## EVIDENCE FOR A CHANGE OF SLOPE IN LARGE-t ELASTIC PROTON-PROTON SCATTERING AT $\sqrt{s}$ = 53 GeV

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## ABSTRACT

New experimental results are presented on proton-proton elastic scattering in the range of momentum transfer  $4 \text{ GeV}^2 < -t < 10 \text{ GeV}^2$  at the centre-of-mass energy of  $\sqrt{s} = 53 \text{ GeV}$ . The data have been obtained using the Split-Field Magnet detector at the CERN Intersecting Storage Rings. We observe another change of slope of the differential cross-section near  $-t = 6.5 \text{ GeV}^2$ .

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We have extended our previous measurements [1] of large-t elastic proton-proton scattering in the range  $4 \text{ GeV}^2 < -t < 10 \text{ GeV}^2$  in order to investigate with better sensitivity the open question whether further diffraction-like minima exist at large values of the momentum transfer. The data have been obtained using the Split-Field Magnet detector (SFM) at the CERN Intersecting Storage Rings (ISR). We report here results obtained at a centre-of-mass energy of  $\sqrt{s} = 53 \text{ GeV}$  for an integrated luminosity of  $4.87 \times 10^9 / \text{mb}$ . We observe a new break in the t-dependence near  $t = -6.5 \text{ GeV}^2$  followed by an exponential decrease of the cross-section with a slope of  $b = (0.88 \pm 0.13) \text{ GeV}^{-2}$  in the t-range from 6 to 10 GeV<sup>2</sup>.

The SFM detector has been described before [2]. It consists of two forward telescopes, each equipped with 12 multiwire proportional chambers, 1 m high, 2 m wide, and having 2 mm wire spacing. Each chamber has a vertical and a horizontal wire plane. The average magnetic field is 1.0 T, providing a momentum resolution of  $\Delta p/p = \pm 7\%$  for scattered protons. Events are selected by requiring two charged tracks (one in each telescope) at the trigger level, and by a cut on the deviation from collinearity, after having performed the geometrical reconstruction and a kinematical fit with four constraints. In terms of  $\chi^2$  this cut corresponds to accepting events at the 1% confidence level.

The value of the differential cross-section at  $t=-10~{\rm GeV}^2$  is about  $10^{10}$  times smaller than at t=0. A careful study of possible contamination by background is therefore required. We have investigated the collinearity distribution, expressed in terms of  $\chi^2$ , as a function of momentum transfer. A scatter diagram for events selected by the criteria described above is shown in Fig. 1. Attributing events with  $\chi^2 > 9$  to background, we estimate a residual contamination of less than 5% for accepted events with  $\chi^2 < 9$  at all t-values. This residual background has been neglected.

We have determined absolute differential cross-sections by correcting for the t-dependent acceptance of the detector and by applying an over-all normalization factor. The detailed procedure is described in Ref. 2; the acceptance at large t has been improved over the previous experiment [1] and has now a value larger than 40% for -t > 3 GeV<sup>2</sup>.

The momentum transfer t is measured with a resolution of  $\sigma_t$  = ±0.04 GeV<sup>2</sup> at t = -4 GeV<sup>2</sup> increasing linearly to  $\sigma_t$  = ±0.10 GeV<sup>2</sup> at t = -10 GeV<sup>2</sup>.

The over-all normalization factor is obtained by collecting monitor counts simultaneously with data taking. The monitoring telescope has been calibrated using the van der Meer method [3]; we estimate the systematic uncertainty of this monitor to be ±5%.

The results\*) are shown in Fig. 2. Describing the differential cross-section by a single exponential with slope b we observe a succession of changing slopes;  $b_1 = (10.3 \pm 0.02) \text{ GeV}^{-2} \text{ for } 0.25 < -t < 0.6 \text{ GeV}^2$  \*\*), followed by a narrow minimum at  $-t = 1.34 \text{ GeV}^2$  and a second maximum at  $-t = 1.8 \text{ GeV}^2$ ;  $b_2 = (1.81 \pm 0.02) \text{ GeV}^{-2}$  for  $2.2 < -t < 6 \text{ GeV}^2$ , followed by a break near  $-t = 6.5 \text{ GeV}^2$  and a third slope  $b_3 = (0.88 \pm 0.13) \text{ GeV}^{-2}$  for  $6 < -t < 10 \text{ GeV}^2$  \*\*\*).

Several authors have predicted a second diffraction-like minimum in the t-range near 4  $\text{GeV}^2$  and an associated change in slope, by using optical models [4], constituent models [5] or exchange models [6]. To our knowledge none has predicted this set of slope values and breaks [7]. Wakaizumi and Tanimoto [8] have shown that a model of multiple scattering of constituents with suitably chosen form factors and wave functions can describe this data. Van Hove [9], using a quark-glue model, predicts a continuous change of slope with b = 1.4  $\text{GeV}^{-2}$  at t = -8  $\text{GeV}^2$ .

<sup>\*)</sup> Tables of the differential cross-sections are available from the authors.

<sup>\*\*)</sup> This value has been obtained at  $\sqrt{s}$  = 23 GeV and 62 GeV, see Ref. 2.

<sup>\*\*\*)</sup> Restricting the maximum likelihood fit to the range  $7 < -t < 10 \text{ GeV}^2$ , we find  $b_3 = (0.39 \pm 0.23) \text{ GeV}^{-2}$ .

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## Figure captions

- Fig. 1 : Scatter diagram of collinearity in terms of  $\chi^2$  as a function of momentum transfer. Events with  $\chi^2$  < 9 are selected. The residual background contamination is less than 5% at all t-values.
- Fig. 2  $^{*}$ : Differential cross-sections as a function of momentum transfer t. The error bars represent statistical and systematic errors. An overall scale uncertainty of  $\pm 5\%$  has to be added.

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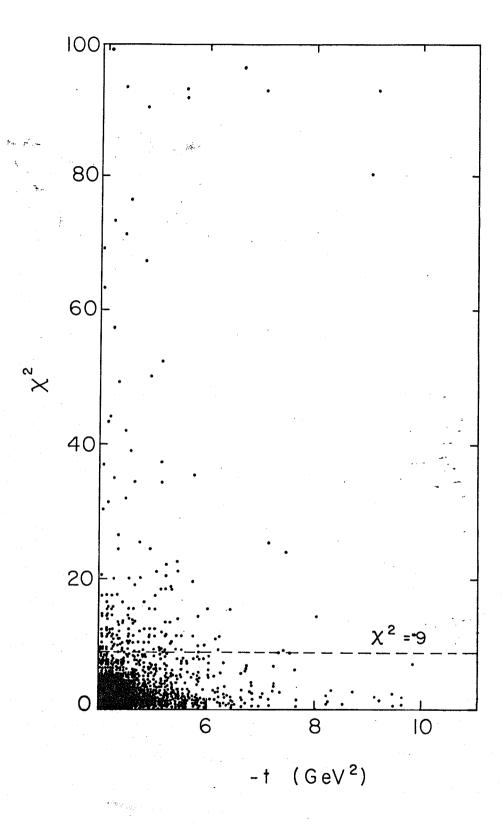


Fig. 1

