

# Invariant mass distribution of di-Higgs production at HL-LHC in the 2HDM

Kateryna Radchenko

in collaboration with Francisco Arco, Sven Heinemeyer and Margarete Mühlleitner

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# The 2HDM model

[Santos, Barroso: [arXiv:hep-ph/9701257](https://arxiv.org/abs/hep-ph/9701257)]

- **CP conserving** 2HDM with two complex doublets:  $\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{v_1 + \rho_1 + i\eta_1}{\sqrt{2}} \end{pmatrix}, \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{v_2 + \rho_2 + i\eta_2}{\sqrt{2}} \end{pmatrix}$
- **Softly broken  $\mathbb{Z}_2$  symmetry** ( $\Phi_1 \rightarrow \Phi_1; \Phi_2 \rightarrow -\Phi_2$ ) entails 4 Yukawa types (Types I and II were analyzed).

**$h$  ( $m_h = 125$  GeV),  $H$  - CP even,  $A$  - CP odd,  $H^+, H^-$**

- Potential: 
$$V_{2\text{HDM}} = m_{11}^2 (\Phi_1^\dagger \Phi_1) + m_{22}^2 (\Phi_2^\dagger \Phi_2) - m_{12}^2 (\Phi_1^\dagger \Phi_2 + \Phi_2^\dagger \Phi_1) + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} ((\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2),$$

- Free parameters:  $m_h, m_A, m_H, m_{H^\pm}, m_{12}^2, v, \cos(\beta - \alpha), \tan\beta$

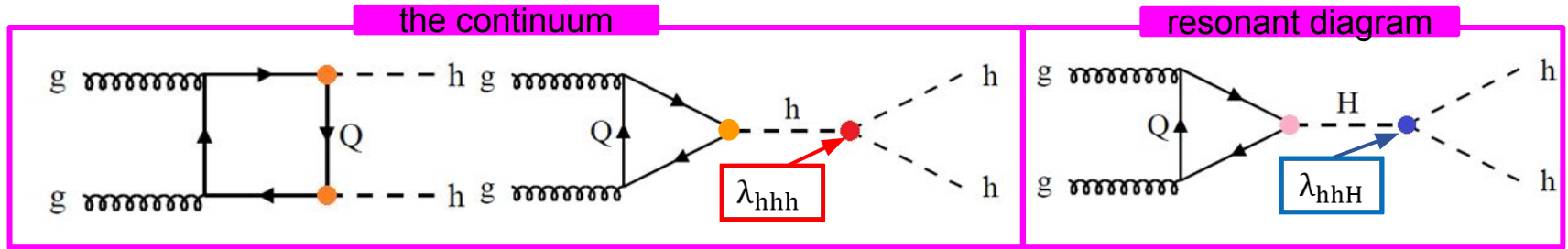
$$\begin{aligned} \tan\beta &= v_2/v_1 \\ v^2 &= v_1^2 + v_2^2 \sim (246 \text{ GeV})^2 \end{aligned}$$

- **Phenomenological implications** can originate from:
  - deviations in **couplings** to fermions and gauge bosons
  - contributions of the **heavy scalars** in the loops

# Di-Higgs production ( $gg \rightarrow hh$ )

[Plehn, Spira, Zerwas : [arXiv:hep-ph/9603205](https://arxiv.org/abs/hep-ph/9603205)]

- Triple Higgs couplings can be accessed through Higgs pair production
- The dominant process at a hadron collider is gluon fusion involving a quark loop



Diagrams that exist in the SM:  
They have a negative interference

$$\sigma_{\text{SM}} \sim 38 \text{ fb at NLO}$$

Diagrams that are sensitive  
to triple Higgs couplings

- We will study the **invariant mass distribution** of two 125 GeV in the final state
- All calculations were done using a modified version of the code **HPAIR**

[Abouabid, Arhrib, Azevedo, El Falaki, Ferreira, Mühlleitner, Santos: [arXiv:hep-ph/2112.12515](https://arxiv.org/abs/hep-ph/2112.12515)]

# Triple Higgs Couplings

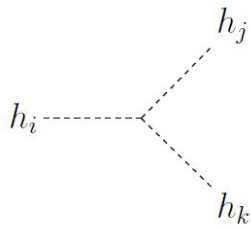
[ATLAS-CONF-2022- 050]

- Can have **large deviations** from SM predictions in BSM:

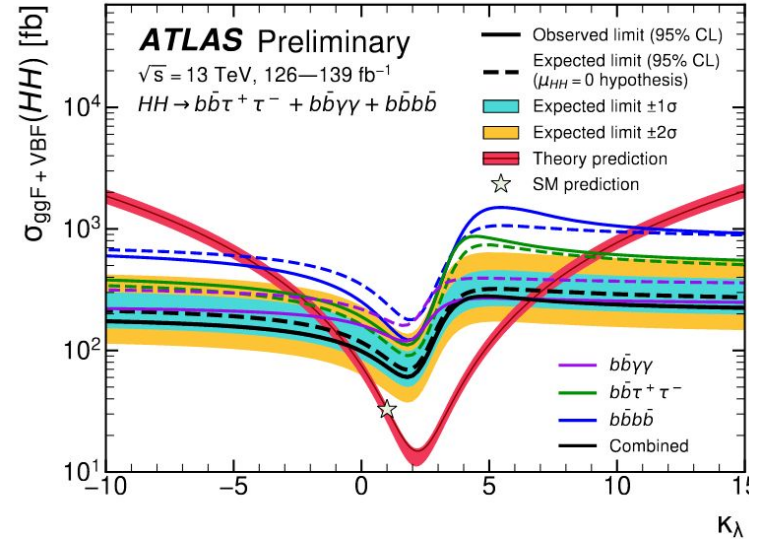
$$\kappa_\lambda = \lambda_{hhh} / \lambda_{hhh}^{SM}$$

$$\lambda_{hhh}^{SM} = \frac{m_h^2}{2v^2} \cong 0.129$$

$[-0.4 < \kappa_\lambda < 6.3]$  (95% CL at LHC Run II)

- Notation:  $h_i$    $= -i v n! \lambda_{h_i h_j h_k}$   
 $n = \text{number of identical Higgses}$

Large  $\lambda$ 's were found: far from the alignment limit and/or for large masses.



## Type I

$$\kappa_\lambda = [-0.5, 1.3]$$

$$\lambda_{hhh} = [-1.7, 1.6]$$

[Arco, Heinemeyer, Herrero: [arXiv2003.12684](https://arxiv.org/abs/2003.12684)]

# Invariant mass distribution: effects of deviations in $\kappa_\lambda$

BP: Type I,  $\cos(\beta - \alpha) = 0.1$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$ ,  $m_H = m_A = m_{H^\pm}$

Prediction for SM couplings:

$$\kappa_\lambda = 1, \lambda_{hhH} = 0$$

Prediction for BSM couplings:

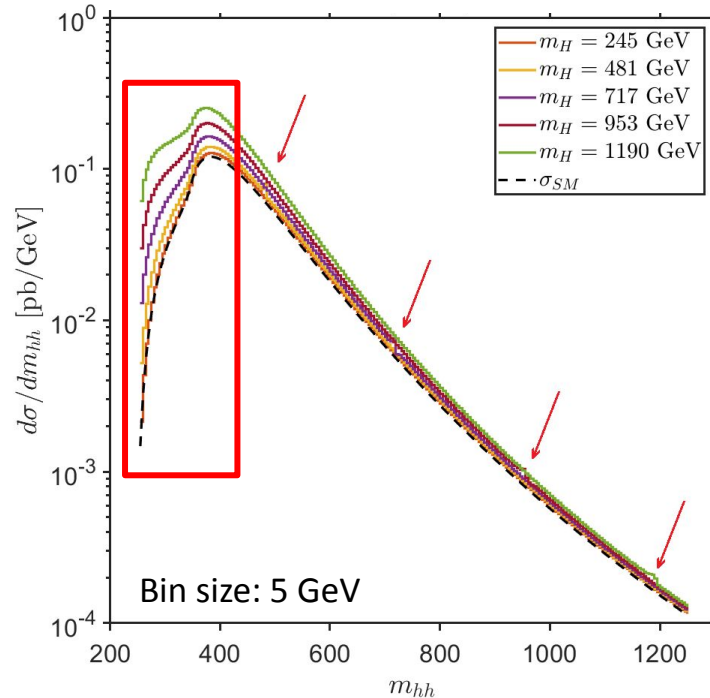
$$\kappa_\lambda = 0.97, \lambda_{hhH} = 0.05$$

$$\kappa_\lambda = 0.85, \lambda_{hhH} = 0.19$$

$$\kappa_\lambda = 0.67, \lambda_{hhH} = 0.42$$

$$\kappa_\lambda = 0.41, \lambda_{hhH} = 0.74$$

$$\kappa_\lambda = 0.08, \lambda_{hhH} = 1.15$$



- Larger sensitivity to  $\kappa_\lambda$  in the low  $m_{hh}$  region

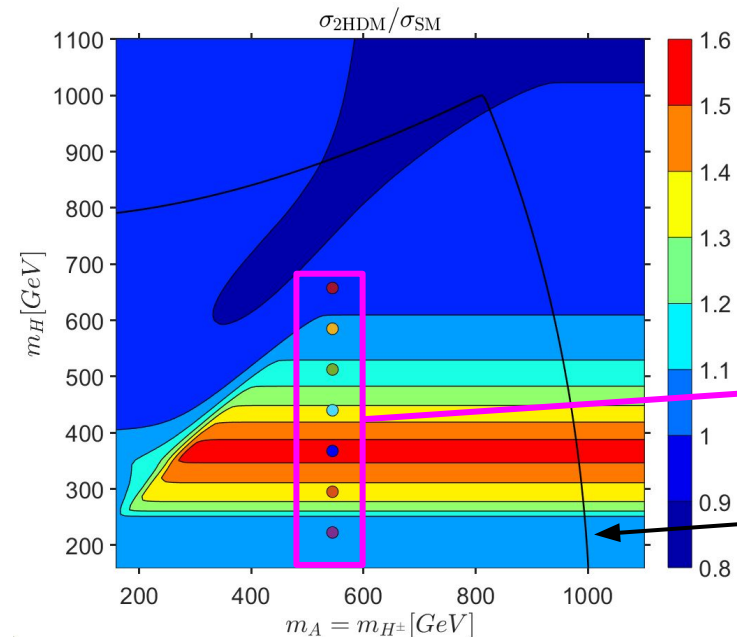
- Resonant contribution very suppressed due to very small top Yukawa  $\xi_H^t \sim 10^{-4}$

$$\xi_H^t = \cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan(\beta)$$

# Effect of the mass of the heavy Higgs

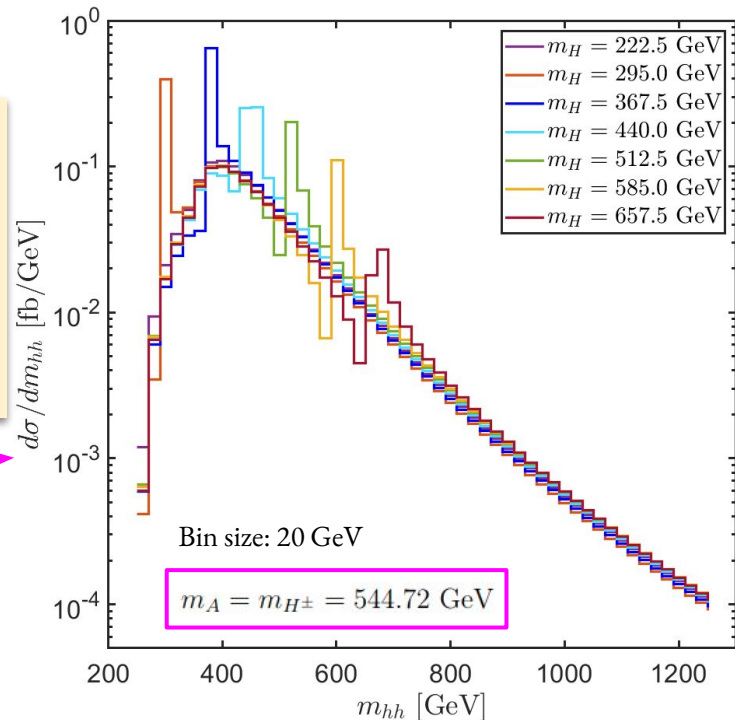
- We vary the mass of the heavy Higgs boson leaving the rest of the parameters of the model fixed.

BP: Type I,  $\cos(\beta - \alpha) = 0.2$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$



Enhancement in the total cross section is resonance dominated. **Location** of the resonance is related to the mass of **H**

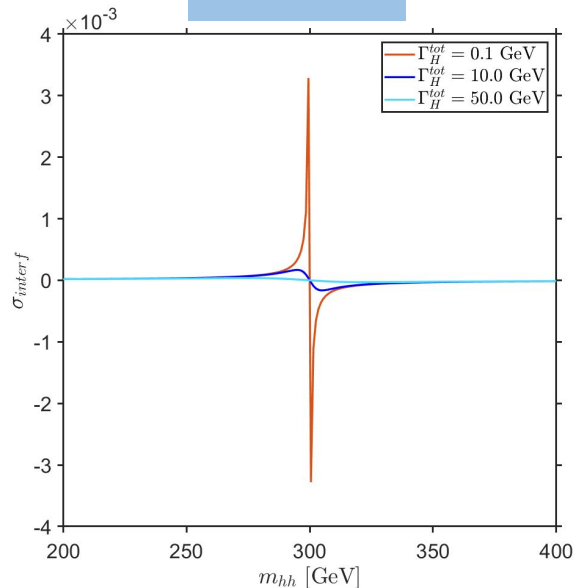
Allowed region inside the black contour.



# Effect of the total decay width

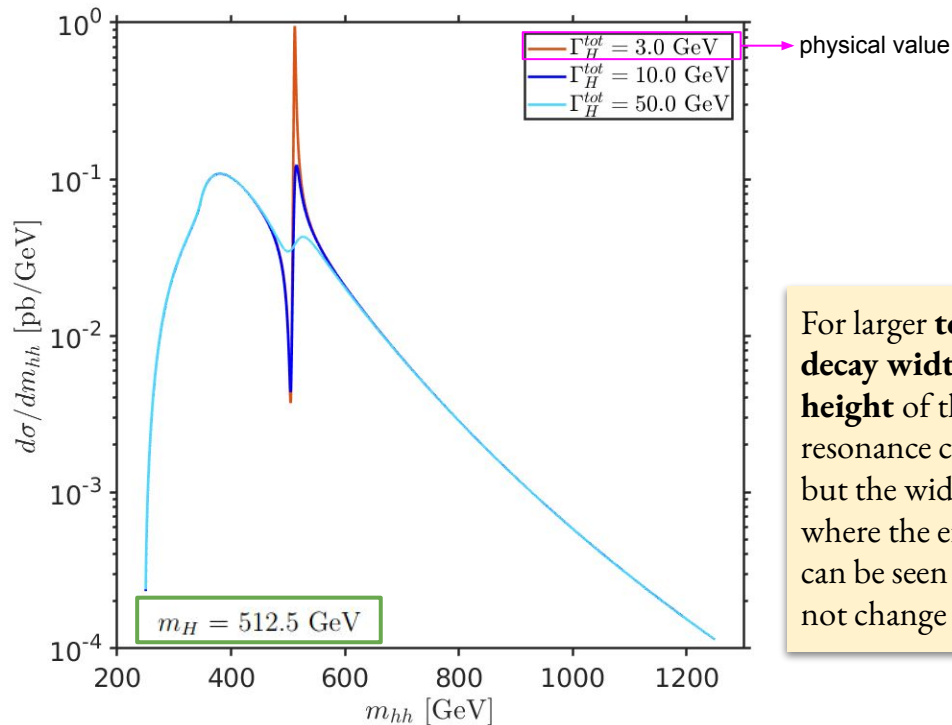
$$\frac{1}{Q^2 - M_{h/H}^2 + i\Gamma_{h/H}M_{h/H}}$$

Toy model



$$\sigma_{\text{interf}} \propto \frac{Q^2 - m_H^2}{(Q^2 - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

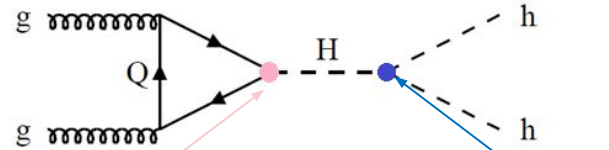
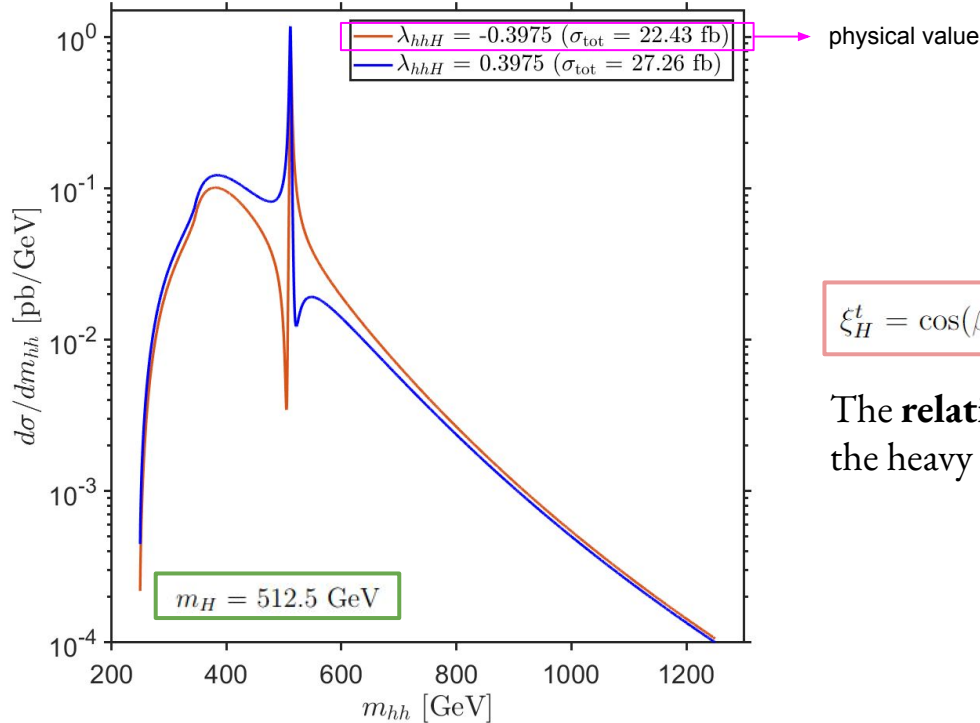
- For the green point of the previous benchmark plane we artificially change the total decay width of the heavy Higgs H:



For larger **total decay widths** the **height** of the resonance changes but the width where the effect can be seen does not change

# Effect of the couplings

- What is the effect of the couplings involved in the resonant diagram on the invariant mass distributions ?



$$\xi_H^t = \cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan(\beta) = 0.104$$

$$\lambda_{hhH}$$

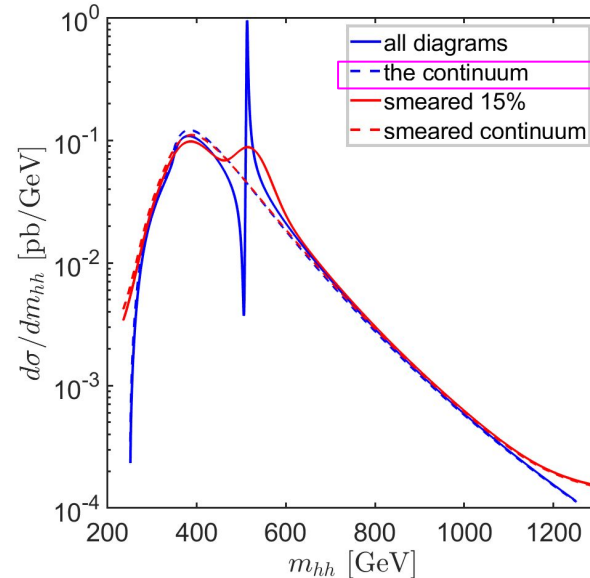
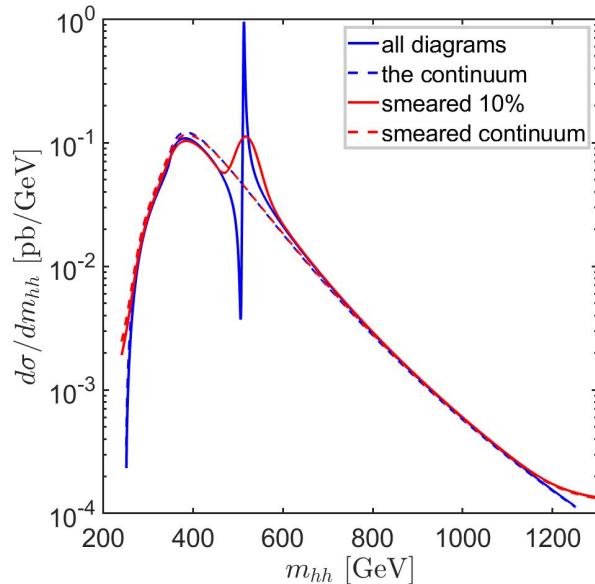
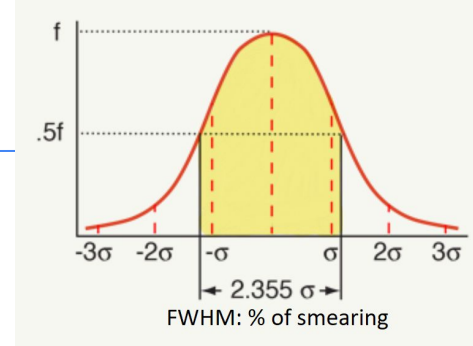
The **relative sign** of the top Yukawa and the BSM coupling to the heavy Higgs gives a **structure** to the resonance:

$\text{sign}(\lambda_{hhH} \cdot \xi_H^t)$	structure
+	peak-dip
-	dip-peak



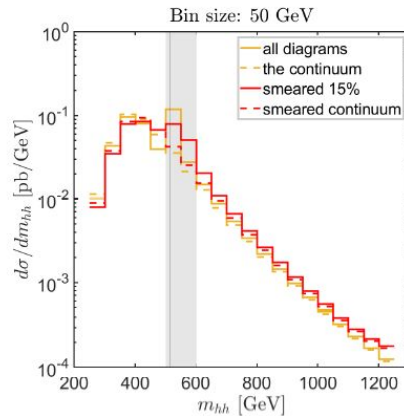
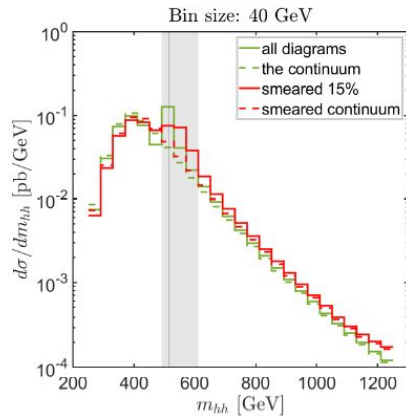
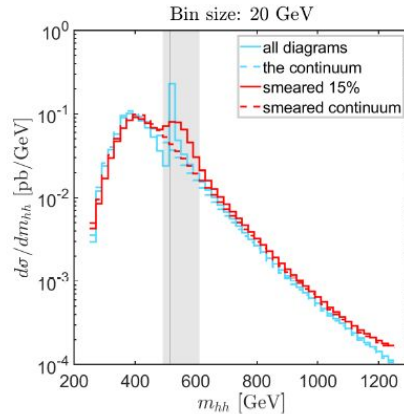
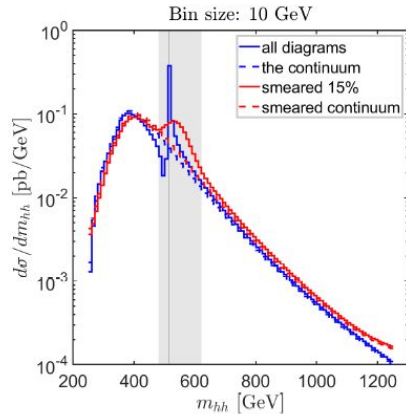
# Experimental challenges: smearing

- Differential cross section measurements are affected by the finite resolution of particle detectors → observed spectrum is “**smear**ed”
- We try to mimic this effect by artificially smearing the theoretical prediction introducing **Gaussian uncertainties** in the invariant mass



box diagram + SM-like Higgs exchange

# Experimental challenges: binning (15 % smearing)



- We define a value for the ‘significance of the signal’ according to the excess of the number of events. Assuming:  $\mathcal{L} = 6000 \text{ fb}^{-1}$

events below resonant  
smeared contribution

events below continuum  
smeared contribution

$$R := \frac{\sum_i (N^R - N^C)}{\sqrt{\sum_i N^C}}$$

Bin size	R
10 GeV	3.79
20 GeV	3.99
40 GeV	3.87
50 GeV	4.04

- Window definition:

$$(N^R - N^C) > (\text{bin size})/50$$

→ Smearing dilutes  
more the resonance  
than binning

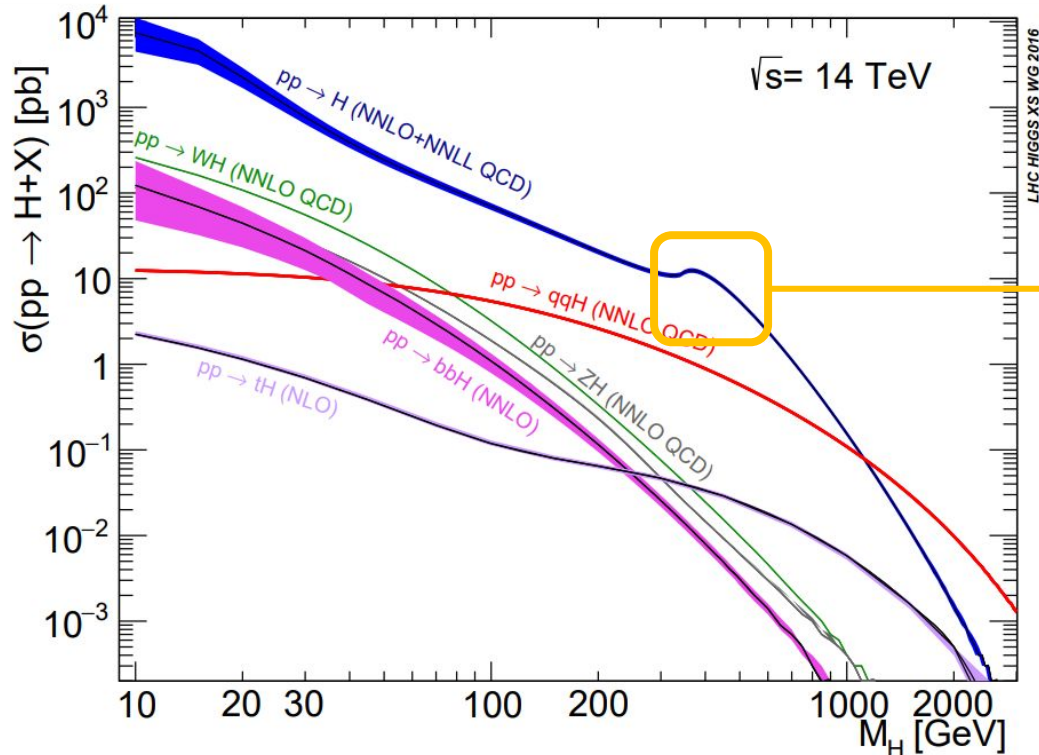
# Conclusion

**Invariant mass distributions** give information about:

1. Deviations of  $\kappa_\lambda$  that can be seen in the low  $m_{hh}$  threshold
2. **resonant production** that can be embedded in BSM models:
  - **mass** of the intermediate Higgs boson  $\rightarrow$  **position** of the resonance
  - **total decay width** of the resonance  $\rightarrow$  **height** of the resonance
  - relative sign of the **couplings**  $\rightarrow$  **structure** of the resonance

These effects may be (partially) washed out by experimental precision (**smearing**)

# Backup: Single Higgs production

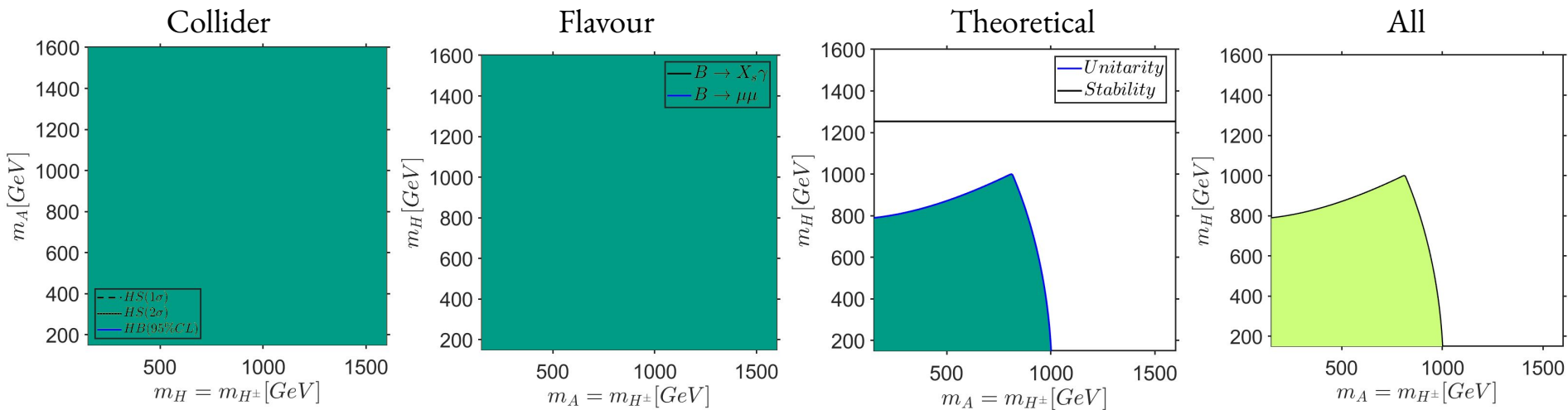


Top pair threshold  $\rightarrow$  gives a hint on the results for Higgs pair production

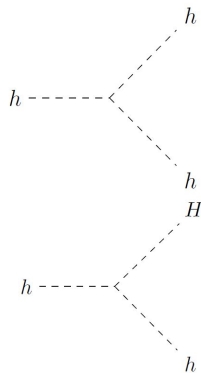
[LHC Higgs Working Group:  
[CERN Yellow Report 4](#)]

# Backup: Constraints

BP: Type I,  $\cos(\beta - \alpha) = 0.2$ ,  $\tan \beta = 10$ ,  $m_{12}^2 = m_H^2 \cos^2 \alpha / \tan \beta$



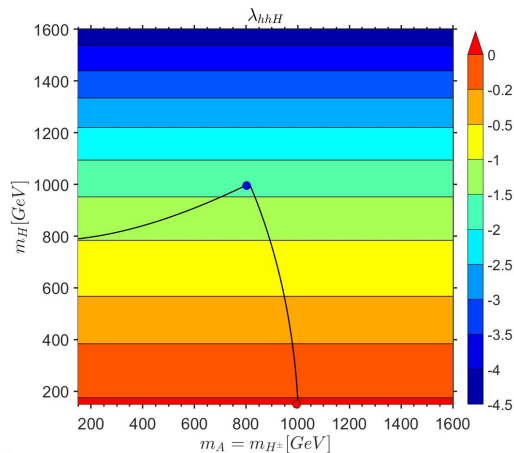
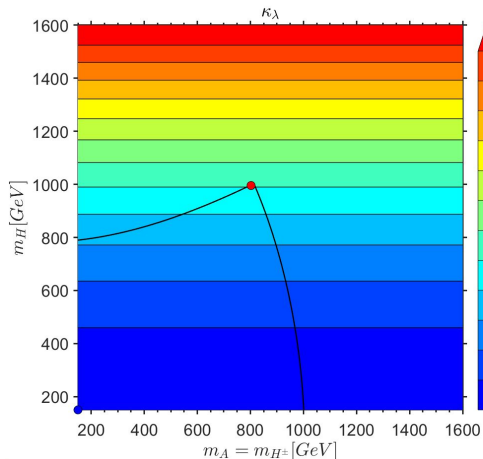
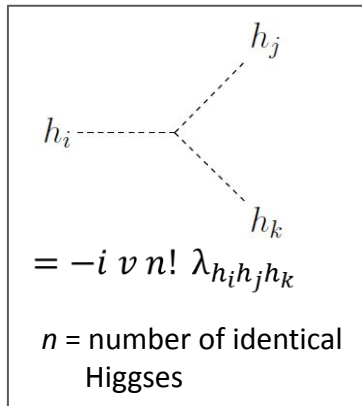
# Backup: Feynman rules for triple Higgs couplings



$$\lambda_{hhh} = \frac{1}{2v^2} \left\{ m_h^2 s_{\beta-\alpha}^3 + (3m_h^2 - 2\bar{m}^2) c_{\beta-\alpha}^2 s_{\beta-\alpha} + 2 \cot 2\beta (m_h^2 - \bar{m}^2) c_{\beta-\alpha}^3 \right\}$$

$$\lambda_{hhH} = \frac{-c_{\beta-\alpha}}{2v^2} \left\{ (2m_h^2 + m_H^2 - 4\bar{m}^2) s_{\beta-\alpha}^2 + 2 \cot 2\beta (2m_h^2 + m_H^2 - 3\bar{m}^2) s_{\beta-\alpha} c_{\beta-\alpha} - (2m_h^2 + m_H^2 - 2\bar{m}^2) c_{\beta-\alpha}^2 \right\}.$$

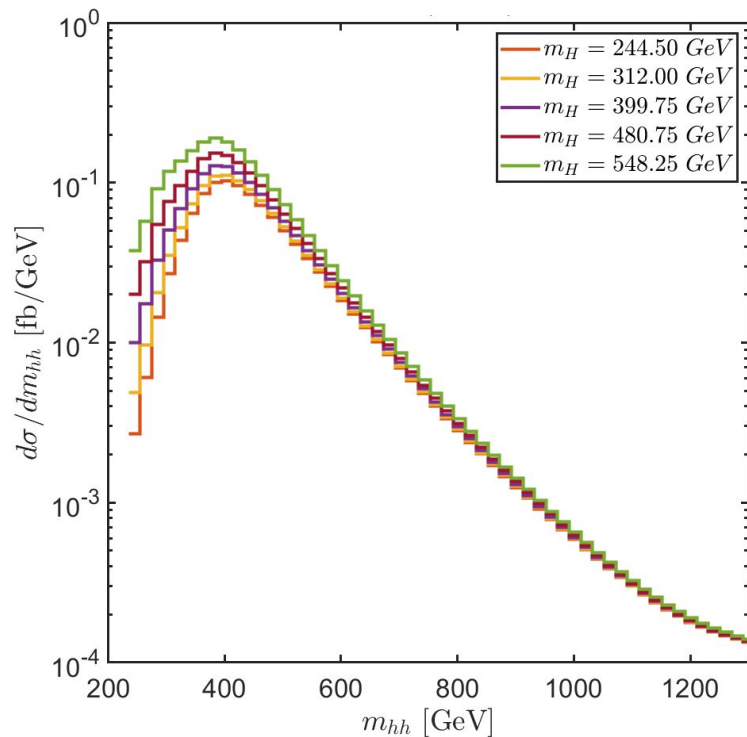
Notation:



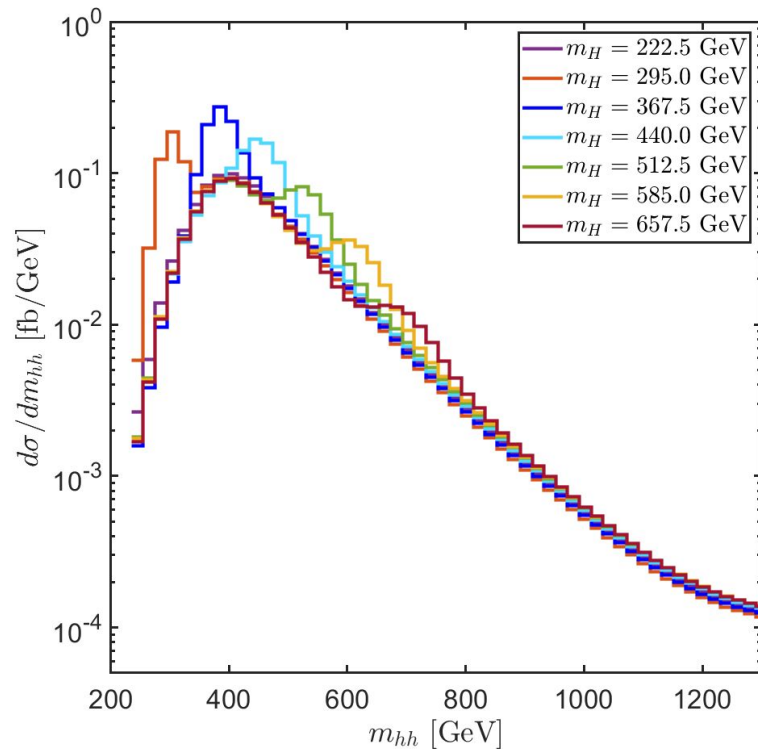
$$\bar{m}^2 = \frac{m_{12}^2}{\sin(\beta)\cos(\beta)}$$

# Smearing applied on the invariant mass distributions

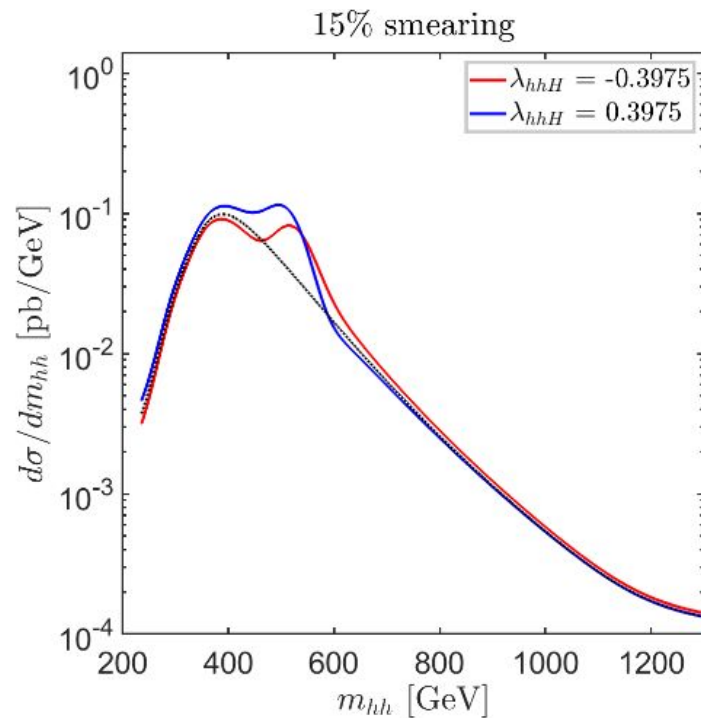
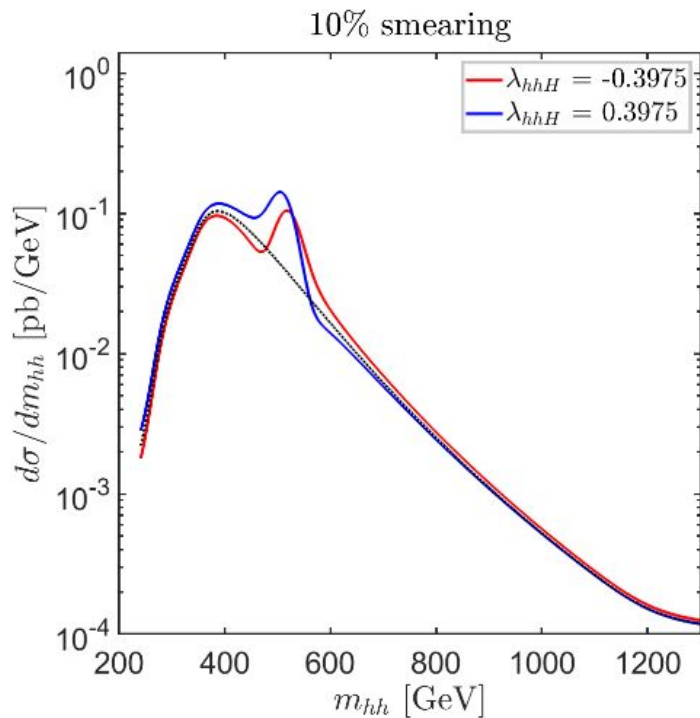
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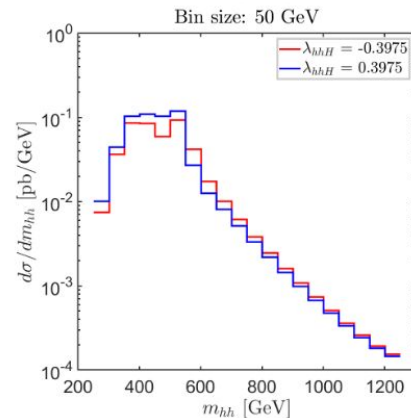
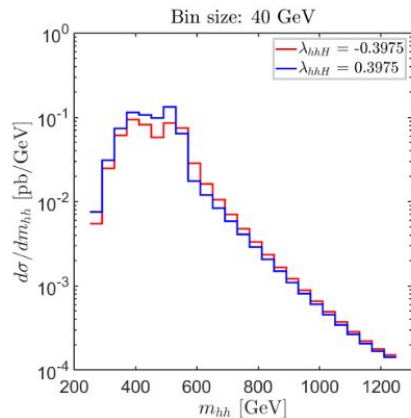
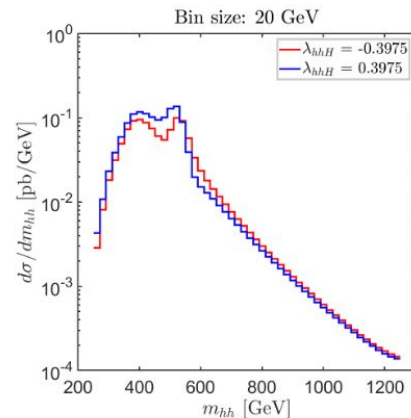
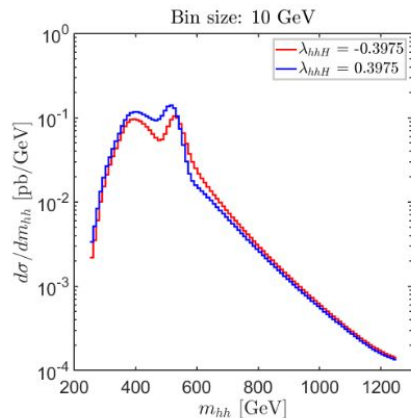


# Smearing applied on the structure of the resonance





# Binning applied on the structure of the resonance



15 % smearing