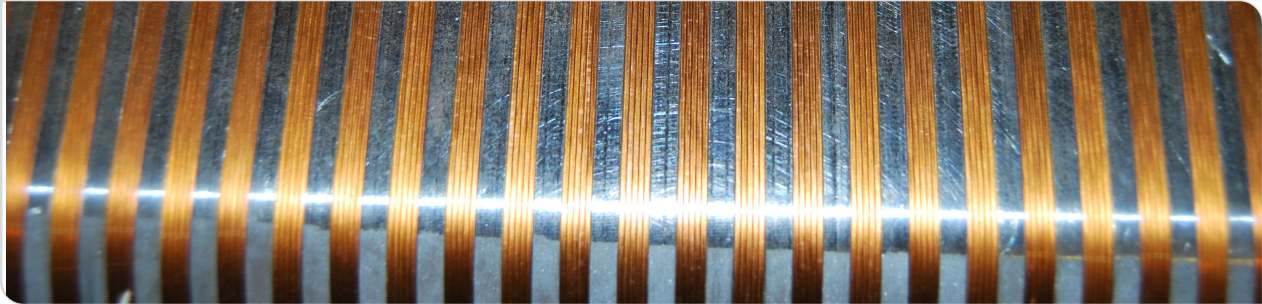
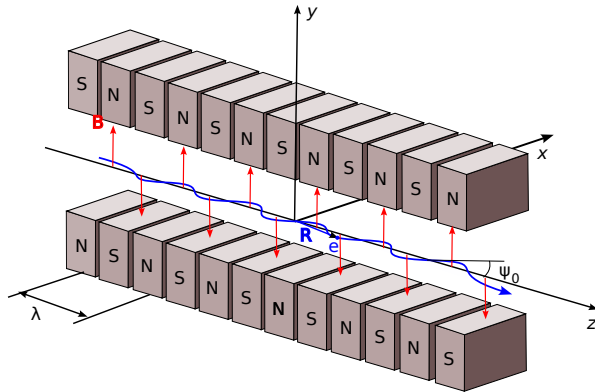


Insertion Devices

Axel Bernhard | 2023-12-01

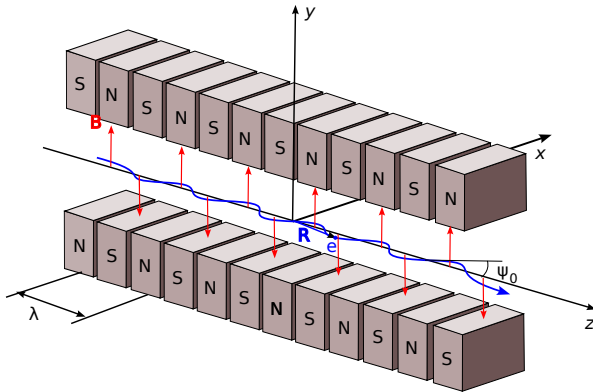


The basic idea of insertion devices



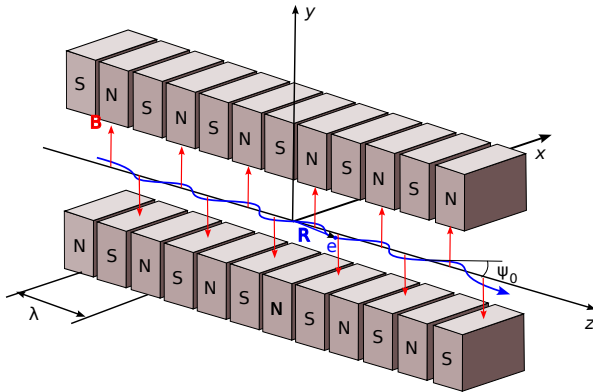
- Magnet in an insertion
- not being part of the lattice

The basic idea of insertion devices



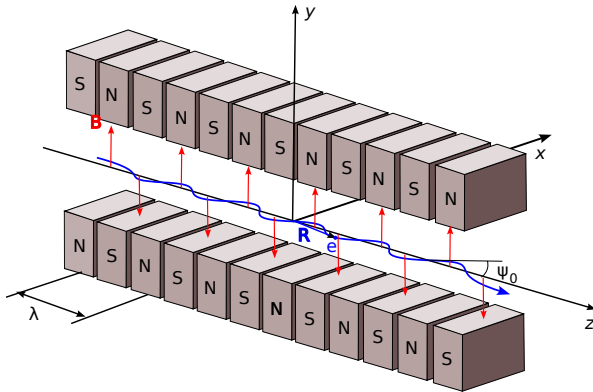
- Magnet in an insertion
- not being part of the lattice
- typically a periodic array with
 - period length λ_u
 - number of periods N_u
 - length $L_u = N_u \lambda_u$
 - zero net deflection: $x'(-L_u/2) = x'(L_u/2)$
 - zero net displacement: $x(-L_u/2) = x(L_u/2)$

The basic idea of insertion devices



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 - zero net displacement: $x(-L_u/2) = x(L_u/2)$
- aim: generate enhanced radiation by multiple deflection

The basic idea of insertion devices

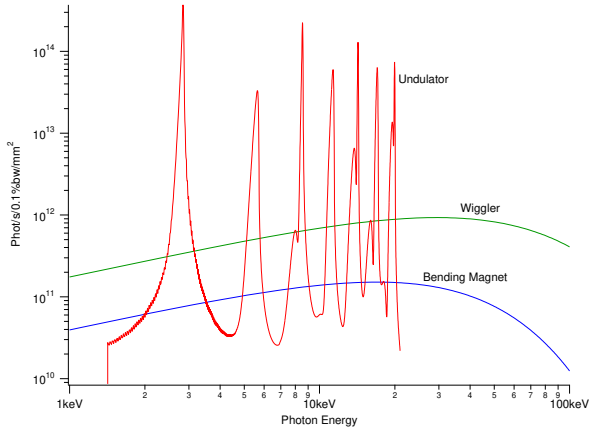


Classification with respect to undulator parameter

$$K_u \approx \frac{\psi_0}{1/\gamma}$$

- $K_u \gg 1$: *Wiggler*
- $K_u \lesssim 1$: *Undulator*

The basic idea of insertion devices



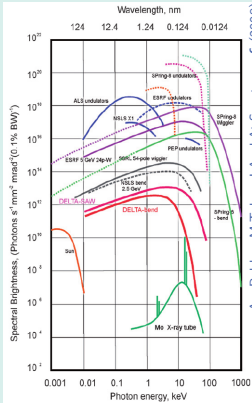
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The purpose of IDs

Generation of brilliant synchrotron radiation



A.v. Bohlen, M. Tolan, J.-Anal. At-Spectrom. 6 (2008)

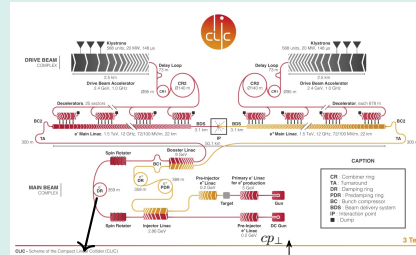


<https://www.esrf.fr>

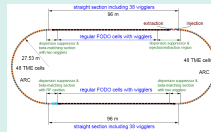


www.xfel.eu

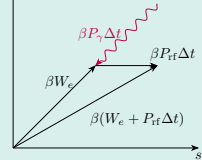
Radiation damping



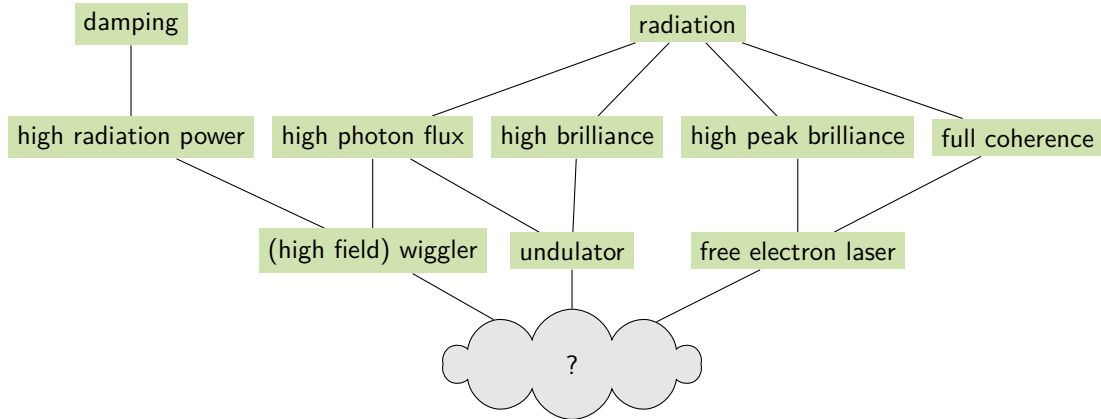
<https://cds.cern.ch/record/2655160>



M. Aichele (ed.), CLIC conceptual design report (2012)



From use case to design requirements



Outline

1. Understanding undulator radiation
2. Understanding the FEL mechanism — Phase requirements
3. Technological boundary conditions for ID magnetic design
4. Aspects of ID 3D and technical design

Outline

1. Understanding undulator radiation

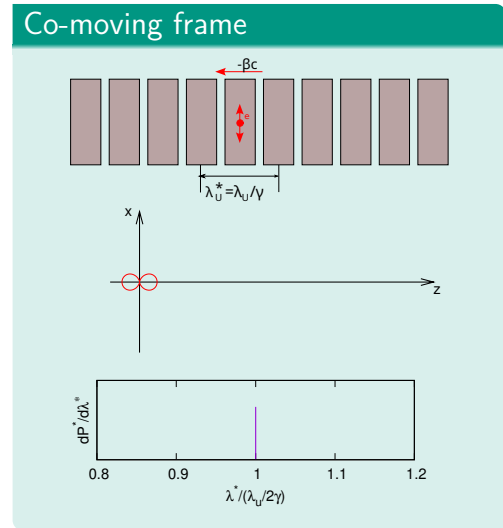
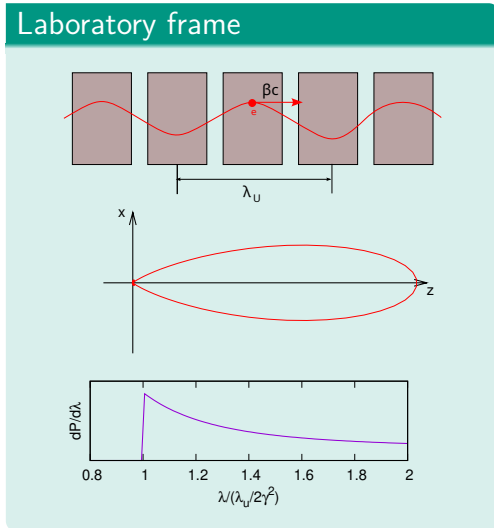
- Planar undulator and wiggler radiation
- Helical undulator radiation

2. Understanding the FEL mechanism — Phase requirements

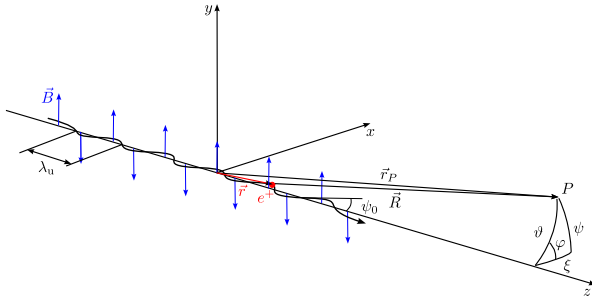
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4. Aspects of ID 3D and technical design

The basic principle



The particle trajectory I



Ideal planar undulator:

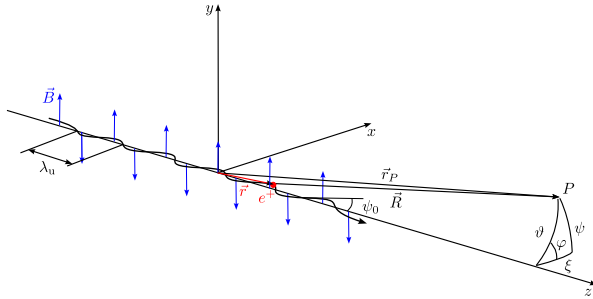
$$\vec{B}(x, 0, z) = B_0(0, \cos k_u z, 0)$$

with $k_u = \frac{2\pi}{\lambda_u}$

Equations of motion: Lorentz force

$$\vec{F} = m\gamma \begin{pmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \end{pmatrix} = e(\vec{v} \times \vec{B}) = eB_0 \begin{pmatrix} -\cos(k_u z)\dot{z} \\ 0 \\ \cos(k_u z)\dot{x} \end{pmatrix}$$

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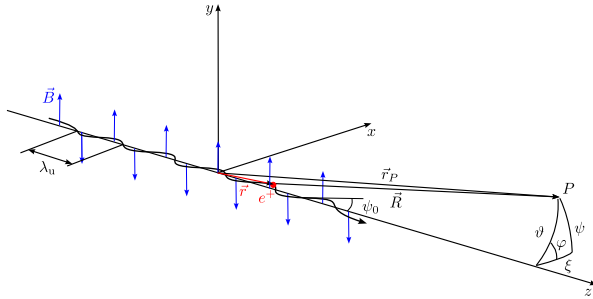
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Integration for x component:

$$\dot{x} = -\frac{eB_0}{m\gamma k_u} \sin(k_u z).$$

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z-component: energy conservation:

$$\dot{x}^2 + \dot{z}^2 = \beta^2 c^2 \Rightarrow \dot{z} = \beta c \sqrt{1 - \frac{\dot{x}^2}{\beta^2 c^2}}$$

The particle trajectory II

Trajectory as function of time

$$x(t') = \frac{K_u}{\beta\gamma k_u} \cos(\Omega_u t')$$

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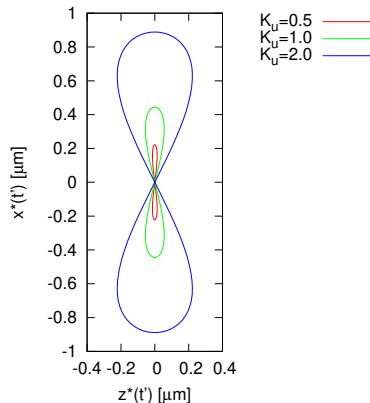
$$z(t') = \beta^* c t' + \frac{K_u^2}{8\beta^2 \gamma^2 k_u} \sin(2\Omega_u t')$$

with

$$\beta^* c = \beta c \left(1 - \frac{K_u^2}{4\beta^2 \gamma^2} \right) < \beta c!$$

$$2\Omega_u = 2k_u \beta^* c$$

$$K_u := \frac{eB_0}{mck_u}$$



Particle trajectory in the co-moving frame with $v = \beta^* c$; numbers given for $\lambda_u = 14$ mm, $W_e = 2.5$ GeV.

The radiation field

Trajectory + Liénard Wiechert potentials:

$$\vec{E}(\vec{r}_P, t) = \frac{e}{4\pi\epsilon_0} \left[\frac{1}{c} \frac{\vec{n} \times [(\vec{n} - \vec{\beta}) \times \dot{\vec{\beta}}]}{(1 - \vec{\beta} \cdot \vec{n})^3 R} \right]_{\text{ret}}$$

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Consequence of reduced mean drift velocity and figure-8 motion:

complex periodic motion \Rightarrow line spectrum with harmonics h of fundamental wavelength

$$\lambda_1 = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K_u^2}{2} + \gamma^2 \vartheta^2 \right)$$

The undulator equation.

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Spectral angular power distribution

$$\frac{d^2 P_h}{d\Omega d\omega} = P_u \gamma^{*2} [F_{h\sigma}(\vartheta, \varphi) + F_{h\pi}(\vartheta, \varphi)] f_N(\Delta\omega_h)$$

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Total power

$$P_u = \frac{2}{3} \frac{e^4 c^3}{4\pi\epsilon_0} \frac{\langle B^2 \rangle W_e^2}{(mc^2)^4}$$

The radiation field

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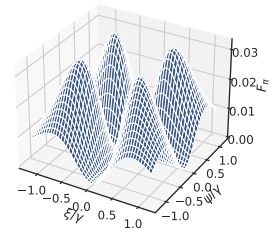
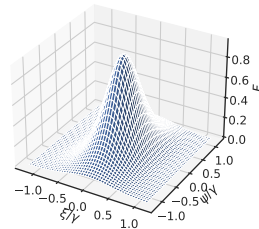
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Normalized angular power distribution 1st harmonic for σ and π polarization, respectively:



The radiation field

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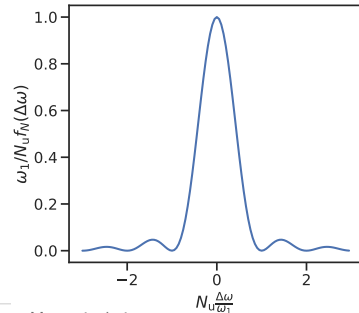
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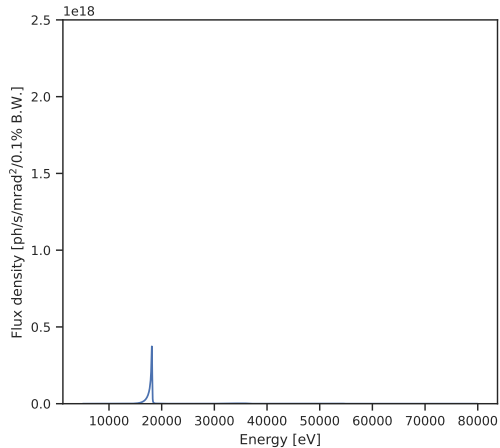
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Spectral function 1st harmonic:

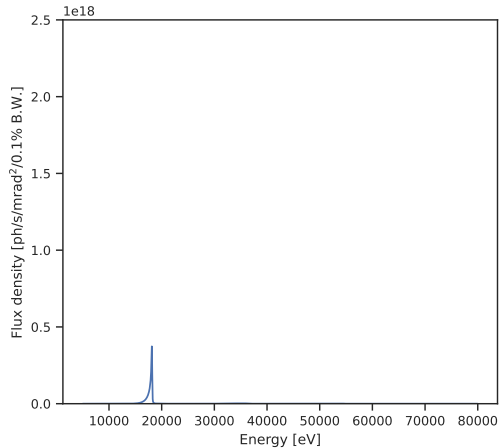


Undulator spectra and tuning



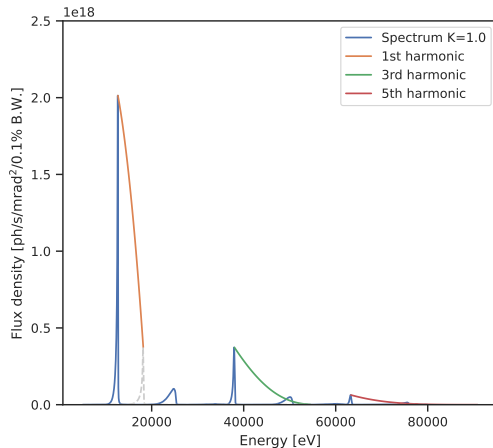
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Undulator spectra and tuning



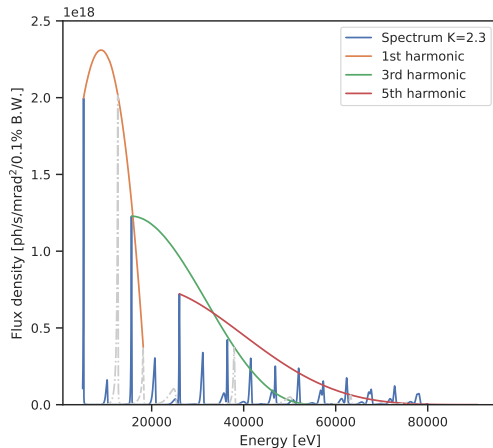
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Undulator spectra and tuning



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Undulator spectra and tuning



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- $K_u = 1.0$: higher harmonics appear, intensity increases (power $\propto B_0^2$), lines are shifted to longer wavelengths (lower energies)
- $K_u = 2.3$: tuning ranges of 1st and 3rd harmonic overlap

Helical undulators

Undulator field and trajectory

On axis field of an ideal helical undulator:

$$\vec{B} = (B_{x0} \cos(k_u z - \phi), B_{y0} \cos(k_u z), 0)$$

⇒ elliptic motion in the transverse plane with

$$\beta_x = -\frac{K_y}{\gamma} \cos(k_u z), \quad \beta_y = -\frac{K_x}{\gamma} \cos(k_u z - \phi)$$

and, in particular, for $K_x = K_y, \phi = \frac{\pi}{2}$

$$\beta_z = \beta - \frac{K_y^2}{4\beta\gamma^2} - \frac{K_x^2}{4\beta\gamma^2} = \text{const.}$$

Radiation field

Similar solutions as for the planar undulator with following main differences:

- modified undulator equation

$$\lambda_1 = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K_x^2}{2} + \frac{K_y^2}{2} + v^2\gamma^2 \right)$$

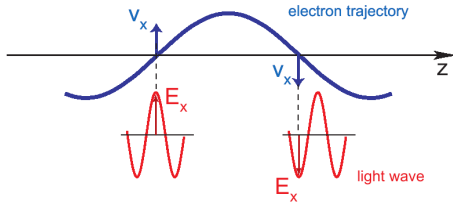
- elliptically, for $K_x = K_y, \phi = \frac{\pi}{2}$ circularly polarized light on axis
- constant $\beta_z \Rightarrow$ virtually no higher harmonics

Outline

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The FEL mechanism: resonance condition

Energy transfer particle ↔ light



P.Schmüser et al., Springer Tracts in Modern Physics (2014)

- oscillatory motion
- $v_x \parallel \vec{E}_{\text{light}}$
- energy transfer possible periodically

Continuous energy transfer

- requires fixed phase relation between electron motion and field oscillation
- that is fulfilled for:

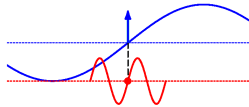
$$\lambda = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K_u^2}{2} \right)$$

$$\Rightarrow \gamma_r = \sqrt{\frac{\lambda_u}{2\lambda} \left(1 + \frac{K_u^2}{2} \right)}$$

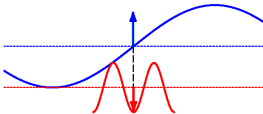
with γ_r the resonance energy

The high gain FEL mechanism: microbunching

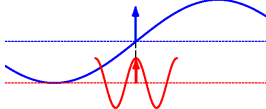
Reference particle: $\psi_0 = -\pi/2$
 zero energy transfer between electron and light wave



Laser-acceleration: $\psi_0 = -\pi$
 energy transfer from light wave to electron



FEL case: $\psi_0 = 0$
 energy transfer from electron to light wave



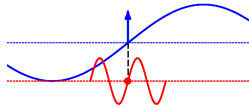
In a long bunch, given the FEL resonance condition is fulfilled:

- all relative phases present
- periodic particle energy modulation
- periodic drift velocity modulation
- periodic particle density modulation

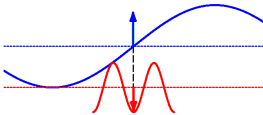
graphics: P.Schmüser et al., ibid.

The high gain FEL mechanism: microbunching

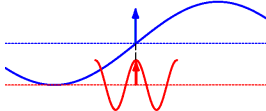
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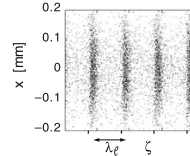
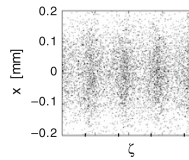
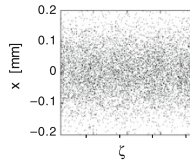
FEL case: $\psi_0 = 0$
 energy transfer from electron to light wave



In a long bunch, given the FEL resonance condition is fulfilled:

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- periodic drift velocity modulation
- periodic particle density modulation

→ Microbunching:



- all particles radiate coherently

graphics: P.Schmüser et al., ibid.

ID field quality: The phase error

Phase relation electron — photon

Phase slip due to different travelling times (over half period):

$$\omega_1(t_{\text{electron}} - t_{\text{photon}}) = \omega_1 \left(\frac{\lambda_u}{2\bar{v}_z} - \frac{\lambda_u}{2c} \right) = \pi$$

Here,

$$\bar{v}_z = \beta^* c = c \left(1 - \frac{1 + \frac{K_u^2}{2}}{2\gamma^2} \right).$$

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Phase errors

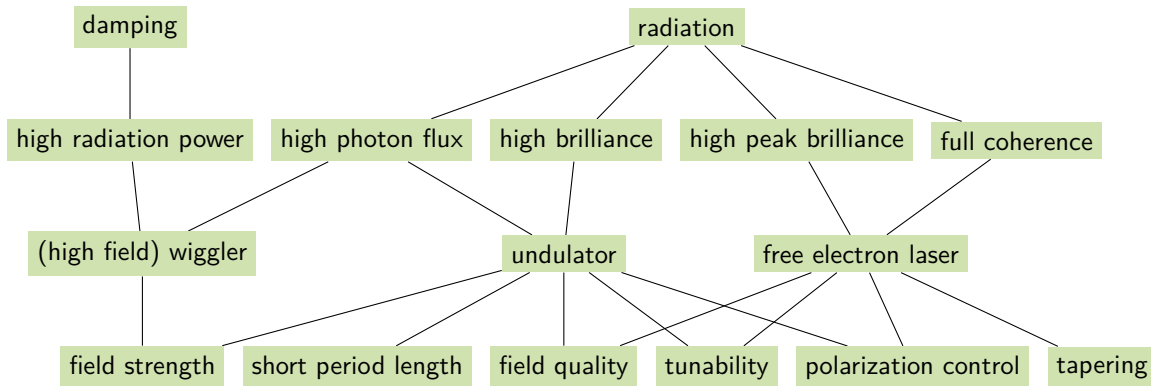
- magnetic field errors \Rightarrow deviations of K_u and \bar{v}_z \Rightarrow local phase deviations ψ_i
- phase error at pole n : accumulated phase deviations of all preceding periods

$$\phi_n = \sum_{i \leq n} \psi_i$$

- usual figure of merit for undulator field quality: rms phase error σ_ϕ^2
- target: few degrees

From use case to design requirements

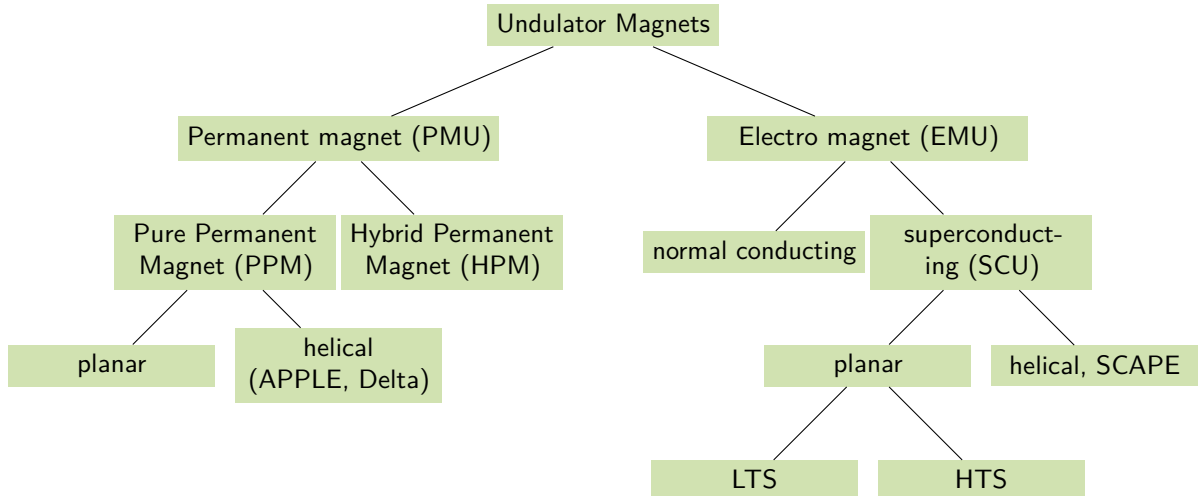
A more complete picture



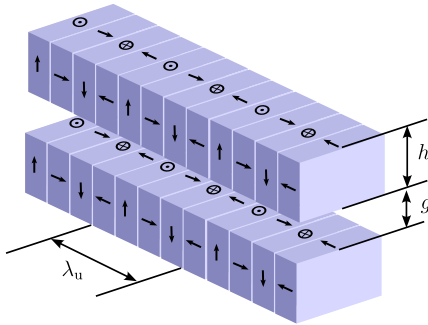
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Magnet technology choices for IDs



Pure Permanent Magnet (PPM)



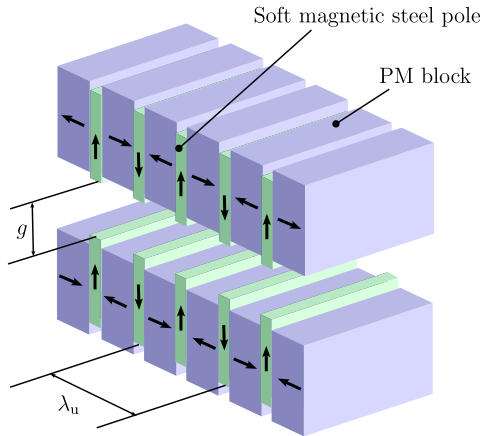
Design principle (K. Halbach NIM 187 (1981))

- permanent magnet blocks
- rotation of easy axis in M steps per period (minimum 4)
- practical scaling law ($M = 4, h = \lambda_u/2$):
 $B_{y0} = 1.72B_r e^{-\pi g/\lambda_u}$
- amplitude determined by B_r, g and λ_u (H_c)
- tuning by variation of g

Analytic field description (on axis)

$$B_y = -2B_r \sum_{i=0}^{\infty} \cos\left(\frac{2n\pi z}{\lambda_u}\right) \cosh\left(\frac{2n\pi y}{\lambda_u}\right) \frac{\sin(n\pi\epsilon/M)}{n\pi/M} e^{-n\pi g/\lambda_u} (1 - e^{-2n\pi h/\lambda_u}) \quad \text{with } n = 1 + iM$$

Hybrid Permanent Magnet



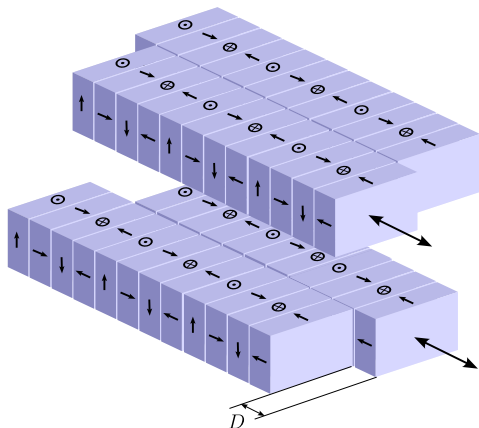
Design principle (K. Halbach (1983))

- soft magnetic poles magnetized by permanent magnet blocks
- non-linear material, numeric field calculation required (FEM)
- practical scaling law (Halbach, J. Phys. Col 44 (1983)):

$$B = a \exp \left[b \left(\frac{g}{\lambda_u} \right) + c \left(\frac{g}{\lambda_u} \right)^2 \right]$$

a, b, c depending on material choice, compilation to be found in F. Nguyen et al., XLS-Report-2019-004 (2019)

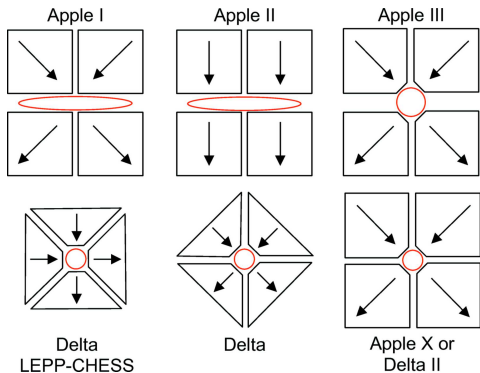
Permanent magnet variable polarization (helical) designs



Design principles

- 4 magnet arrays for 2 orthogonal field components
- always fixed phase shift of $\frac{\pi}{2}$ between B_x and B_y
- amplitude ratio varies by longitudinal shifting of arrays

Permanent magnet variable polarization (helical) designs

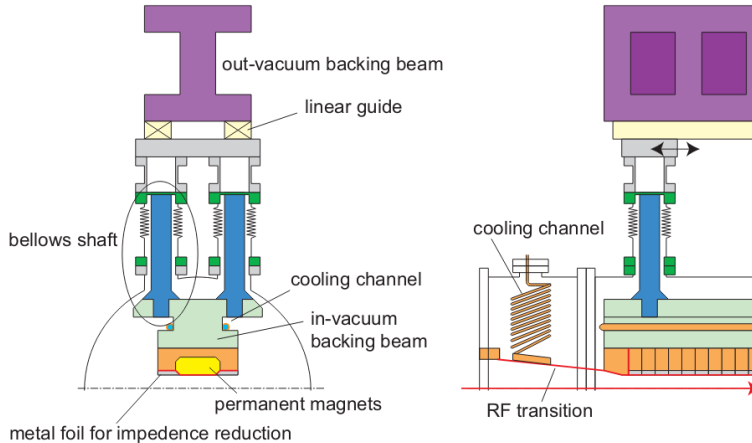


M. Calvi et al., *J. Synchrotron Rad.* **24** (2017)

Design principles

- 4 magnet arrays for 2 orthogonal field components
- always fixed phase shift of $\frac{\pi}{2}$ between B_x and B_y
- amplitude ratio varies by longitudinal shifting of arrays
- design variations, particularly for round beams
- Scaling law as for planar PMUs, a, b, c additionally depend on helical undulator design type (see also XLS-Report-2019-004)

Strategies for increasing B

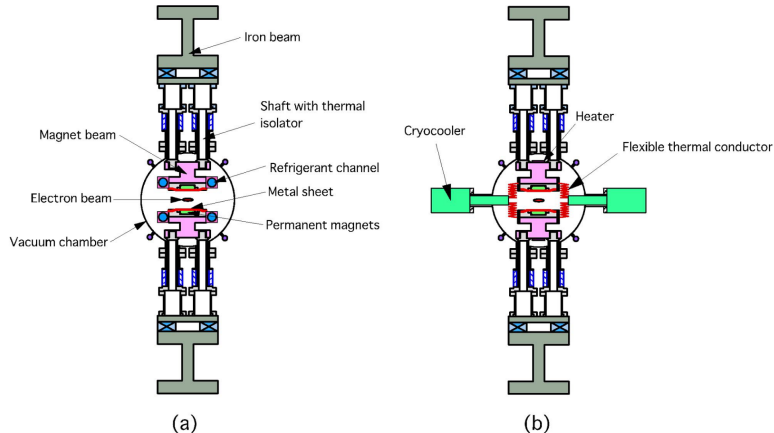


Advanced PMU designs

- In vacuum undulators (IVU)
 - reduce magnetic gap
 - increase ratio B_0/λ_u

T.Tanaka et al., Proc. FEL 2005

Strategies for increasing B

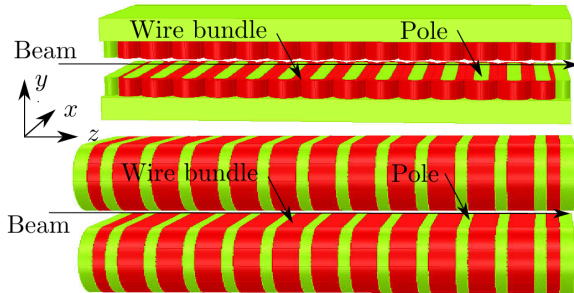


T.Hara et al., PRSTAB 7 (2004)

Advanced PMU designs

- In vacuum undulators (IVU)
 - reduce magnetic gap
 - increase ratio B_0/λ_u
- Cooled PMUs (CPMU)
 - in addition:
 - increase remanent field
 - increase coercive force

Electromagnets: Planar SCUs and SCWs



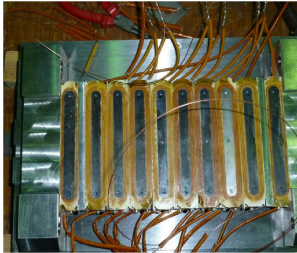
Aiming at very high fields or short periods

Approximating with dipole field, K_u scales like

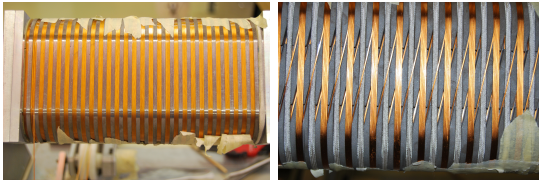
$$K_u = 2.35 \times 10^{-4} NI \frac{\lambda_u}{g} \quad (1)$$

- ⇒ generally high current densities are required
- ⇒ SCUs are the mainly relevant EMUs

Electromagnets: Planar SCUs and SCWs



N.Mezentsev et al., CLIC DW Tech. Report (2016)

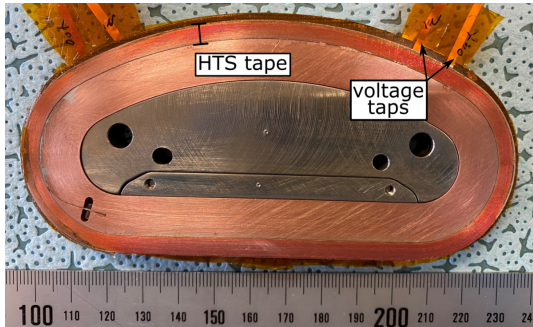


Undulator short model 'Undine', built in-house

State of the art

- low temperature SC, Nb-Ti
- horizontal racetrack
 - simple coil production (good)
 - many splices (not so good, but manageable)
- vertical racetrack
 - single wire
 - enabling very short periods

Electromagnets: Planar SCUs and SCWs



S. Richter, Dissertation, KIT (2023)

HTS vertical racetrack demonstrator coil (1 period)

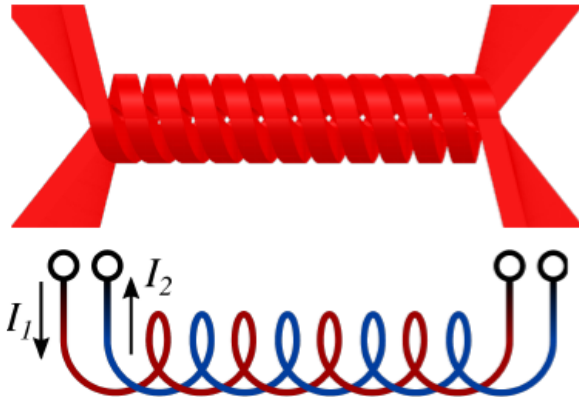
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On the horizon

- high temperature SC, ReBCO tape
- unprecedentedly high current densities and field amplitudes

SC helical and variable polarization

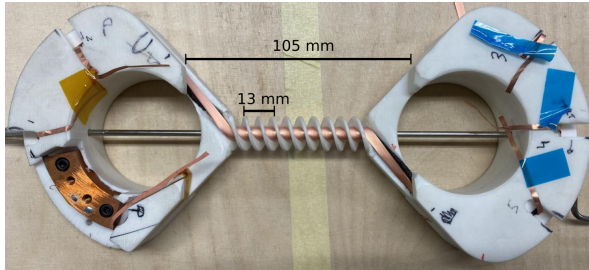


S. Richter, Dissertation, KIT, 2023

Bifilar helix

- bifilar helical coil around beam pipe
- fixed helicity (not changed upon current inversion)
- extremely resource efficient

SC helical and variable polarization

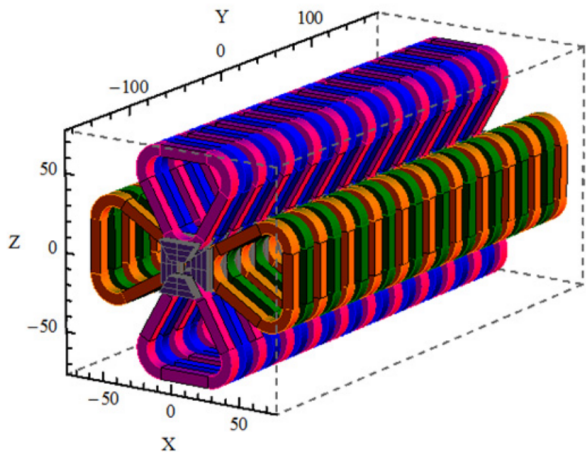


S. Richter, Dissertation, KIT, 2023

Bifilar helix

- bifilar helical coil around beam pipe
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- extremely resource efficient
- demonstrated for LTS and HTS

SC helical and variable polarization



Y.Ivanyushenkov et al., Proc.IPAC 2017

Bifilar helix

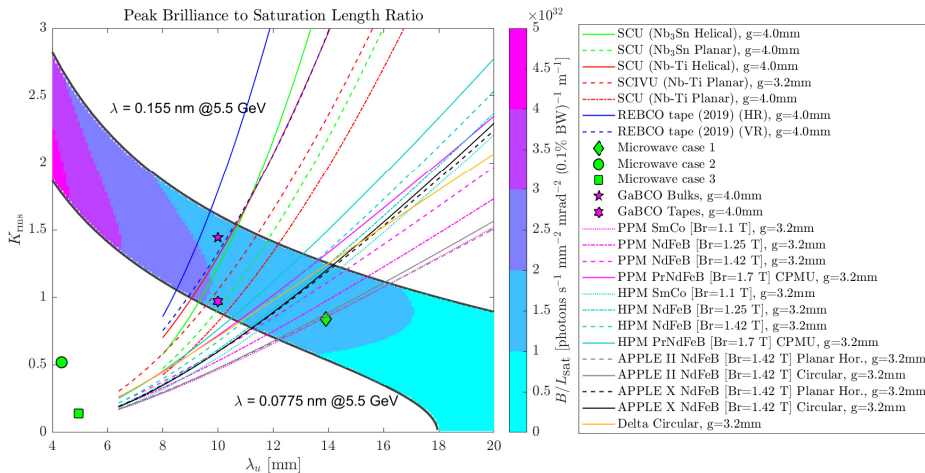
- bifilar helical coil around beam pipe
- fixed helicity (not changed upon current inversion)
- extremely resource efficient
- demonstrated for LTS and HTS

SCAPE

- SC arbitrary polarizing emitter
- full polarization control through two independently powered coil pairs

Comparison of undulator technologies

Example from the CompactLight design study



F. Nguyen et al., XLS Report 2019-004

Outline

1. Understanding undulator radiation
2. Understanding the FEL mechanism — Phase requirements
3. Technological boundary conditions for ID magnetic design
- 4. Aspects of ID 3D and technical design**

Aspects of 3D and technical design

a non-exhaustive list

- **field termination**
3D design and technical measures to ensure the beam-optical transparency
- **forces, structural mechanics**
mechanics for taking up and moving against magnetic forces (variable gap, polarization control), thereby maintaining field quality
- **cooling**
cryo-engineering for CPMUs and SCUs
- **SC magnet protection**
quench detection and magnet protection, simplified compared to beam transport magnets due to the low field energy
- **fiducialization**
- **field quality measurement and control**

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Field termination

Transparency requirement:

$$x\left(\frac{-L_u}{2}\right) = x\left(\frac{L_u}{2}\right) \quad \text{and} \quad x'\left(\frac{-L_u}{2}\right) = x'\left(\frac{L_u}{2}\right)$$

Field termination

Transparency requirement:

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This is equivalent with

$$I_1 := \int_{-\frac{L_u}{2}}^{\frac{L_u}{2}} B_y dz = 0$$

$$\text{and } I_2 := \int_{-\frac{L_u}{2}}^{\frac{L_u}{2}} dz \int_{-\frac{L_u}{2}}^z B_y dz' = 0,$$

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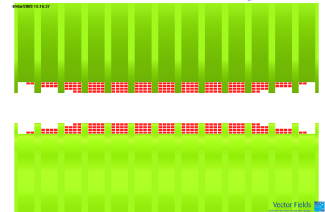
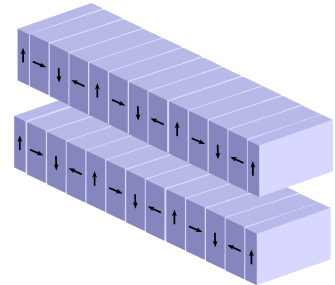
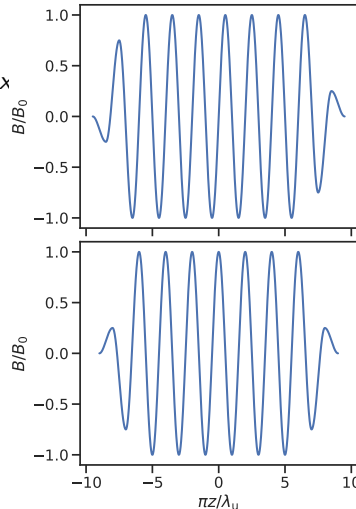
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Practically achieved with

$$B_y(z) = -B_y(-z) \quad \text{and} \quad I_2 = 0$$

$$\text{or } B_y(z) = B_y(-z) \quad \text{and} \quad I_1 = 0.$$



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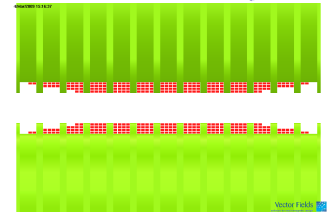
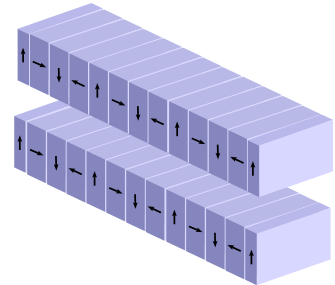
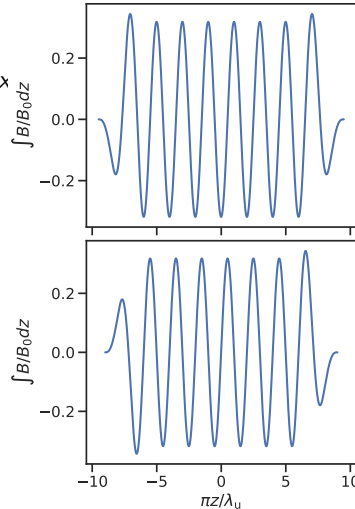
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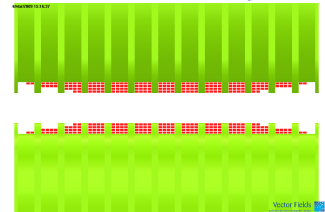
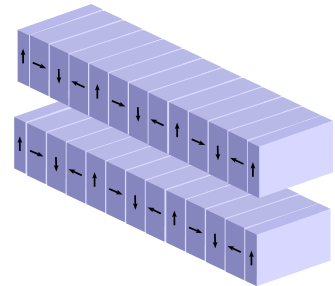
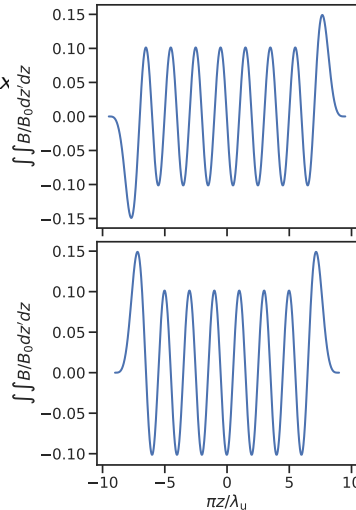
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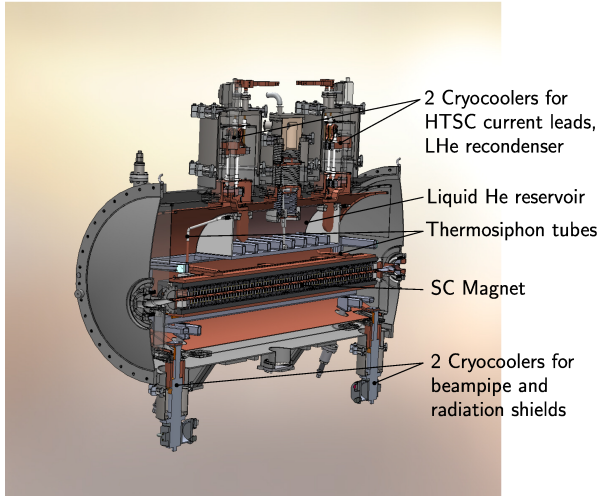
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N.Mezentsev et al., CLIC DW Tech. Report (2016)

Requirements and techniques

- **SCU**: operating temperature ~ 4 K
- **CPMU**: operating temperature ~ 150 K
- **HTSCU**: operating temperature 4 K to 10 K, still to be investigated and optimized
- Operation in light sources
 - light sources: stand-alone, closed cycle preferred
 - bath or contact cooling with liquids or/and cryocoolers
- Operation in FELs and damping rings
 - forced flow cooling
 - central cryoplant more efficient than local cryocoolers

Field quality measurement and control

Field quality control techniques

- measurement techniques in general
 - Hall probe (, pulsed wire)
→ local field, phase error
 - flipping coil, stretched wire, pulsed wire
→ field integrals

Field quality measurement and control

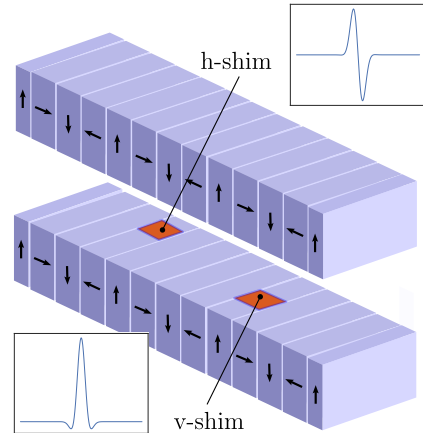
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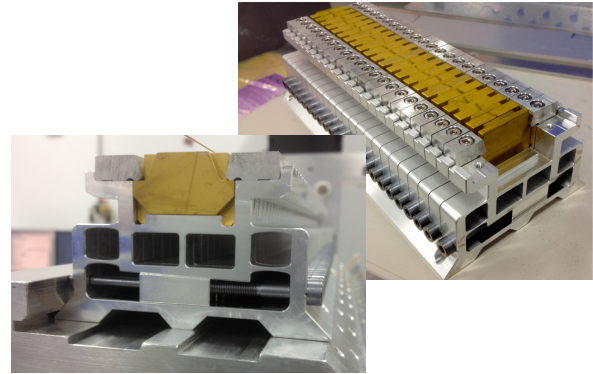


Basic idea of local field correction by shimming

Field quality measurement and control

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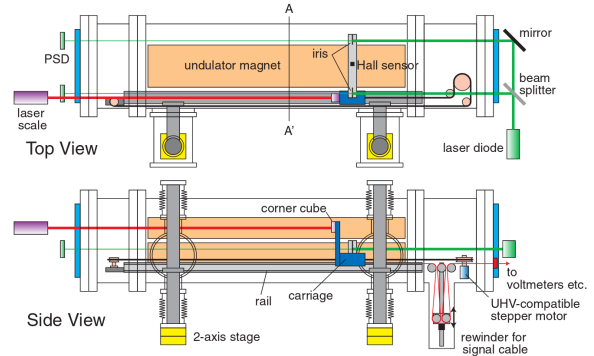
C.J. Milne et al., *Appl. Sci.* 7 (2017)

Magnet assembly and adjustment system of the Swiss-FEL ARAMIS undulators

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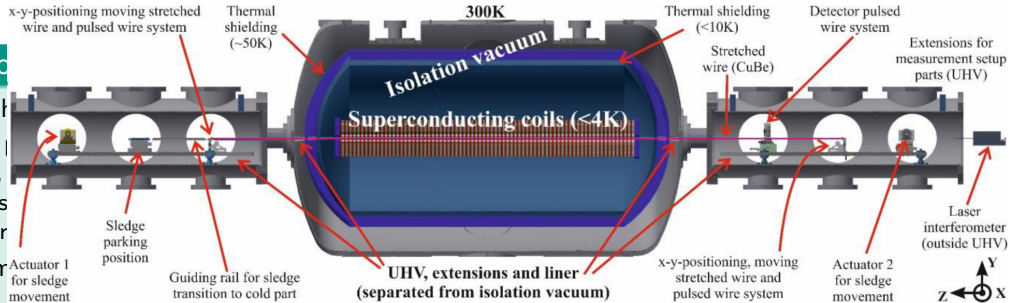
T. Tanaka et al. PRSTAB 12 (2009)

Self-aligned field analyzer with laser instrumentation (SAFALI), UHV compatible.

Field quality measurement and control

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 - field integr
- phase error minimization
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 - pole/block adjustment
- in vacuum and CPMU: specialized Hall mapping techniques, e.g. SAFALI
- SCUs specialized cryogenic and in-situ set-ups



A. Grau et al., Proc. IPAC 2019

In-vacuum, cryogenic magnetic measurement setup for magnetic qualification of SCUs in the final cryostat

Summary

- Insertion device design is mainly driven by synchrotron light users' demands in terms of
 - Photon flux and brilliance
 - Polarization control
 - Coherence
 - but also: availability, sustainability
- general trend: towards shorter periods, yet maintaining the tunability range up to $K = 2$
- Permanent magnets are the state of the art, still improving
- Superconducting magnets as an alternative are becoming increasingly relevant, in particular in view of the trend towards more compact light sources
- HTS technology bears a large potential for future short-period high-field insertion devices