

# **Giant Gravitons and non-conformal vacua in Twisted Holography**

Kasia Budzik

CERN, June 2023

Based on arxiv:2106.14859, arxiv:2211.01419 with Davide Gaiotto



# Twisting Supersymmetric QFTs

- Twisting SQFT is the procedure of passing to a cohomology of a supercharge:

$$[Q, \phi] = 0 \quad (Q\text{-closed})$$

$$\phi \sim \phi + [Q, \psi] \quad (\text{modulo } Q\text{-exact})$$

# Twisting Supersymmetric QFTs

- Twisting SQFT is the procedure of passing to a cohomology of a supercharge:

$$\begin{aligned} [Q, \phi] &= 0 && (Q\text{-closed}) \\ \phi &\sim \phi + [Q, \psi] && (\text{modulo } Q\text{-exact}) \end{aligned}$$

- Physics motivation:

- ▶ Produces a consistent subsector of SQFT

# Twisting Supersymmetric QFTs

- Twisting SQFT is the procedure of passing to a cohomology of a supercharge:

$$\begin{aligned} [Q, \phi] &= 0 && (Q\text{-closed}) \\ \phi &\sim \phi + [Q, \psi] && (\text{modulo } Q\text{-exact}) \end{aligned}$$

- Physics motivation:

- ▶ Produces a consistent subsector of SQFT
- ▶ Restricts to protected (BPS) quantities

# Twisting Supersymmetric QFTs

- Twisting SQFT is the procedure of passing to a cohomology of a supercharge:

$$\begin{aligned} [Q, \phi] &= 0 && (Q\text{-closed}) \\ \phi &\sim \phi + [Q, \psi] && (\text{modulo } Q\text{-exact}) \end{aligned}$$

- Physics motivation:

- ▶ Produces a consistent subsector of SQFT
- ▶ Restricts to protected (BPS) quantities
- ▶ Correlation functions independent of some coordinates:

$$\{Q, \tilde{Q}\} \sim P$$

Eg. **topological** or **holomorphic** twist

# Twisting Supersymmetric QFTs

- Twisting SQFT is the procedure of passing to a cohomology of a supercharge:

$$\begin{aligned} [Q, \phi] &= 0 && (Q\text{-closed}) \\ \phi &\sim \phi + [Q, \psi] && (\text{modulo } Q\text{-exact}) \end{aligned}$$

- Physics motivation:

- ▶ Produces a consistent subsector of SQFT
- ▶ Restricts to protected (BPS) quantities
- ▶ Correlation functions independent of some coordinates:

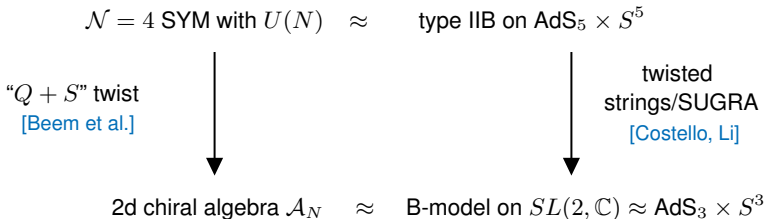
$$\{Q, \tilde{Q}\} \sim P$$

Eg. **topological** or **holomorphic** twist

- Twisted holography: holographic duals of these twists

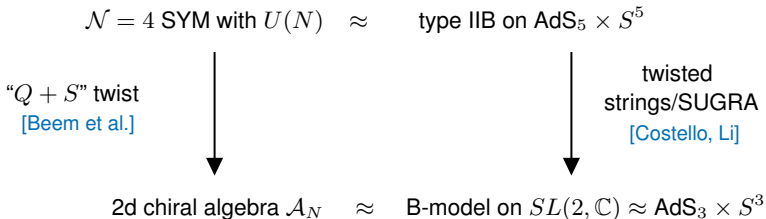
# Twisted Holography

Example: protected subsector of  $\text{AdS}_5/\text{CFT}_4$  [Costello, Gaiotto '18]:



# Twisted Holography

**Example:** protected subsector of  $\text{AdS}_5/\text{CFT}_4$  [Costello, Gaiotto '18]:



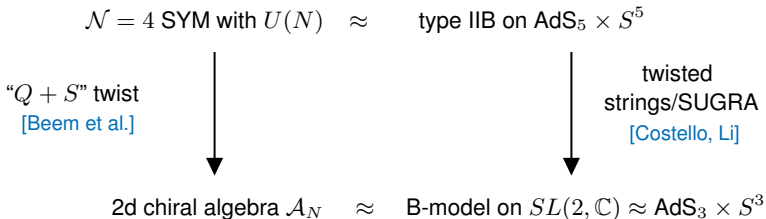
## **Motivation:**

- Many simplifications occur



# Twisted Holography

**Example:** protected subsector of  $\text{AdS}_5/\text{CFT}_4$  [Costello, Gaiotto '18]:



## Motivation:

- Many simplifications occur
- Connections to math

# In this talk

- 1 Review the duality

[Costello, Gaiotto '18]

2d chiral algebra  $\mathcal{A}_N \leftrightarrow$  topological B-model on  $SL(2, \mathbb{C})$

# In this talk

- 1 Review the duality

[Costello, Gaiotto '18]

2d chiral algebra  $\mathcal{A}_N \leftrightarrow$  topological B-model on  $SL(2, \mathbb{C})$

- 2 Correspondence between **determinants** and **Giant Gravitons**
  - ▶ Match saddles of determinant correlation functions with D1-brane configurations
  - ▶ Spectral curve construction

# In this talk

- 1 Review the duality

[Costello, Gaiotto '18]

2d chiral algebra  $\mathcal{A}_N \leftrightarrow$  topological B-model on  $SL(2, \mathbb{C})$

- 2 Correspondence between determinants and Giant Gravitons
  - ▶ Match saddles of determinant correlation functions with D1-brane configurations
  - ▶ Spectral curve construction
- 3 Extend the duality to non-conformal vacua of the chiral algebra

# In this talk

- 1 Review the duality

[Costello, Gaiotto '18]

2d chiral algebra  $\mathcal{A}_N \leftrightarrow$  topological B-model on  $SL(2, \mathbb{C})$

- 2 Correspondence between **determinants** and **Giant Gravitons**
  - ▶ Match saddles of determinant correlation functions with D1-brane configurations
  - ▶ Spectral curve construction
- 3 Extend the duality to **non-conformal vacua** of the chiral algebra
- 4 Future directions:
  - ▶ “Bootstrapping” to  $AdS_5 \times S^5$
  - ▶ LLM type geometries
  - ▶ Holomorphic twist of  $\mathcal{N} = 4$  SYM

# Twisted Holography

Twisted holography example:

[Costello, Gaiotto '18]

chiral algebra  $\mathcal{A}_N$   
gauged  $\beta\gamma$  system in adj. of  $U(N)$   
(large  $N$  expansion of)

$\leftrightarrow$

B-model on  $SL(2, \mathbb{C})$   
(with coupling  $N^{-1}$ )

# Twisted Holography

Twisted holography example:

[Costello, Gaiotto '18]

chiral algebra  $\mathcal{A}_N$   
gauged  $\beta\gamma$  system in adj. of  $U(N)$   
(large  $N$  expansion of)

$\leftrightarrow$

B-model on  $SL(2, \mathbb{C})$   
(with coupling  $N^{-1}$ )

## Simplifications:

- Dependence on t'Hooft coupling drops out

# Twisted Holography

Twisted holography example:

[Costello, Gaiotto '18]

chiral algebra $\mathcal{A}_N$ gauged $\beta\gamma$ system in adj. of $U(N)$ (large $N$ expansion of)	$\leftrightarrow$	B-model on $SL(2, \mathbb{C})$ (with coupling $N^{-1}$ )
---	-------------------	---

## Simplifications:

- Dependence on t'Hooft coupling drops out
- (Almost) **free field** theory computations in the chiral algebra  $\mathcal{A}_N$



# Twisted Holography

Twisted holography example:

[Costello, Gaiotto '18]

chiral algebra $\mathcal{A}_N$ gauged $\beta\gamma$ system in adj. of $U(N)$ (large $N$ expansion of)	$\leftrightarrow$	B-model on $SL(2, \mathbb{C})$ (with coupling $N^{-1}$ )
---	-------------------	---

## Simplifications:

- Dependence on t'Hooft coupling drops out
- (Almost) **free field** theory computations in the chiral algebra  $\mathcal{A}_N$
- D1-branes are **holomorphic curves** in  $SL(2, \mathbb{C})$

# Chiral algebra $\mathcal{A}_N$

- Any 4d  $\mathcal{N} = 2$  SCFT contains a 2d chiral algebra subsector

[Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]

# Chiral algebra $\mathcal{A}_N$

- Any 4d  $\mathcal{N} = 2$  SCFT contains a 2d chiral algebra subsector  
[Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]
- The chiral algebra of  $\mathcal{N} = 4$  SYM is a gauged  $\beta\gamma$  system:

$$X_b^a(z)Y_d^c(0) \sim \delta_d^a \delta_b^c \frac{1}{N} \frac{1}{z}$$
$$Q_{\text{BRST}} \sim N \oint \text{Tr} \left( c[X, Y] + \frac{1}{2} b[c, c] \right)$$

# Chiral algebra $\mathcal{A}_N$

- Any 4d  $\mathcal{N} = 2$  SCFT contains a 2d chiral algebra subsector  
[Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]
- The chiral algebra of  $\mathcal{N} = 4$  SYM is a gauged  $\beta\gamma$  system:

$$X_b^a(z)Y_d^c(0) \sim \delta_d^a \delta_b^c \frac{1}{N} \frac{1}{z}$$
$$Q_{\text{BRST}} \sim N \oint \text{Tr} \left( c[X, Y] + \frac{1}{2} b[c, c] \right)$$

- Free theory computations in the chiral algebra (for BRST closed operators)

# Chiral algebra $\mathcal{A}_N$

- Any 4d  $\mathcal{N} = 2$  SCFT contains a 2d chiral algebra subsector  
[Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]
- The chiral algebra of  $\mathcal{N} = 4$  SYM is a gauged  $\beta\gamma$  system:

$$X_b^a(z)Y_d^c(0) \sim \delta_d^a \delta_b^c \frac{1}{N} \frac{1}{z}$$
$$Q_{\text{BRST}} \sim N \oint \text{Tr} \left( c[X, Y] + \frac{1}{2} b[c, c] \right)$$

- Free theory computations in the chiral algebra (for BRST closed operators)
- For the future, define a linear combination:

$$Z(u; z) \equiv X(z) + uY(z)$$

# Chiral algebra $\mathcal{A}_N$

- Any 4d  $\mathcal{N} = 2$  SCFT contains a 2d chiral algebra subsector  
[Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]
- The chiral algebra of  $\mathcal{N} = 4$  SYM is a gauged  $\beta\gamma$  system:

$$X_b^a(z)Y_d^c(0) \sim \delta_d^a \delta_b^c \frac{1}{N} \frac{1}{z}$$
$$Q_{\text{BRST}} \sim N \oint \text{Tr} \left( c[X, Y] + \frac{1}{2} b[c, c] \right)$$

- Free theory computations in the chiral algebra (for BRST closed operators)
- For the future, define a linear combination:

$$Z(u; z) \equiv X(z) + uY(z)$$

- We will be interested in correlation functions of determinant operators

$$\det(m + Z(u; z))$$

# Topological B-model

- The spacetime theory is Kodaira-Spencer (BCOV) theory on 3d CY
  - ▶ Holomorphic

# Topological B-model

- The spacetime theory is Kodaira-Spencer (BCOV) theory on 3d CY
  - ▶ Holomorphic
  - ▶ Depends only on the complex structure (part of metric) of  $\mathcal{X}$



# Topological B-model

- The spacetime theory is Kodaira-Spencer (BCOV) theory on 3d CY
  - ▶ Holomorphic
  - ▶ Depends only on the complex structure (part of metric) of  $\mathcal{X}$
  - ▶ Fields are poly-vectorfields

$$PV^{j,i}(\mathcal{X}) = \Omega^{(0,i)}(\mathcal{X}, \wedge^j T\mathcal{X}) \quad (\text{locally } f_{m\dots}^{n\dots} d\bar{z}_n \dots \partial_{z_m} \dots)$$

# Topological B-model

- The spacetime theory is Kodaira-Spencer (BCOV) theory on 3d CY
  - ▶ Holomorphic
  - ▶ Depends only on the complex structure (part of metric) of  $\mathcal{X}$
  - ▶ Fields are poly-vectorfields

$$PV^{j,i}(\mathcal{X}) = \Omega^{(0,i)}(\mathcal{X}, \wedge^j T\mathcal{X}) \quad (\text{locally } f_m^{n\dots} d\bar{z}_n \dots \partial_{z_m} \dots)$$

For example,

$$\beta \in PV^{1,1}(\mathcal{X})$$

is a Beltrami differential which deforms the complex structure

# Topological B-model

- The spacetime theory is Kodaira-Spencer (BCOV) theory on 3d CY
  - ▶ Holomorphic
  - ▶ Depends only on the complex structure (part of metric) of  $\mathcal{X}$
  - ▶ Fields are poly-vectorfields

$$PV^{j,i}(\mathcal{X}) = \Omega^{(0,i)}(\mathcal{X}, \wedge^j T\mathcal{X}) \quad (\text{locally } f_m^{n\dots} d\bar{z}_n \dots \partial_{z_m} \dots)$$

For example,

$$\beta \in PV^{1,1}(\mathcal{X})$$

is a Beltrami differential which deforms the complex structure

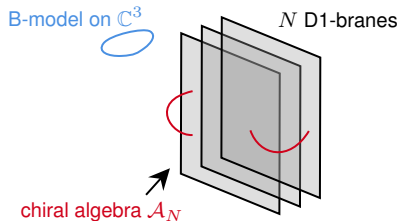
- D-branes wrap holomorphic submanifolds eg. D1-branes are holomorphic complex lines

# Twisted Holography

- 1 Consider topological B-model in flat space  $\mathcal{X} = \mathbb{C}^3$

# Twisted Holography

- 1 Consider topological B-model in flat space  $\mathcal{X} = \mathbb{C}^3$
- 2 The chiral algebra  $\mathcal{A}_N$  is supported by  $N$  D1-branes wrapping  $\mathbb{C} \subset \mathbb{C}^3$



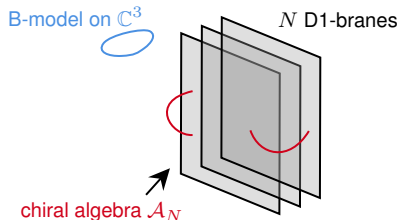
# Twisted Holography

1 Consider topological B-model in flat space  $\mathcal{X} = \mathbb{C}^3$

2 The chiral algebra  $\mathcal{A}_N$  is supported by  $N$  D1-branes wrapping  $\mathbb{C} \subset \mathbb{C}^3$

3 The stack of branes sources a **Beltrami differential** which deforms the complex structure:

$$\mathbb{C}^3 \setminus \mathbb{C} \rightarrow SL(2, \mathbb{C})$$

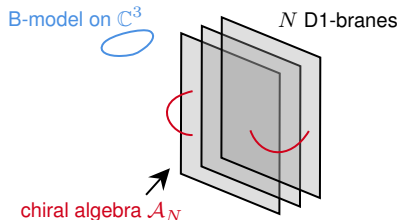


# Twisted Holography

1 Consider topological B-model in flat space  $\mathcal{X} = \mathbb{C}^3$

2 The chiral algebra  $\mathcal{A}_N$  is supported by  $N$  D1-branes wrapping  $\mathbb{C} \subset \mathbb{C}^3$

3 The stack of branes sources a **Beltrami differential** which deforms the complex structure:



$$\mathbb{C}^3 \setminus \mathbb{C} \rightarrow SL(2, \mathbb{C})$$

[Costello, Gaiotto '18]

$$\text{B-model on } \mathbb{C}^3 + N \text{ D1-branes} \longrightarrow \text{B-model on } SL(2, \mathbb{C}) \approx \text{AdS}_3 \times S^3$$

$$\begin{array}{c} \uparrow \\ \mathcal{A}_N \end{array}$$

# Giant Gravitons

- Determinant operators in  $\mathcal{N} = 4$  SYM are dual to D3-branes wrapping

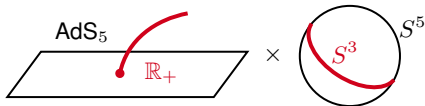
$$\mathbb{R}_+ \times S^3 \text{ inside } \text{AdS}_5 \times S^5$$



# Giant Gravitons

- Determinant operators in  $\mathcal{N} = 4$  SYM are dual to D3-branes wrapping

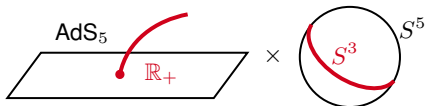
$$\mathbb{R}_+ \times S^3 \text{ inside } \text{AdS}_5 \times S^5$$



# Giant Gravitons

- Determinant operators in  $\mathcal{N} = 4$  SYM are dual to D3-branes wrapping

$$\mathbb{R}_+ \times S^3 \text{ inside } \text{AdS}_5 \times S^5$$



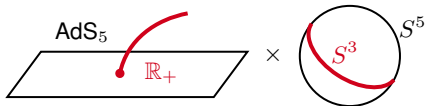
- Determinant operators in chiral algebra  $\mathcal{A}_N$  are dual to D1-branes wrapping

$$\mathbb{C}^* \equiv \mathbb{R}_+ \times S^1 \text{ inside } SL(2, \mathbb{C}) \equiv \text{AdS}_3 \times S^3$$

# Giant Gravitons

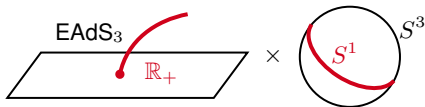
- Determinant operators in  $\mathcal{N} = 4$  SYM are dual to D3-branes wrapping

$$\mathbb{R}_+ \times S^3 \text{ inside } \text{AdS}_5 \times S^5$$



- Determinant operators in chiral algebra  $\mathcal{A}_N$  are dual to D1-branes wrapping

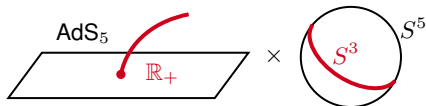
$$\mathbb{C}^* \equiv \mathbb{R}_+ \times S^1 \text{ inside } SL(2, \mathbb{C}) \equiv \text{AdS}_3 \times S^3$$



# Giant Gravitons

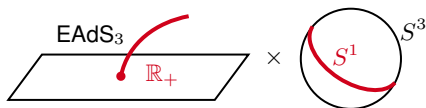
- Determinant operators in  $\mathcal{N} = 4$  SYM are dual to D3-branes wrapping

$$\mathbb{R}_+ \times S^3 \text{ inside } \text{AdS}_5 \times S^5$$



- Determinant operators in chiral algebra  $\mathcal{A}_N$  are dual to D1-branes wrapping

$$\mathbb{C}^* \equiv \mathbb{R}_+ \times S^1 \text{ inside } SL(2, \mathbb{C}) \equiv \text{AdS}_3 \times S^3$$



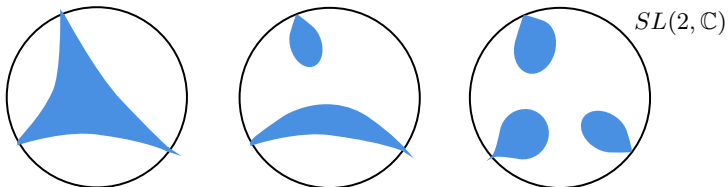
- Determinant operator:

$$\det(m + Z(u; z)), \quad Z(u; z) = X(z) + uY(z)$$

- $z$  = position at the boundary of  $\text{AdS}_3$
- $u$  controls orientation of  $S^1 \subset S^3$
- $m$  controls size of  $S^1 \subset S^3$

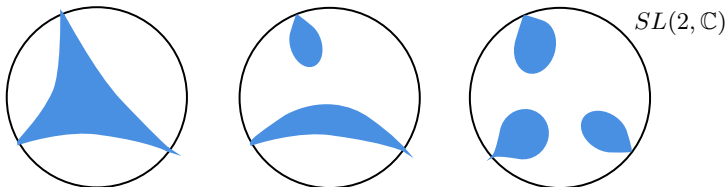
# Giant Gravitons

- Many possible brane configurations with the same boundary behaviour



# Giant Gravitons

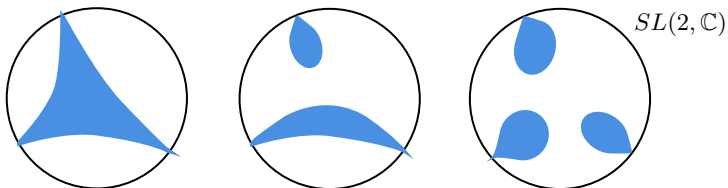
- Many possible brane configurations with the same boundary behaviour



- We will match saddles  $\rho^*$  of correlation functions of determinants with brane configurations

# Giant Gravitons

- Many possible brane configurations with the same boundary behaviour



- We will match saddles  $\rho^*$  of correlation functions of determinants with brane configurations
  - ▶  $m_i, u_i, z_i$  control boundary behaviour
  - ▶ Saddles  $\rho$  will control the shape in the bulk

# Determinant correlation functions

[Jiang, Komatsu, Vescovi '19]

- Fermionize determinants

$$\det(m + Z(u; z)) = \int [d\bar{\psi}d\psi] e^{\bar{\psi}(m+Z(u,z))\psi}$$



# Determinant correlation functions

[Jiang, Komatsu, Vescovi '19]

- Fermionize determinants

$$\det(m + Z(u; z)) = \int [d\bar{\psi}d\psi] e^{\bar{\psi}(m+Z(u,z))\psi}$$

- Rewrite correlation function using auxiliary bosonic variables  $\rho_j^i$  for  $i \neq j$ ,  
 $\rho_i^i \equiv m_i$

$$\left\langle \prod_i^k \det(m_i + Z(u_i; z_i)) \right\rangle \sim \int [d\rho] e^{N S[\rho]}$$

with action

$$S[\rho] = \frac{1}{2} \sum_{i \neq j} \frac{z_i - z_j}{u_i - u_j} \rho_j^i \rho_i^j + \log \det \rho$$

# Saddles and branes

- Saddle point equations in the matrix form:

$$[\zeta, \rho] + [\mu, \rho^{-1}] = 0$$

where

$$\zeta = \begin{pmatrix} z_1 & & \\ & \ddots & \\ & & z_k \end{pmatrix}, \quad \mu = \begin{pmatrix} u_1 & & \\ & \ddots & \\ & & u_k \end{pmatrix}, \quad \rho = \begin{pmatrix} m_1 & & ? \\ & \ddots & \\ ? & & m_k \end{pmatrix}$$

# Saddles and branes

- Saddle point equations in the matrix form:

$$[\zeta, \rho] + [\mu, \rho^{-1}] = 0$$

where

$$\zeta = \begin{pmatrix} z_1 & & \\ & \ddots & \\ & & z_k \end{pmatrix}, \quad \mu = \begin{pmatrix} u_1 & & \\ & \ddots & \\ & & u_k \end{pmatrix}, \quad \rho = \begin{pmatrix} m_1 & & ? \\ & \ddots & \\ ? & & m_k \end{pmatrix}$$

- We will match saddles  $\rho$  to classical brane configurations in B-model on  $SL(2, \mathbb{C})$

# Saddles and branes

- Saddle point equations in the matrix form:

$$[\zeta, \rho] + [\mu, \rho^{-1}] = 0$$

where

$$\zeta = \begin{pmatrix} z_1 & & \\ & \ddots & \\ & & z_k \end{pmatrix}, \quad \mu = \begin{pmatrix} u_1 & & \\ & \ddots & \\ & & u_k \end{pmatrix}, \quad \rho = \begin{pmatrix} m_1 & & ? \\ & \ddots & \\ ? & & m_k \end{pmatrix}$$

- We will match saddles  $\rho$  to classical brane configurations in B-model on  $SL(2, \mathbb{C})$
- For each  $\rho$  we will define a spectral curve  $S_\rho$  in  $SL(2, \mathbb{C})$

# Saddles and branes

- Saddle point equations in the matrix form:

$$[\zeta, \rho] + [\mu, \rho^{-1}] = 0$$

where

$$\zeta = \begin{pmatrix} z_1 & & \\ & \ddots & \\ & & z_k \end{pmatrix}, \quad \mu = \begin{pmatrix} u_1 & & \\ & \ddots & \\ & & u_k \end{pmatrix}, \quad \rho = \begin{pmatrix} m_1 & & ? \\ & \ddots & \\ ? & & m_k \end{pmatrix}$$

- We will match saddles  $\rho$  to classical brane configurations in B-model on  $SL(2, \mathbb{C})$
- For each  $\rho$  we will define a spectral curve  $S_\rho$  in  $SL(2, \mathbb{C})$
- We check it matches dual Giant Graviton brane

# Saddles and branes

- Saddle point equations in the matrix form:

$$[\zeta, \rho] + [\mu, \rho^{-1}] = 0$$

where

$$\zeta = \begin{pmatrix} z_1 & & \\ & \ddots & \\ & & z_k \end{pmatrix}, \quad \mu = \begin{pmatrix} u_1 & & \\ & \ddots & \\ & & u_k \end{pmatrix}, \quad \rho = \begin{pmatrix} m_1 & & ? \\ & \ddots & \\ ? & & m_k \end{pmatrix}$$

- We will match saddles  $\rho$  to classical brane configurations in B-model on  $SL(2, \mathbb{C})$
- For each  $\rho$  we will define a spectral curve  $S_\rho$  in  $SL(2, \mathbb{C})$
- We check it matches dual Giant Graviton brane
- $k > 2$  would be hard in  $AdS_5 \times S^5$

# Spectral curve

For each saddle  $\rho$  we define a spectral curve  $S_\rho$ :

- Define **commuting matrices**:

$$B(a) = a\mu - \rho, \quad C(a) = a\zeta + \rho^{-1}, \quad D(a) = a\zeta\mu + \rho^{-1}\mu - \zeta\rho,$$

# Spectral curve

For each saddle  $\rho$  we define a spectral curve  $S_\rho$ :

- Define **commuting matrices**:

$$B(a) = a\mu - \rho, \quad C(a) = a\zeta + \rho^{-1}, \quad D(a) = a\zeta\mu + \rho^{-1}\mu - \zeta\rho,$$

- Define spectral curve:

$$\mathcal{S}_\rho = \{(a, b, c, d) \\ \text{s.t. } b, c, d \text{ are **simultaneous eigenvalues** of } B(a), C(a), D(a)\}$$



# Spectral curve

For each saddle  $\rho$  we define a spectral curve  $S_\rho$ :

- Define **commuting matrices**:

$$B(a) = a\mu - \rho, \quad C(a) = a\zeta + \rho^{-1}, \quad D(a) = a\zeta\mu + \rho^{-1}\mu - \zeta\rho,$$

- Define spectral curve:

$$\mathcal{S}_\rho = \{(a, b, c, d) \\ \text{s.t. } b, c, d \text{ are **simultaneous eigenvalues** of } B(a), C(a), D(a)\}$$

- The matrices are defined so that:

- ▶ They commute when  $\rho$  satisfies the saddle point equations
- ▶ They satisfy

$$aD(a) - B(a)C(a) = 1$$

- ▶ They have the expected boundary behavior

# Holographic checks

- Boundary behaviour  $a \rightarrow \infty$ :

$$\frac{B(a)}{a} = \left( \begin{array}{c} u_1 - \frac{m_1}{a} \\ \dots \\ u_k - \frac{m_k}{a} \end{array} \right) + \dots, \quad \frac{C(a)}{a} = \left( \begin{array}{c} z_1 \\ \dots \\ z_k \end{array} \right) + \dots$$

# Holographic checks

- Boundary behaviour  $a \rightarrow \infty$ :

$$\frac{B(a)}{a} = \left( \begin{array}{ccc} u_1 - \frac{m_1}{a} & & \\ & \ddots & \\ & & u_k - \frac{m_k}{a} \end{array} \right) + \dots, \quad \frac{C(a)}{a} = \left( \begin{array}{ccc} z_1 & & \\ & \ddots & \\ & & z_k \end{array} \right) + \dots$$

$\implies$  spectral curve  $S_\rho$  goes to the boundary of  $SL(2, \mathbb{C})$  in  $k$  points

# Holographic checks

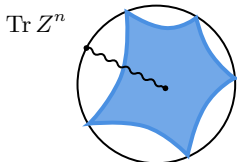
- Boundary behaviour  $a \rightarrow \infty$ :

$$\frac{B(a)}{a} = \left( \begin{array}{ccc} u_1 - \frac{m_1}{a} & & \\ & \ddots & \\ & & u_k - \frac{m_k}{a} \end{array} \right) + \dots, \quad \frac{C(a)}{a} = \left( \begin{array}{c} z_1 \\ \ddots \\ z_k \end{array} \right) + \dots$$

$\implies$  spectral curve  $S_\rho$  goes to the boundary of  $SL(2, \mathbb{C})$  in  $k$  points

- Various holographic checks:

- ▶ Correlation functions of determinants with a single-trace [Jiang, Komatsu, Vescovi '19]
- ▶ Action  $S[\rho]$  vs  $S[\text{brane}]$
- ▶ Modifications of determinants / excitations of the brane



# Determinant modifications

- For example

$$\det X \rightarrow \frac{1}{N!} \varepsilon \varepsilon(X, X, X, \dots, Y^2)$$

---

\* There are also 3 other types of generators but we focus on one tower.

# Determinant modifications

- For example

$$\det X \rightarrow \frac{1}{N!} \varepsilon \varepsilon (X, X, X, \dots, Y^2)$$

- Employ the global symmetry algebra<sup>\*</sup> of  $\mathcal{A}_N$ :

$$\oint z^k \operatorname{Tr} Z^{(i_1} Z^{i_2} \dots Z^{i_n)}, \quad 0 \leq k \leq n-2$$

---

<sup>\*</sup> There are also 3 other types of generators but we focus on one tower.

# Determinant modifications

- For example

$$\det X \rightarrow \frac{1}{N!} \varepsilon \varepsilon (X, X, X, \dots, Y^2)$$

- Employ the global symmetry algebra<sup>\*</sup> of  $\mathcal{A}_N$ :

$$\oint z^k \operatorname{Tr} Z^{i_1} Z^{i_2} \dots Z^{i_n}, \quad 0 \leq k \leq n-2$$

- For example, the lowest modes are the  $su(2)$  generators:

$$\oint \operatorname{Tr} X X, \quad \oint \operatorname{Tr} X Y, \quad \oint \operatorname{Tr} Y Y$$

---

<sup>\*</sup> There are also 3 other types of generators but we focus on one tower.

# Determinant modifications

- For example

$$\det X \rightarrow \frac{1}{N!} \varepsilon \varepsilon(X, X, X, \dots, Y^2)$$

- Employ the global symmetry algebra\* of  $\mathcal{A}_N$ :

$$\oint z^k \operatorname{Tr} Z^{i_1} Z^{i_2} \dots Z^{i_n}, \quad 0 \leq k \leq n-2$$

- For example, the lowest modes are the  $su(2)$  generators:

$$\oint \operatorname{Tr} X X, \quad \oint \operatorname{Tr} X Y, \quad \oint \operatorname{Tr} Y Y$$

- Create modifications by acting with the modes, eg.

$$\oint \operatorname{Tr} Y^4(z) \det X(0) \sim \varepsilon \varepsilon(X, X, X, \dots, Y^3) + \dots$$

---

\* There are also 3 other types of generators but we focus on one tower.



# Determinant modifications

- We computed 2-pt functions of determinant modifications

$$\left\langle [J_{p',q'}^{(n')}, \det Y(\infty)] [J_{p,q}^{(n)}, \det X(0)] \right\rangle \Big|_{N \rightarrow \infty}$$

# Determinant modifications

- We computed 2-pt functions of determinant modifications

$$\left\langle [J_{p',q'}^{(n')}, \det Y(\infty)] [J_{p,q}^{(n)}, \det X(0)] \right\rangle \Big|_{N \rightarrow \infty}$$

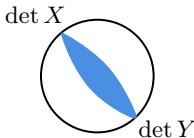
- There are only two types of determinant modifications:

$$\begin{aligned} J_{p,p-1}^n &: \det X \longrightarrow n \varepsilon \varepsilon(X, X, \dots, Y^{1-2p}) \\ J_{p,p+1}^n &: \det X \longrightarrow n \varepsilon \varepsilon(X, X, \dots, Y^{-2p-1} \partial X) \\ &\quad + n \varepsilon \varepsilon(X, X, \dots, \partial^2 Y^{-2p-3}) \end{aligned}$$

# Brane excitations

- $\langle \det Y(\infty) \det X(0) \rangle$  has a single nontrivial saddle corresponding to the brane:

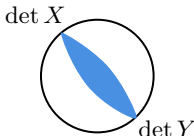
$$\begin{pmatrix} a & 0 \\ 0 & 1/a \end{pmatrix} \in SL(2, \mathbb{C})$$



# Brane excitations

- $\langle \det Y(\infty) \det X(0) \rangle$  has a single nontrivial saddle corresponding to the brane:

$$\begin{pmatrix} a & 0 \\ 0 & 1/a \end{pmatrix} \in SL(2, \mathbb{C})$$



- Holographic global symmetry algebra acts by

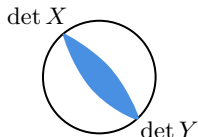
[Costello, Gaiotto]

holomorphic divergence-free vector fields on  $SL(2, \mathbb{C})$

# Brane excitations

- $\langle \det Y(\infty) \det X(0) \rangle$  has a single nontrivial saddle corresponding to the brane:

$$\begin{pmatrix} a & 0 \\ 0 & 1/a \end{pmatrix} \in SL(2, \mathbb{C})$$



- Holographic global symmetry algebra acts by [Costello, Gaiotto]  
holomorphic divergence-free vector fields on  $SL(2, \mathbb{C})$

- There are only two types of brane excitations:

$$J_{p,p-1}^{(n)} : \begin{pmatrix} a & 0 \\ 0 & 1/a \end{pmatrix} \longrightarrow \begin{pmatrix} a & \delta b \\ 0 & 1/a \end{pmatrix}, \quad \delta b \sim +na^{1-2p}$$

$$J_{p,p+1}^{(n)} : \begin{pmatrix} a & 0 \\ 0 & 1/a \end{pmatrix} \longrightarrow \begin{pmatrix} a & 0 \\ \delta c & 1/a \end{pmatrix}, \quad \delta c \sim -na^{-1-2p}$$

# Coulomb branch geometries

- Duality can be extended to non-conformal vacua of the chiral algebra  $\mathcal{A}_N$

# Coulomb branch geometries

- Duality can be extended to non-conformal vacua of the chiral algebra  $\mathcal{A}_N$
- Twisted analog of

Coulomb branch of  $\mathcal{N} = 4$  SYM  $\longleftrightarrow$  multi-center solutions

# Coulomb branch geometries

- Duality can be extended to non-conformal vacua of the chiral algebra  $\mathcal{A}_N$
- Twisted analog of

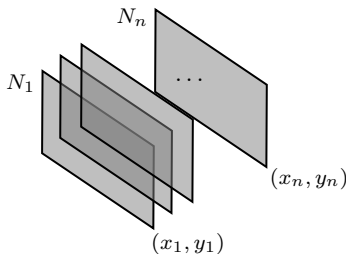
Coulomb branch of  $\mathcal{N} = 4$  SYM  $\longleftrightarrow$  multi-center solutions

- Backreact stack of non-coincident branes
- Dual Calabi-Yau geometries are deformations of  $SL(2, \mathbb{C})$

$$z_I - z_{I'} = + \frac{N_i/N}{(x - x_i)(y - y_i)}$$

For standard  $SL(2, \mathbb{C})$  geometry:

$$z_0 - z_\infty = \frac{1}{xy}$$





# Coulomb branch geometries

- Duality can be extended to non-conformal vacua of the chiral algebra  $\mathcal{A}_N$
- Twisted analog of

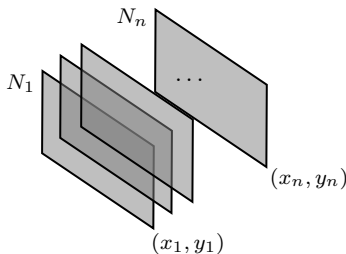
Coulomb branch of  $\mathcal{N} = 4$  SYM  $\longleftrightarrow$  multi-center solutions

- Backreact stack of non-coincident branes
- Dual Calabi-Yau geometries are deformations of  $SL(2, \mathbb{C})$

$$z_I - z_{I'} = + \frac{N_i/N}{(x - x_i)(y - y_i)}$$

For standard  $SL(2, \mathbb{C})$  geometry:

$$z_0 - z_\infty = \frac{1}{xy}$$



- Holographic check:
  - ▶ Determinant correlation functions (with a single-trace) and dual Giant Graviton branes

# Future directions

- Spectral curve construction in other examples of twisted or free field holography

# Future directions

- Spectral curve construction in other examples of twisted or free field holography
- Find SUSY D3-branes in  $\text{AdS}_5 \times S^5$  that correspond to our B-model D1-branes

# Future directions

- Spectral curve construction in other examples of twisted or free field holography
- Find SUSY D3-branes in  $\text{AdS}_5 \times S^5$  that correspond to our B-model D1-branes
- Mathematical question:

Solutions of matrix equations $[\zeta, \rho] + [\mu, \rho^{-1}] = 0$	$\iff$	Holomorphic curves in $SL(2, \mathbb{C})$
---	--------	--

- $\Rightarrow$  Spectral curve construction
- $\Leftarrow$  For genus  $g = 0$ , we can go back
- $\Leftarrow$  For genus  $g > 0$ , we don't know

## Future directions

- Consider  $\mathcal{O}(N^2)$  operators eg.  $(\det Z)^N$ , which are dual to backreacted geometries

# Future directions

- Consider  $\mathcal{O}(N^2)$  operators eg.  $(\det Z)^N$ , which are dual to backreacted geometries

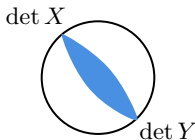
New holomorphic coordinates:

$$\tilde{a}_1 = a e^{-\frac{\bar{c}}{b} \frac{1}{|b|^2 + |c|^2}}, \quad b \neq 0$$

$$\tilde{a}_2 = a e^{\frac{\bar{b}}{c} \frac{1}{|b|^2 + |c|^2}}, \quad c \neq 0$$

with transition:

$$\frac{\tilde{a}_1}{\tilde{a}_2} = e^{\frac{1}{bc}}$$



# Future directions

- Consider  $\mathcal{O}(N^2)$  operators eg.  $(\det Z)^N$ , which are dual to backreacted geometries

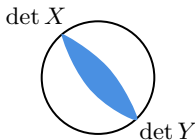
New holomorphic coordinates:

$$\tilde{a}_1 = a e^{-\frac{\bar{c}}{b} \frac{1}{|b|^2 + |c|^2}}, \quad b \neq 0$$

$$\tilde{a}_2 = a e^{\frac{\bar{b}}{c} \frac{1}{|b|^2 + |c|^2}}, \quad c \neq 0$$

with transition:

$$\frac{\tilde{a}_1}{\tilde{a}_2} = e^{\frac{1}{bc}}$$



- Holographic dual of the holomorphic twist of  $\mathcal{N} = 4$  SYM

# Future directions

- Consider  $\mathcal{O}(N^2)$  operators eg.  $(\det Z)^N$ , which are dual to backreacted geometries

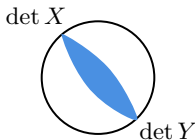
New holomorphic coordinates:

$$\tilde{a}_1 = a e^{-\frac{\bar{c}}{b} \frac{1}{|b|^2 + |c|^2}}, \quad b \neq 0$$

$$\tilde{a}_2 = a e^{\frac{\bar{b}}{c} \frac{1}{|b|^2 + |c|^2}}, \quad c \neq 0$$

with transition:

$$\frac{\tilde{a}_1}{\tilde{a}_2} = e^{\frac{1}{bc}}$$



- Holographic dual of the holomorphic twist of  $\mathcal{N} = 4$  SYM
  - ▶ Proposed to be topological B-model on  $\mathbb{C}^5$ , in the presence of a certain background field [Costello, Li '16]



# Future directions

- Consider  $\mathcal{O}(N^2)$  operators eg.  $(\det Z)^N$ , which are dual to backreacted geometries

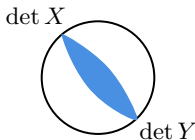
New holomorphic coordinates:

$$\tilde{a}_1 = a e^{-\frac{\bar{c}}{b} \frac{1}{|b|^2 + |c|^2}}, \quad b \neq 0$$

$$\tilde{a}_2 = a e^{\frac{\bar{b}}{c} \frac{1}{|b|^2 + |c|^2}}, \quad c \neq 0$$

with transition:

$$\frac{\tilde{a}_1}{\tilde{a}_2} = e^{\frac{1}{bc}}$$



- Holomorphic dual of the holomorphic twist of  $\mathcal{N} = 4$  SYM
  - ▶ Proposed to be topological B-model on  $\mathbb{C}^5$ , in the presence of a certain background field [Costello, Li '16]
  - ▶ Maybe useful for "non-multigraviton" cohomology classes?

# Future directions

- Consider  $\mathcal{O}(N^2)$  operators eg.  $(\det Z)^N$ , which are dual to backreacted geometries

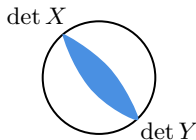
New holomorphic coordinates:

$$\tilde{a}_1 = a e^{-\frac{\bar{c}}{b} \frac{1}{|b|^2 + |c|^2}}, \quad b \neq 0$$

$$\tilde{a}_2 = a e^{\frac{\bar{b}}{c} \frac{1}{|b|^2 + |c|^2}}, \quad c \neq 0$$

with transition:

$$\frac{\tilde{a}_1}{\tilde{a}_2} = e^{\frac{1}{bc}}$$



- Holomorphic dual of the holomorphic twist of  $\mathcal{N} = 4$  SYM
  - ▶ Proposed to be topological B-model on  $\mathbb{C}^5$ , in the presence of a certain background field [Costello, Li '16]
  - ▶ Maybe useful for "non-multigraviton" cohomology classes?

Thank you!

# Holographic checks

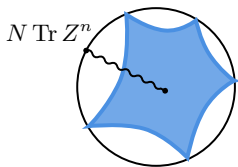
Correlation function of determinants with a single-trace

[Jiang, Komatsu, Vescovi '19]:

$$\left\langle \prod_i \mathcal{D}(m_i; u_i; z_i) N \operatorname{Tr} Z^n \right\rangle \Big|_{N \rightarrow \infty} \approx -e^{NS[\rho^*]} N \operatorname{Tr}_{k \times k} \left( -\rho \frac{\mu - u}{\zeta - z} \right)^n \Big|_{\rho = \rho^*}$$

We can rewrite it to a form

$$\int_{S_{\rho^*}} \partial^{-1} \alpha,$$



where  $\alpha$  is a Kodaira-Spencer field sourced by  $N \operatorname{Tr} Z^n$ :

$$\alpha = \partial \left( (b - ua)^n \delta_{\frac{c}{a} = z} + (za - c)^{-n} \delta_{\frac{b}{a} = u} \right)$$