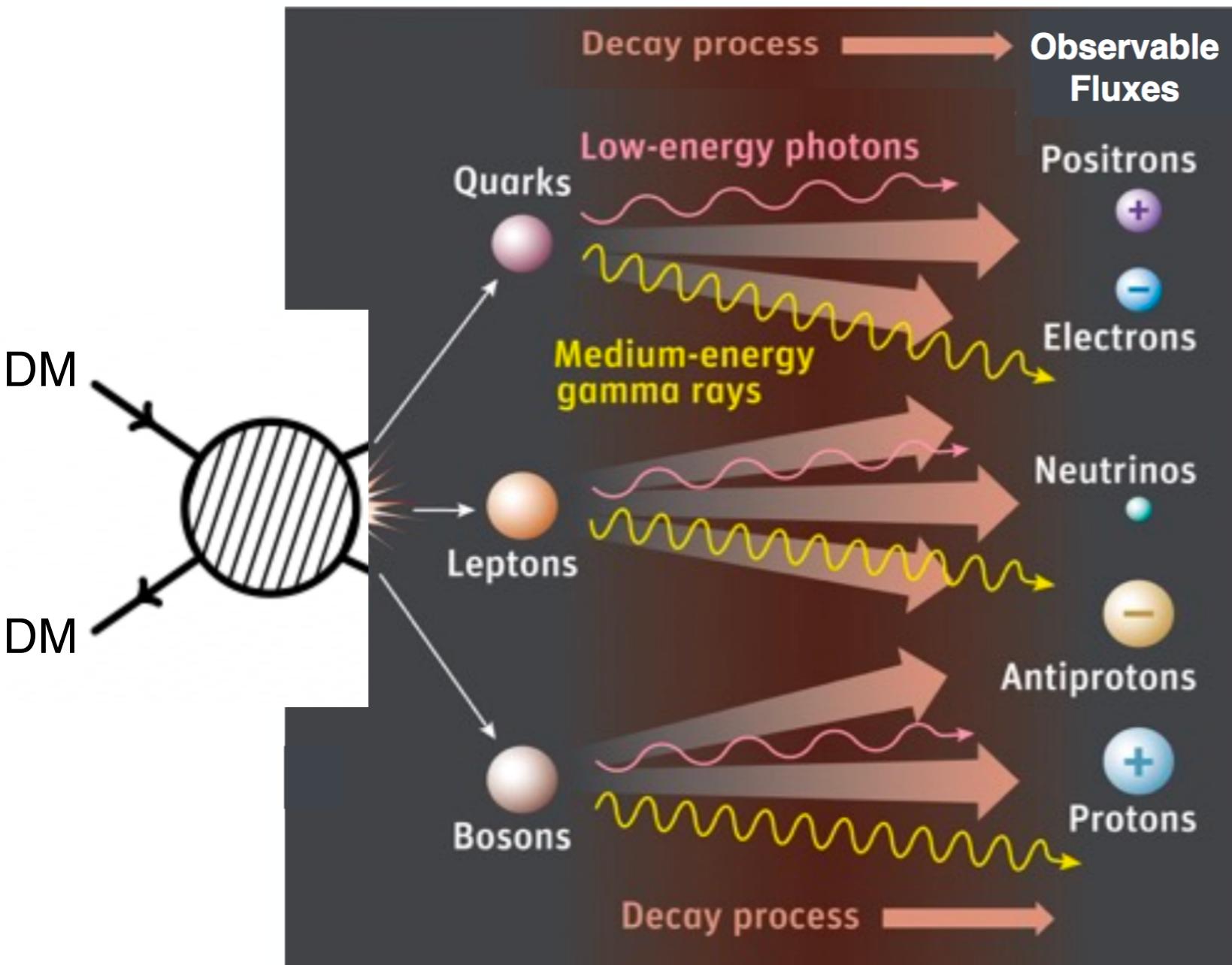


Indirect detection



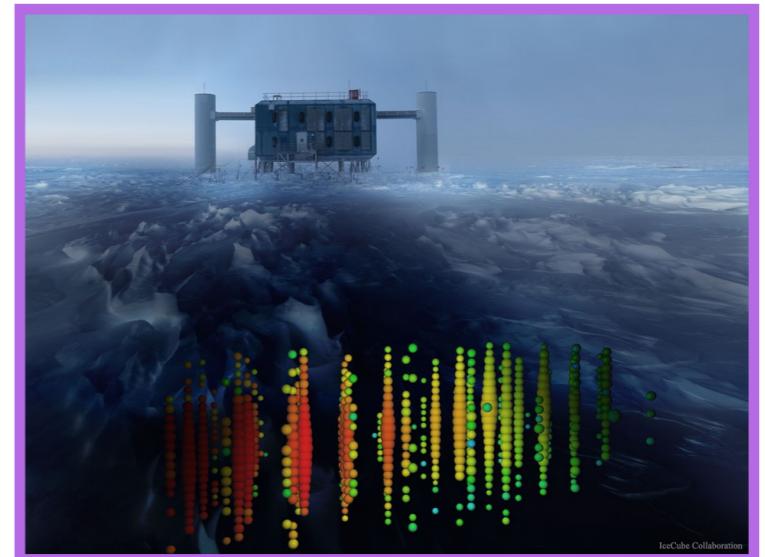
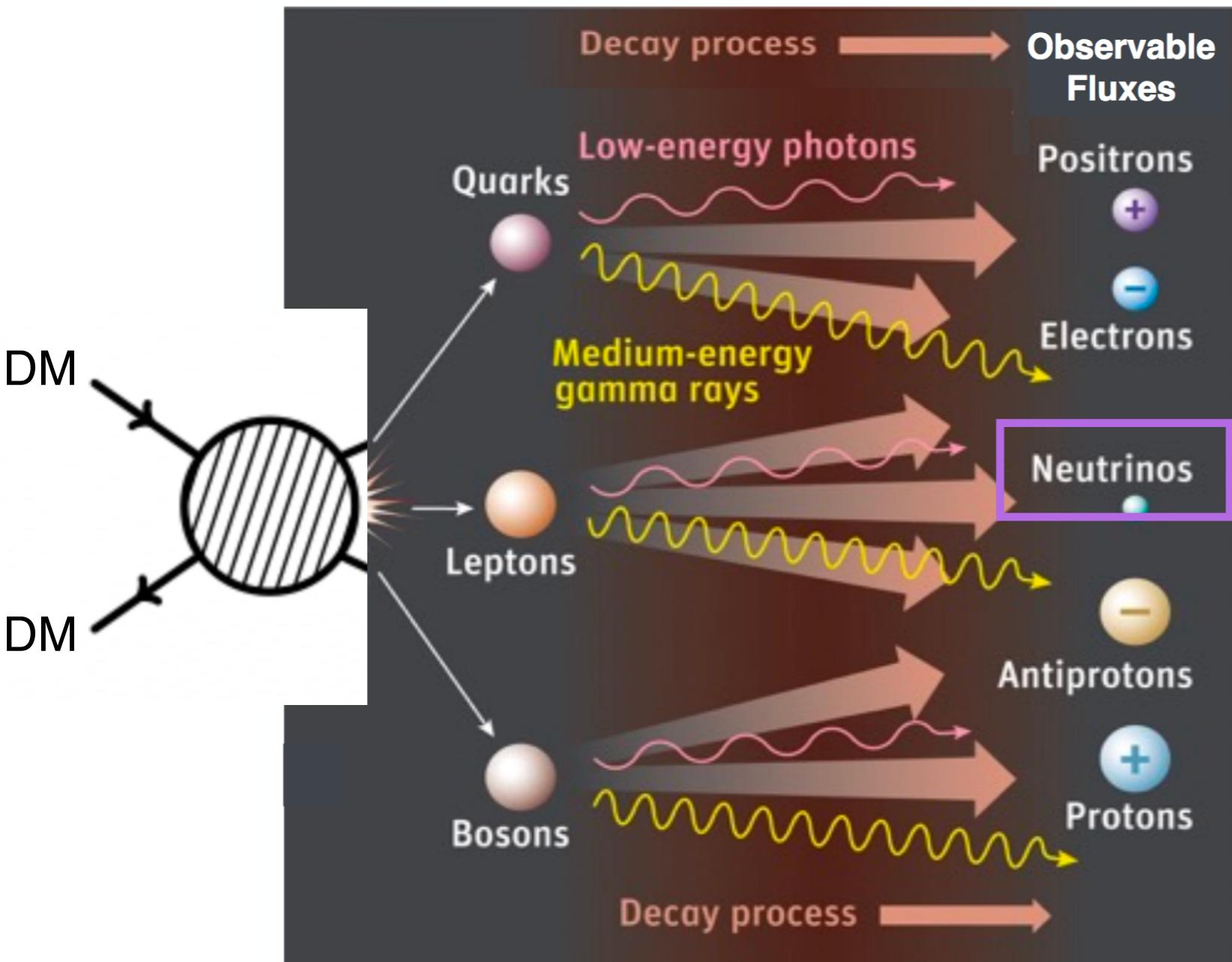
Indirect detection

Prompt flux



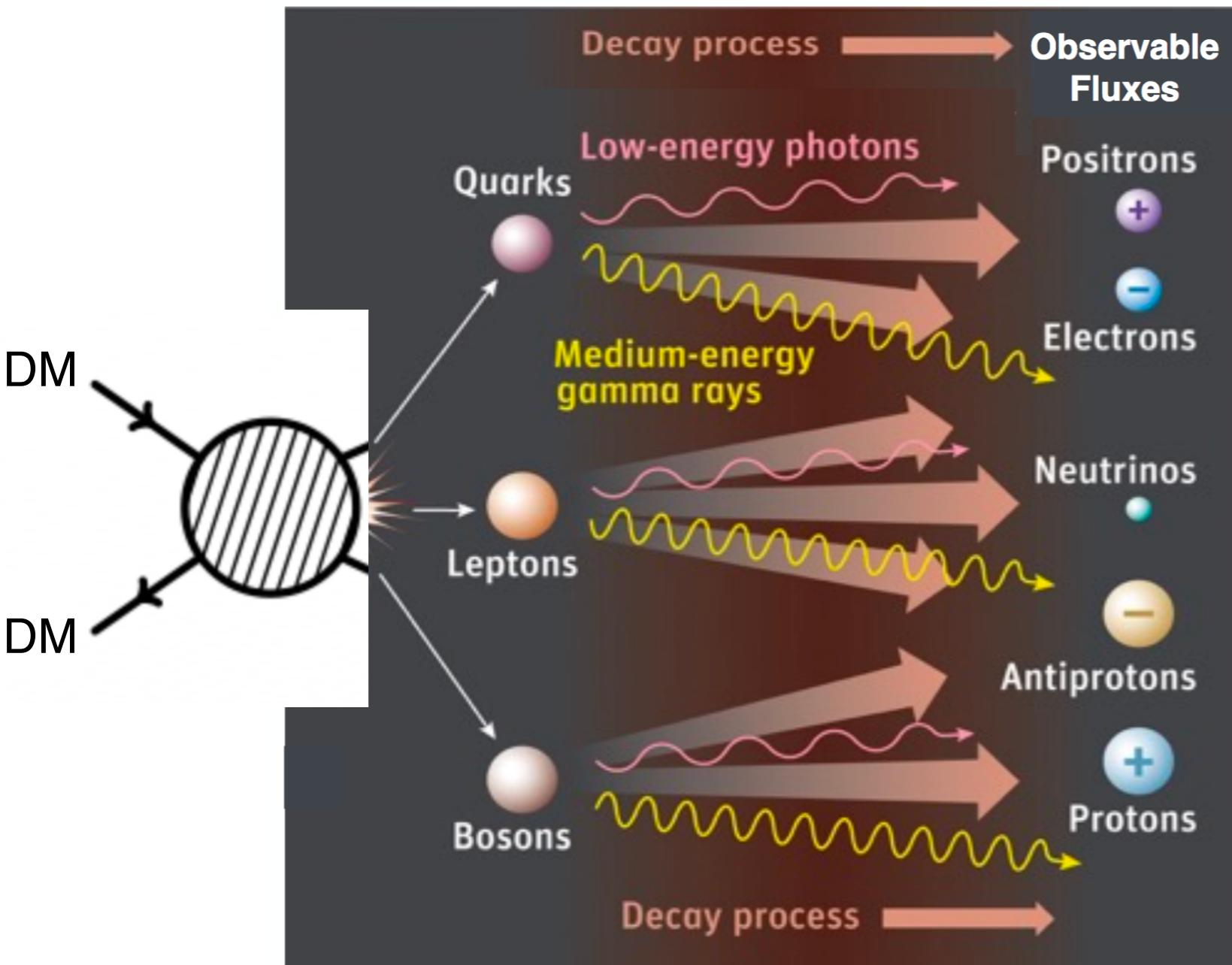
Indirect detection

Prompt flux



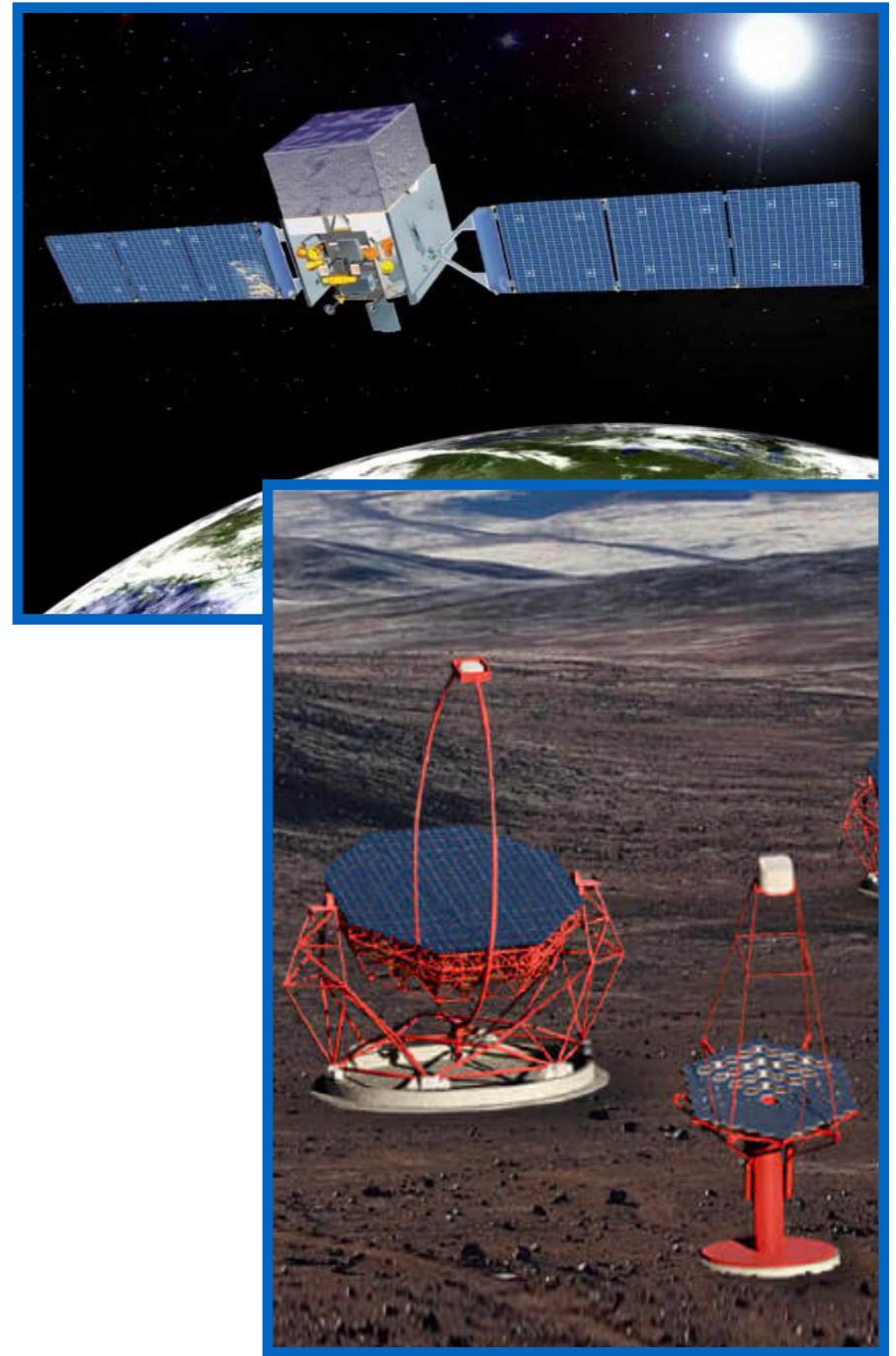
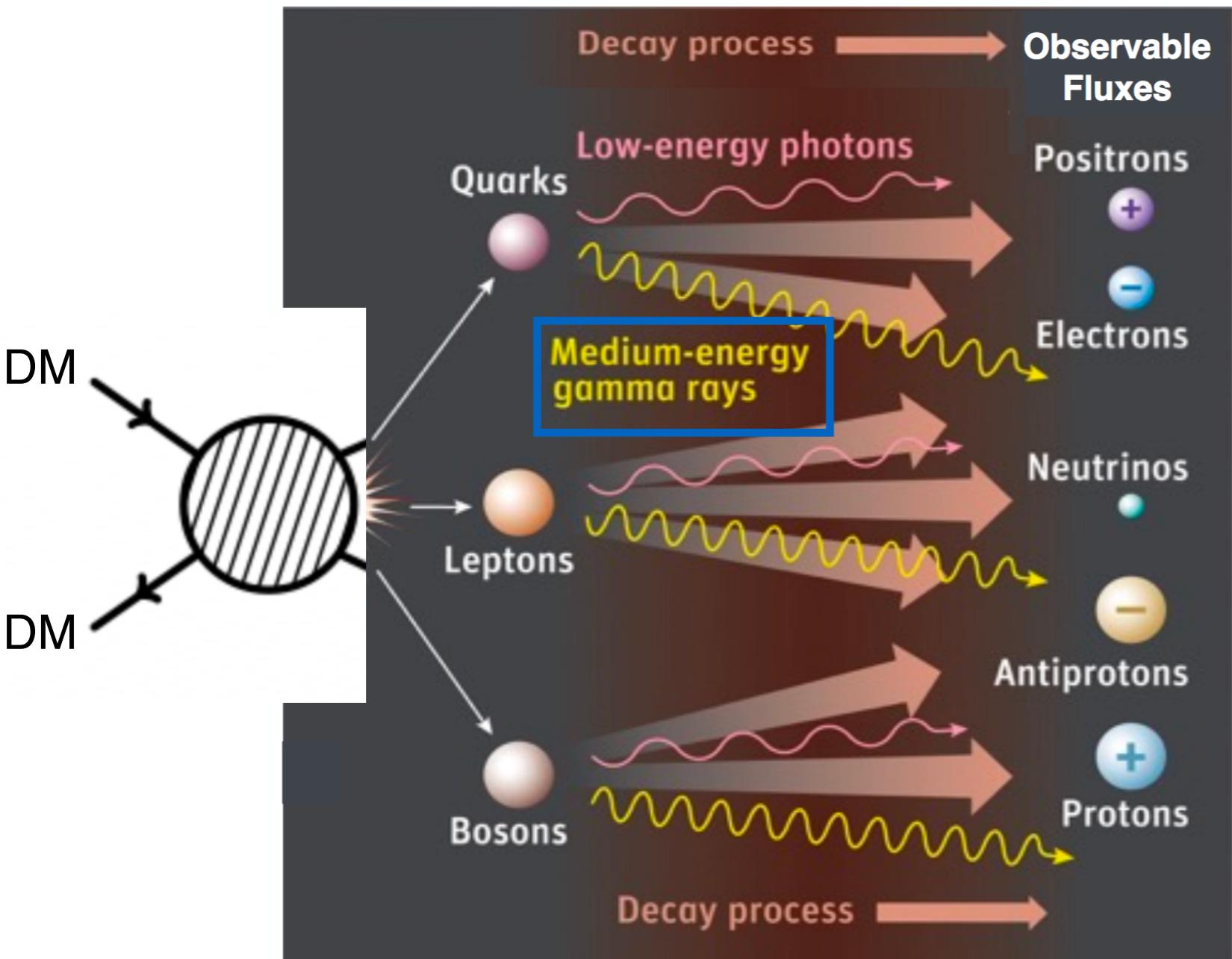
Indirect detection

Prompt flux



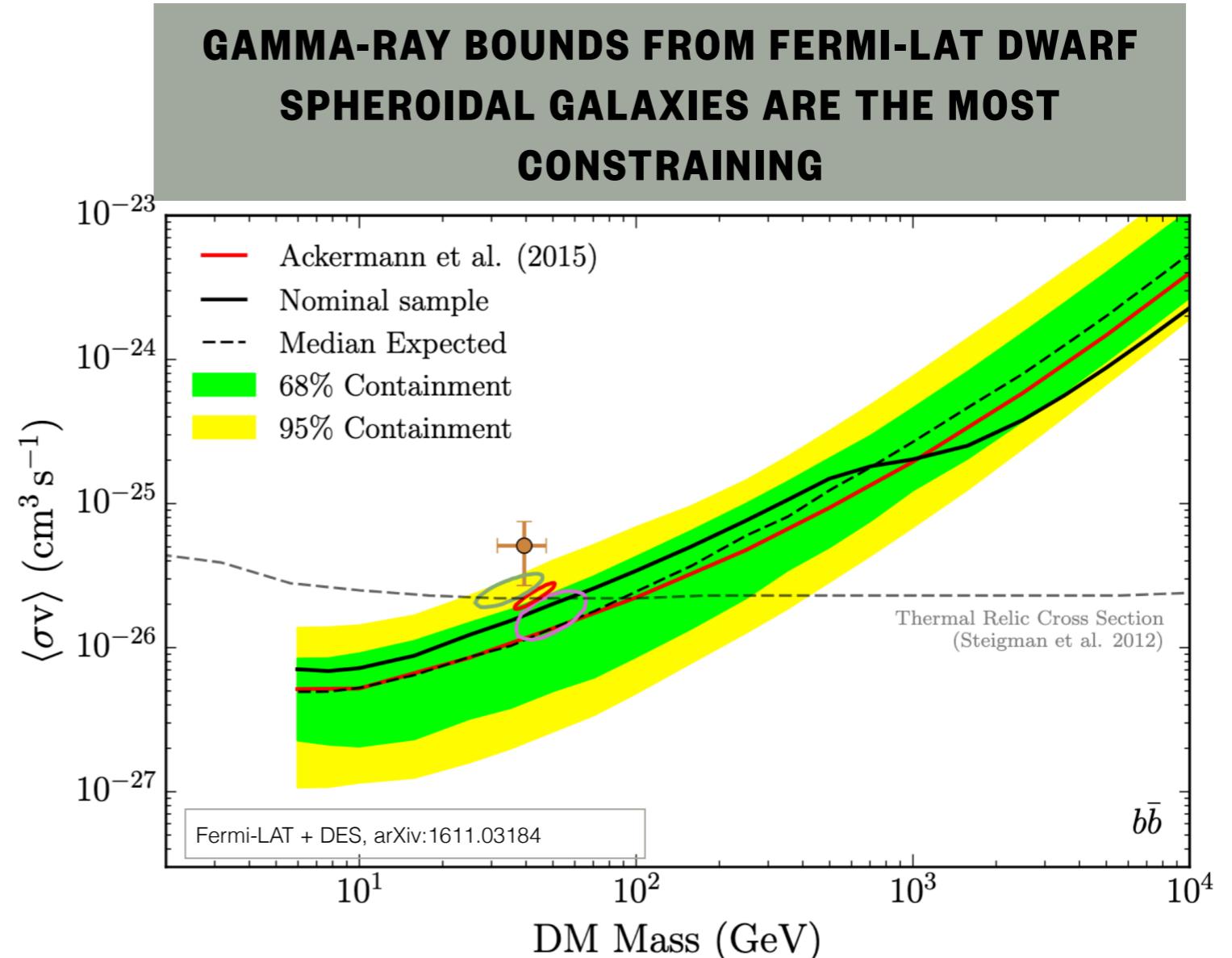
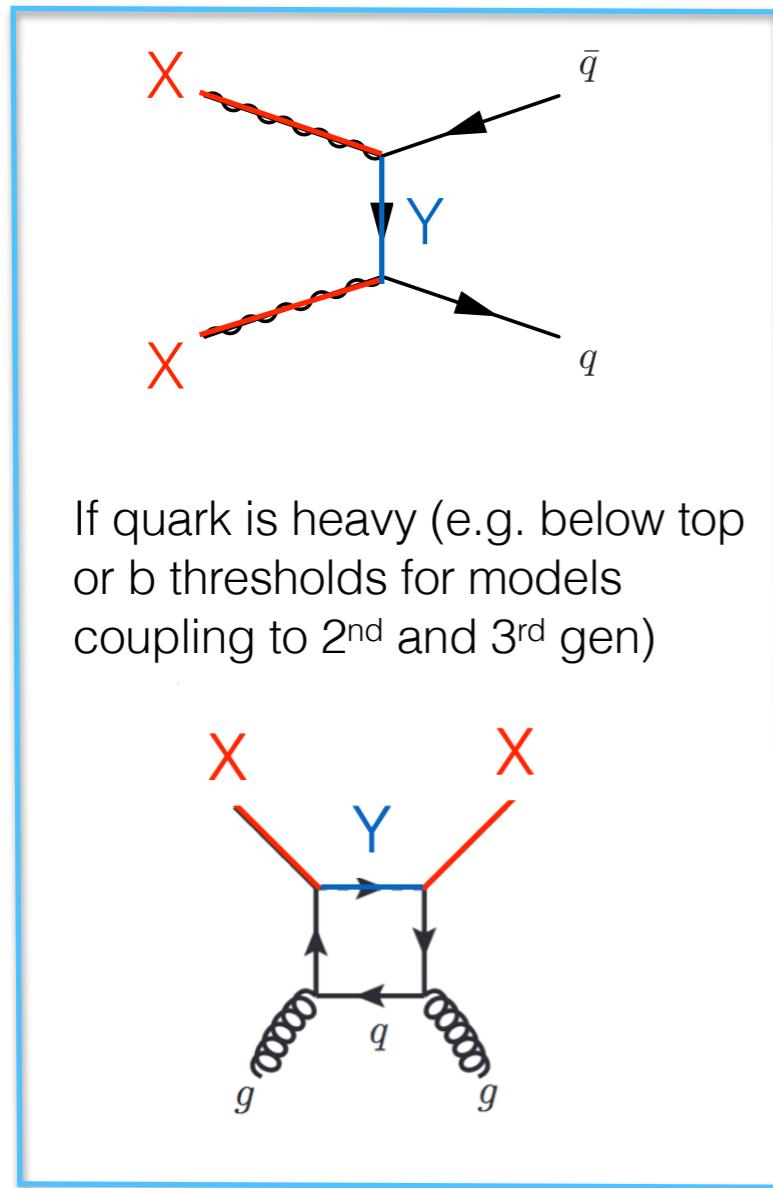
Indirect detection

Prompt flux



Dark matter annihilation in galactic halos at present time

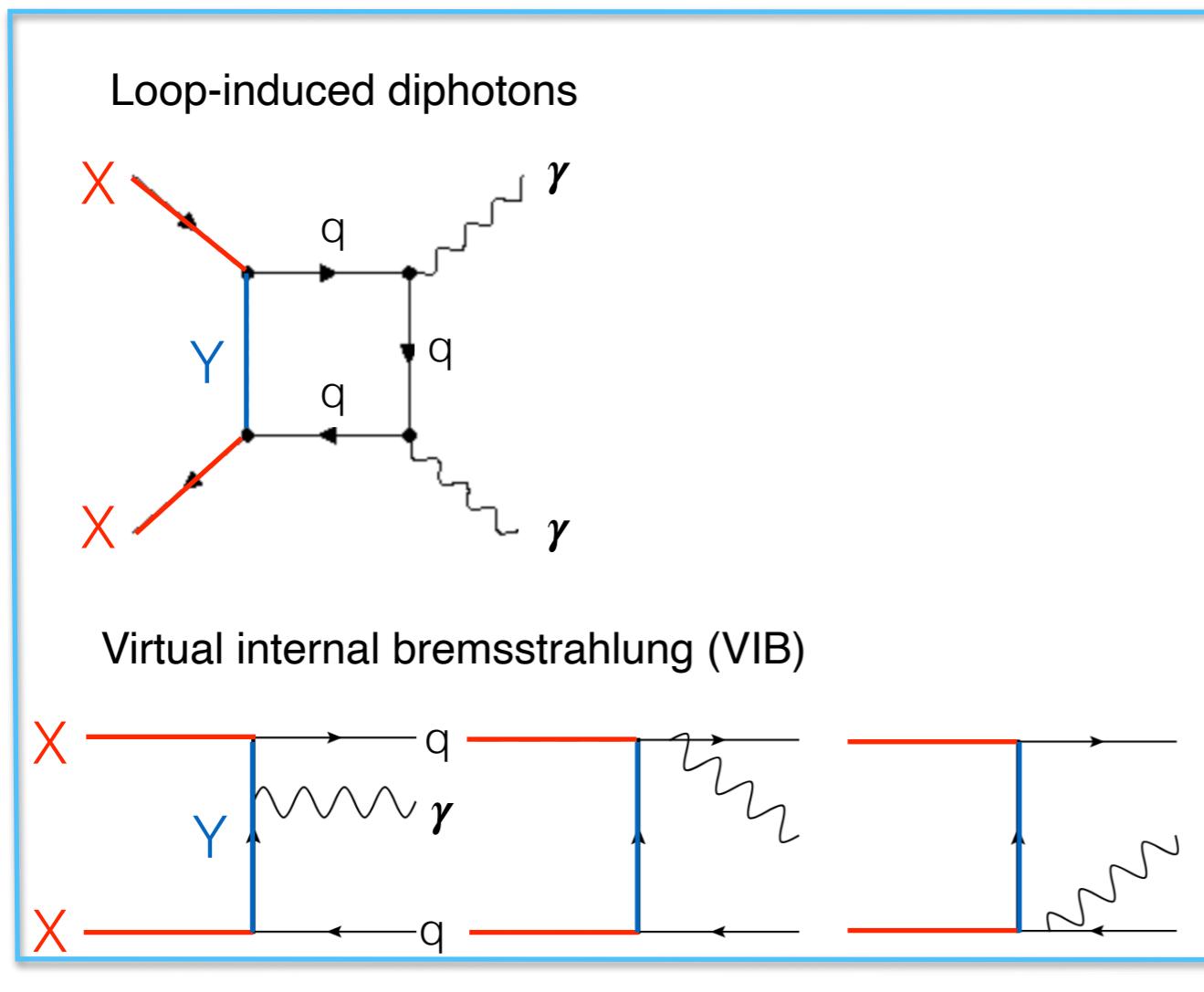
LO ANNIHILATION (DIRAC AND VECTORIAL DARK MATTER)



BOUNDS FROM FERMI-LAT DWARF SPHEROIDAL GALAXIES CAN BE COMPUTED AUTOMATICALLY WITHIN MADDM

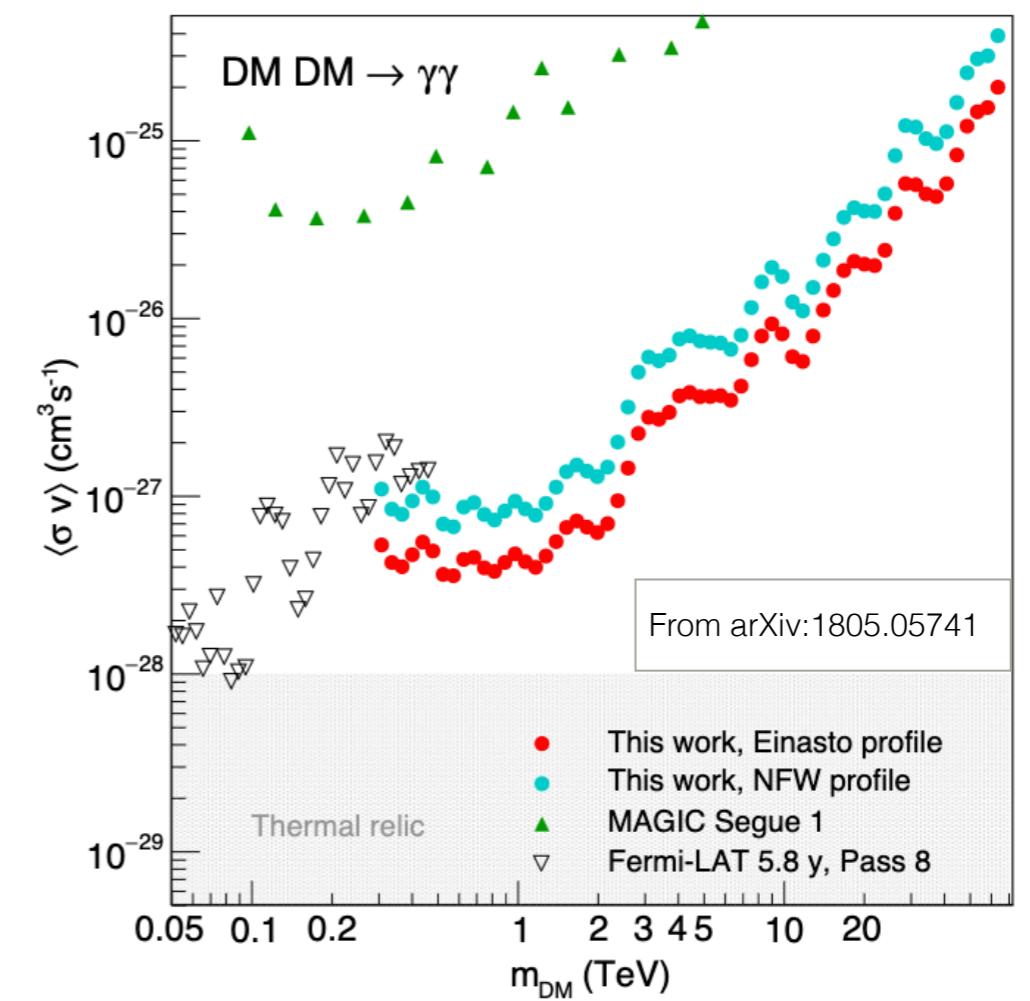
Dark matter annihilation in galactic halos at present time

- LO ANNIHILATION IS P-WAVE OR D-WAVE SUPPRESSED (MAJORANA OR SCALAR DARK MATTER)
- NLO PROCESSES UPLIFT THE SUPPRESSION AND PRODUCE A SHARP FEATURE IN THE GAMMA-RAY ENERGY SPECTRUM



CHIARA ARINA, 16/05/2024

GAMMA LINE SEARCHES ARE THE MOST STRINGENT BOUNDS

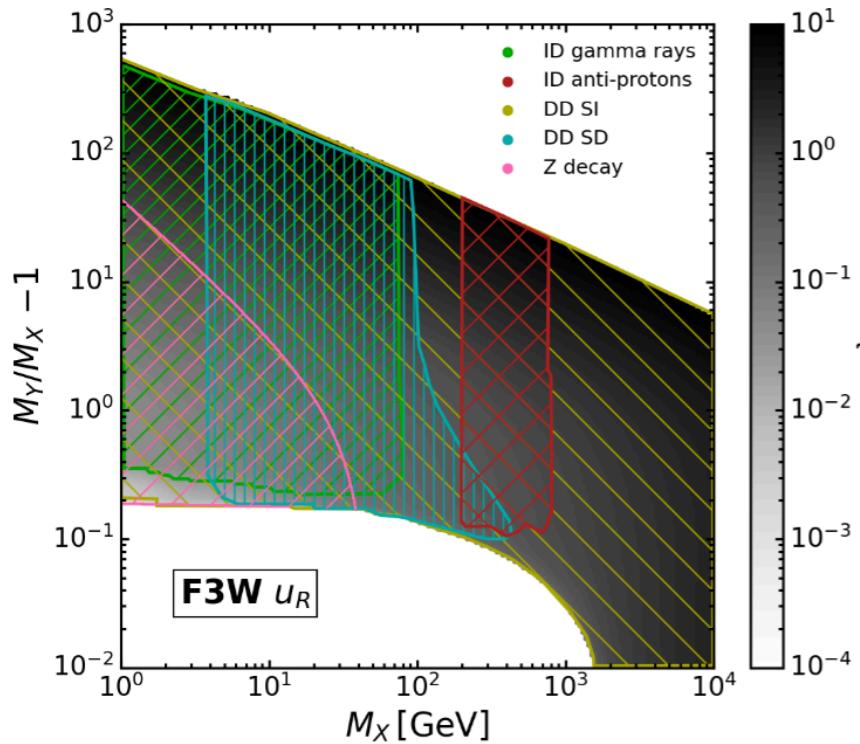


Benchmark 1: coupling to u_R

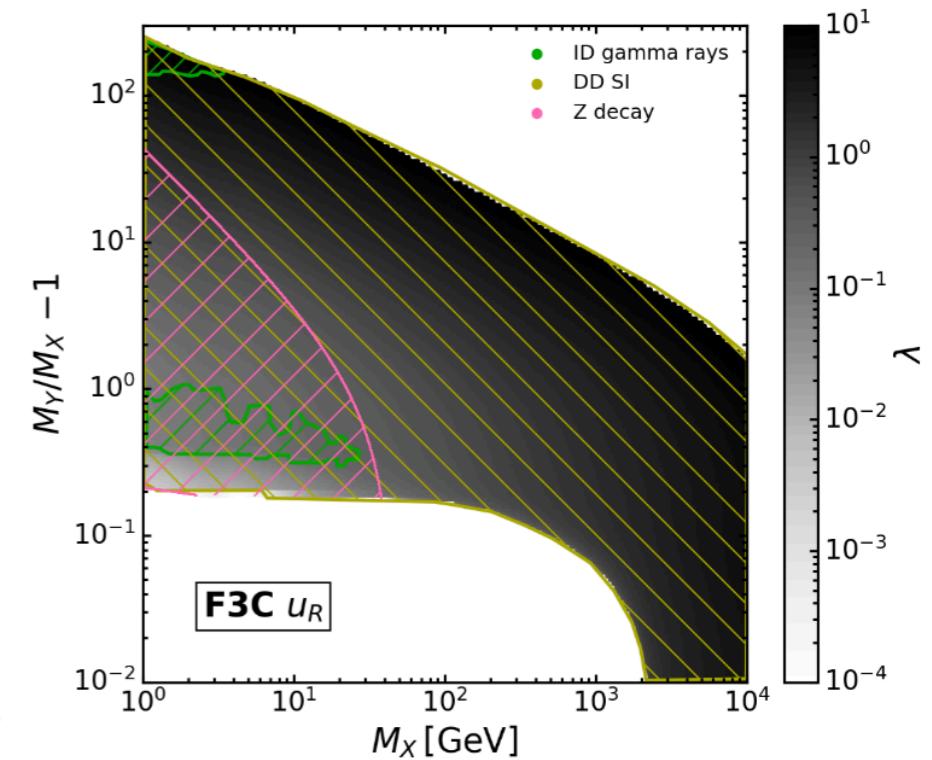
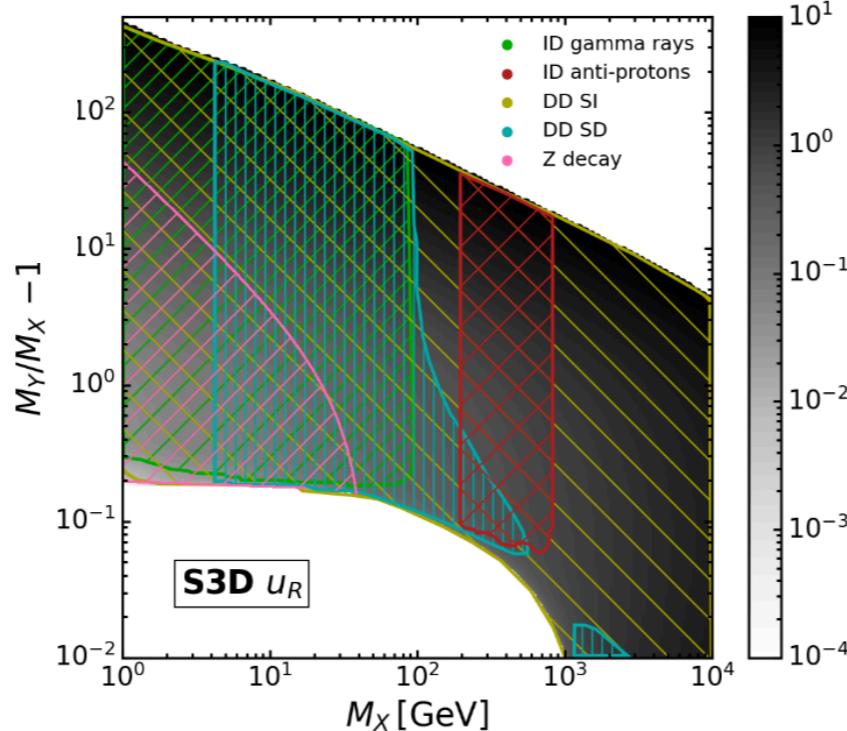
Name	DM	Mediators	Parameters	We consider all models coupling to u_R quark 6 in total (M_χ, M_Y, λ)
S3M_uni	$\tilde{\chi}$	$\varphi_{Q_f}, \varphi_{u_f}, \varphi_{d_f}$		
S3D_uni	χ			
S3M_3rd	$\tilde{\chi}$	$\varphi_{Q_3}, \varphi_{u_3}, \varphi_{d_3}$	$M_\varphi, M_\chi, \lambda_\varphi$	
S3D_3rd	χ			
S3M_uR	$\tilde{\chi}$	φ_{u_1}		Fermionic DM (Majorana and Dirac) and scalar mediator
S3D_uR	χ			
F3S_uni	\tilde{S}	$\psi_{Q_f}, \psi_{u_f}, \psi_{d_f}$		
F3C_uni	S			
F3S_3rd	\tilde{S}	$\psi_{Q_3}, \psi_{u_3}, \psi_{d_3}$	$M_S, M_\psi, \hat{\lambda}_\psi$	
F3C_3rd	S			
F3S_uR	\tilde{S}	ψ_{u_1}		Scalar DM (real and complex) and fermionic mediator
F3C_uR	S			
F3V_uni	\tilde{V}_μ	$\psi_{Q_f}, \psi_{u_f}, \psi_{d_f}$		
F3W_uni	V_μ			
F3V_3rd	\tilde{V}_μ	$\psi_{Q_3}, \psi_{u_3}, \psi_{d_3}$	$M_V, M_\psi, \hat{\lambda}_\psi$	
F3W_3rd	V_μ			
F3V_uR	\tilde{V}_μ	ψ_{u_1}		Vector DM (real and complex) and fermionic mediator
F3W_uR	V_μ			

Benchmark 1: coupling to u_R complex case

From arXiv:2307.10367

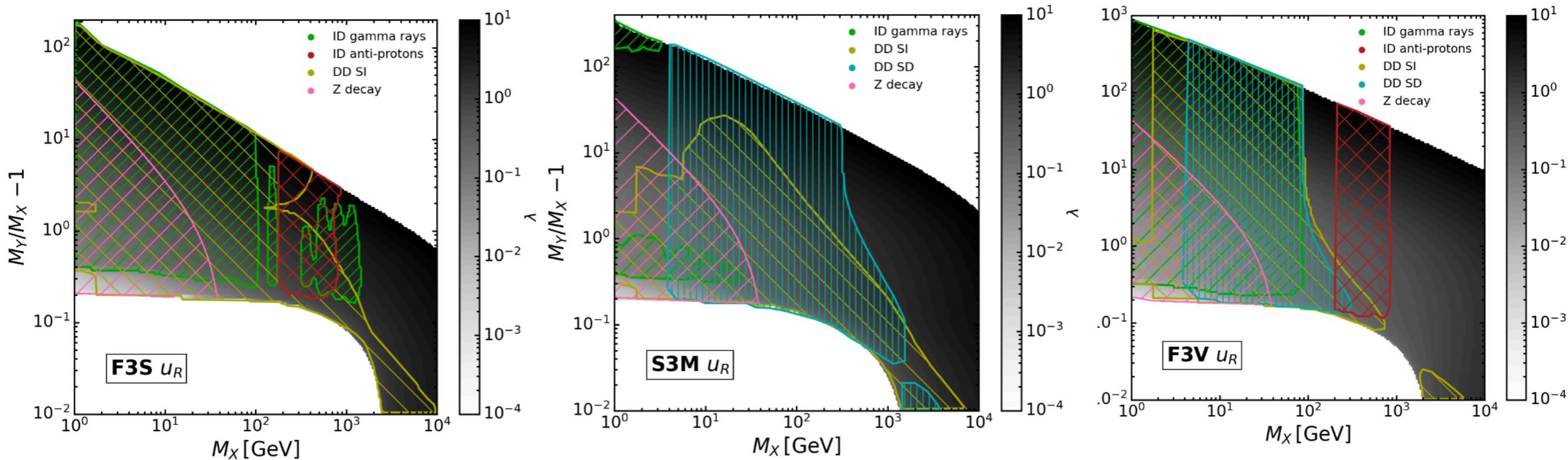


**ASSUMING WIMP SCENARIO
AND CORRECT RELIC
DENSITY VIA FO ALL
COMPLEX MODELS ARE
EXCLUDED UP TO 10 TEV BY
DIRECT DETECTION !**



**FREEZE-IN REGION AND
COMPRESSED SPECTRA TO BE
INVESTIGATED STILL...**

Benchmark 1: coupling to u_R real case



**ASSUMING WIMP SCENARIO
AND CORRECT RELIC
DENSITY VIA FO COSMOLOGY
EXCLUDES A LOT BUT STILL
VIABLE BENCHMARKS**

**FREEZE-IN REGION AND
COMPRESSED SPECTRA TO BE
INVESTIGATED STILL...**

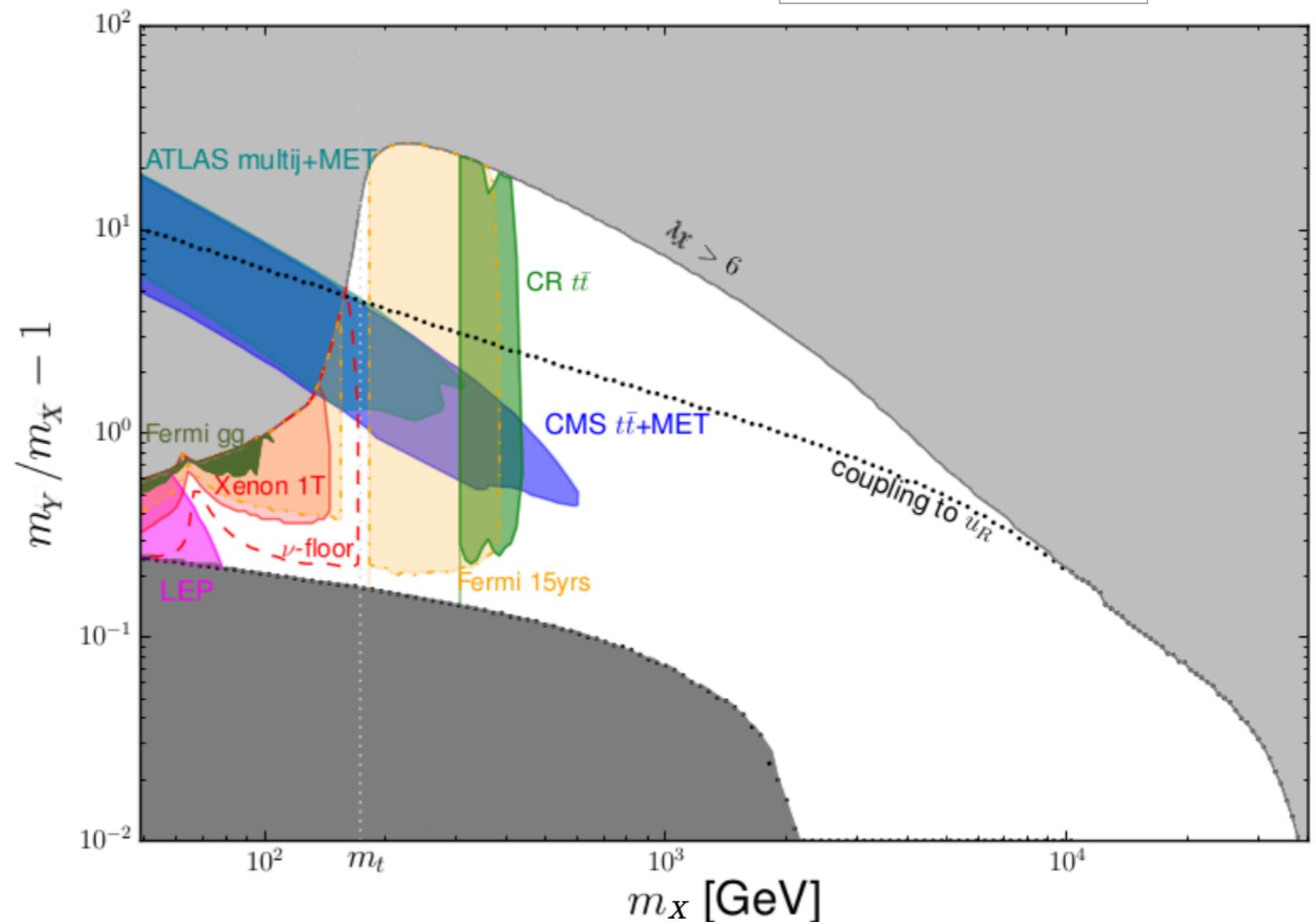
From arXiv:2307.10367

Benchmark 2: coupling to t_R

Real scalar particle as X and fermionic
particle as Y
(M_X, M_Y, λ)

- BEHAVIOR IS SIMILAR TO THE CASE OF U_R
- NON NEGIGIBLE MASS OF THE TOP IS RELEVANT AND NEED A CAREFUL TREATMENT
- COMPLEMENTARITY OF ASTROPARTICLE AND COLLIDER SEARCHES

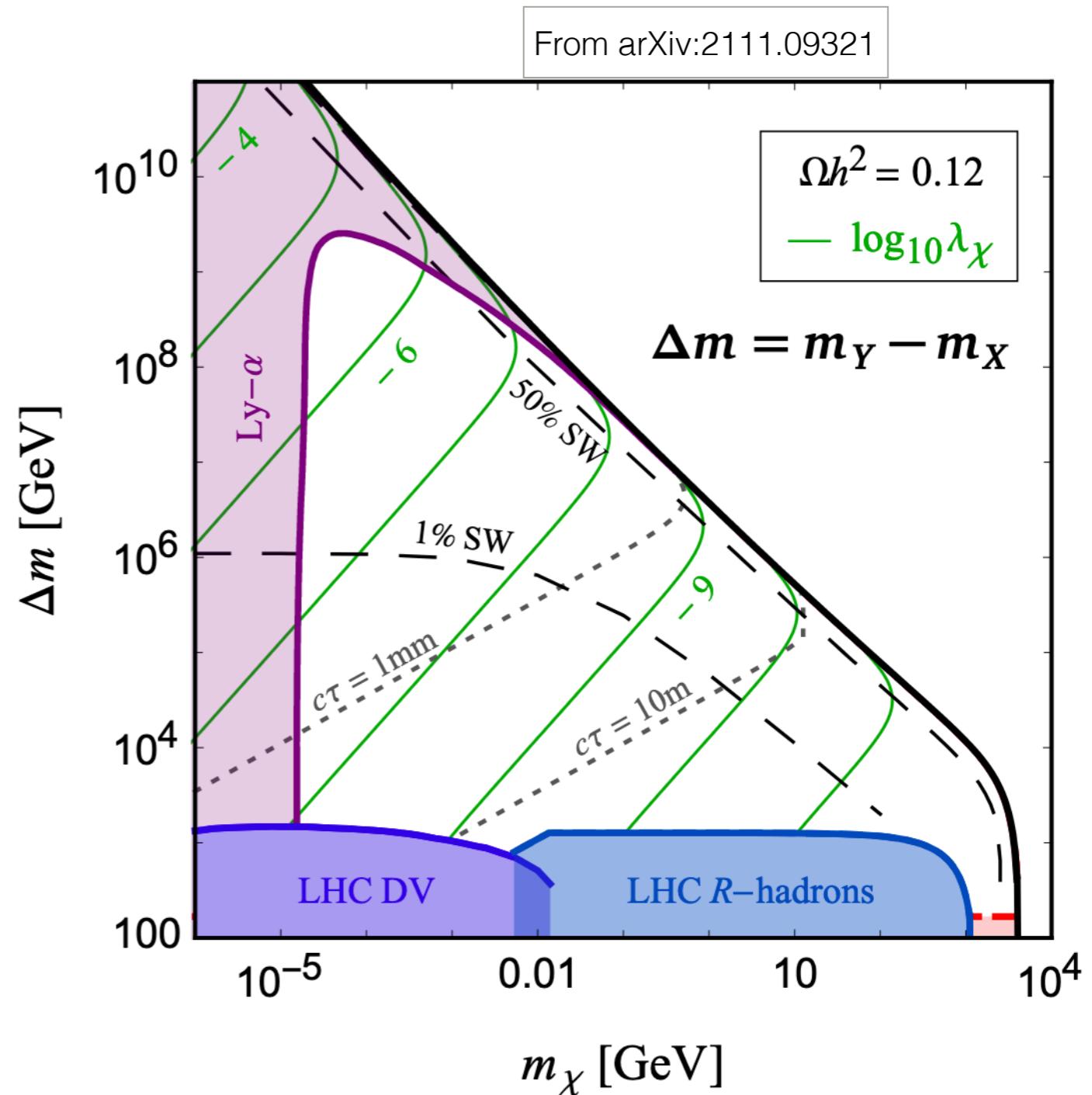
From arXiv:1804.05068



Benchmark 2: coupling to t_R for freeze-in and superWIMP

Majorana particle as X and scalar
particle as Y
(M_X, M_Y, λ)

- WIDE RANGE OF DM MASSES CAN BE CONSIDERED
- LOW MASSES ARE CONSTRAINED BY COSMOLOGY
- COMPRESSED SPECTRA CONSTRAINED BY DISPLACED VERTICES AND R-HADRONS
- NICE COMPLEMENTARITY
- WIDE VIABLE REGION (NO DD OR ID BOUND)



Summary

DARK MATTER CANDIDATES ARISING FROM SIMPLIFIED AND MINIMAL T-CHANNEL MODELS ARE PARTICULAR ATTRACTIVE BECAUSE OF THEIR PREDICTIVITY:

- IF CONSIDERED AS WIMPS THEY ARE PRETTY CONSTRAINED BY INDIRECT AND DIRECT SEARCHES OF DARK MATTER;
 - MODELS COUPLING TO U AND D ARE MORE CONSTRAINED (COMPLEX CANDIDATES ARE BASICALLY EXCLUDED)
 - MODELS COUPLING TO 2ND AND 3RD GEN ARE LESS CONSTRAINED BY DIRECT DETECTION
 - PHENO OF RELIC AND INDIRECT DETECTION DO NOT STRONGLY DEPEND ON THE QUARK FLAVOR (BESIDES MASS EFFECTS)
-
- Models beyond the minimal version and/or theoretical complete models feature enlarged parameter space, specific signatures, ...;
 - Freeze-in, superWIMP, conversion driven freeze-out are other mechanisms that shape LLP regions and/or very compressed spectra: different regions and masses to explore yet;
 - A non-standard cosmological history can change the model parameter space and open up new regions;
 - Much more in the white paper to appear, stay tuned!

Back up slides

Bound states and Sommerfeld enhancement (from arXiv:2203.04326)

SOMMERFELD ENHANCEMENT

$$\sigma_{\text{SE},[\mathbf{R}]} v_{\text{rel}} = c_{[\mathbf{R}]} S_{0,[\mathbf{R}]} \sigma_0,$$

$$\sigma_{\mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow gg} v_{\text{rel}} = \sigma_{\mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow gg,0} \left(\frac{2}{7} S_{0,[1]} + \frac{5}{7} S_{0,[8]} \right),$$

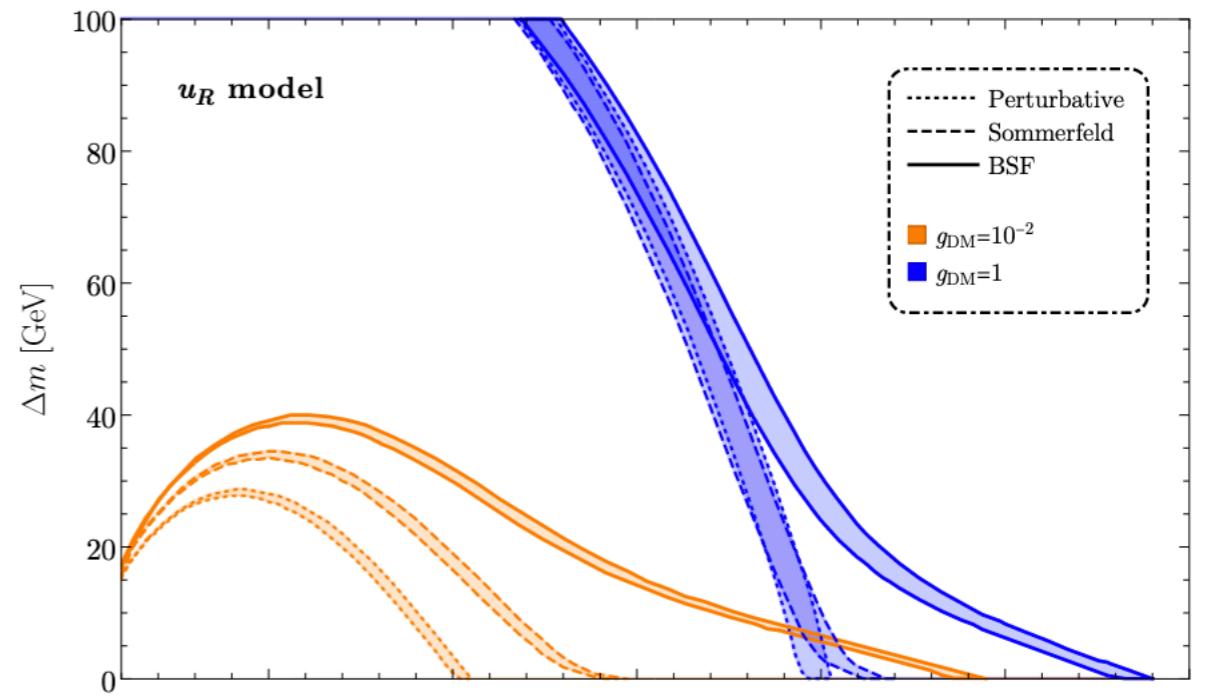
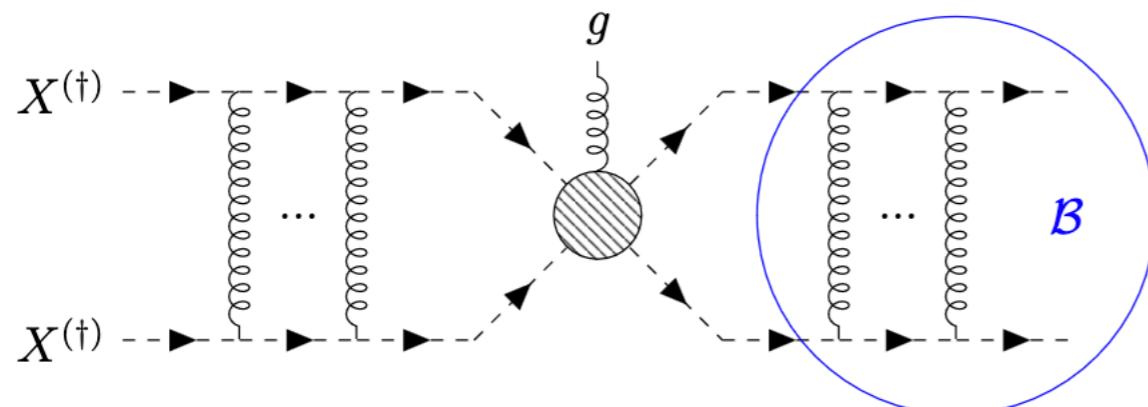
$$\sigma_{\mathbf{3} \otimes \bar{\mathbf{3}} \rightarrow q\bar{q}} v_{\text{rel}} = \sigma_{\mathbf{3} \otimes \bar{\mathbf{3}},0} (f_{[1]}(g_s, g_{\text{DM}}) S_{0,[1]} + f_{[8]}(g_s, g_{\text{DM}}) S_{0,[8]}),$$

$$\sigma_{\mathbf{3} \otimes \mathbf{3} \rightarrow qq} v_{\text{rel}} = \sigma_{\mathbf{3} \otimes \mathbf{3} \rightarrow qq,0} S_{0,[6]},$$

$$\sigma_{\mathbf{3}_i \otimes \mathbf{3}_j \rightarrow q_i q_j} = \sigma_{\mathbf{3}_i \otimes \mathbf{3}_j \rightarrow q_i q_j,0} \left(\frac{1}{3} S_{0,[\bar{\mathbf{3}}]} + \frac{2}{3} S_{0,[\mathbf{6}]} \right).$$

$$V(r)_{\mathbf{3} \otimes \bar{\mathbf{3}}} = \begin{cases} -\frac{4}{3} \frac{\alpha_s}{r} & [1] \\ +\frac{1}{6} \frac{\alpha_s}{r} & [8] \end{cases}; \quad V(r)_{\mathbf{3} \otimes \mathbf{3}} = \begin{cases} -\frac{2}{3} \frac{\alpha_s}{r} & [\bar{\mathbf{3}}] \\ +\frac{1}{3} \frac{\alpha_s}{r} & [\mathbf{6}] \end{cases}$$

EXAMPLE OF BOUND STATE FORMATION



Lyman-alpha bounds

**FI AND SW (NON-THERMAL) CAN BE PRODUCED
WITH A CERTAIN BOOST AND AFFECT STRUCTURE
FORMATION**

**BOUNDS FROM WARM DARK
MATTER (THERMAL) CAN BE
TRANSLATED INTO FI AND SW
TERMS**

$$m_X \gtrsim \begin{cases} 15 \text{ keV} \times \left(\frac{106.75}{g_*(T_{\text{FI}})} \right)^{1/3} & \text{for FI through decays,} \\ 3.8 \text{ GeV} \times (R_\Gamma^{\text{SW}} / 10^{-12})^{-1/2} \times \left(\frac{106.75}{g_*(T_{\text{SW}})} \right)^{1/3} & \text{for SW,} \end{cases}$$