

Quark Mass Effects in Gluon Fusion at NNLO in QCD

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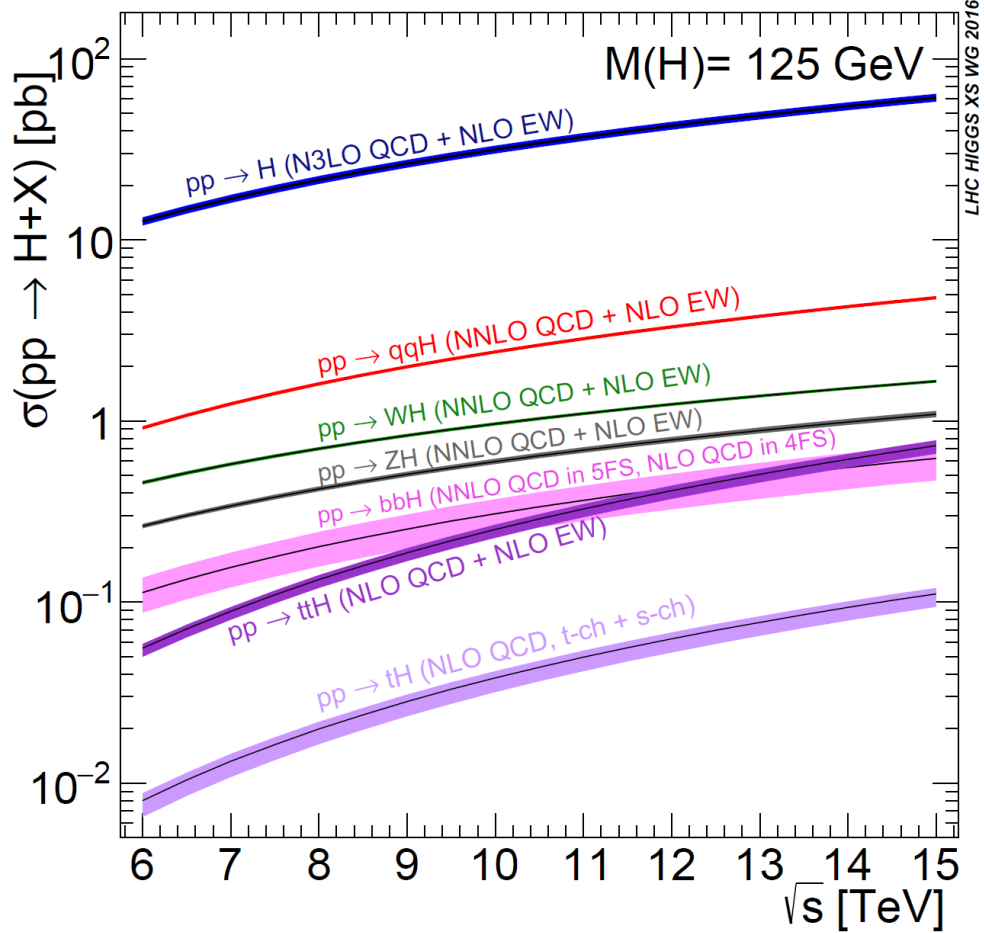
Max-Planck-Institute for Physics Munich

Based on [PRL 127 \(2021\), 162002](#)

Motivation

- Gluon fusion is the predominant Higgs-boson production mode at the LHC
- Higgs-boson plays unique role in the SM:
 - Only scalar particle
 - Only particle with Yukawa interactions to fermions

Handbook of LHC Higgs cross sections:
 4. Deciphering the nature of the Higgs sector
 Report of the LHC Higgs Cross Section Working Group `16

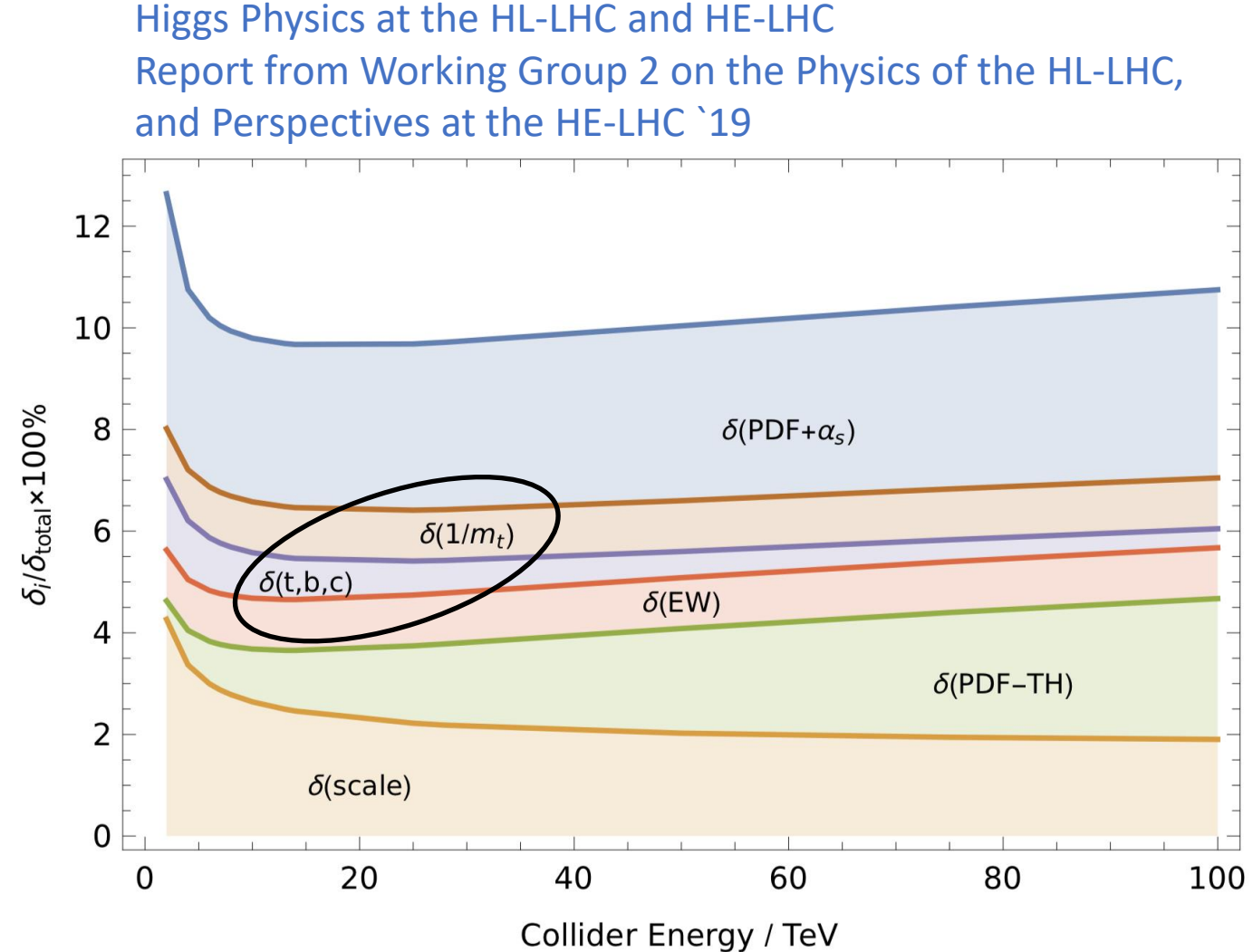


LHC @13 TeV

$$\sigma = 48.58 \text{ pb}^{+2.22 \text{ pb} (+4.56\%)}_{-3.27 \text{ pb} (-6.72\%)} \text{ (theory)} \pm 1.56 \text{ pb} (3.20\%) \text{ (PDF}+\alpha_s)$$

Theory uncertainties

- $\delta(\text{scale})$ and $\delta(\text{PDF-TH})$ due to missing higher-order terms in $\hat{\sigma}$ and PDFs [Anastasiou, et al. `15](#)
- $\delta(\text{trunc})$ has been removed [Mistlberger `18](#)
- $\delta(\text{EW})$ was addressed recently
[Bonetti, Melnikov, Tancredi `18](#)
[Anastasiou, del Duca, et al. `19](#)
[Becchetti, Bonciani, et al. `21](#)
- $\delta(t,b,c)$ and $\delta(1/m_t)$ related to quark mass effects

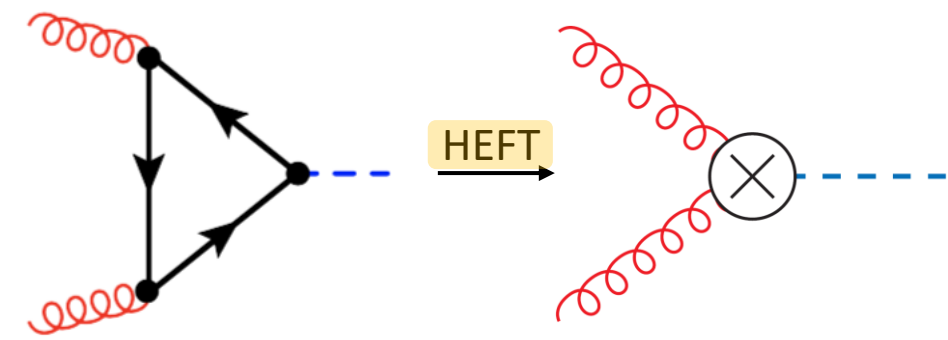


$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
+0.10 pb -1.15 pb	± 0.18 pb	± 0.56 pb	± 0.49 pb	± 0.40 pb	± 0.49 pb
+0.21% -2.37%	$\pm 0.37\%$	$\pm 1.16\%$	$\pm 1\%$	$\pm 0.83\%$	$\pm 1\%$

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Contributions to σ_{tot}

48.58 pb =	16.00 pb	(+32.9%)	(LO, rEFT)
	+ 20.84 pb	(+42.9%)	(NLO, rEFT)
	- 2.05 pb	(-4.2%)	((t, b, c), exact NLO)
	+ 9.56 pb	(+19.7%)	(NNLO, rEFT)
	+ 0.34 pb	(+0.7%)	(NNLO, $1/m_t$)
	+ 2.40 pb	(+4.9%)	(EW, QCD-EW)
	+ 1.49 pb	(+3.1%)	(N ³ LO, rEFT)



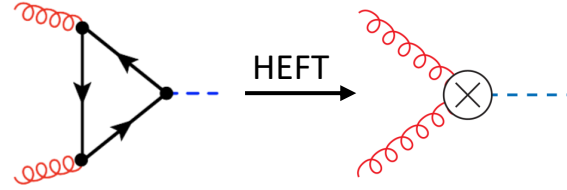
"Born-improved" total cross section:

$$\sigma_{\text{HEFT}}^{\text{HO}} = \left(\frac{\sigma^{\text{HO}}}{\sigma^{\text{LO}}} \right)_{M_t \rightarrow \infty} \sigma^{\text{LO}}$$

- Gluon-fusion is induced by quark loops
 - NLO result available for arbitrary quark masses [Graudenz, Spira, Zerwas '93](#)
 - Radiative corrections beyond NLO restricted to top-loop induced terms [Anastasiou, Melnikov '02](#)
[Harlander, Kilgore '02](#)
[Ravindran, Smith, van Neerven '03](#)
[Marzani, Ball, Del Duca, et al. '08](#)
[Harlander, Mantler, Marzani, et al. '09](#)
[Pak, Rogal, Steinhauser '09](#)
- Dominant effect of top-loop induced terms can be accounted for in HEFT approximation

HEFT

- Introduce effective Higgs-gluon vertex
 → reduce number of **loops** by one
 → reduce number of **scales** by one



- Very good agreement with exact result at NLO

→ Remarkable, because

- M_t being the largest scale is invalid over large range of $\sqrt{\hat{s}}$
- $M_t \rightarrow \infty$ is applied to more than 50% of total cross section
- HEFT fails to capture top-mass effects for partonic quark channels

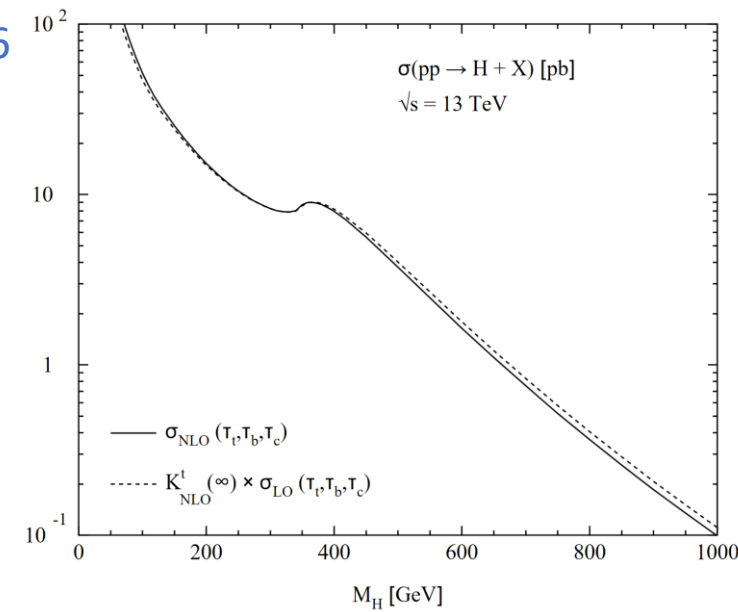
- Qualitative explanation:

- Suppression of large- \hat{s} region by PDFs
- Dominance of the soft region

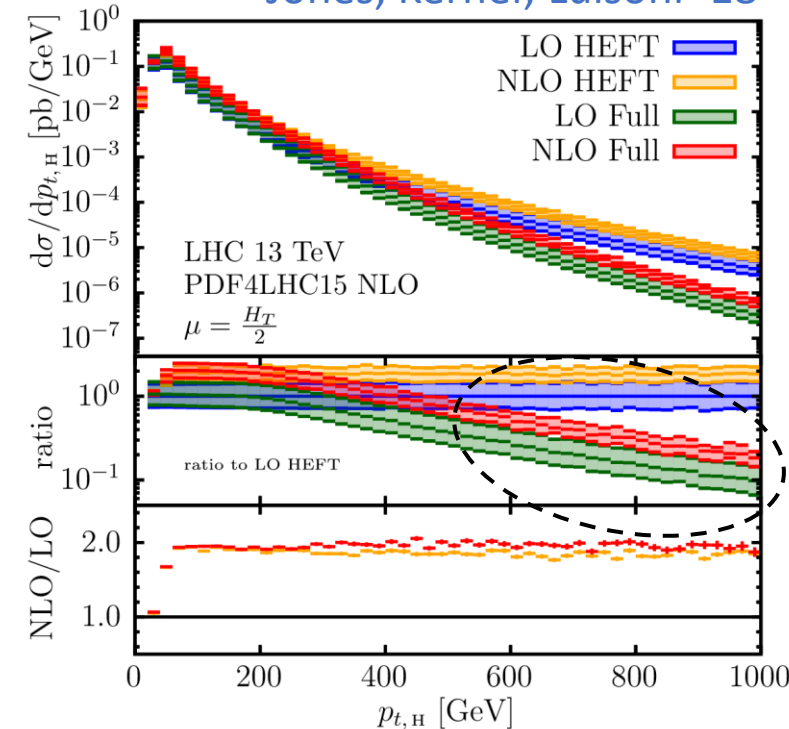
- Only estimate of top-mass effects beyond HEFT at NNLO based on combination of $1/M_t$ -expansion with leading terms in large- \hat{s} limit

→ Eliminate this uncertainty with exact calculation of top-quark mass effects

Spira '16



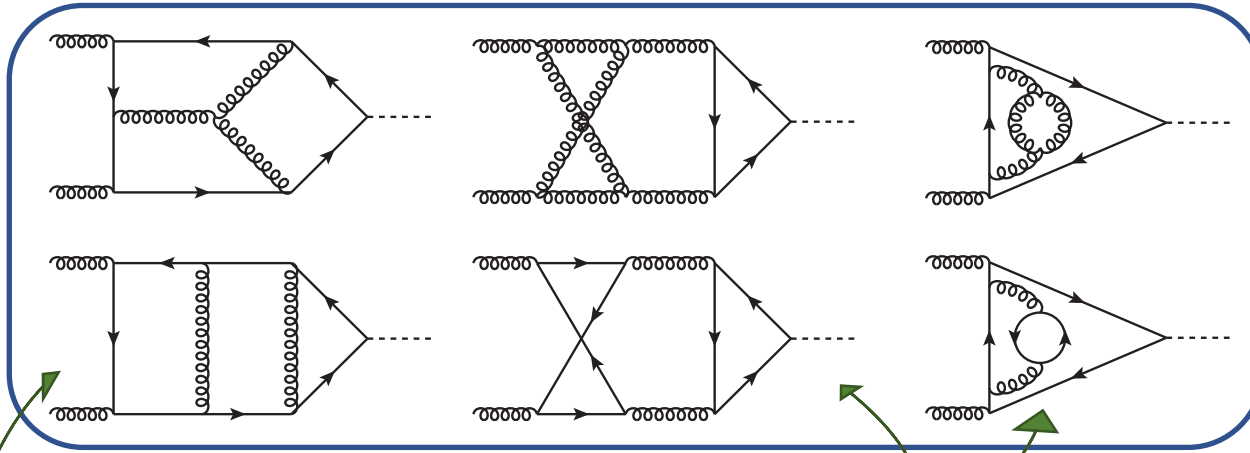
Jones, Kerner, Luisoni '18



(also Bonciani, Del Duca, Frellesvig, et al. '22)

Ingredients for gluon fusion at NNLO

Ingredients – Double Virtual



Light-fermion contribution (analytically):

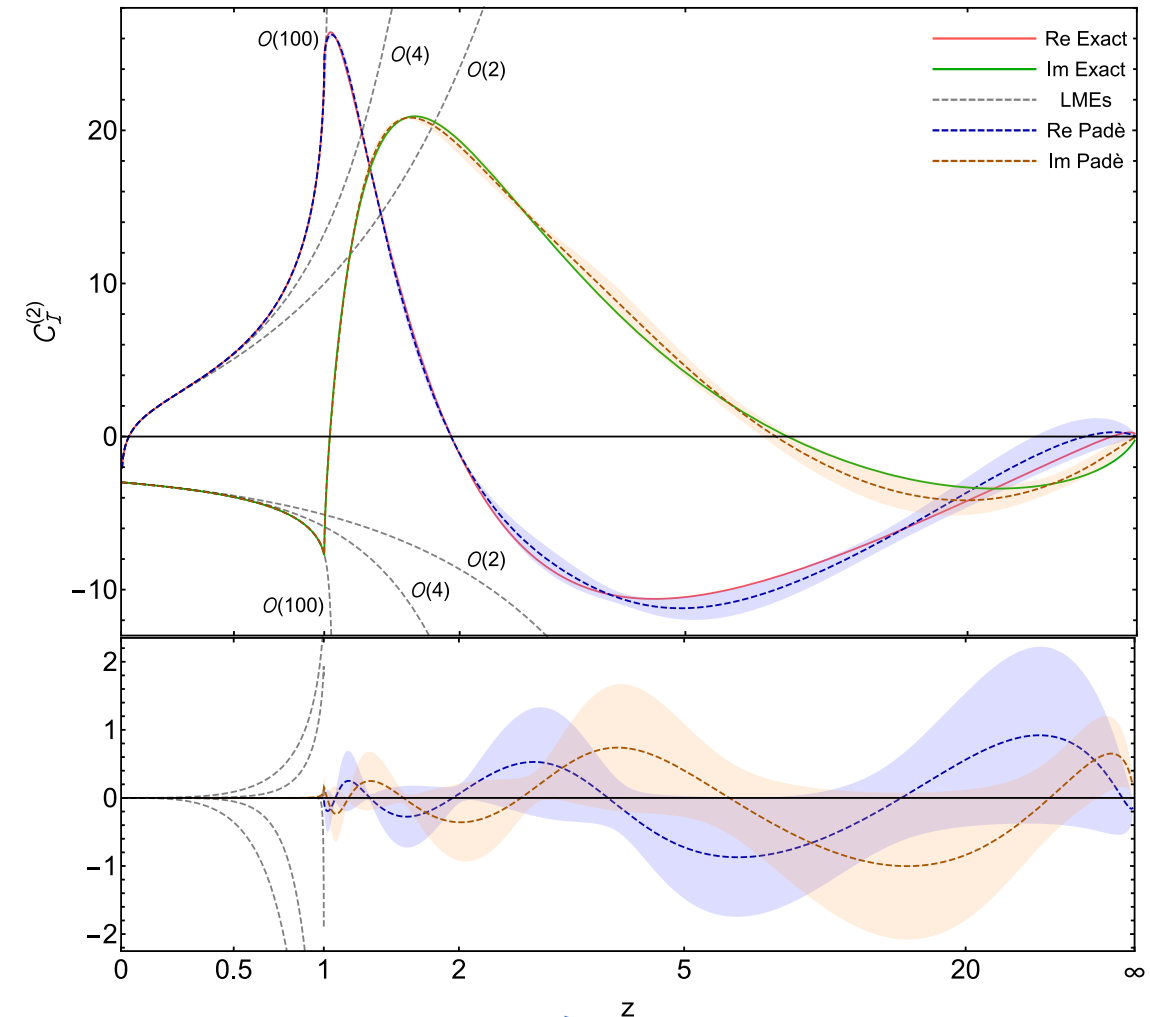
Harlander, Prusa, Usovitsch `19

Leading color contribution (analytically):

Prusa, Usovitsch `20

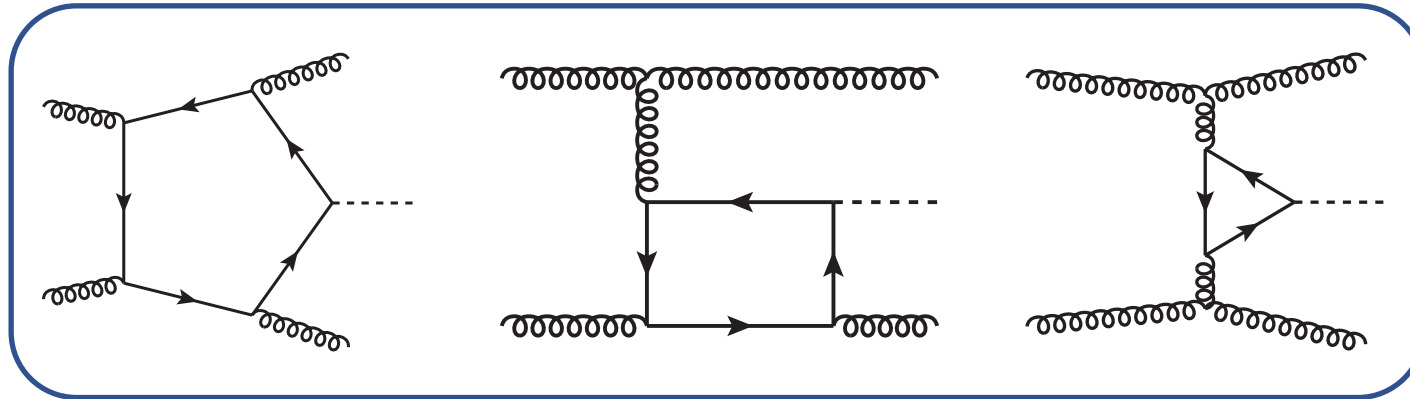
Padé approximation: Davies, Gröber, Maier et al. `19

Numerically exact: Czakon, MN `20



$$z = m_H^2 / 4m_q^2$$

Ingredients – Double Real



+ quark channels with possibly different quark flavors

Evaluation:

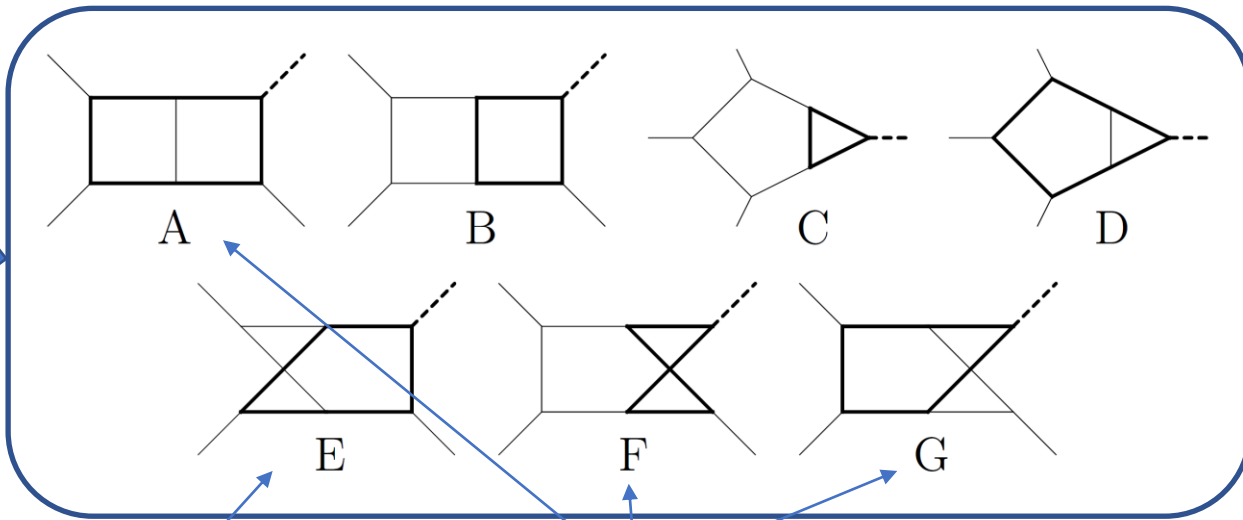
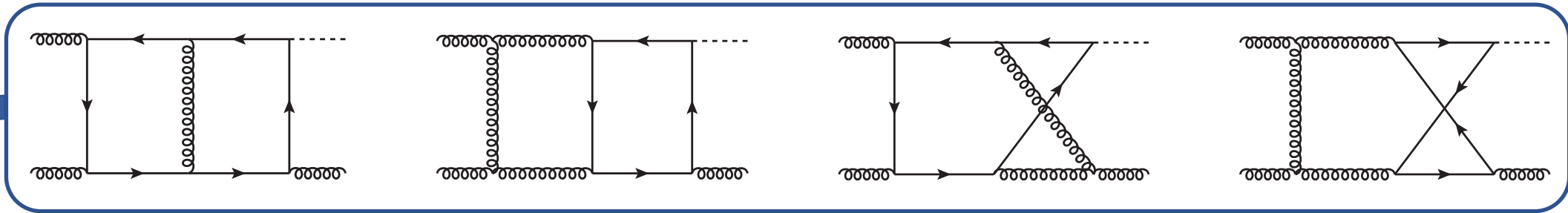
Analytically: [Del Duca, Kilgore, Oleari, et al. '01](#)

OpenLoops 2: [Buccioni, Lang, Lindert, et al. '19](#)



Analytically (more compact): [Budge, Campbell, De Laurentis, et al. '20](#)

Ingredients – Real-Virtual



vanishing color factor

Elliptic sector

A,B,C,D: [Bonciani, Del Duca, Frellesvig, et al. `16](#)

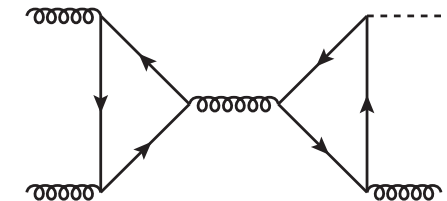
F: [Bonciani, Del Duca, Frellesvig, et al. `19](#)

G: [Frellesvig, Hidding, Maestri, et al. `19](#)

→ leading to H+jet at NLO with top and bottom massive:

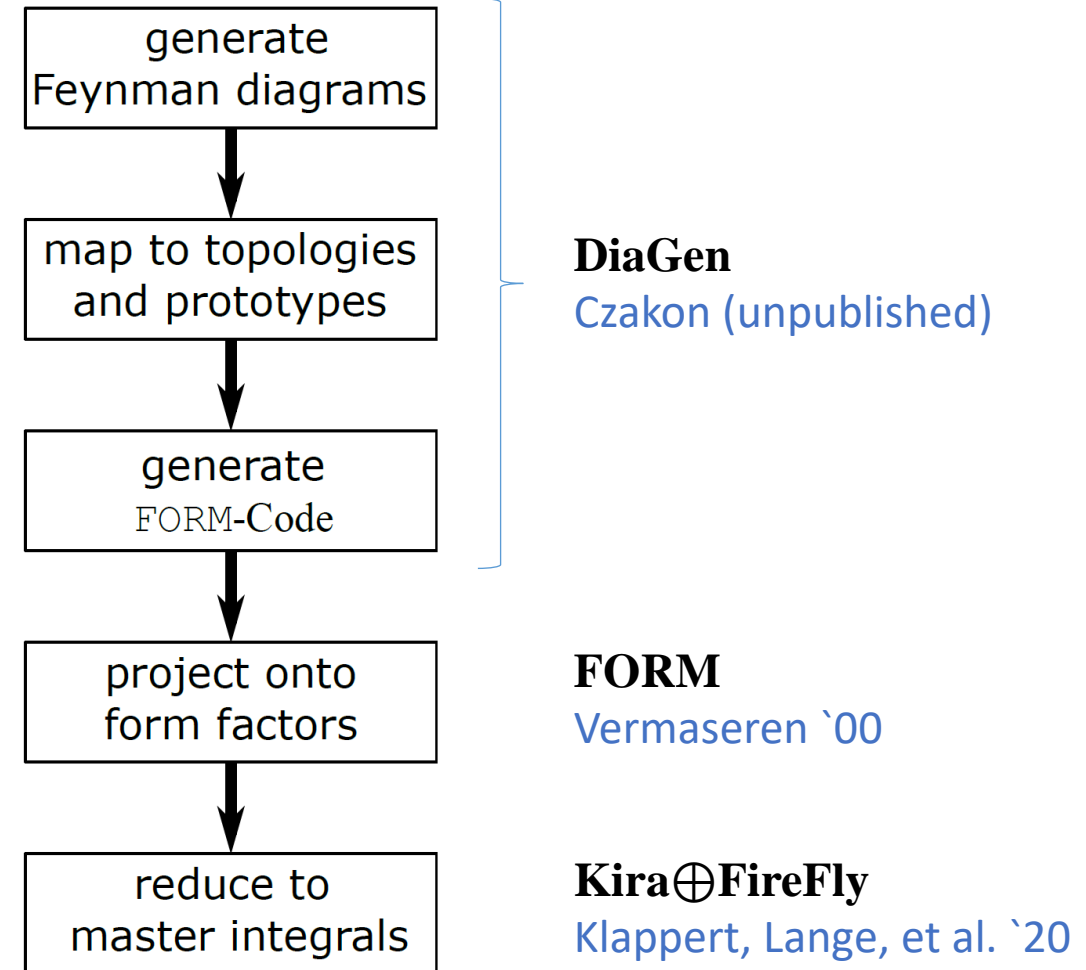
[Bonciani, Del Duca, Frellesvig, et al. `22](#)

Contributions with two closed fermion chains are always factorizable:



Workflow of the computation

- Get rid of tensor/colour structure to end up with a linear combination of scalar integrals with rational function coefficients in front
- Reduce the scalar integrals to a linearly independent set of master integrals (MI) (447 master integrals for $gg \rightarrow Hg$)
- Reduction is highly non-trivial since rational coefficients depend on 5 variables!
→ Use finite fields to reconstruct symbolic coefficients from numerical probes of the system of equations



Computation of the MIs

Parametrization

- Variables: \hat{s} , \hat{t} , \hat{u} , m_H^2 , m_t^2
- Introduce dimensionless variables and fix ratio m_t^2/m_H^2
 - z parametrizes **soft** limit
 - λ parametrizes **collinear** limit

$$z = 1 - m_H^2/\hat{s}$$

$$\lambda = \hat{t}/(\hat{t} + \hat{u})$$

$$m_t^2/m_H^2 = 23/12$$

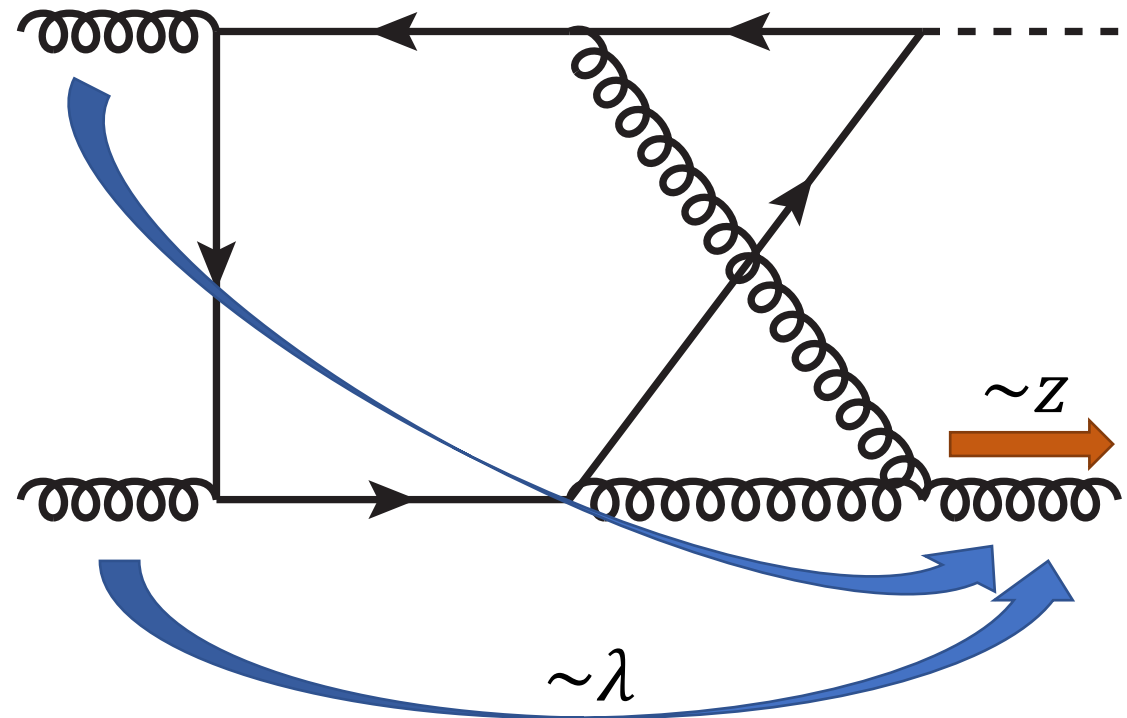
- Range of parameters:
- $\lambda \in (0,1)$
 - $z \in (0,1)$

$$\hat{t}/\hat{s} = z \lambda$$

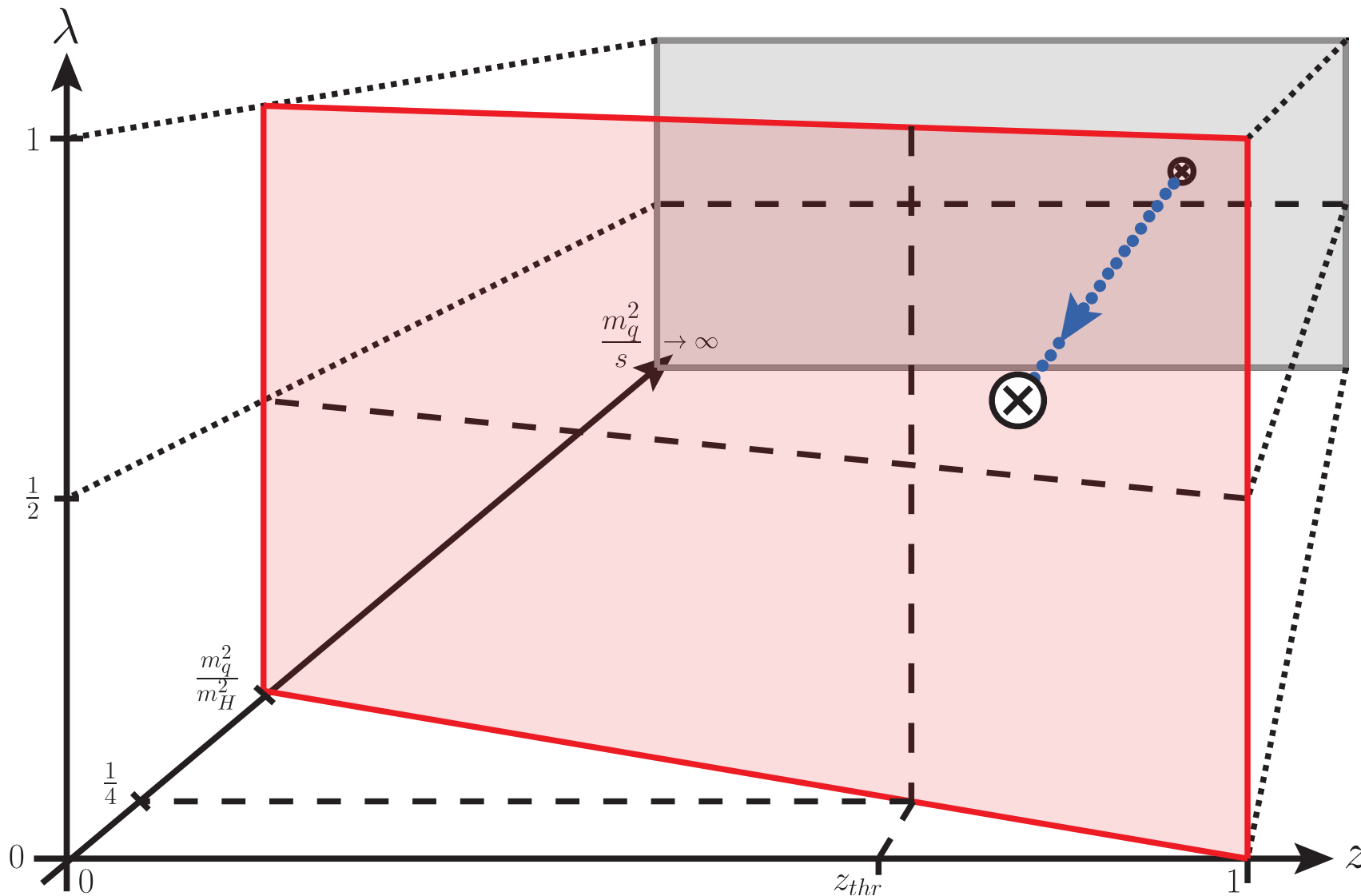
$$\hat{u}/\hat{s} = z (1-\lambda)$$

$$z = 1 - m_H^2/\hat{s}$$

$$\lambda = \hat{t}/(\hat{t} + \hat{u})$$



Evolution of differential equations



$$z = 1 - m_H^2 / \hat{s}$$

$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$m_t^2 / m_H^2 = 23/12$$

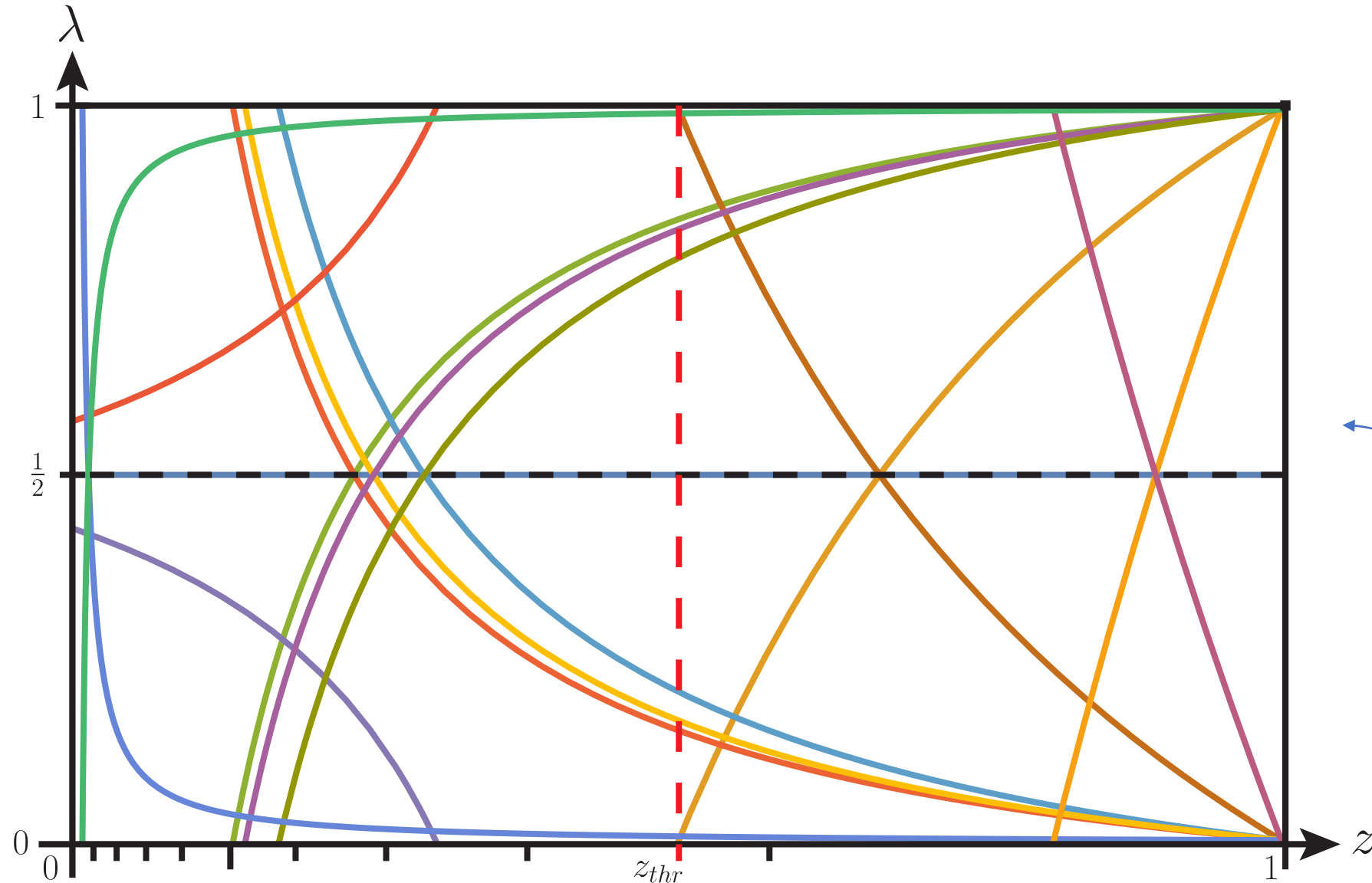
Range of parameters:

- $\lambda \in (0, 1)$
- $z \in (0, 1)$

$$z_{thr} = 1 - \frac{m_H^2}{4m_q^2}$$

For $m_t^2 / m_H^2 = 23/12$:
 $z_{thr} = 20/23 \approx 0.87$

Evolution in the (z, λ) -plane



$$z = 1 - m_H^2 / \hat{s}$$

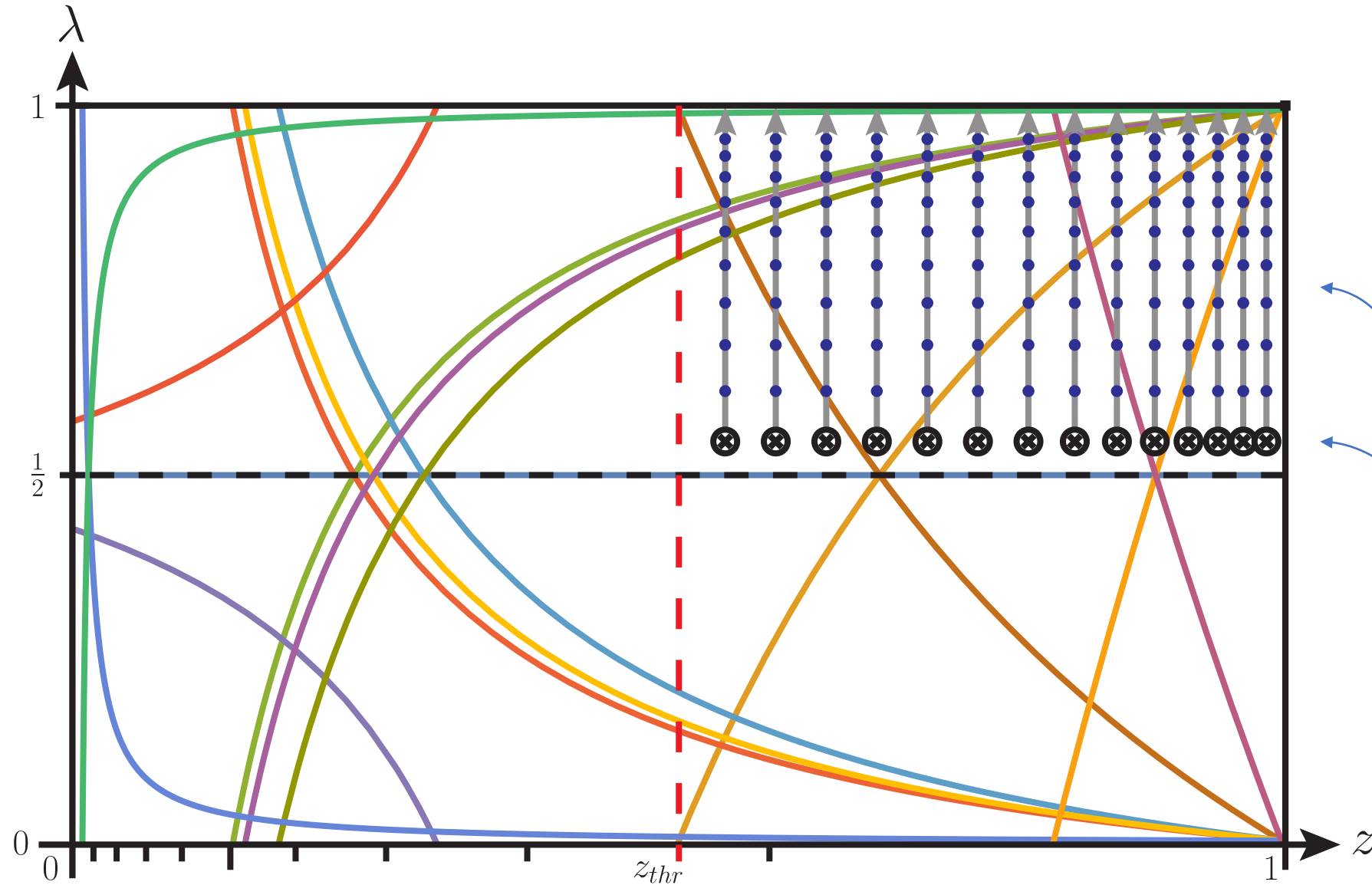
$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$m_t^2 / m_H^2 = 23/12$$

- Range of parameters:
- $\lambda \in (0, 1)$
 - $z \in (0, 1)$

Poles of differential equations in λ

Evolution in the (z, λ) -plane



$$z = 1 - m_H^2 / \hat{s}$$

$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$m_t^2 / m_H^2 = 23/12$$

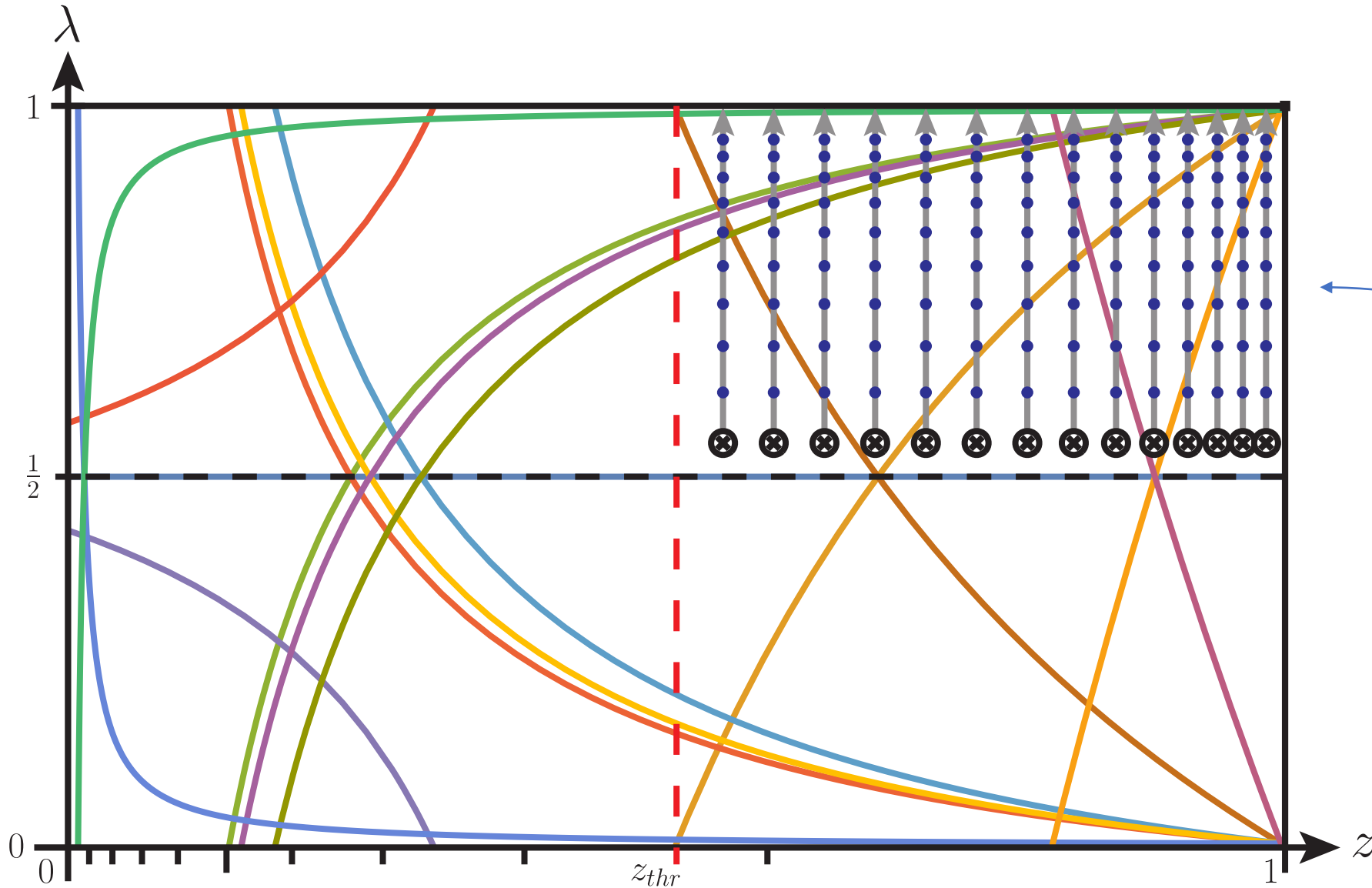
Range of parameters:

- $\lambda \in (0, 1)$
- $z \in (0, 1)$

Collect numerical samples for MI along straight integration contours

Boundaries for numerical integration in the mass

Evolution in the (z, λ) -plane

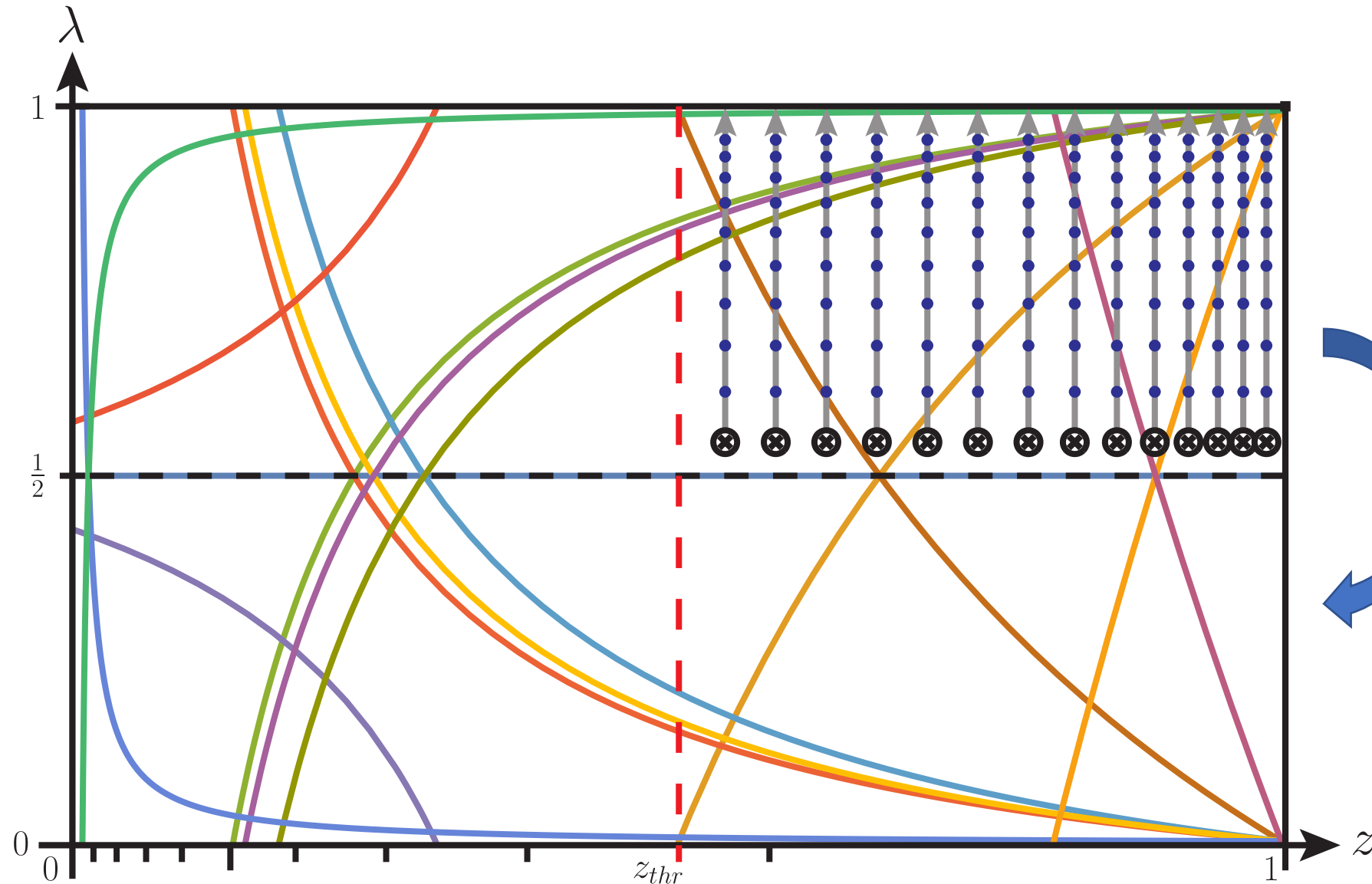


$$z = 1 - m_H^2 / \hat{s}$$
$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$
$$m_t^2 / m_H^2 = 23/12$$

- Range of parameters:
- $\lambda \in (0, 1)$
 - $z \in (0, 1)$

- 302 integration contours at different z
- Collected more than 1.5×10^6 samples

Evolution in the (z, λ) -plane



$$z = 1 - m_H^2 / \hat{s}$$

$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

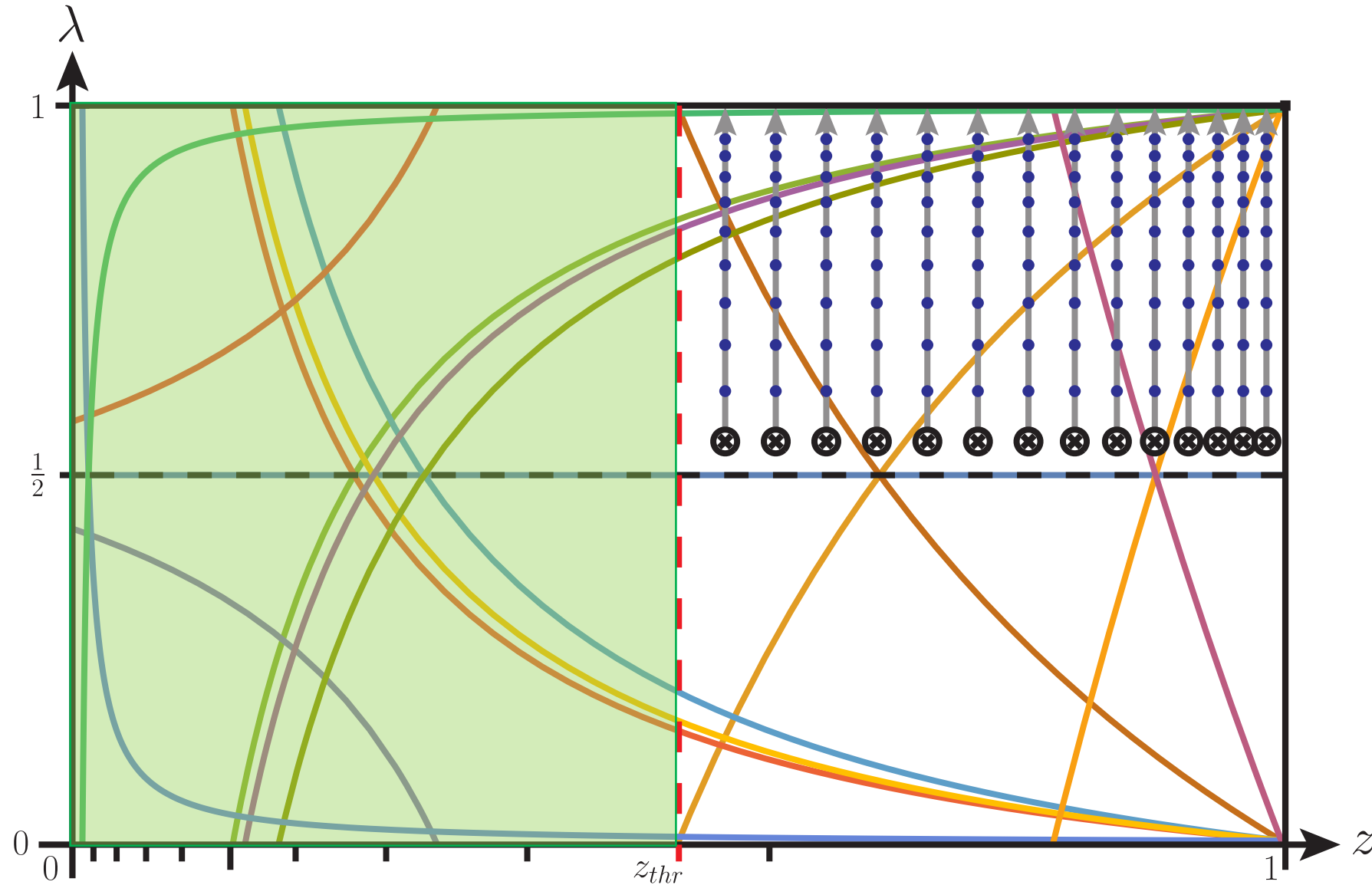
$$m_t^2 / m_H^2 = 23/12$$

Range of parameters:

- $\lambda \in (0, 1)$
- $z \in (0, 1)$

Exploit symmetry of the problem

Evolution in the (z, λ) -plane



$$z = 1 - m_H^2 / \hat{s}$$

$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$m_t^2 / m_H^2 = 23/12$$

- Range of parameters:
- $\lambda \in (0, 1)$
 - $z \in (0, 1)$

Region below threshold covered by LME

LME
 $\mathcal{O}((1/m_q^2)^{40})$

Subtraction of IR limits

Subtraction for $gg \rightarrow gH$

$$z = 1 - m_H^2/\hat{s}$$

$$\lambda = \hat{t}/(\hat{t} + \hat{u})$$

$$m_t^2/m_H^2 = 23/12$$

- Interested in finite results
 - Soft and collinear divergences
 - Amplitudes have to be regulated

Higgs-gluon form factor in HEFT is exact at leading order $|F_{\text{exact}}^{(1)}|^2 = |F_{\text{HEFT}}^{(1)}|^2$

- Directly evaluate difference between HEFT and exact result:

$$\text{NLO: } \langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(1)} \rangle_{\text{regulated}} \equiv \langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(1)} \rangle - \left[\langle M_{\text{HEFT}}^{(1)} | M_{\text{HEFT}}^{(1)} \rangle \right]$$

$$\text{NNLO: } \langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(2)} \rangle_{\text{regulated}} \equiv \langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(2)} \rangle - \left[\langle M_{\text{HEFT}}^{(1)} | M_{\text{HEFT}}^{(2)} \rangle + \frac{8\pi\alpha_s}{\hat{t}} \left\langle P_{gg}^{(0)} \left(\frac{\hat{s}}{\hat{s} + \hat{u}} \right) \right\rangle \langle F^{(1)} | (F_{\text{exact}}^{(2)} - F_{\text{HEFT}}^{(2)}) \rangle \right]$$

Free of divergences!

- Better numerical stability!
- Appearance of divergences delayed!
 - Simple counterterms!

Splitting-function regulates residual soft and collinear divergences

One of the reasons for the smallness of top-quark mass effects beyond HEFT

Subtraction for $gg \rightarrow gH$

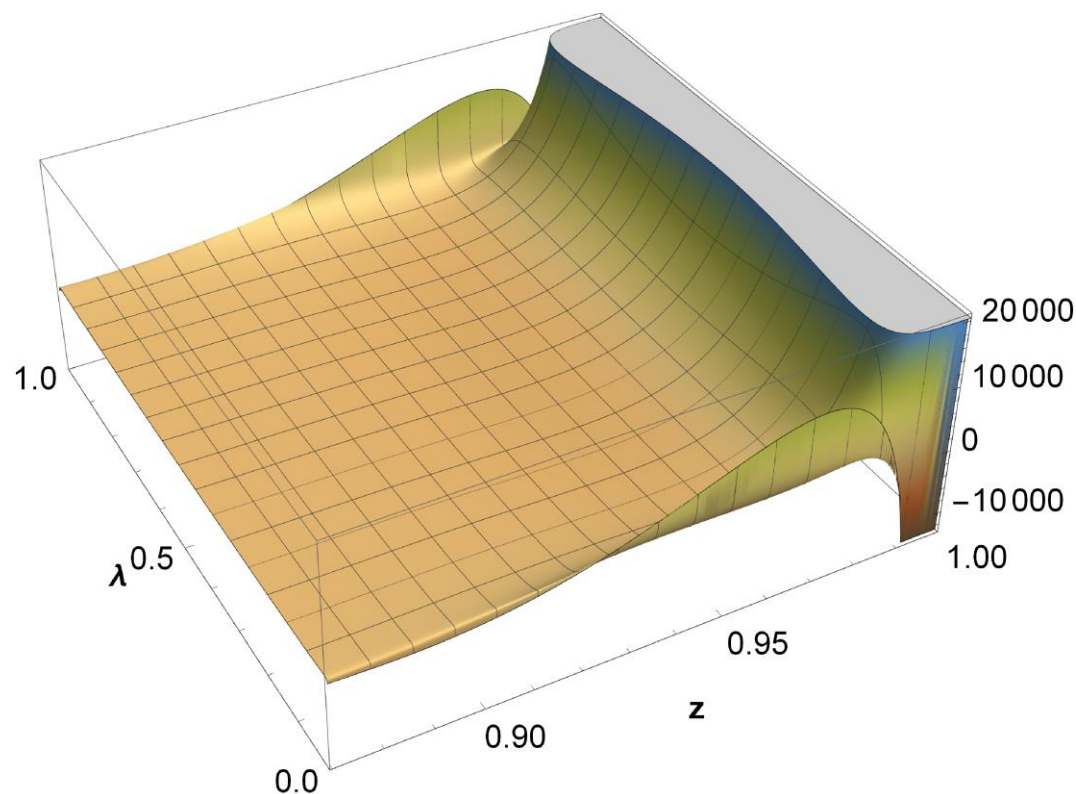
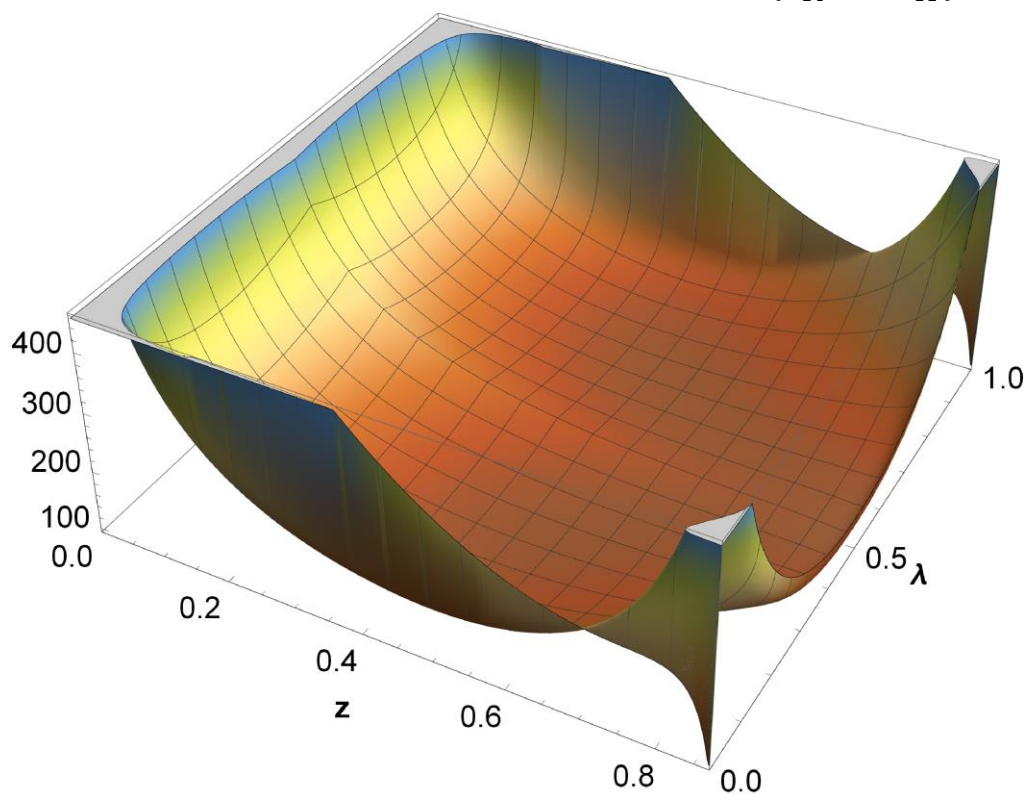
$$z = 1 - m_H^2 / \hat{s}$$

$$\lambda = \hat{t} / (\hat{t} + \hat{u})$$

$$m_t^2 / m_H^2 = 23/12$$

$$\langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(2)} \rangle \Big|_{\text{regulated}} \equiv \langle M_{\text{exact}}^{(1)} | M_{\text{exact}}^{(2)} \rangle - \left[\langle M_{\text{HEFT}}^{(1)} | M_{\text{HEFT}}^{(2)} \rangle + \frac{8\pi\alpha_s}{\hat{t}} \left\langle P_{gg}^{(0)} \left(\frac{\hat{s}}{\hat{s} + \hat{u}} \right) \right\rangle \langle F^{(1)} | (F_{\text{exact}}^{(2)} - F_{\text{HEFT}}^{(2)}) \rangle \right]$$

- Real part of the regulated quantity at $\mu_R = m_H/2$:



- Integrate in λ and convolute with PDFs to obtain contribution to σ_{tot}
- Subtraction term and other contributions are computed with Monte Carlo methods using **Stripper** [Czakon \(unpublished\)](#) 21

Results

- Effects of a finite top-quark mass on the total hadronic Higgs-boson production cross section for the LHC
 - PDF set: NNPDF31_nnlo_as_0118
 - $\mu_R = \mu_F = m_H/2$
 - $M_H = 125 \text{ GeV} \Rightarrow M_t \approx 173.055 \text{ GeV}$

channel	$\sigma_{\text{HEFT}}^{\text{NNLO}} [\text{pb}]$ $\mathcal{O}(\alpha_s^2) + \mathcal{O}(\alpha_s^3) + \mathcal{O}(\alpha_s^4)$	$(\sigma_{\text{exact}}^{\text{NNLO}} - \sigma_{\text{HEFT}}^{\text{NNLO}}) [\text{pb}]$ $\mathcal{O}(\alpha_s^3) \quad \mathcal{O}(\alpha_s^4)$		$(\sigma_{\text{exact}}^{\text{NNLO}} / \sigma_{\text{HEFT}}^{\text{NNLO}} - 1) [\%]$
$\sqrt{s} = 8 \text{ TeV}$				
<i>gg</i>	7.39 + 8.58 + 3.88	+0.0353	+0.0879 ± 0.0005	+0.62
<i>qg</i>	0.55 + 0.26	-0.1397	-0.0021 ± 0.0005	-18
<i>qq</i>	0.01 + 0.04	+0.0171	-0.0191 ± 0.0002	-4
total	7.39 + 9.15 + 4.18	-0.0873	+0.0667 ± 0.0007	-0.10
$\sqrt{s} = 13 \text{ TeV}$				
<i>gg</i>	16.30 + 19.64 + 8.76	+0.0345	+0.2431 ± 0.0020	+0.62
<i>qg</i>	1.49 + 0.84	-0.3696	-0.0115 ± 0.0010	-16
<i>qq</i>	0.02 + 0.10	+0.0322	-0.0501 ± 0.0006	-15
total	16.30 + 21.15 + 9.70	-0.3029	+0.1815 ± 0.0023	-0.26

Results

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channel	$\sigma_{\text{HEFT}}^{\text{NNLO}} [\text{pb}]$ $\mathcal{O}(\alpha_s^2) + \mathcal{O}(\alpha_s^3) + \mathcal{O}(\alpha_s^4)$	$(\sigma_{\text{exact}}^{\text{NNLO}} - \sigma_{\text{HEFT}}^{\text{NNLO}}) [\text{pb}]$ $\mathcal{O}(\alpha_s^3)$ $\mathcal{O}(\alpha_s^4)$		$(\sigma_{\text{exact}}^{\text{NNLO}} / \sigma_{\text{HEFT}}^{\text{NNLO}} - 1) [\%]$
$\sqrt{s} = 8 \text{ TeV}$				
<i>gg</i>	7.39 + 8.58 + 3.88	+0.0353	+0.0879 ± 0.0005	+0.62 (= 0.18 + 0.44)
<i>qg</i>	0.55 + 0.26	-0.1397	-0.0021 ± 0.0005	-18 (= -17.5 - 0.5)
<i>qq</i>	0.01 + 0.04	+0.0171	-0.0191 ± 0.0002	-4 (= +34 - 38)
total	7.39 + 9.15 + 4.18	-0.0873	+0.0667 ± 0.0007	-0.10 (= -0.42 + 0.32)
$\sqrt{s} = 13 \text{ TeV}$				
<i>gg</i>	16.30 + 19.64 + 8.76	+0.0345	+0.2431 ± 0.0020	+0.62 (= 0.08 + 0.54)
<i>qg</i>	1.49 + 0.84	-0.3696	-0.0115 ± 0.0010	-16 (= -15.5 - 0.5)
<i>qq</i>	0.02 + 0.10	+0.0322	-0.0501 ± 0.0006	-15 (= +27 - 42)
total	16.30 + 21.15 + 9.70	-0.3029	+0.1815 ± 0.0023	-0.26 (= -0.64 + 0.38)

Comparison with HEFT $\oplus (1/M_t^n)$

- The impact of subleading terms in $1/M_t^2$ is determined with SusHi [Harlander, Liebler, Mantler '16](#)
 - Include terms up to $1/M_t^4$ at NLO and NNLO and match with high-energy limit
 - Total result of $1/M_t$ approximation very close to exact result

channel	$\sigma_{\text{HEFT}}^{\text{NNLO}}$ [pb] $\mathcal{O}(\alpha_s^2) + \mathcal{O}(\alpha_s^3) + \mathcal{O}(\alpha_s^4)$	$(\sigma_{\text{subl.}}^{\text{NNLO}} - \sigma_{\text{HEFT}}^{\text{NNLO}})$ [pb] $\mathcal{O}(\alpha_s^3)$ $\mathcal{O}(\alpha_s^4)$		$(\sigma_{\text{subl.}}^{\text{NNLO}} / \sigma_{\text{HEFT}}^{\text{NNLO}} - 1)$ [%]
$\sqrt{s} = 8 \text{ TeV}$				
<i>gg</i>	7.39 + 8.58 + 3.88	-0.0104	+0.1088	+0.50
<i>qg</i>	0.55 + 0.26	-0.1265	+0.0142	-14
<i>qq</i>	0.01 + 0.04	+0.0025	-0.0076	-10
total	7.39 + 9.15 + 4.18	-0.1344	+0.1153	-0.09
$\sqrt{s} = 13 \text{ TeV}$				
<i>gg</i>	16.30 + 19.64 + 8.76	-0.1145	+0.3127	+0.44
<i>qg</i>	1.49 + 0.84	-0.3348	+0.0482	-12
<i>qq</i>	0.02 + 0.10	+0.0036	-0.0246	-17
total	16.30 + 21.15 + 9.70	-0.4457	+0.3363	-0.23

Summary

- ✓ The hadronic Higgs production cross section including the full top-quark mass dependence was computed!

- ✓ Slight decrease relative to HEFT at NNLO
 - -0.26% at 13 TeV
 - -0.10% at 8 TeV

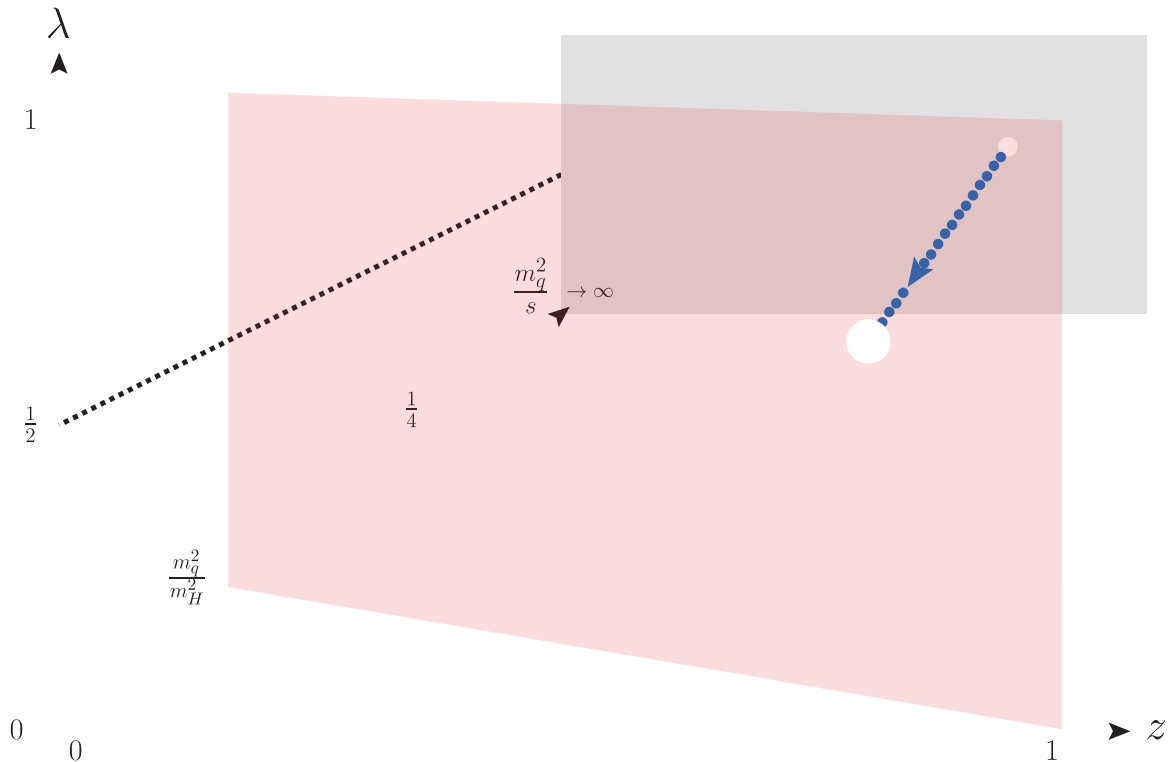
- ✓ The result confirms and eliminates the uncertainty estimate from the lack of knowledge of the exact top-quark mass effects!

- ✓ Amplitudes for $gg \rightarrow gH$ (and quark channels) have been computed numerically for a physical top quark mass by supplementing a deep LME with a dense grid of phase space samples above threshold
 - Computation can be repeated for different values of m_q^2 and m_H^2
 - Now straightforward for $m_t^2/m_H^2 \sim 23/12$

- Same techniques can be applied to compute **bottom quark mass effects...**

Summary

- Same techniques can be applied to compute **bottom quark mass effects**...
- Large hierarchy between m_b^2 and m_H^2 can lead to numerical instabilities when solving the differential equations
- Boundaries at $m_q^2 \rightarrow \infty$ not optimal



Corrections to $gg \rightarrow H$ at three loops for two different massive quark flavors unknown

