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# LTS Conductors

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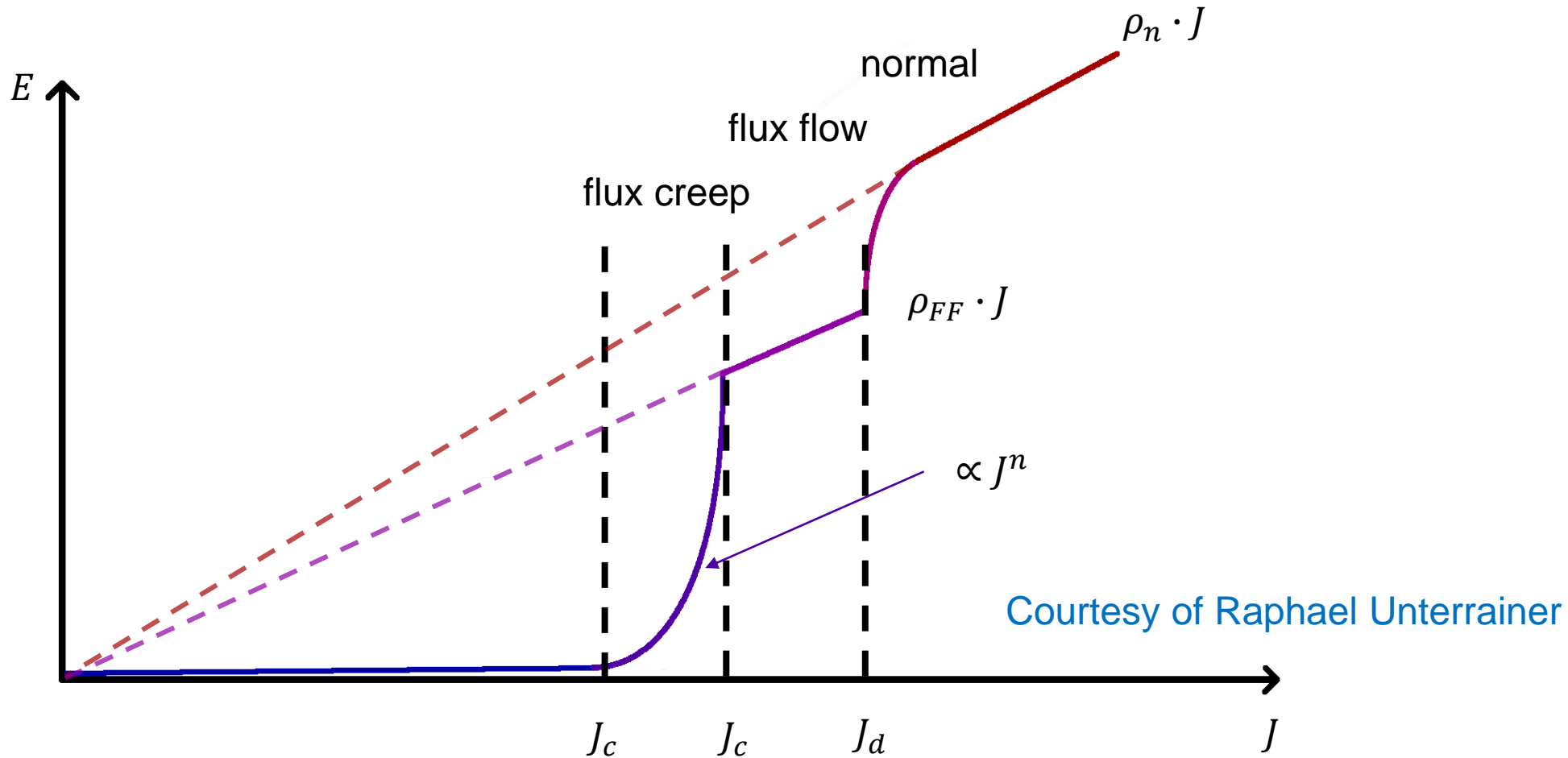


- Technical  $I_c$
- Stabilization
  - Multifilamentary wires
- Superconducting materials used for conductors
  - NbTi
  - Nb<sub>3</sub>Sn
  - (MgB<sub>2</sub>)
- Effect of forces on superconductors



# Technical $I_c$

- Dissipation in superconductors

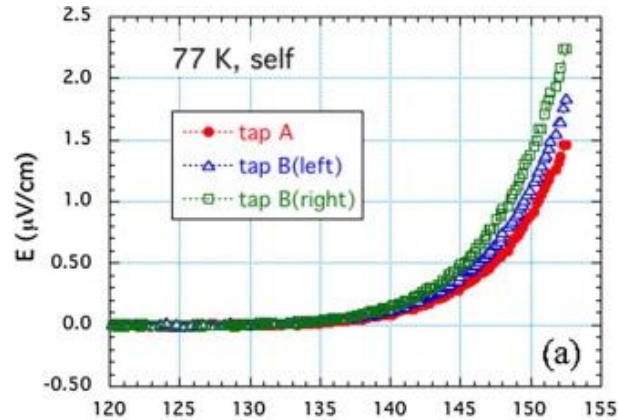


$J_c \neq J_c$ ? Technical  $J_c$  defined by arbitrary electric field criterion, e.g.  $0.1 \mu\text{V}/\text{cm}$  4



# Resistive transition

- Transition can be often approximated by  $E = E_c \left(\frac{J}{J_c}\right)^n$  or  $I \propto E_c \left(\frac{I}{I_c}\right)^n$ 
  - Linear on a double logarithmic scale



Tsuchiya et al., *Cryogenics* **85** (2017) 1

- Two main reasons for a broadening of the transition
  - Intrinsic: thermal activation of vortices out of the pinning potential  $U$

$$n = \frac{U}{k_B T}$$

- Material inhomogeneities  $n = ?$
- “Well-behaved” superconductors:  $n \gtrsim 30$
- Important for applications: predictable (extremely) low loss behavior.



# Requirements on a superconducting wire

***A high current density is a necessary condition for applications, but it is not sufficient.***

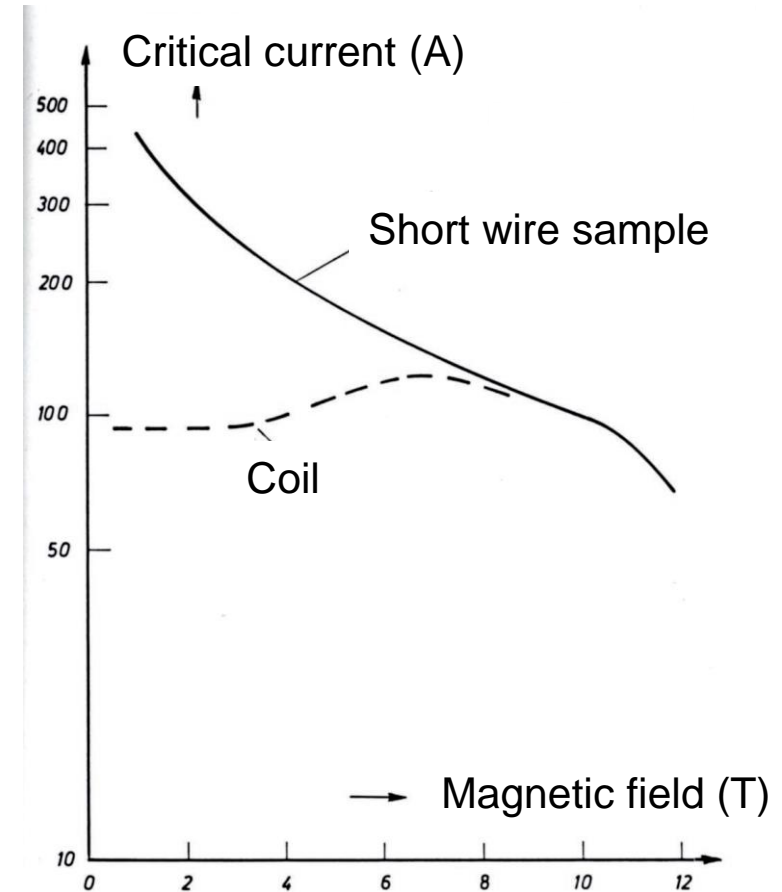
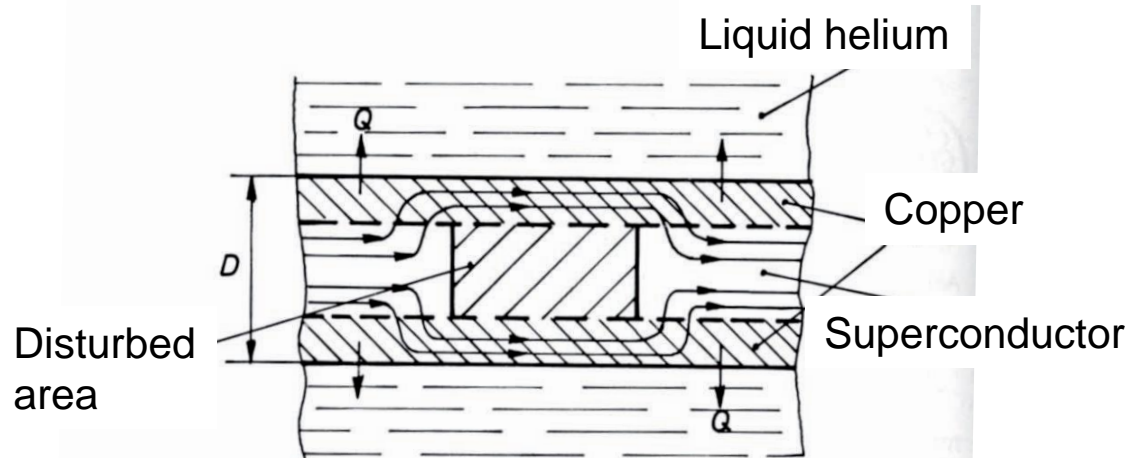
## ***Other crucial requirements:***

- ***Have high tolerance to stress*** *Magnetic forces*
- ***Be safe in case of magnet quench*** *Quench detection, NZPV*
- ***Have low magnetization*** *Applications to NMR, MRI, HEP magnets*
- ***Have a persistent joint technology*** *Applications to NMR, MRI*
- ***Stability against small thermal perturbances***

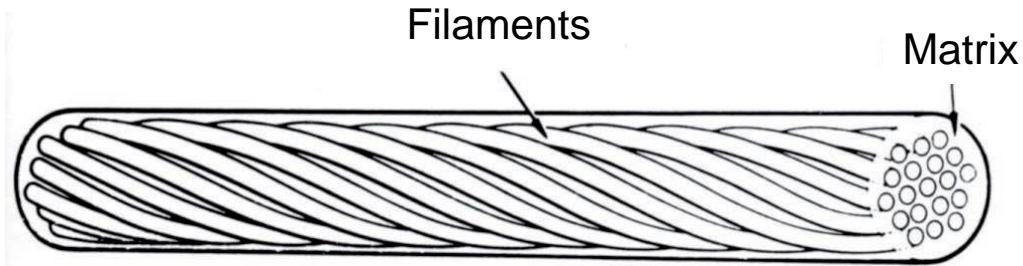


# Stabilization

- A technical conductor needs stabilization otherwise:
  - A small disturbance leads to dissipation
  - The superconductor heats up
  - Dissipation increases
  - Quench
- Idea: Current can bypass the disturbed area
  - Allows the superconductor to recover



# Multifilamentary wires



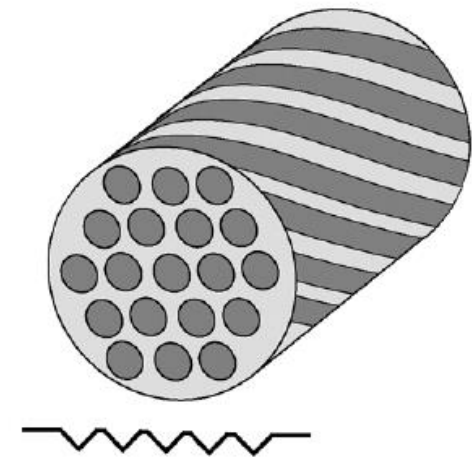
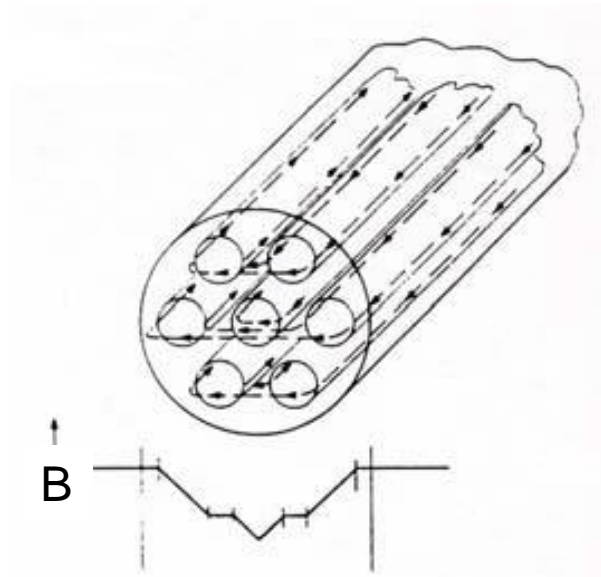
- Thin filaments: Large surface to volume fraction
  - Efficient heat and current transfer
- Matrix
  - Low resistivity
  - Compatible properties
- Twisted filaments
  - Avoid inductions loops





# Multifilamentary wires

- Induction loops
  - store energy that may trigger an instability
  - distort the field quality
- Self screening of filaments cannot be avoided
  - Transport current only flows in a small part of the filament when ramping a magnet at low fields
  - Remaining part shields any field changes
- Current loops can close over the low ohmic matrix driven by the induction voltage (or over bridges)
- Twisting reduces inductive filament coupling efficiently
  - Current loops become smaller (longitudinally and transversal)
  - A local disturbance releases less energy



# LTS materials (typical values)

	$T_c$ (K)	$B_{c2}$ (T) <sup>1</sup>	$J_d$ (MA/cm <sup>2</sup> ) <sup>1</sup>	Advantages	Disadvantages
NbTi	9.5	11	40	<ul style="list-style-type: none"> <li>• easy to handle</li> <li>• optimized</li> <li>• cheap</li> </ul>	<ul style="list-style-type: none"> <li>• low <math>T_c</math></li> <li>• small <math>B_{c2}</math></li> </ul>
Nb <sub>3</sub> Sn	18	27	200	<ul style="list-style-type: none"> <li>• higher <math>T_c</math></li> <li>• high <math>B_{c2}</math></li> </ul>	<ul style="list-style-type: none"> <li>• brittle</li> <li>• in-situ process</li> <li>• wind and react</li> <li>• expensive</li> <li>• under optimization</li> </ul>
MgB <sub>2</sub>	39	10	130	<ul style="list-style-type: none"> <li>• “High“ <math>T_c</math></li> <li>• cheap?</li> </ul>	<ul style="list-style-type: none"> <li>• small <math>B_{c2}</math></li> <li>• wind and react?</li> <li>• expensive</li> <li>• not optimized</li> </ul>

<sup>1</sup>at 4.2 K



# NbTi

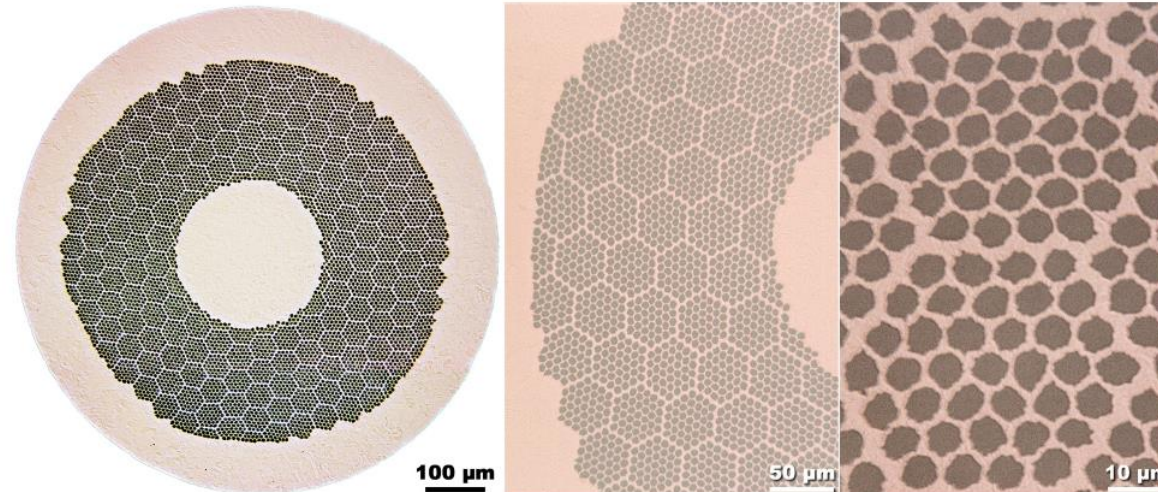
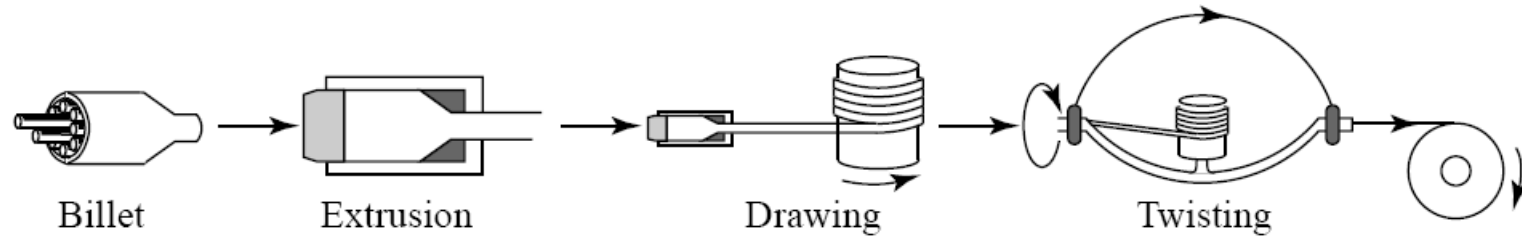
Life is easy



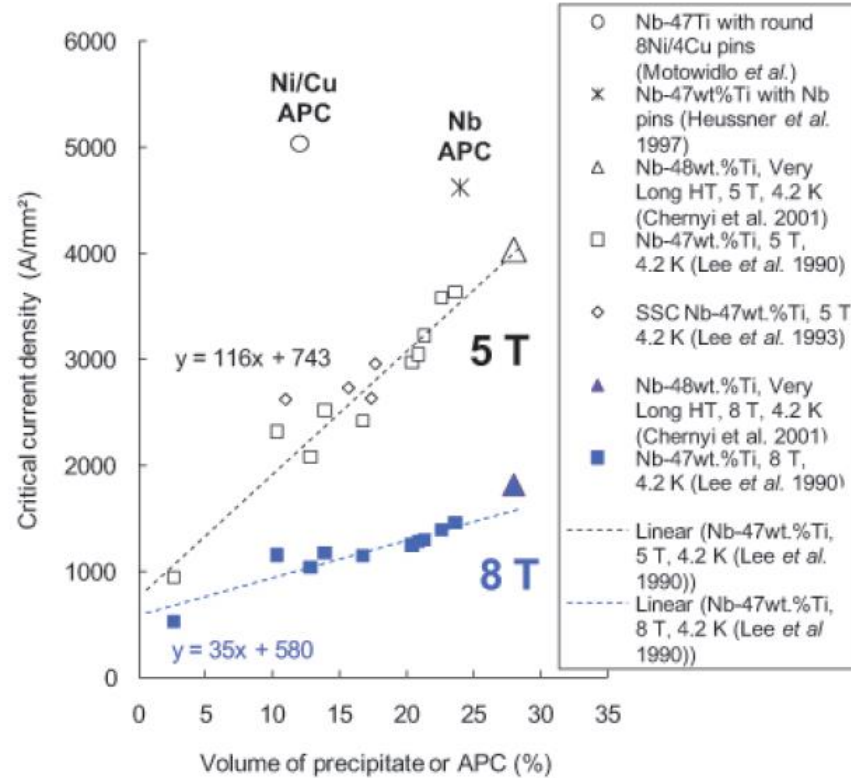
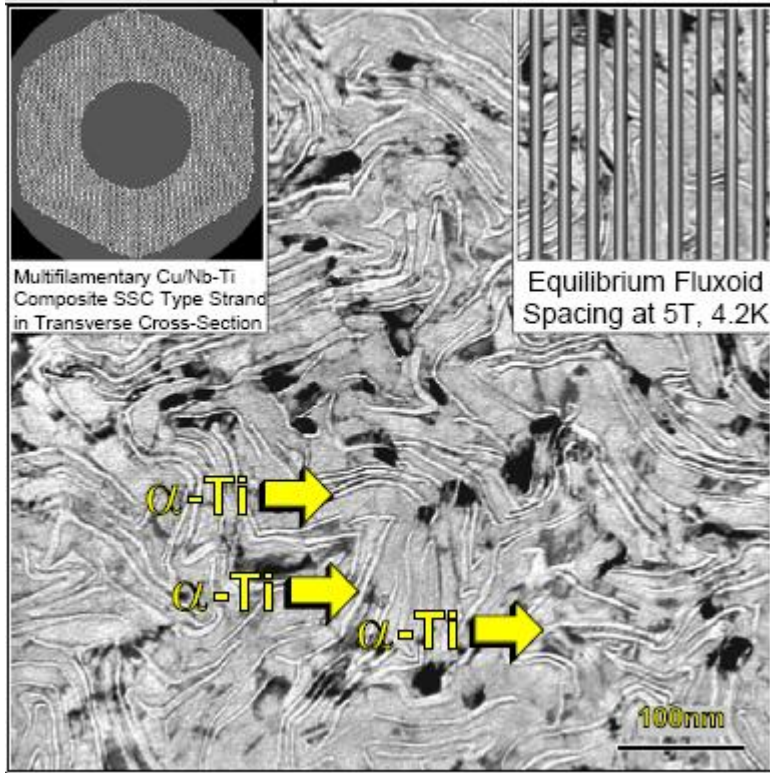
# Nb-Ti : the King of the Hill

Type II	$T_c$ [K]	$\mu_0 H_{c2}^*$ [T]
Nb (metals)	9.5	0.2*
NbTi (alloys)	9.8	10.5†

- Enabling technology for the large diffusion of MRI (a 4'000 M€ market!)
- 1200+ tonnes of Nb-Ti in LHC



# Nb-47wt%Ti : How to get high $J_c$



$\alpha$ -Ti  $\rightarrow$  hcp  
 $\beta$ -Ti  $\rightarrow$  bcc

**FIGURE 11.15:** TEM image of the microstructure (transverse cross-section) of the first 3700 A/mm<sup>2</sup> (5 T, 4.2 K) multifilamentary strand from a US manufacturer (OST). This previously unpublished image taken on September 5<sup>th</sup> 1986, shows the dense array of folded  $\alpha$ -Ti ribbons (lighter contrast) that create the strong vortex pinning.

$\alpha$ -Ti precipitates are adjusted to the proper dimensions in order to pin vortices

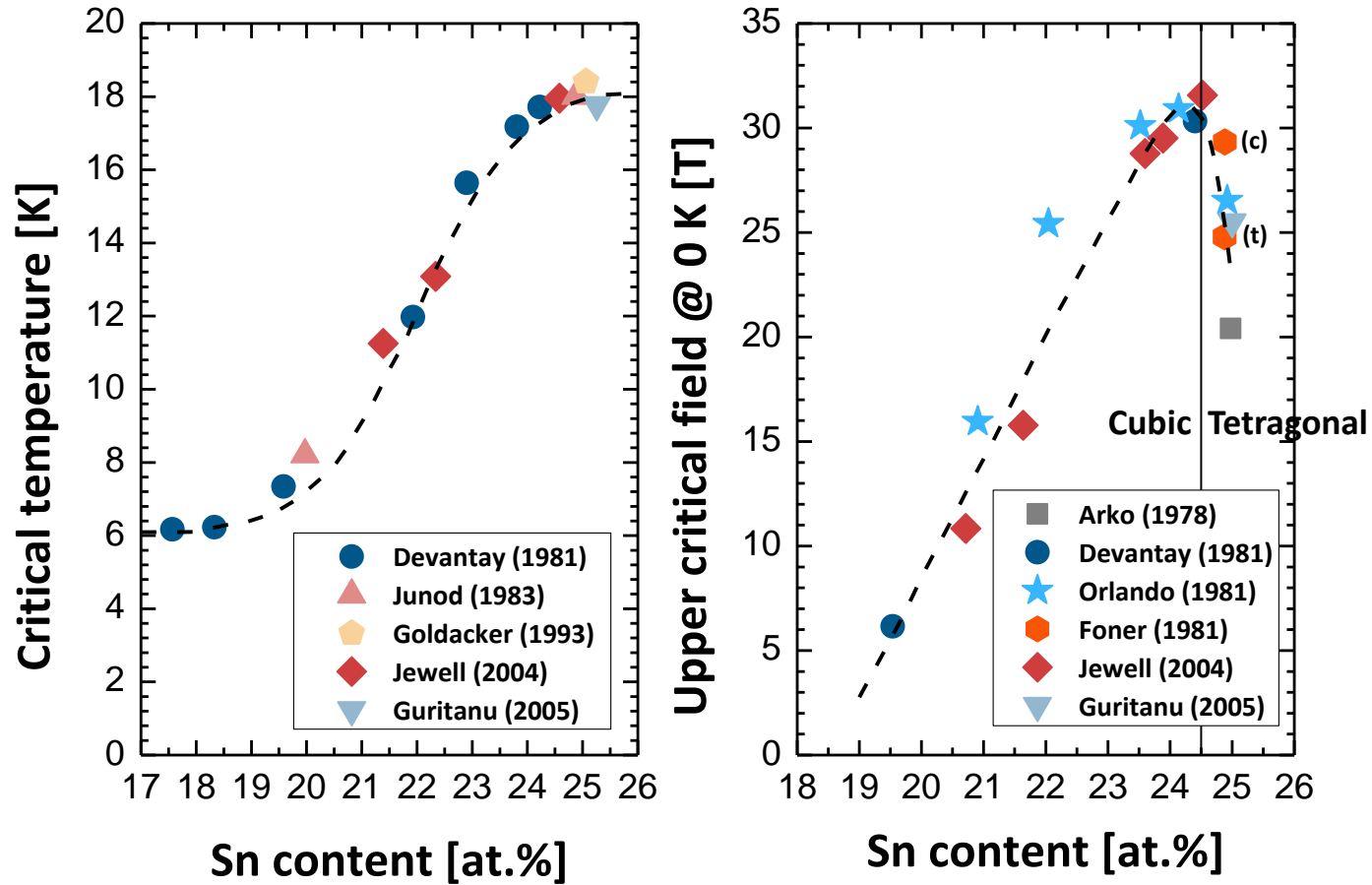


# $\text{Nb}_3\text{Sn}$

It's a hard life



## Influence of the Sn content on T<sub>c</sub> and B<sub>c2</sub>



Nb<sub>3+x</sub>Sn<sub>1-x</sub> remains superconducting when it deviates from stoichiometry



# How to rise $H_{c2}$ – Let's play it dirty

Upper critical field of a superconductor

$$H_{c2} = \frac{\Phi_0}{2\pi\xi^2}$$

Disorder reduces the electron mean free path  $l$ , which in turn leads to decrease of  $\xi$

$$\frac{1}{\xi(l)} = \frac{1}{\xi(\infty)} + \frac{1}{l}$$

An useful expression of  $H_{c2}$  in the dirty limit

$$H_{c2}(T=0) \cong \frac{k_B e}{\mu_0} N(E_F) \rho_n T_c \propto \gamma \rho_n T_c$$

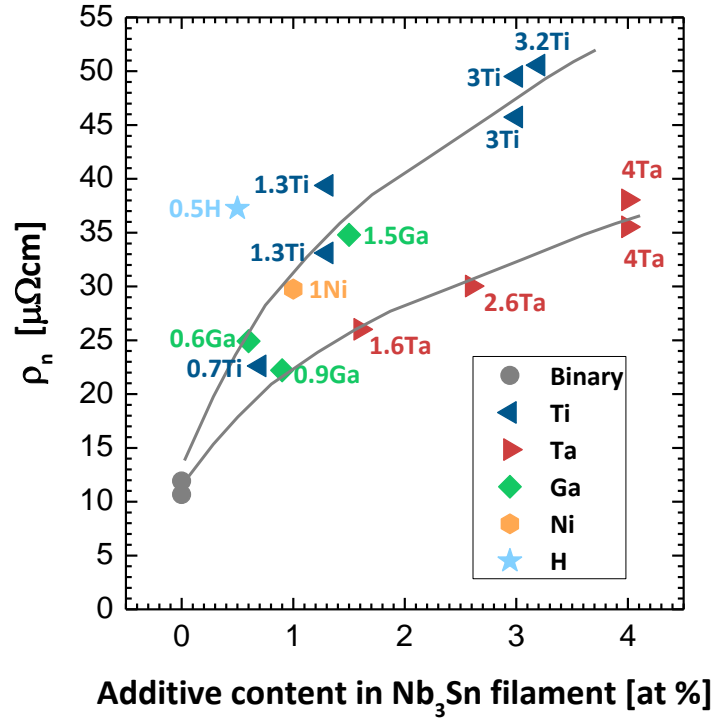




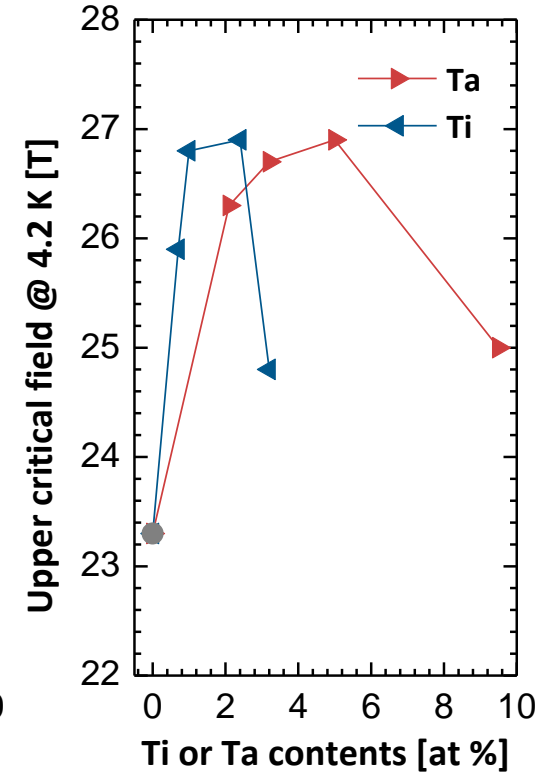
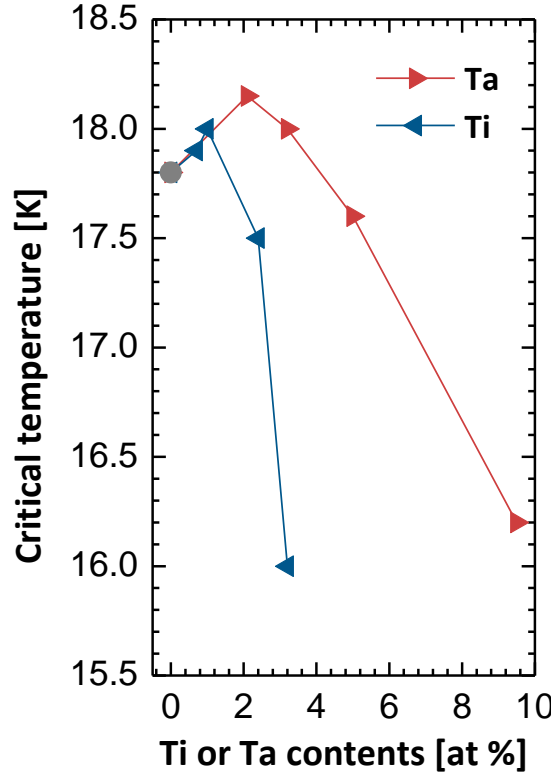
# Alloying (doping) Nb<sub>3</sub>Sn to rise H<sub>c2</sub>

The additions of Ta and Ti are particularly beneficial

$$H_{c2} \propto \gamma \rho_n T_c$$



R. Flükiger et al., *Cryogenics* 48 (2008) 293



Ti substitutes Nb

Ta substitutes both Nb and Sn

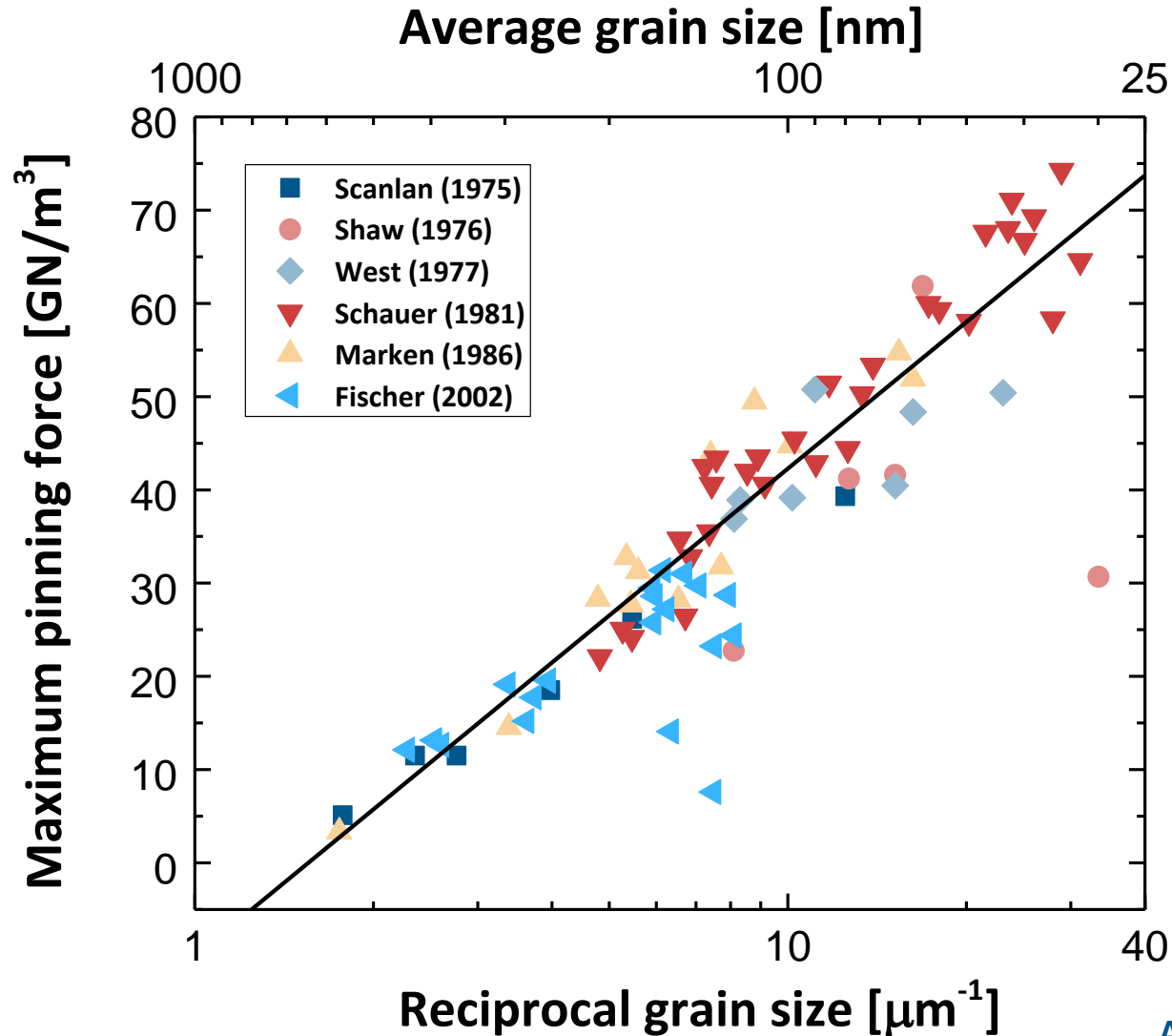
$$\frac{H_{c2}(4.2 \text{ K})}{H_{c2}(0 \text{ K})} = 0.89$$

M. Suenaga et al., *JAP* 59 (1986) 840

S.M. Heald et al., *Sci. Rep.* 8 (2018) 4798

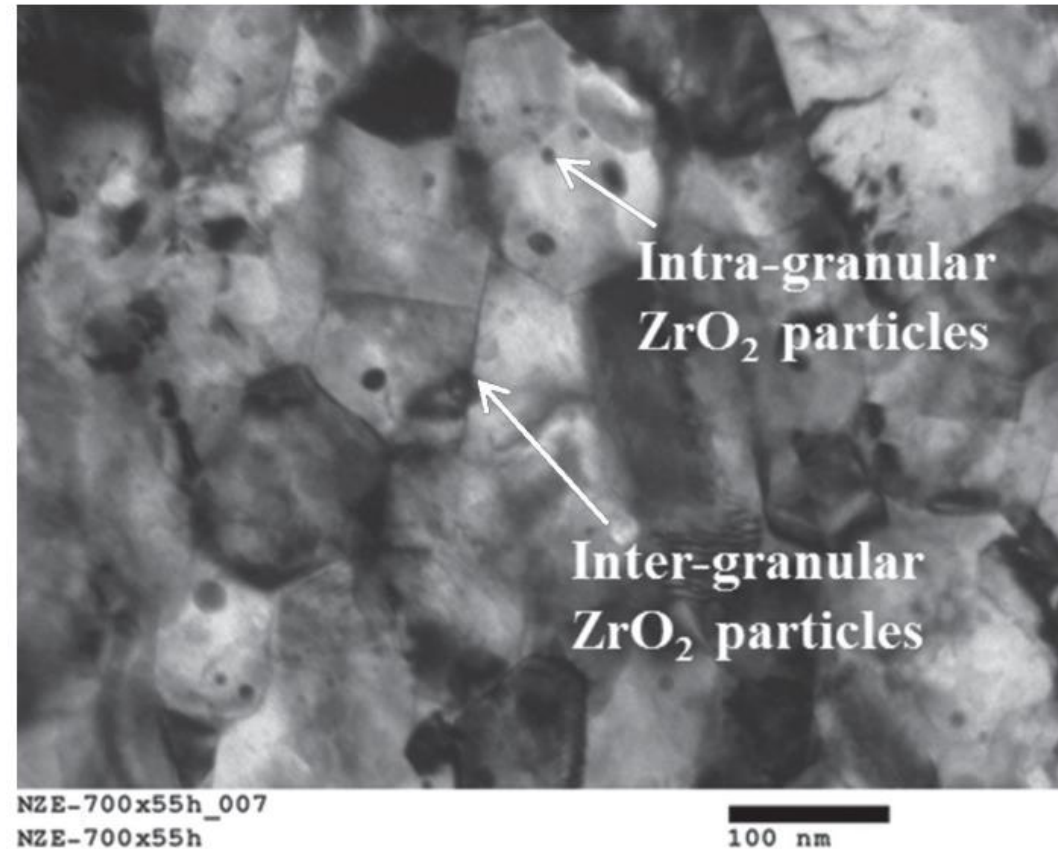


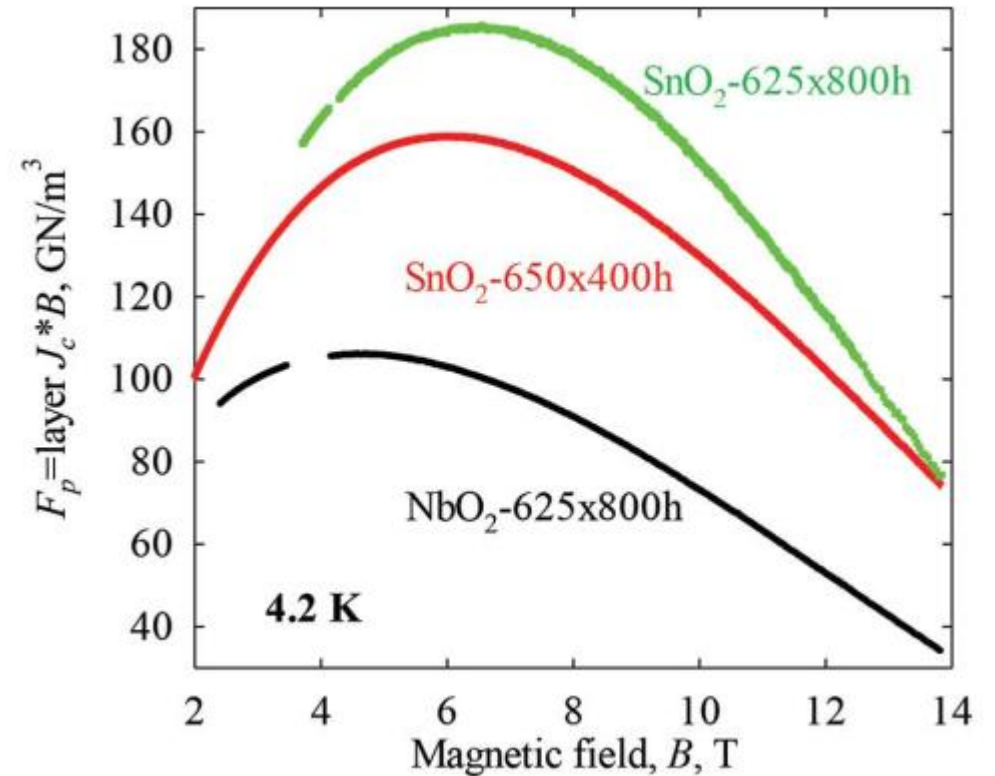
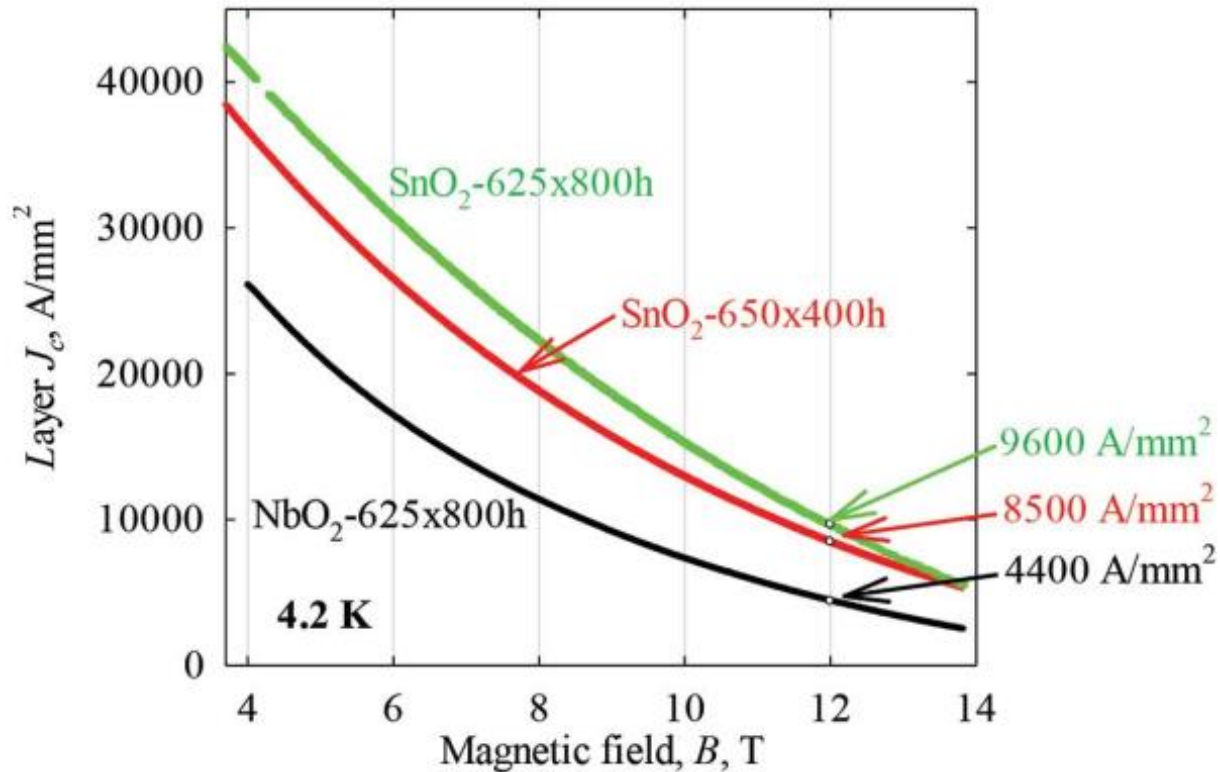
## Grain boundaries impede vortex motion to increase $J_c$



## Introduction of ZrO<sub>2</sub> nanoparticles

- Alloy Zr to Nb
- Add an oxygen source  
(e.g. NbO<sub>2</sub> or SnO<sub>2</sub>)
- ZrO<sub>2</sub> forms during heat treatment  
(„internal oxidation“)





$J_c$  increases due to both a grain refinement and pinning on nanoparticles.

Not commercially available yet.

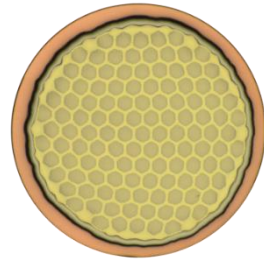


# $\text{Nb}_3\text{Sn}$

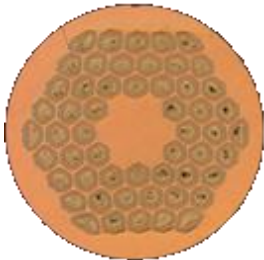
## Production routes



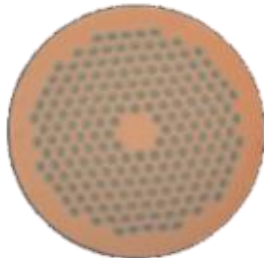
Three technologies have been developed at industrial scale



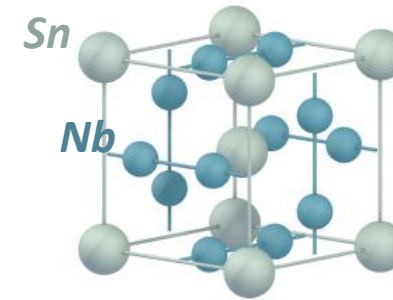
- *Bronze route*



- *Internal Sn diffusion*



- *Powder-In-Tube (PIT)*

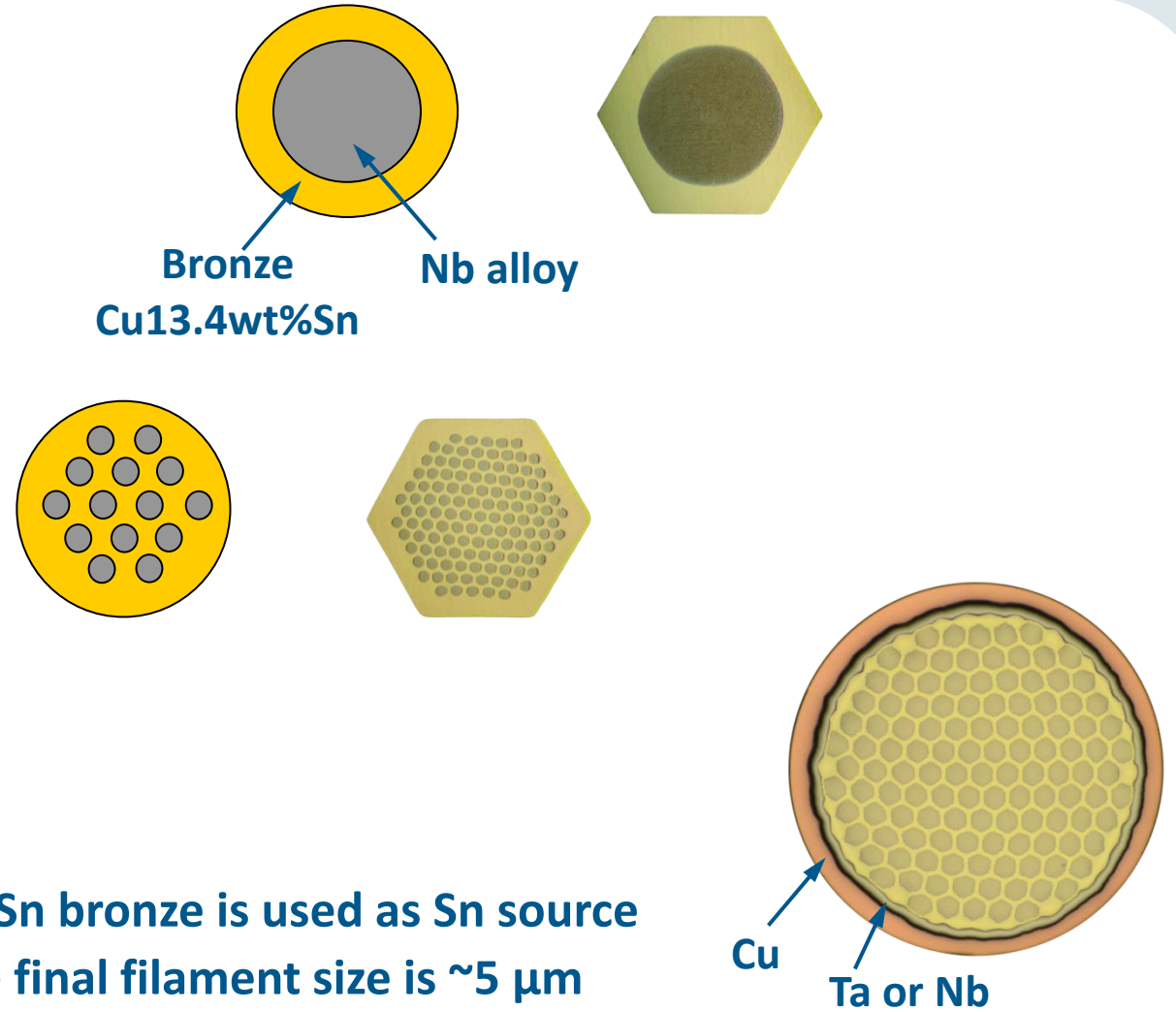
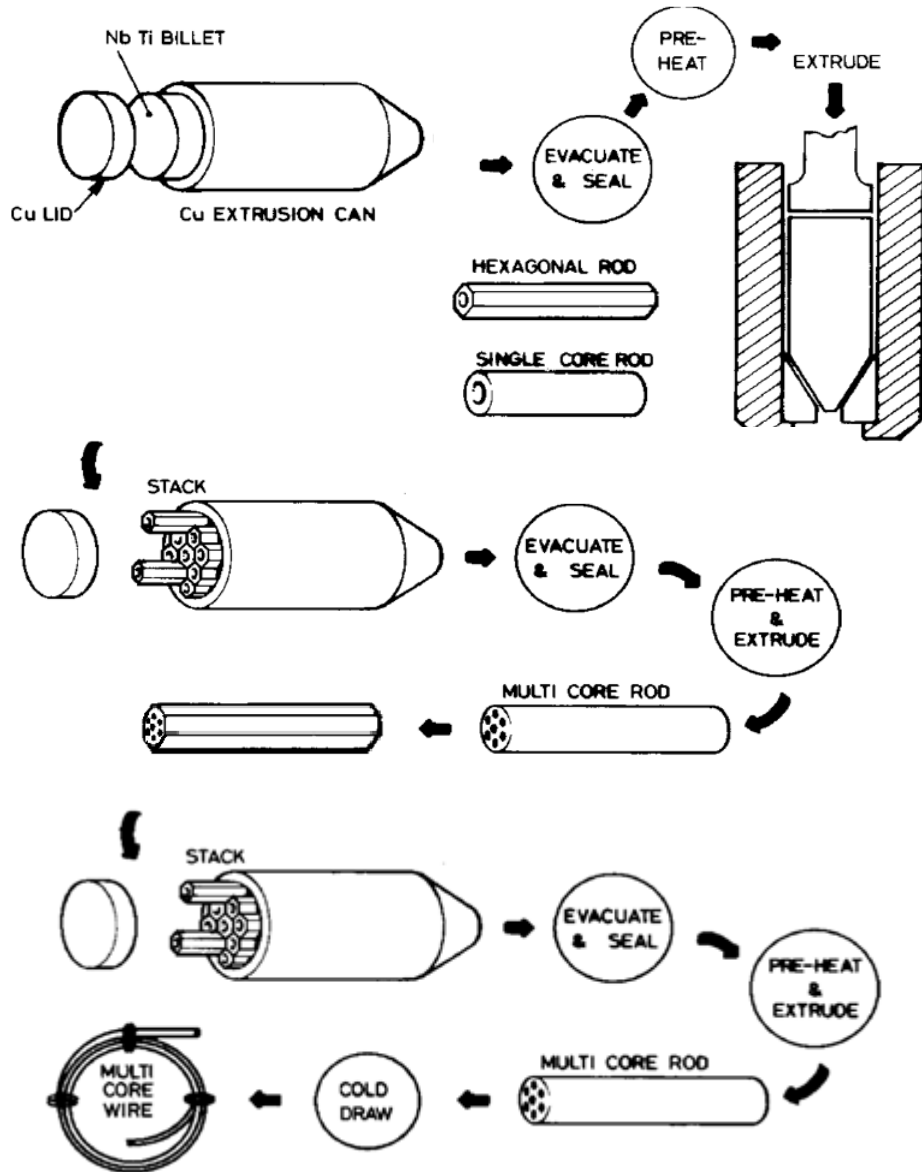


*The Sn source is the main difference*

Presently produced by

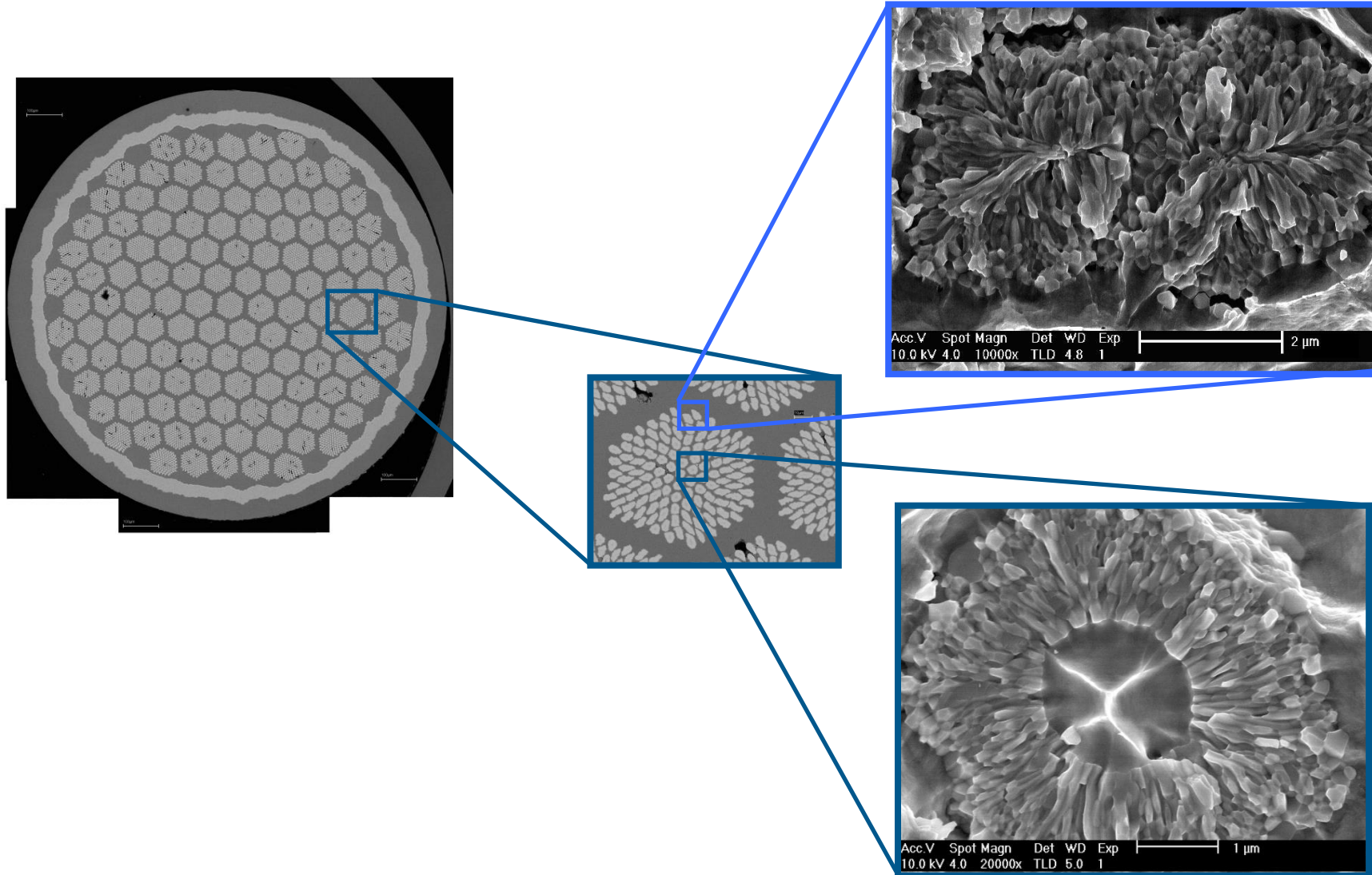


# The Bronze Process for multifilamentary Nb<sub>3</sub>Sn wires



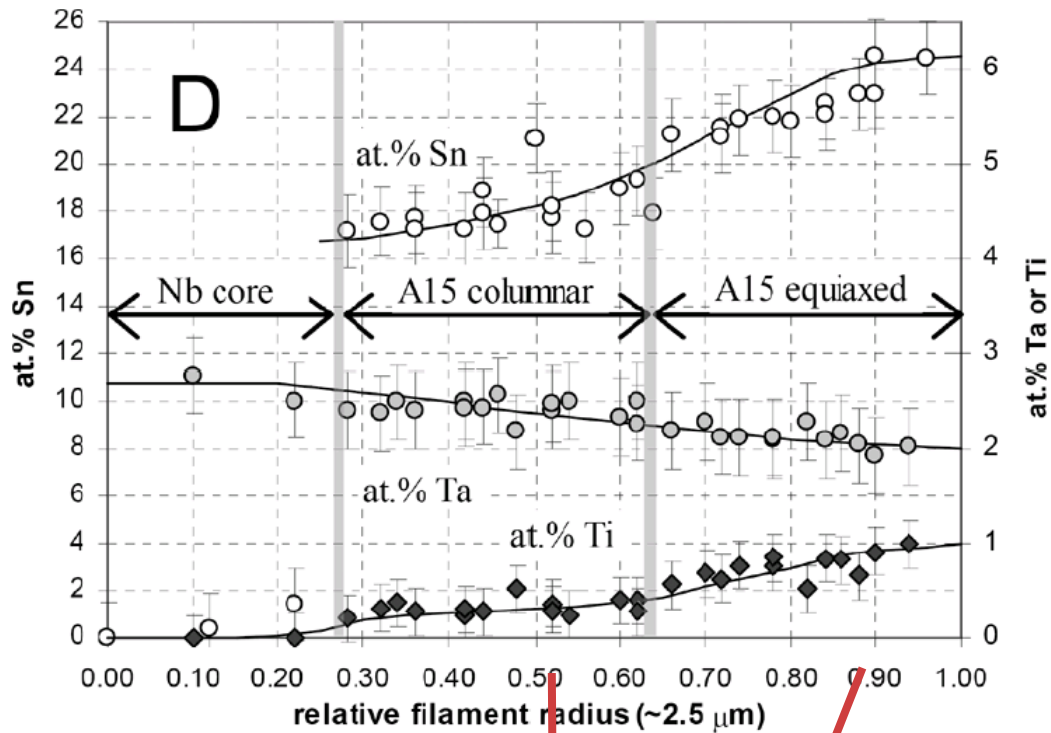
Cu-Sn bronze is used as Sn source  
 The final filament size is ~5 μm  
 Wires are then reacted at ~650°C  
 for >100 hours to form Nb<sub>3</sub>Sn

# Bronze Route $Nb_3Sn$ wires, after reaction





# Sn gradient over the filament radius: Bronze Route



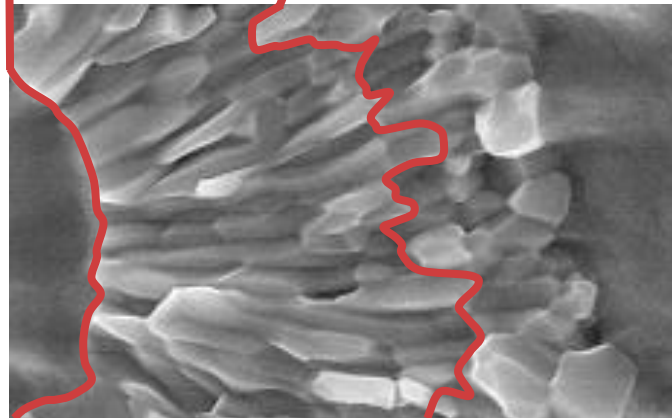
**Correlation between Sn content and grain morphology**

**Equiaxed grains: 21-25 at.% Sn**

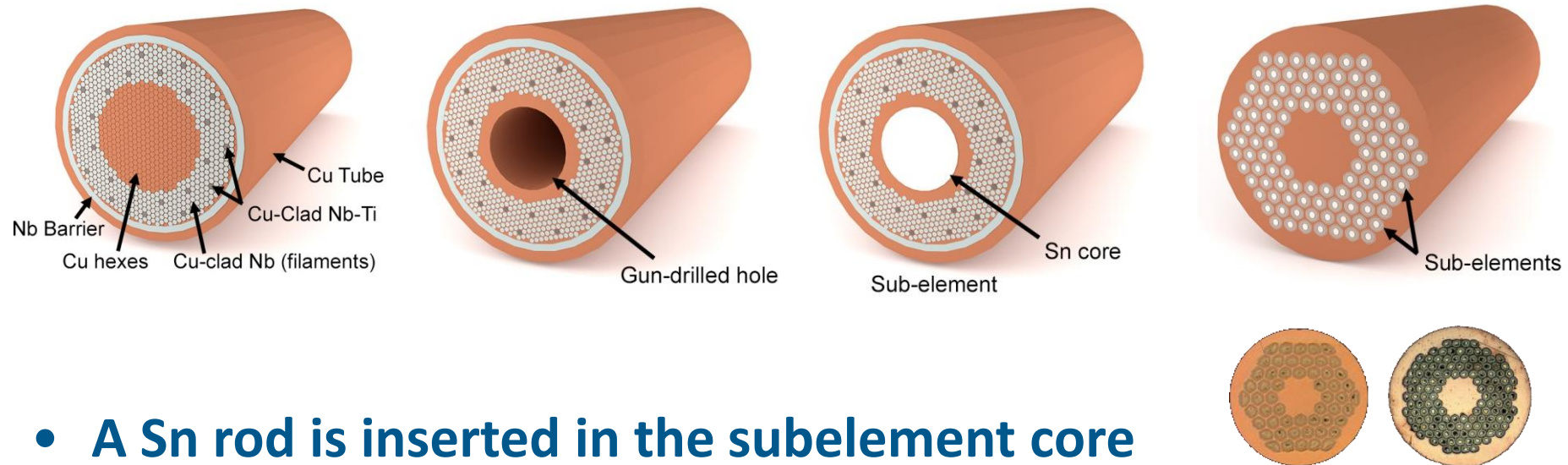
**Columnar grains: 18-21 at.%**

**Equiaxed grain size ~150 nm**

**Columnar grain size up to 400 nm**



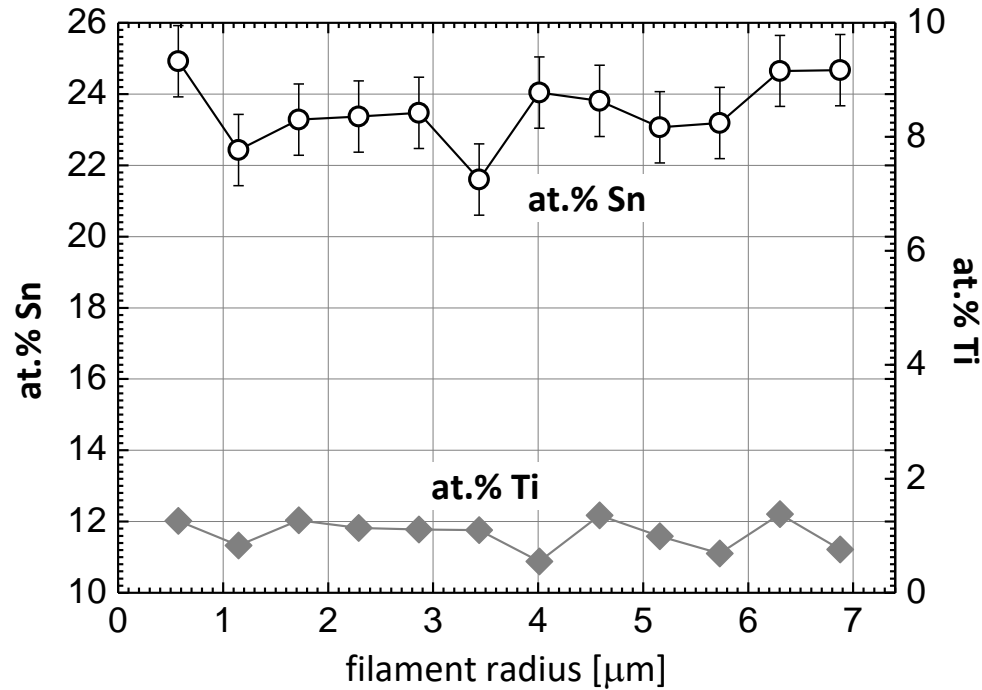
# The Internal Sn diffusion process



- A Sn rod is inserted in the subelement core
- After the insertion of Sn, only cold deformations are possible
- Subelement size ranges between 20 and 100  $\mu\text{m}$
- A long-duration multistep reaction schedule is required to form  $\text{Nb}_3\text{Sn}$

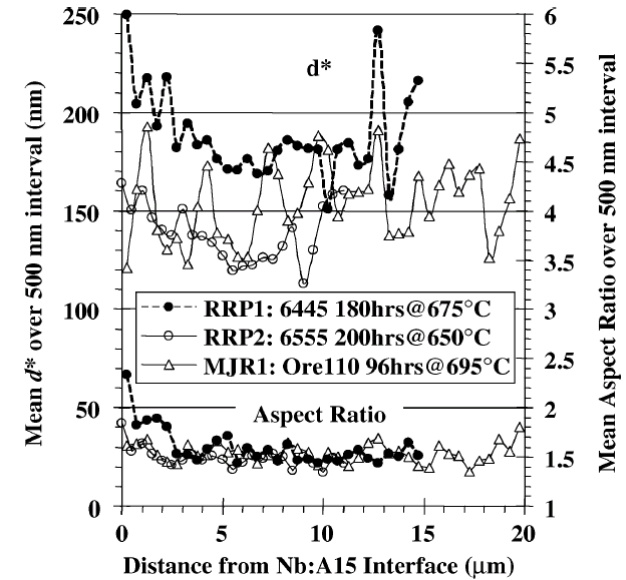
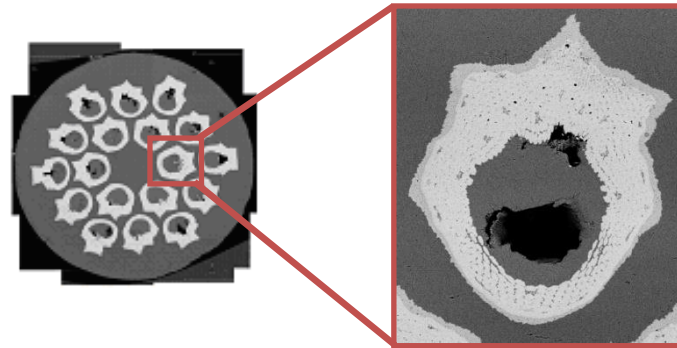


# Sn gradient across the filament radius: Internal Sn

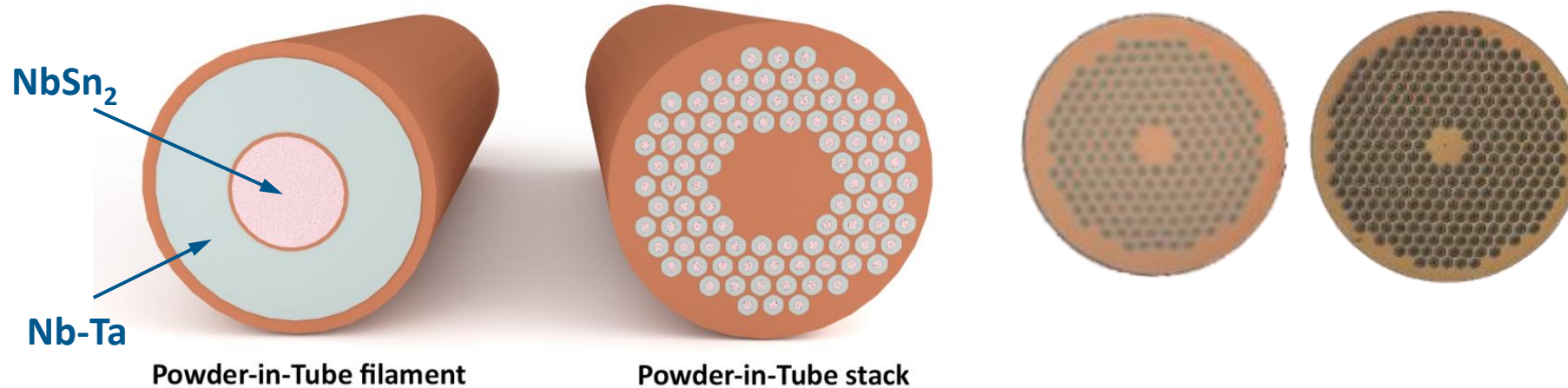


D. Uglietti, PhD Thesis (2006), UNIGE

**All the grains are equiaxed and almost stoichiometric !!**



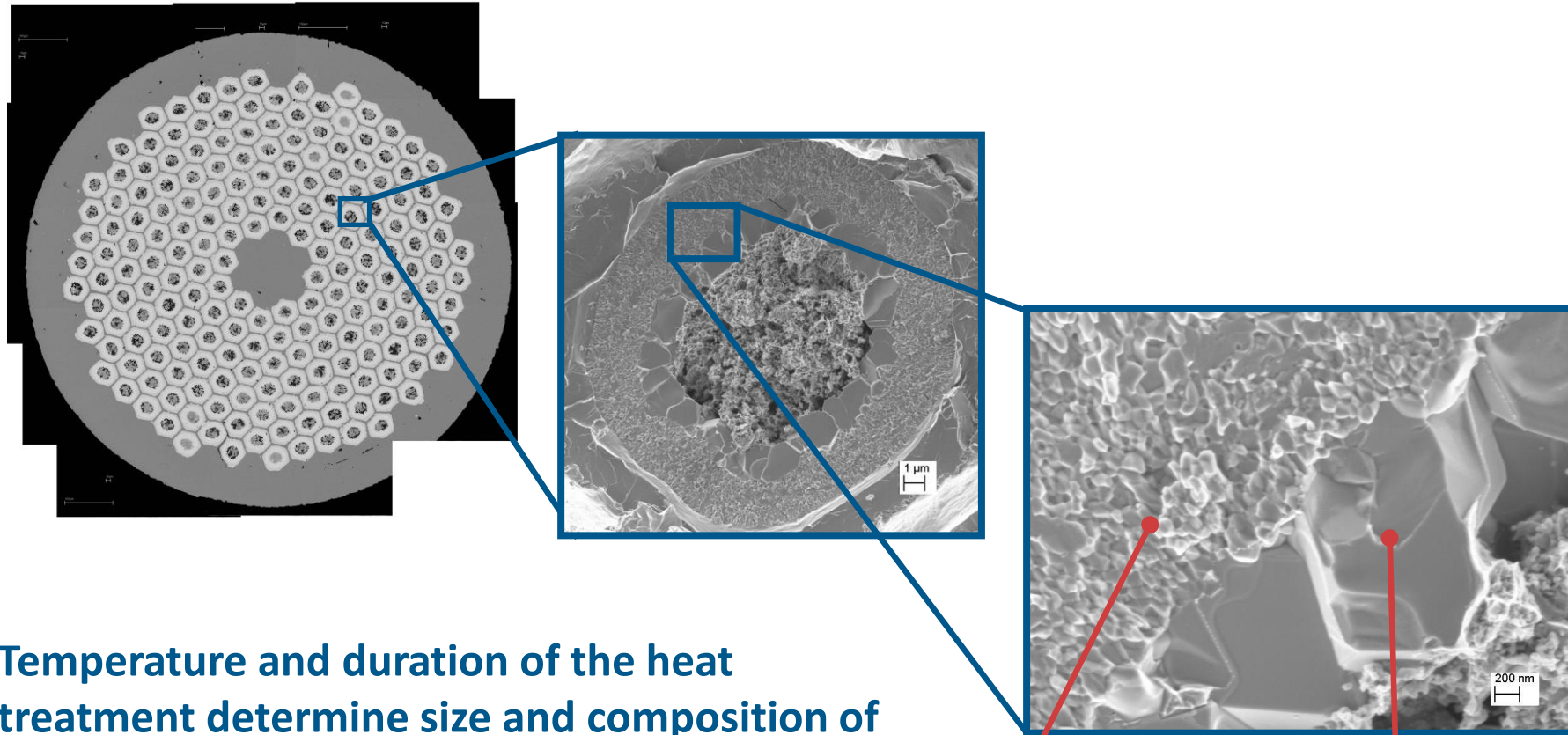
# The Powder-In-Tube method



- **NbSn<sub>2</sub> and Sn powders are used as Sn source**
- **Subelement size ranges between 20 and 100 μm**
- **A long-duration multistep reaction schedule is required to form Nb<sub>3</sub>Sn**



# Formation and Microstructure of the A15 phases

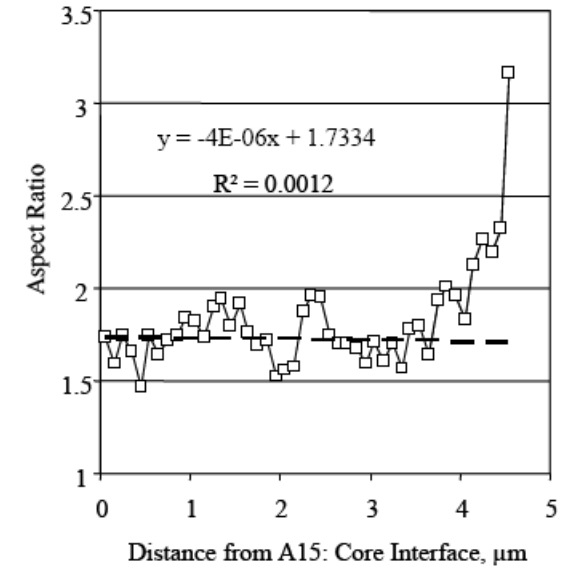
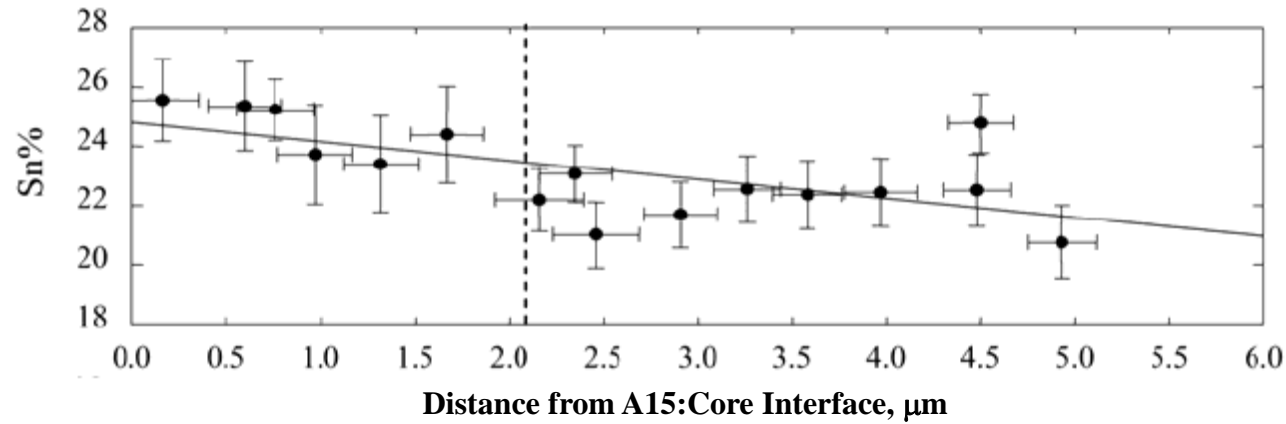


Temperature and duration of the heat treatment determine size and composition of the two A15 regions

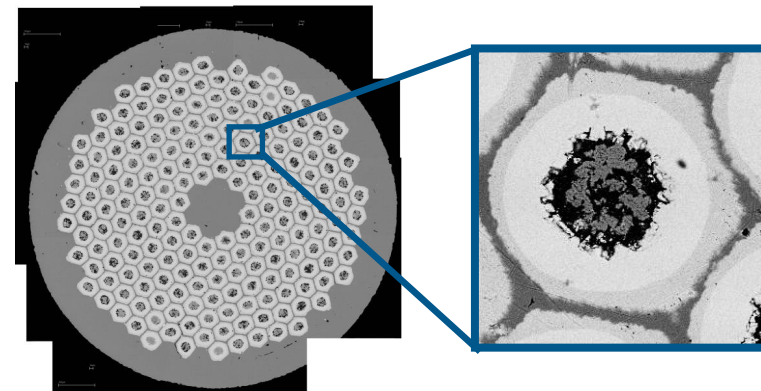
Fine grains (~ 200 nm)  
~ 23 at.% Sn

Large grains (> 1 μm)  
~ 25 at.% Sn

# Sn gradient over the filament radius: PIT

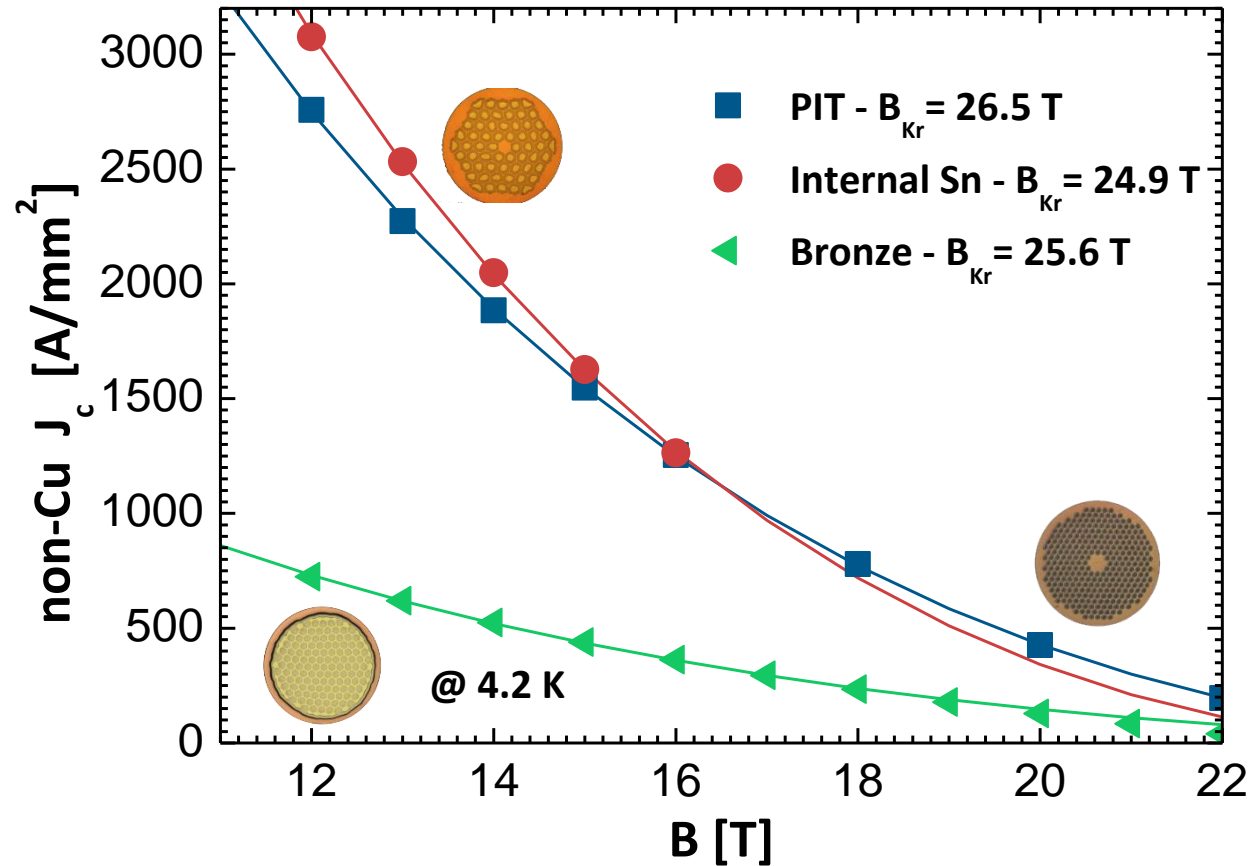


- Sn content decreases linearly along the filament radius
- Both fine and large grains are (almost) equiaxed



# Critical current density vs. magnetic field

Best performance achieved so far in industrial wires



[T. Boutboul et al., IEEE TASC 19 \(2009\) 2564](#)

[J. Parrell et al., AIP Conf. Proc. 711 \(2004\) 369](#)

[V. Abächerli et al., IEEE TASC 17 \(2007\) 2564](#)



# Nb<sub>3</sub>Sn

It's a hard life (stress and strain)

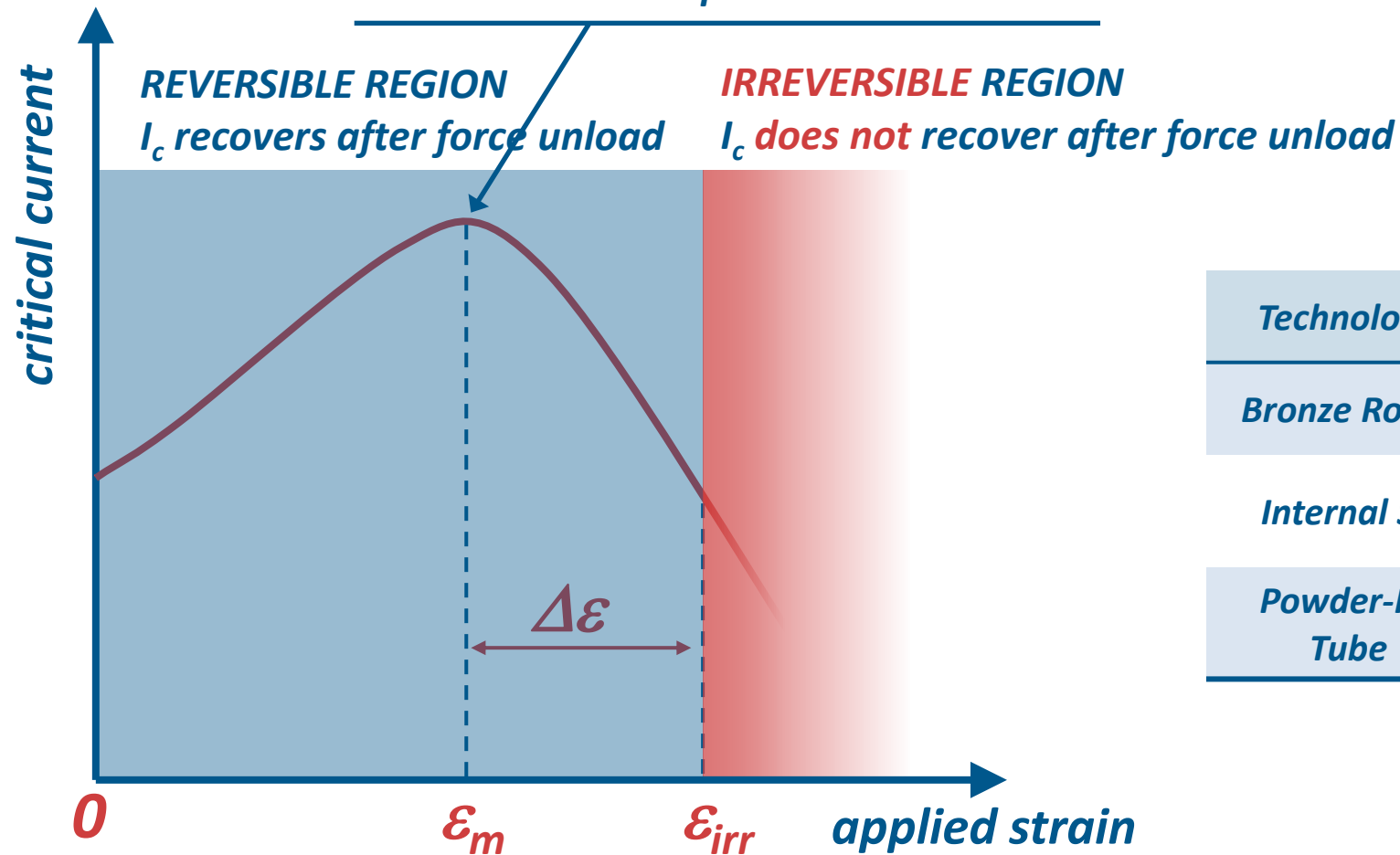




# Strain-induced changes in the critical current

## Effects of the longitudinal strain

The applied force reduces the effects of the thermal precompression and the critical current increases up to a maximum

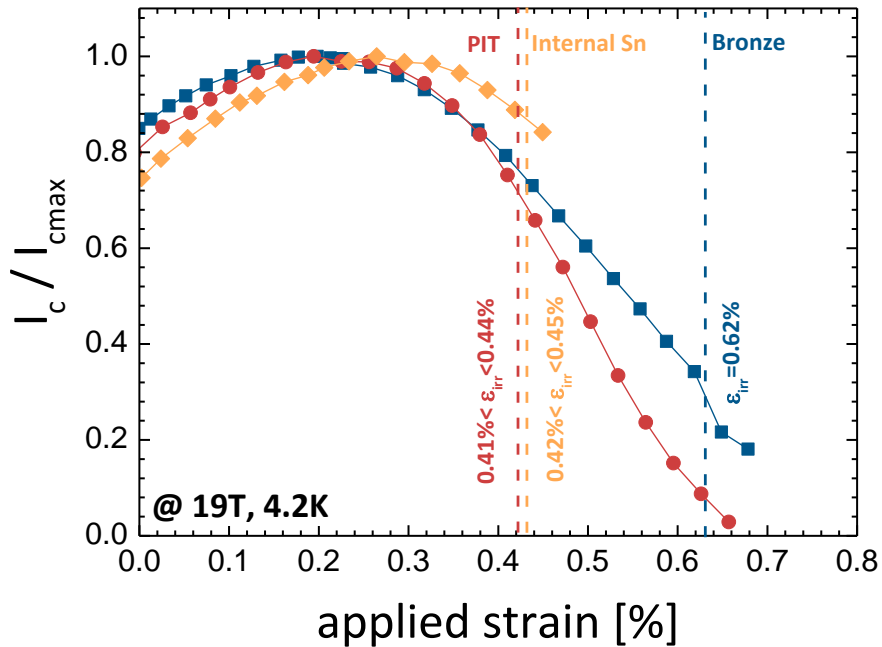


Technology	$\Delta\epsilon = \epsilon_m - \epsilon_{irr}$ [%]
Bronze Route	0.4 – 0.6
Internal Sn	0.05 – 0.2
Powder-In-Tube	0.15 – 0.3

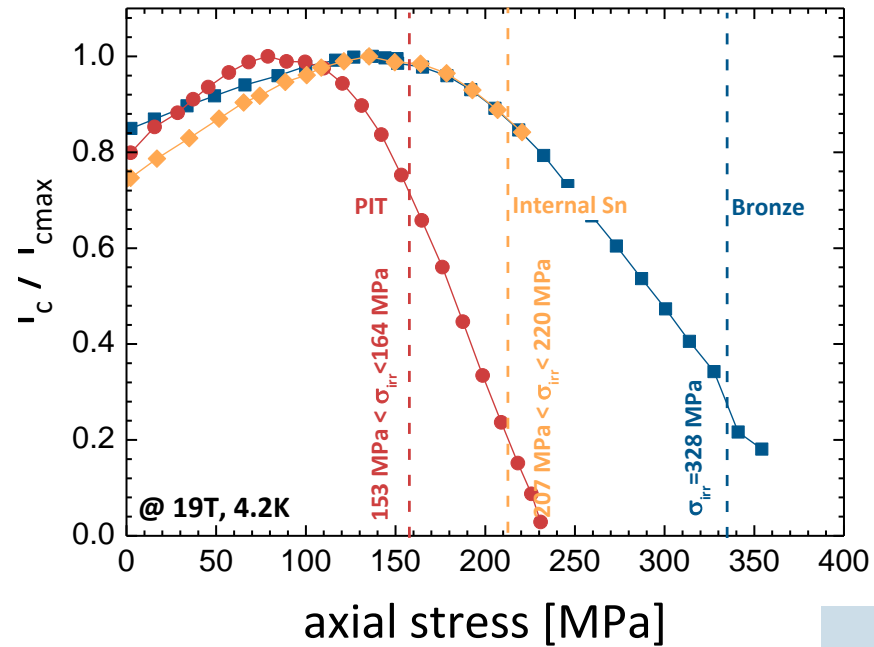


## Reversible behaviour and irreversible limit

### $I_c$ vs. axial strain



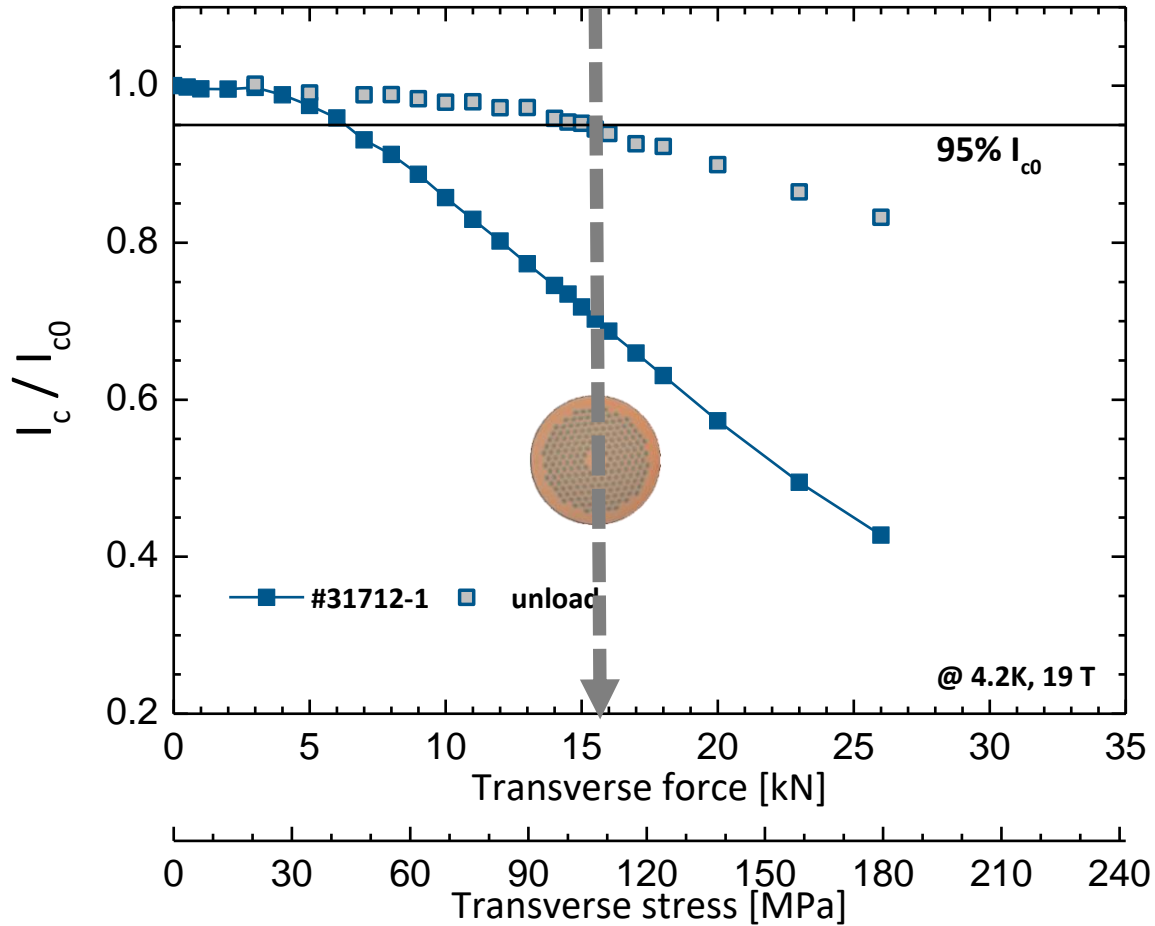
### $I_c$ vs. axial stress



Technology	$\sigma_{irr}$
Bronze Route	330 MPa
Internal Sn	210 MPa
Powder-In-Tube	150 MPa



# $I_c$ vs. transverse stress



The irreversible limit is defined at the force level leading to a 95% recovery of the initial  $I_c$  after unload

Here

$$F_{irr} = 16 \text{ kN}$$

The corresponding irreversible stress limit is

$$\sigma_{irr} = 110 \text{ MPa}$$

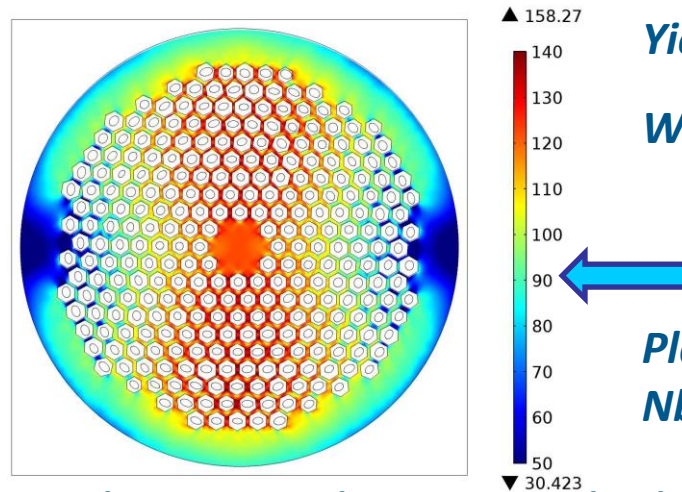
where

$$\text{Stress} = \frac{\text{Force}}{\text{groove length} \times \text{groove width}}$$



## Two irreversible phenomena play together

- **Plastic deformation of the Cu matrix**



*PIT Nb<sub>3</sub>Sn wire under transverse load*  
*Stress map of the Cu matrix*

*Yield strength of Cu  $\sigma_y \sim 90$  MPa*

*What is the residual stress on Cu after unload ?*

$$\sigma_{res}^{Cu} = \max(\sigma_{Mises}^{Cu} - \sigma_y, 0)$$

*Plastically deformed Cu imposes a stress on Nb<sub>3</sub>Sn after force unload (= lattice deformation)*

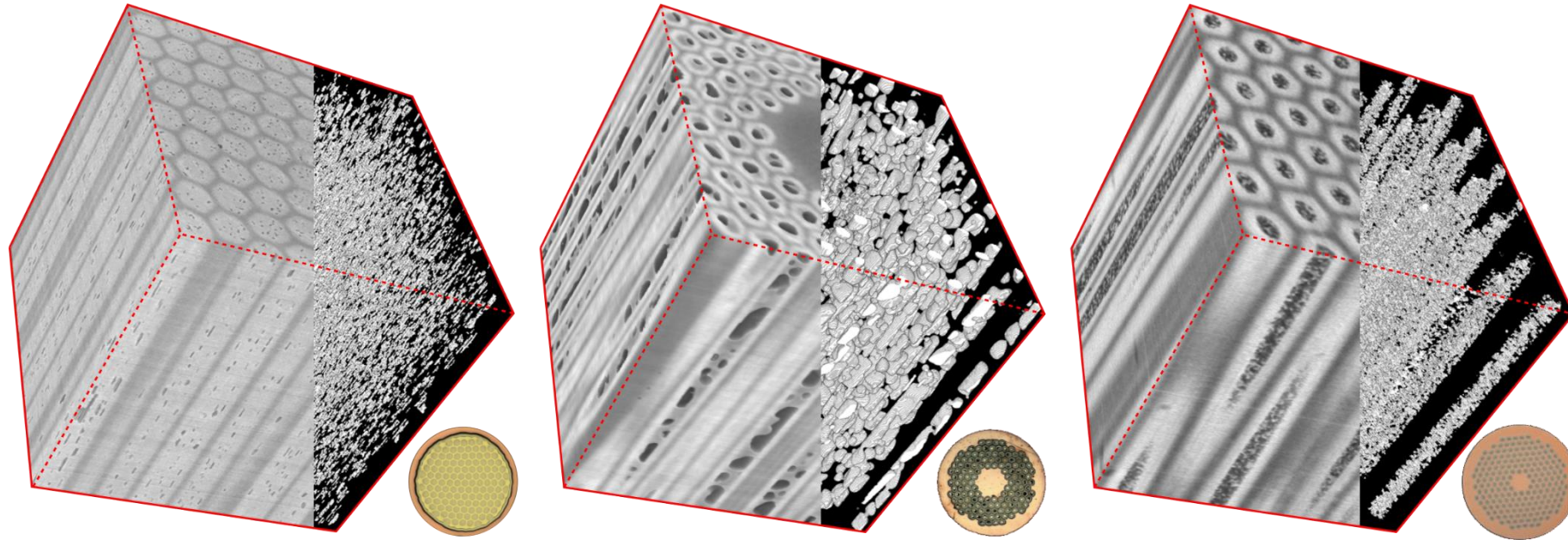
- **Crack formation in Nb<sub>3</sub>Sn**

*Nb<sub>3</sub>Sn is a brittle material and is characterized by a strong propensity to fracture*

*Voids formed during the reaction cause localized stress concentrations where cracks nucleate*



## XRD microtomography reconstruction



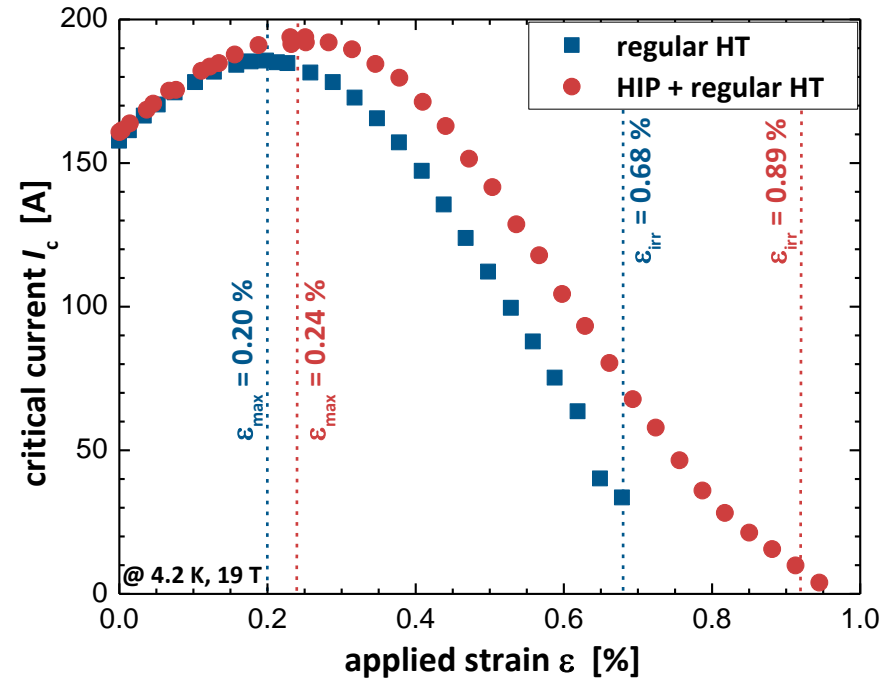
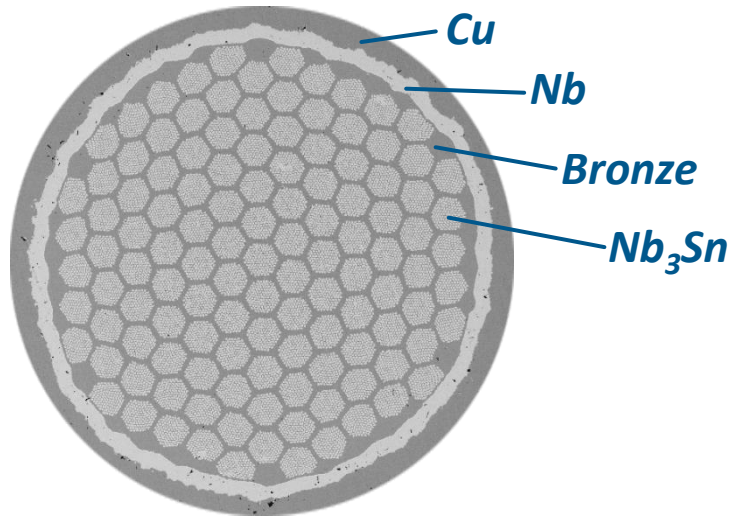
**Bronze Route**  
121 x 121 filaments

**Internal Sn**  
132/169 subelements

**PIT**  
192 filaments

*Can we quantify the impact of voids on the electromechanical limits?*

# A case study on Bronze Route $Nb_3Sn$ wires



manufacturer	UNIVERSITÉ DE GENÈVE FACULTÉ DES SCIENCES
wire diameter	1.25 mm
# of filaments	121 x 121
filament size	4.5 $\mu m$

**Regular HT: 600°C/100h + 670°/150h**

$$\epsilon_c = \epsilon_{irr} - \epsilon_{max} = 0.48 \%$$

**HIP 550°C/1h/200MPa + Regular HT**

$$\epsilon_c = \epsilon_{irr} - \epsilon_{max} = 0.65 \%$$

**With HIP treatment  $\epsilon_c$  increases by +0.17 %**



# Summary

- A superconducting wire for magnet application should consist of tiny twisted filaments in a highly conductive matrix.
- NbTi is close to an ideal superconductor if cooling is not an issue and the magnetic field remains below about 8 T
- Nb<sub>3</sub>Sn is the actual LTS high field conductor
  - It is brittle (wind and react)
  - It has to be produced in-situ
    - Imperfect, inhomogeneous microstructure
    - Stoichiometry issues
    - Voids increase the strain sensitivity
  - Strain is a crucial parameter for the design of superconducting magnets



**Rogalla & Kes**  
**100 Years of Superconductivity**  
**Chapter 3 Section 7**  
**Chapter 11 Section 2 (Nb-Ti)**  
**Section 3 (Nb<sub>3</sub>Sn)**

**Poole, Farach, Creswick, Prozorov**  
**Superconductivity (2<sup>nd</sup> edition)**  
**Chapter 3**

**Papers cited in the slides**

