



Heat transfer, cryostat, (conduction-)cooling (II of II)

CAS – Austria 2023

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With contributions from P. Borges de Sousa (CERN), T. Kottig (CERN), F. Ferrand (CERN)

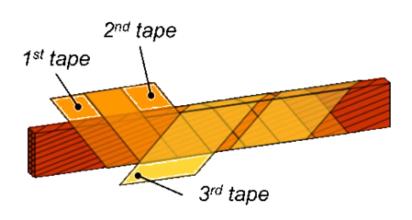


Outline

- Fully immersed in HeII:
 - LHC dipoles and HL-LHC cable-stacks
 - Bayonet HX cooling scheme
 - Numerical tool & application to HL-LHC Nb₃Sn quadrupoles
- Helium availability and typical LHC reliance (fully immersed magnets)
- Food for thought of operating fully helium immersed accelerator magnets
- Some points to remember
- References



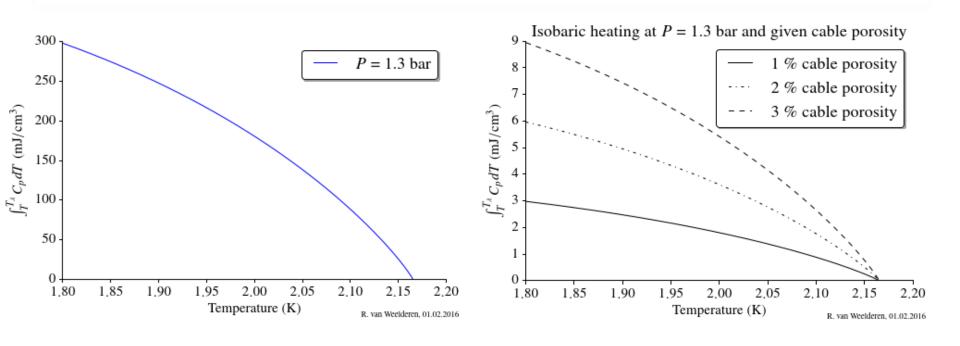
Fully immersed in HeII: LHC-dipoles



NbTi - cable is:

- non-impregnated
- Electrically insulated by partially overlapping layers of Kapton
- --> porous to helium

Fully immersed in HeII: LHC-dipoles

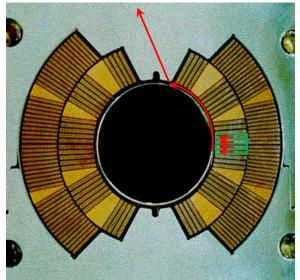


The high thermal capacity of the helium inside the porous NbTi - cable contributes to the stability

The order of magnitude is ~2 mJ/cm³ per % of He volume



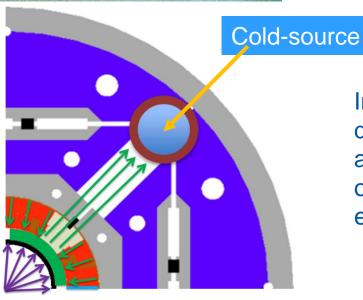
Cold-source Heat sources and heat-sinks



(aka "cold-sources")

In *fully immersed* magnets the heat generated in the coil-pack must find its way out to the coldsource via helium path-ways kept clear in the cold-mass construct

In the example top left, of the LHC main dipole, heat flows from the coil pack into the annular space between beam-pipe and coil-pack and out via space between the collar laminations



In the example bottom left, of the HL-LHC MQXF quadrupole, heat flows from the coil pack into the annular space between beam-pipe and coil- pack and out via dedicated passages (8 mm diameter holes every 50 mm along the length of the magnet)

Heat sources and heat-sinks

(aka "cold-sources")

Mechanical Concept CERN, CNAO, INFN and MedAUSTRON on novel ion gantry concepts: "dry" conduction

cooled magnet

Epoxy-impregnated 2-layer coils with inter-layer splice, wound with 34-strand 8.75 mm Nb-Ti cable with braided glass insulation

2 Stiff austenitic steel collars with 0.15..0.2 mm thick spacers on one side to follow the coil curvature

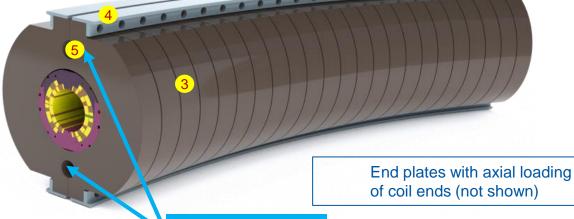
Horizontally split laminated iron yoke made of 1-mm-thick Si-steel with b-staged resin coating.
Yoke sectors machined out of cured lamination stacks.

Courtesy M. Karpinnen & Ch. Kokkinos

In a "dry" magnets the heat generated in the coil-pack must find its way out to the cold-source via solid conduction and solid-liquid interface

Yoke assembly clamps mounted under yoking press

5 Thermalisation at 4.5 K



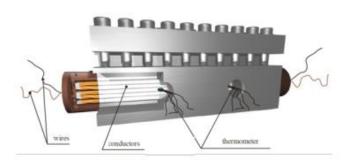
Heat sink(s)

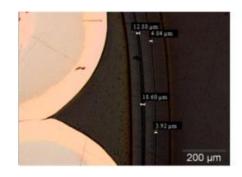


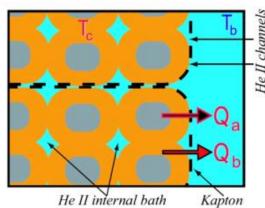
Fully immersed in HeII



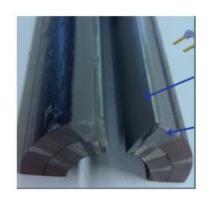
Cable T excursions w.r.t. helium bath measured as function of power deposit

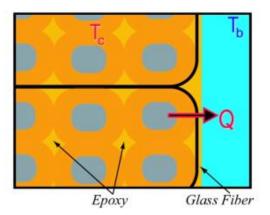






NbTi cables showing the porousity w.r.t. helium



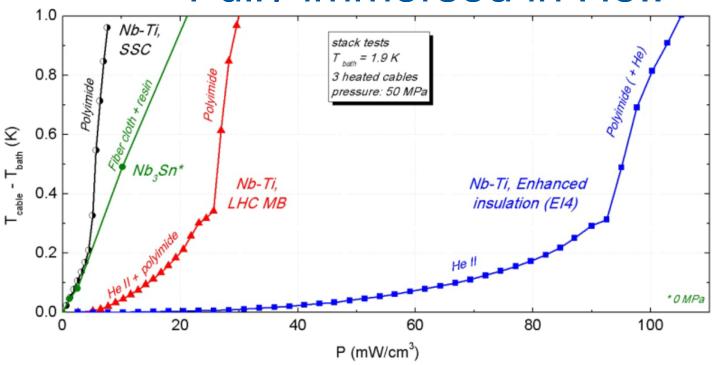


Nb3Sn cables, fully impregnated, only conduction through solids

courtesy H. R. Correia Rodrigues



Fully immersed in Hell

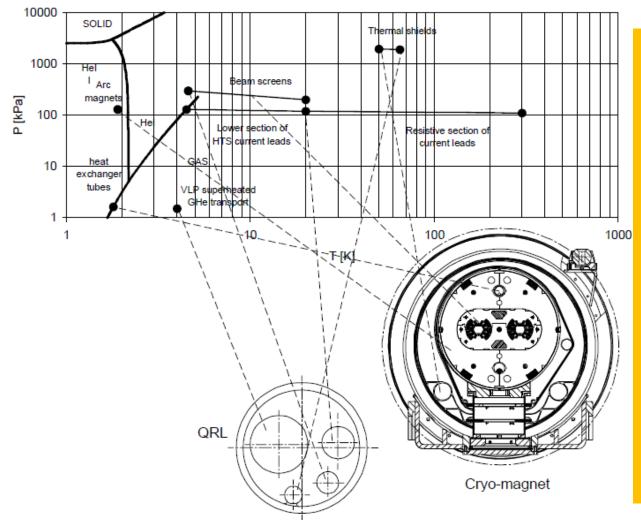


Stack measurement results, all faces open, showing the stark difference between porous and fully impregnated cables

Nb₃Sn measurements have since, the last 5 years, been addressed in more detail (see [3])



Fully immersed in HeII: LHC-dipoles

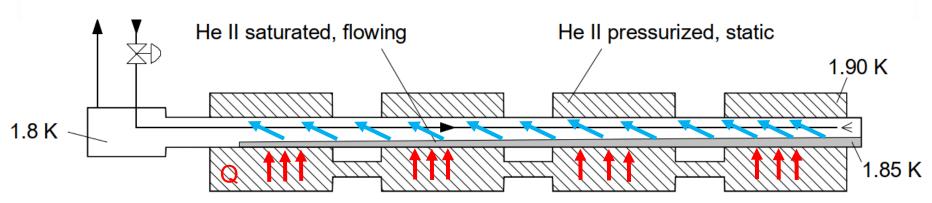


Addressing yesterday's question:

Even if the coil, collar and yoke are in pressurized Hell at 1.3 bar and ~ 1.9 K

The heat-sink (upper hole), where the heat is taken up by vapourizing liquid helium (~1.85 K, ~15 mbar) is in a copper-pipe protruding over the length of several magnets (107 m)

Fully immersed in HeII: LHC-dipoles



All the extracted heat has eventually to flow across the tube wall, where it encounters three thermal impedances in series:

- the limited solid conduction across the metal constitutive of the wall,
- the Kapitza resistances produced by the refraction of thermal phonons at the metal-to-helium interfaces.

The overall transverse impedance was measured on fully wetted test samples at varying temperature, so that the different temperature dependence of the solid conduction and Kapitza terms enabled to resolve them. For tubes with a wall thickness up to about 1 mm, the Kapitza resistance largely dominates below 2 K, and the use of high-purity, cryogenic-grade copper is not required.

See ref[1, 2]



Evaporation of two-phase, very low pressure, helium flow

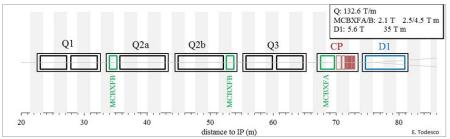


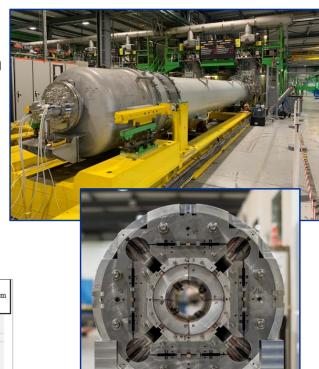
Fully immersed in HeII: MQXF

(HL-LHC Nb₃Sn quadrupoles)

- In the context of HL-LHC at CERN, a robust multi-region CFD numerical framework for the modelling of heat transfer in complex cryogenic system geometries involving He II has been developed
- Numerical tool allowed to provide thermal design requirements early enough in the magnets' design phase
- Now, a fully upgraded and consolidated tool is used to reassess previous results, and for systematic analysis of heat extraction pathways in complex He II – composite solid geometries





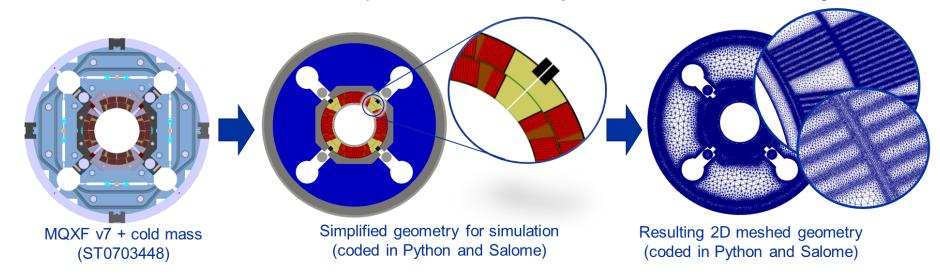


Following slides: courtesy P. Borges de Sousa | Revised estimates of temperature margins in MQXF, see also ref[x]



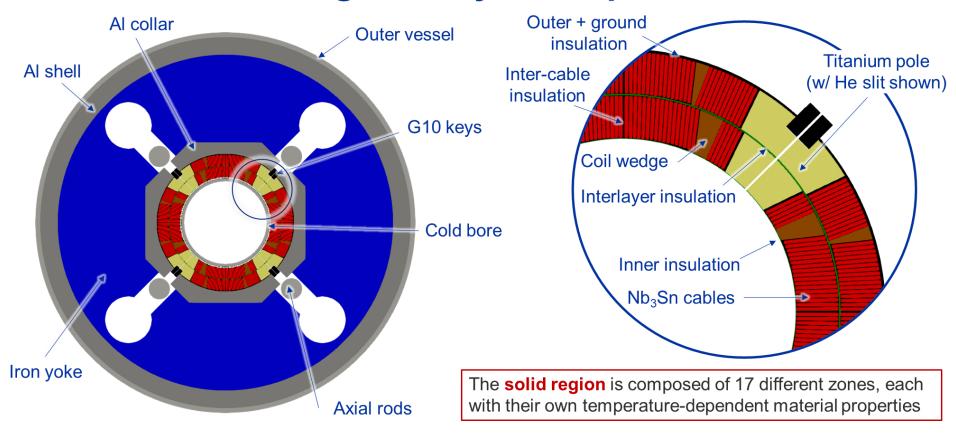
Numerical code applied to full-scale cold masses

- Libraries and solver implemented in OpenFOAM v7, geometry created via Python API for Salome, easily adjustable for parametric studies
- Currently handles 2D geometries and static He II conditions (no flow) and He II → He I transition
- Power deposition data can be mapped directly onto the geometry, or direct heat input to any part
- Coil described in detail; mesh composed of 3.5M cells, special refinement on thin layers



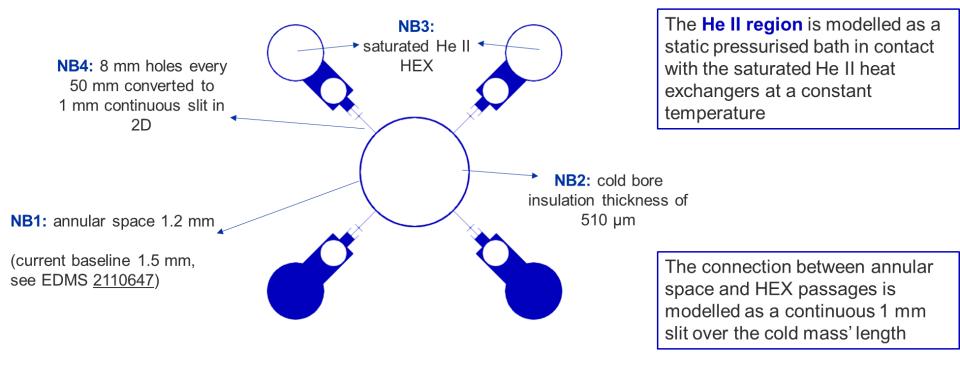


MQXF cold mass geometry - composite solid



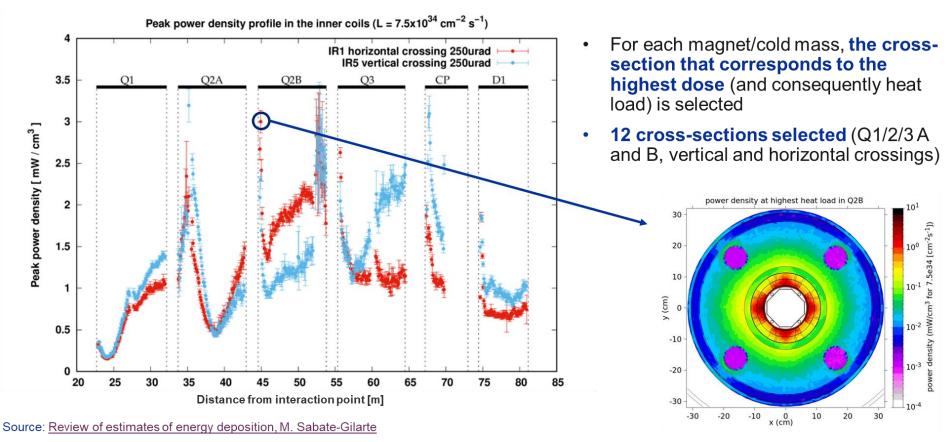


MQXF cold mass geometry – He II region





Input for heat deposition from dose calculations (I)

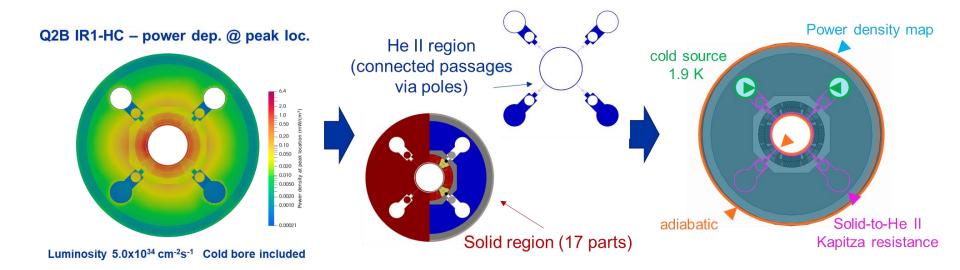




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Input for heat deposition from dose calculations (II)

- For each cold mass, the 2D **power density** (mW/cm³) **at the peak** (maximum) **location** for nominal luminosity (7 TeV, 5.0x10³⁴ cm⁻²s⁻¹) is mapped onto the mesh
- Mesh is split into two regions, a He II region and a solid region
- Set of boundary conditions is chosen; power density map changes for every analysed cross-section



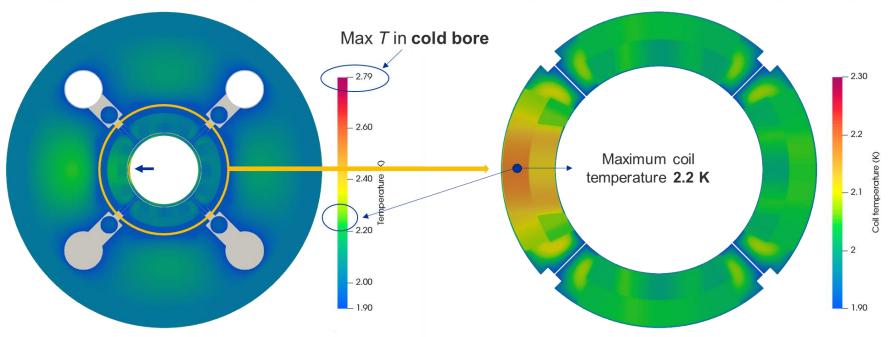


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Results: example Q2B IR1-HC @ nominal luminosity

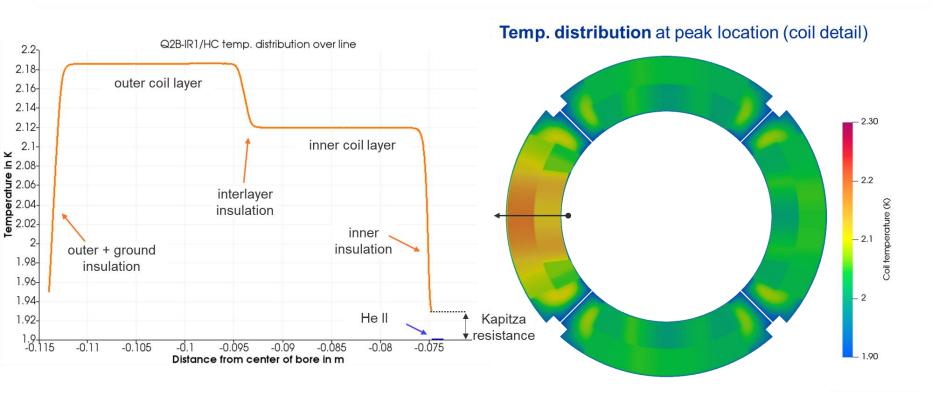
Temperature distribution at peak location

Temp. distribution at peak location (coil detail)



NB: Cold bore included, luminosity = 5.0×10^{34} cm⁻²s⁻¹

Results: example Q2B IR1-HC @ nominal luminosity



With fully impregnated cables, the thermal gradients are mainly determined by the interfaces (in contrast with NbTi-helium permeable cables)



With fully impregnated cables, the thermal gradients are mainly determined by the interfaces (in contrast with NbTi-helium permeable cables)

→ Going from fully helium-immersed coils to "dry" conduction-cooled coils is for fully impregnated cables, apart from an additional DT, a similar situation.

For future magnet development based on fully impregnated coils the temperature gradient ΔT across the coil must be kept as low as possible:

- Control the interfaces
 - Electrical insulation (Kapton foils p.e.)
 - Quench heaters (external, CLIQ, inductive...)
 - New impregnation materials (can we enhance the conductivity & Cp?)
- New conductors
 - Anisotropy of HTS tapes
 - Electrically insulated or non-insulated HTS tapes
 - •
- Conductivity & Cp data of new magnet designs to be measured
 - Accurate data greatly detailed thermal analysis



Key takeaways

- A robust, easily adaptable numerical framework has been consolidated to evaluate heat extraction mechanisms in complex cryogenic system geometries involving He II
- We are now able to produce results in a timeframe that allows for parametric investigation of geometry, operational *T*, and power deposition on magnet systems cooled by static, pressurised He II → this enables a systematic approach, and can be used as a tool in magnet design w.r.t. heat extraction
- Numerical simulations directly benefit from the material data obtained from the experimental test campaign, allowing for more realistic, accurate calculations
- The heat extraction pathways for the HL-LHC inner triplet magnets have been validated for the latest geometry, considering a steady-state power deposition at peak location for nominal operating conditions (5.0x10³⁴ cm⁻²s⁻¹, 7 TeV), at 1.9 K cold source



ASC 2022 paper submitted to IEEE Transactions on Applied Superconductivity (TAS)

Numerical Assessment of the Inhomogeneous Temperature Field and the Quality of Heat Extraction of Nb₃Sn Impregnated Magnets for the High Luminosity Upgrade of the LHC

Abstract—The High Luminosity apprade of the Large Hadron visibler (ILL-HC) at CERN forevers the installation of Mysher control of the CERN forevers the installation of the structure. The previous based logs of and magnet thermal charac-retictic and heat cutraction performance in response to power promision during both monital and utilinate conditions in exer-ual in determining safe operating margins. A 20 numerical granureouth has been developed to systematically assess the temper-

Index Terms— cryogenics, numerical simulation/model, power leposition, Nb3Sn magnets, Helium II, HL-LHC

In the context of the High-Luminosity upgrade of the LHC at CERN, a robust multi-region CFD numerical toolkit using open-source software for the modeling of heat transfer in complex cryogenic system geometries involving He II has been developed, aiming to provide thermal design requirements [1] early enough in the design phase of the inner triplet magnets, such that they could be implemented in the mechani-

sign that was originally proposed.

The work presented here is a continuation of previous studies [2]-[4], implementing updated MQXF geometry, power density maps and experimentally-obtained material properties, as well as a fully upgraded numerical toolkit that enables a systematic analysis of multiple magnet cross-sections and al

He I phase-transitions and converges more than an order of

II. GEOMETRY AND NUMERICAL TOOLKIT

tank (CERN) A. Geometry of MQXF cold mass rule, Baden

A 2D cross-section of a MQXF quadrupole magnet cold more of the farmers in this paper are available unline mass representative of the cold mass' straight section, along with the cold bore (see Fig. 1), was created according to the



[7]) The individual con-directly from ROXIE*

ole CAD files. The main

eated using a Python* API a* 9.7, an open-source sci-decause the geometry is traceable coordinates, any magazed to the rest of the

read in Fig. 2. The coil is steed in Fig. 2. The coil is ble level, including inter-meet-layer invulation. The seel vessel and the magnet, a aluminum shell, the iron r, and the coils. The cold lation, is also present. The as well as alignment slots

every 50 mm along the magnet axis and that connect the inner He annular space between the cold bore and the inner radius of the coll to the cooling channels, a continuous 1 mm slit along the intamium pole and G10 alignment key was created to ena-ble a 2D geometry with equivalent longinodinal cross-section (can dead in Eq. 2 (5)). (see detail in Fig. 2 (a))

(see detail in Fig. 2 (a)).
Each component of the cold mass is assigned appropriate temperature-dependent material properties such as density, specific best, and thermal conductivity. Wherever possible, material properties syntacted from experimentally obtained data on MQXF coil samples [11] were used rather than literature values; the interface contact resistance between solid materials and He II was also experimentally measured on a MQXF coil

B. Mesh creation and refinement

The geometry was methed using a Python* API for the Meshing module of Salome* 9.7, with special refinement on the coil pack and this layers such as insulation and belium passages. The final mesh used for the numerical calculations is composed of 3.5 million cells.

C. Mismerical positive descriptions.

The numerical about used in this work: was implemented in Oppard Double vib. may open-course C++* insolvine fine the development of the property of the with its own material properties, and a pressurized He II re-gion. The two regions interact by means of a baffle-type boundary condition that implements experimentally measured [11] interface contact resistance between He II and solid mate nals, referred to as Kapitza resistance

A. Energy deposition from collision debris

Debris from proton-proton collisions at the interactions points induce energy deposition in the superconducting mag-nets and respective cold masses, especially those in the inser-tion regions (RR [13] The power density deposited in the coils leads to an increase in temperature that is non-uniform

coun seaso of an increase in temperature tant is non-minoral throughout the cold mass structure.

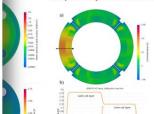
The temperature increase due to debris-penezated energy deposition needs to be evaluated to provide a more realistic operating temperature field for the magnetis in the insertion ra-pions. The energy deposition in the inner triplet magnetis in the insertion replons IR1 (ATLAS) and IR3 (CMS) used as input for the present work was estimated by FLUXA for a baseline

, for proton-proton col. He II here exchangers is set at a fixed temperature of 19 K 14 TeV, for HL-HC operation of the set of the magnets that FLUKA (for a luminosity of 50,00° cm 2 of 50,00° cm 2 of m proton-18 Q/2a, Q/b, Q/b, and proton-18 Q/2a, Q/b, Q/b, and proton-18 Q/2a (for a luminosity of 14 TeV) is spond vertical crisisness gained to the entire cross-section of the MQXT cold and cold the cold of the Q/2a Cold and cold the cold of the Q/2a Cold and cold the Cold of the Cold mass: this is the only input that varies with each of the 12 modeled magnet cross-sections.

IV RESULTS

Results consist of a temperature map for each selected mag-Results consist of a temperature map for each selected mag-net cross-section; an example of power density (a) and result-ing temperature profile (b) is shown in Fig. 3 for the most ex-posed cross-section of Q2B at insertion region 1 (horizontal crossing). There is a localized hotspot in the cold bore, reach-ing 2.5 K, where the power density is highest. ore (chosen to allow adering a continuous 1 mm nament keys. The outside he inside of the cold bore

Fig. 4 (a) shows the same temperature profile as Fig. 3, fo-cused only on the coil strelf. The coil shows a fairly homogeneous temperature distribution, with a maximum calculated temmerature of 2.22 K on the outer layer i.e. a 0.32 K eradien with respect to the cooling source in the saturated He II heat



the temperature distribu-tion of a shown by the general control of the strength of the shown of the shown of the temperature margin for each analyzed magner cross-section general control of the strength of the shown of the shown of the strength of the shown of the shown of the shown of the shown of the strength of the shown of perature discontinu-te resistance between margin above 5 K. considering an otherwise ideal conductor.

A robust, easily adaptable numerical toolkit has been con-solidated to produce results in a timeframe that allows for paruction through the alumi-ysis we conclude that the e coil is dominated most-ed by the interlayer insu-

at 12 T and 4.2 K [15])

Well the court multi-index of the court multi-index of the court multi-net conclude that the presence of the court of t stance to the helium bath, ed He II bath (not shown) g that the heat extraction sugn are not a bottleneck tration of the magnets. value for each quantity is

amre of 1.9 K. For the cross-sections of each magnet at peak power density, a maximum temperature of 2.22 K was calculated on the coil pack (on Q3a, IRS), while the maximum temperature for the rest of the coid mass is invariably located on the coil down, reaching 2.79 K on Q3b, IRS. These results compare positively against the conservative assessment made during the cold mass design phase [2]. Considering the calculated temperature distribution, for the

design B. J. and cold source at 1.9 K, on an otherwise idea. magnet, a minimum temperature margin of 4.78 K was found for a single conductor, nearest to the pole, in Q3a, IR5. location-dependent, being lensity is the highest. The by ROXIE* assume $f_{\mu\nu}$ = dependent $B_{\mu\nu}$ uniform a coil, and reference con-

The authors would like to acknowledge the Cryolab team a CERN for their support in the experimental campaign on He II heat transfer to Nb,Sn magnet samples.

On coil pack				On cold mass			
2 4 5	Maximum temperature [K]	Average temperature [K]	Minimum calculated temperature margin (K)	Maximum power den- sity [mW/cm ²]	Average power den- sity [mWicm ³]	Maximum temperature [K]	Average temperature [K]
	2.06	2.01	4.02	1.05	0.13	2.21	2.01
	2.19	2.07	4.84	1.52	0.21	2,38	2.06
	2.15	2.07	4.86	2.96	0.22	2.41	2.04
-	2.19	2.04	4.88	4.77	0.19	2.51	2.02
	2.18	2.07	4.82	4.93	0.23	2.64	2.05
•	2.09	2.02	4.89	3.91	0.16	2.50	2.00
	2.07	2.01	4.92	1.24	0.13	2.23	2.01
•	2.20	2.07	4.84	1.79	0.21	2.44	2.07
-	2.21	2.08	4.33	5.42	0.25	2.74	2.05
-	2.14	2.04	4.86	3.87	0.17	2.51	2.02
-	2.22	2.09	4.78	5.03	6.25	2.67	2.06
п	2.19	2.03	484	6.37	0.16	2.79	2.01



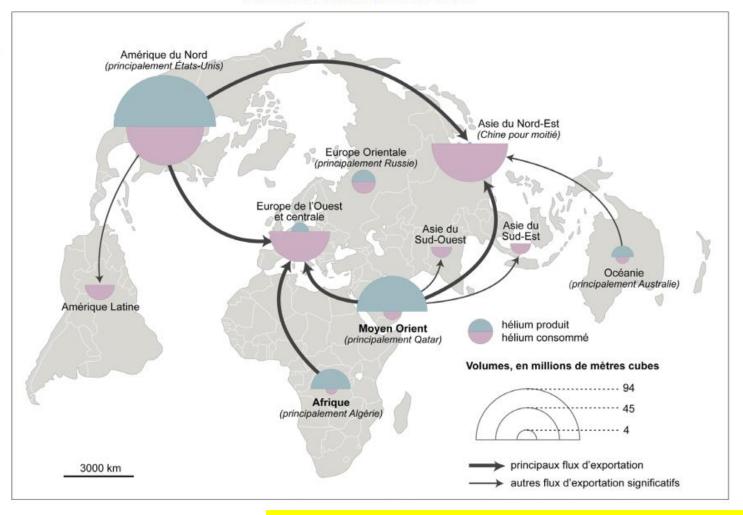
Helium availability and typical LHC reliance (fully immersed magnets)

Next slides source: Frederic Ferrand (CERN Cryogenics) see [7] for the full transcript

- Presentation based on available public information only
- Not an economist nor a seller, only interest is to supply necessary molecules for the laboratory in the coming years
- Helium market is currently (2022-data) evolving really fast and situation may be significantly different in a few months



Carte 1 : production, échanges et consommation d'hélium dans le monde en 2018



Source: Gubler et al. (2019).

Europe (and CERN) are outside supplier dependent

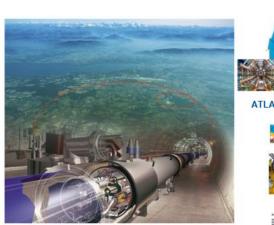
Ref: https://www.ifri.org/fr/publications/briefings-de-lifri/helium-nouvelles-geographies-dune-ressource-critique



Use of helium cryogenics at CERN



LHC & Experiments



LHC 36'000 tonnes of cold mass distributed over the 26.7 km underground accelerator



ATLAS 1'275 tonnes of cold mass



CMS 225 tonnes of cold mass

LHC accelerator cooling of the superconducting magnets at 1.9 K ATLAS cooling at 4.5 K of the superconducting magnetic system CMS cooling at 4.5 K of the superconducting solenoid

non-LHC & Test Facilities





SM18 test Facility

West Area test Facility







HIE Isolde Cryo Modules

SPS BA4 COLDEX

SPS BA6 Crab Cavities







Central Liquefier

North Area

Test Benches for accelerator magnets, wires, RF cavities non-LHC facilities and fixed target physics experiments Central Helium liquefier and Cryogenics laboratory

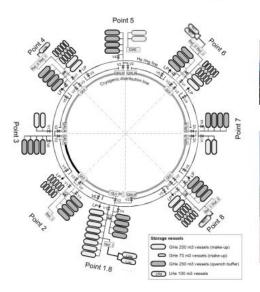
Courtesy Frederic Ferrand



LHC infrastructure



Surface storage & distribution



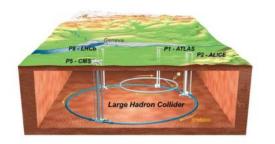


Gaseous helium storage 250m3



Liquid helium storage 250m³

Underground distribution & use



4.5K & 1.8K Cold Boxes

Distribution Valve Boxes

Distribution System







- Total inventory at CERN 175 tonnes including 130 tonnes minimum for LHC
- Distribution through the LHC Helium Ring Line, inter-sites High Pressure helium distribution line for non LHC and trailers







Liquid helium supply **200 tonnes**

to compensate helium losses

Mainly for testing and experiments

For **LHC** machine constant effort to reduce losses showed results over the past 15 years



Helium quantities included:

- LHC needs for Run 3 and refilling of the machine at the end of Long Shutdown 3
- Requirement for new HL LHC cryogenic infrastructure are also included

Helium management 140 tonnes

temporary sent out from CERN to contractual suppliers and returned on demand

This is a necessary extension of CERN storage capacity when LHC machine is partially or completely warm



During Run year end closure up to 20 tonnes are sent out for about 2 months

During Long Shutdowns up to 80 tonnes are sent out for about 18 months

Gaseous helium supply 6 tonnes

for LHC machine pressure tests

High Pressure helium trailers are used to pressurized sectors for safety pressure test



These deliveries are foreseen at the end of Long Shutdown 3

Courtesy Frederic Ferrand





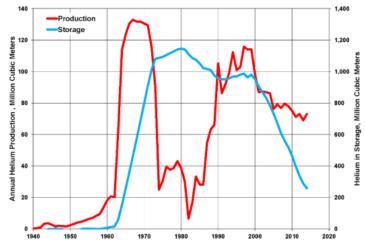
Helium market long term evolution driven by US strategic decisions

Helium act 1925

- Significant increase of the production in the 60's
- Cliffside infrastructure and underground storage of the federal strategic reserve
- Almost monopoly of the US on the worldwide helium market
- Rapid decrease of production in the 70's, reaching peak in 1980
- Managed by the state BLM with significant associated costs
- Confort zone, low prospecting and investment elsewhere

Helium privatization act in 1996

- Federal government expenses pay back by selling 1bcm till 2015
- Large industrial investment in Algeria 1997 and Qatar since 2008
- Production capacity developments in the private sector was too low



Wikipedia, based on US geological survey data

Stewardship Act of 2013

- Regulation of the of the market by selling of the strategic reserve through yearly auctions to the private sector
- Make the bridge with new international extraction capacity development
- Regulation of the US domestic market, but reduction of exports over the years

Good Read (in French): IFRI Briefing

Courtesy Frederic Ferrand





Known parameters affecting helium availability

- Unbalanced supply and demand, helium shortage 2006 & 2013
- Geopolitical stability in the country of extraction, Qatar 2017
- Logistics complexity blockage of Suez Canal in 2020
- Maintenance shutdown or technical event on LNG feed and helium liquefaction plants can impact market





What's happening in 2022?

- BLM closure in January
 - Temporarily closed procedures for safe handling of chemical materials (Gasworld)
 - Outsource of BLM operations to private contractor
- Explosion at AMUR site in Russia production site in January
 - Several events leading to stop of the commissioning (Gasworld);
 - Expecting ramp up of the production postponed
- Commercial limitation with Russia due to conflict in Ukraine
 - Impact on the AMUR project repair and recommissioning of the helium plant
 - Takeover of the engineering from European by a Chinese company
- Planned maintenance during spring period
- Restart of industrial activities after Covid period





Consequences

Current situation

- Market shortage is affecting industrial and scientific customers
- Manufacturing industry contracts are impacted with volume limitations
- Large scientific instrument cannot do so & rely on established industrial partnership

Helium market still at risk in 2023 and for the coming years

- Uncertainty on the effective Russian production capacity and market access
- Algerian gas production transferred using pipeline instead of LNG
- No more back-up from the US federal authorities, Cliffside for sale! (C&en News)





Perspective of new helium production capacities

LNG related

- Algeria debottlenecking of existing production sites
- South Africa → 900Mm3 (lower estimate)
- Qatar IV ?

Nitrogen-based helium production (Physicsworld article)

- Canada (Gasworld <u>article 1</u> & <u>article 2</u>) \rightarrow Announced production objectives 10% ww capacipty by 2030
- Tanzania (The Citizen) 54bcf of primary helium at Lake Rukwa → Production announced in 2025



Food for thought of operating fully helium immersed accelerator magnets

Disclaimer: the LHC has been an immense success and functions marvelously. A 27 km long accelerator of which about 23 km at 1.9 K!

A cryogenic success, on both magnet cooling as well as the whole cryogenic infrastructure which functions with very high availability!

However, there are a few points which are food for thought on whether we can do better for the next machine...

(whichever one that might be)



Safety: avoiding trapped volumes

What happens if one neglects (or even forgets) radial passages!

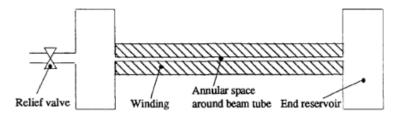
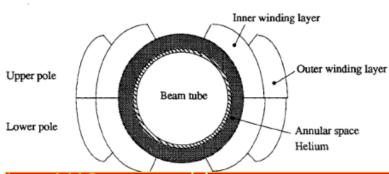
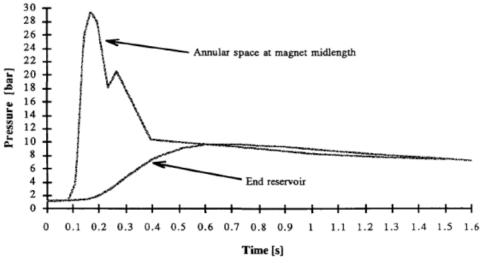


Figure 1.1 Dipole magnet





In a LHC-type cold-mass structure foreseen for a future 10 T coil, a low field (< 5 T) coil was placed, but to ease assembly a teflon sheet was mistakenly inserted between collar and yoke!

All radial escapes were closed, and even with the magnet energy of 5 T field, annular space pressures reached 30 bar.

A higher field coil of 10 T would have squeezed/distroyed the beam-pipe.

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Safety: avoiding trapped volumes

Since the experience with this very early stage LHC-cold mass test the
explicit control of radial and longitudinal helium escape paths has
become an integral part of the fully helium immersed magnet cold-mass
designs and has never been an issue anymore

Safety: cold-mass pressure vessel

- A quench of a magnet-coil transfers about 25 % or more of the stored energy directly to the helium in which it is immersed
 - → pressure built up would be of the order of > 200 bar if not managed by safety relief devices
 - → The rate of pressure rise is coil-specific (porous or not p.e.) but initial adiabatic shock waves of several bar to about 10 bar cannot be avoided.
 - → The pressure rating of the cold-mass due to the helium pressure rise phenomena may be higher than desired because of this (think of machine-detector-interface magnets p.e.)



Safety: containing helium spill consequences

The shear amount of helium inside accelerator magnets has a heavy impact on personnel safety, especially in underground areas, and the infrastructure to deal with this:

- Helium release has to be captured in dedicated cryogenic transfer lines
- Access restrictions for machine maintenance become dependent on the amount of helium in the cold-masses and their powering status
- Should accidents break the cold-mass to insulation vacuum helium will instead of being captured by the cryogenic transfer line spill into the ambient environment → Asphyxiation hazard

Inventory handling and supplier dependence



Some points to remember

<u>Energy efficiency:</u> Incorporate in the magnet design, whenever possible, features to intercept heat at as high T as possible!

Energy efficiency: If possible, use conductors that have a high Tc

<u>Facilitating robust thermal design:</u> Incorporate temperature margin to deal with heat extraction (no magnet operates at an ideal homogeneous temperature)

<u>Inventory handling, supplier dependence, general & personnel safety:</u> Aim for reduced helium content, preferably conduction cooled magnets



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