

# The Compact Muon Solenoid Experiment

Mailing address: CMS CERN, CH-1211 GENEVA 23, Switzerland



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# $B_s$ decay vertex resolution

A. Starodumov $^{a)}$ , Z. Xi $e^{b)}$ 

INFN, Pisa, Italy

#### Abstract

The resolution of the  $B_s$  decay vertex in the channel  $B_s^0 \rightarrow D_s^- \pi^+ \rightarrow \phi \pi^- \pi^+ \rightarrow K^+ K^- \pi^- \pi^+$  is evaluated. Full simulation of the decay channel is performed using CMSIM v.112 package. The CMS vertex detector for low luminosity operation with a first pixel layer at R = 3.8cm is assumed. The parameterisation of the secondary vertex resolution is done versus the transverse momentum of  $B_s$ .

<sup>&</sup>lt;sup>a)</sup> On leave from ITEP, Moscow, Russia

<sup>&</sup>lt;sup>b)</sup> also Scuola Normale Superiore, Pisa, Italy

# **1** Introduction

One of the most promising channel for the  $B_s - \overline{B_s}$  oscillation study [1] is the following:

$$B_s^0 \to D_s^- \pi^+ \to \phi \pi^- \pi^+ \to K^+ K^- \pi^- \pi^+$$
 (1)

The channel is very challenging since only one charged track ( $\pi$ ) comes from the  $B_s$  decay vertex. This vertex can be reconstructed only if  $D_s$  decay vertex is reconstructed as well. The algorithm used in the study minimises  $\chi^2$  of the simultaneous fit of the two vertices. In this Note we study the decay vertex resolution which can be obtained with the CMS tracker.

The note is organised as follows. In Section 2 we describe the event sample used in the study. Results of secondary vertex reconstruction using CMSIM package is shown in Section 3. Parameterisation of the secondary vertex resolution versus transverse momentum of  $B_s$  is done in Section 4. In Section 5 we discuss results and make conclusion.

# 2 Event sample

Proton-proton collisions at centre of mass energy of 14 TeV have been simulated using PYTHIA 5.7 event generator [2]. Both mechanisms of heavy quark production, gluon splitting and fusion, are taken into account. About  $10^6$  $b\bar{b}$  events have been selected and stored (for the details of the simulation procedure see [3]). Simulation of each event is stopped at the parton level.

Each  $b\bar{b}$  event hadronised, fragmented and decayed 5 times to get reasonable statistics of selected  $B_s$  mesons. Decay channels of both B mesons have been forced: one B decays into  $\mu + X$  (in the study a muon tagging technique is used) and  $B_s$  decays according to the equation (1) The collected statistics corresponds to one month of LHC operation assuming integrated luminosity of  $L = 10^4 pb^{-1}$  and total  $b\bar{b}$  cross-section  $\sigma_{b\bar{b}} = 500 \mu b$ .

Fig. 1 shows  $p_t$  distributions of  $B_s$  mesons and final state hadrons after the kinematical selection described in Section 3.2. Because of different initial random number used to hadronise and decay the same event 5 times, no any bias is introduced during this procedure. Fig. 1 illustrates that there is no artificial statistical fluctuations in the transverse momentum spectra.

# **3** Secondary vertex reconstruction

#### **3.1** How to fit $B_s$ decay vertex

To investigate the  $B_s$  oscillation one should measure the proper time of  $B_s$ . This is calculated according to the following formula:

$$t = m \times L/p \tag{2}$$

where t is the proper time, m is the mass, p is the momentum and L is the flight path of  $B_s$ .

The most important term in this formula is the flight path. The precision of the flight path mainly depends on the precision of the  $B_s$  decay vertex reconstruction. In decay (1) only one track ( $\pi$ ) comes from the  $B_s$  itself, while the other tracks come from the  $D_s$ . To reconstruct the  $B_s$  decay vertex one needs to reconstruct  $D_s$  decay vertex as well. In the study two vertices have been reconstructed in space simultaneously by minimising the following expression:

$$\chi^{2} = \sum \frac{d_{i}^{2}(X_{B}, Y_{B}, Z_{B}, a)}{\sigma_{i}^{2}}$$
(3)

Here  $i = 1 \div 4$ ;  $d_i$  is a distance from track *i* to the corresponding decay vertex: for i = 1 it is the  $B_s$  decay vertex, for  $i = 2 \div 4$  it is the  $D_s$  decay vertex;  $\sigma_i$  are impact parameter errors, which come from track fitting and are treated separately in the  $r - \phi$  and r - z planes, *a* is the distance between the  $B_s$  and  $D_s$  decay vertices, which satisfies the following equation:

$$\vec{X}_{D_s} = \vec{X}_{B_s} + a \frac{\vec{P}_{D_s}}{|P_{D_s}|}$$
(4)

Here,  $\vec{P}_{D_s}$  is the reconstructed momentum of  $D_s$ ,  $\vec{X}_{B_s}$  and  $\vec{X}_{D_s}$  three-vectors of  $B_s$  and  $D_s$  decay vertices respectively.



Figure 1:  $p_t$  distribution of  $B_s$  (a) and final state hadrons:  $\pi 1$  (b) is the pion from  $B_s$  decay,  $\pi 2$ (c) is the pion from  $D_s$  decay, K (d) comes from the  $\phi$  decay.

## **3.2** Selection of events

Track reconstruction and vertex fitting have been done using CMSIM v.112 [4]. CMS tracker Version 3 is used in the study with low luminosity option of pixel detectors (GEOM control card is 'PXLX' 23). For track finding Monte-Carlo method (RECO control card is 'TRAK' 1) is used. MC method means that for the track fitting hits from the real simulated track are used. So, possible pattern recognition problems are ignored in this study. The following kinematical and selection criteria have been applied:

- $|\eta| \le 2.4$  for all particles;
- $p_t \ge 1 GeV$  for all hadrons (Fig. 1  $b \div d$ );
- single muon trigger with the threshold of  $p_t \ge 6GeV$ ;
- 2 hits in pixel layers for each track (1 hit/layer)  $\equiv$  good tracks;
- only events with 4 good tracks from  $B_s$  decay are used for vertex fit;
- $\chi^2$  of the vertices fit is less than 20 (Fig. 2 a).



Figure 2:  $\chi^2$  of the vertex fit (a) and residuals of the  $B_s$  decay vertex in **x** (b), **y** (c), **z** (d) directions

# **3.3** Result of the $B_s$ decay vertex fit

The fitting procedure returns reconstructed vertices and  $\chi^2$  of the fit, but not the errors of the fit. To estimate a secondary vertex resolution one can plot the difference between reconstructed and simulated values. Fig.  $2b \div d$  shows residuals of the  $B_s$  decay vertex in the **x**, **y** and **z** directions. The Gaussian fit of these distributions provides secondary vertex resolution which is ~  $60\mu m$  in the **x** and **y** directions and ~  $80\mu m$  in the **z** direction. To evaluate the error of proper time one has to find a projection of the error on the flight path direction. Fig. 3 illustrates flight path, the errors of the secondary vertex along the flight path (in the  $r - \phi$  plane and in space) and in the perpendicular direction. The error along the  $B_s$  flight path is about ~  $77/115\mu m$  in the transverse plane and in space respectively. The error in the direction which is perpendicular to the flight path is quite small ~  $20/36\mu m$  in the transverse plane and space respectively. In Section 4 secondary vertex resolution along the transverse flight path of  $B_s$  will be parametrised versus  $p_t$  of  $B_s$ .

## **4** Parametrisation of the secondary vertex resolution

The difference between generated and reconstructed decay vertex of  $B_s$  divided by flight path (relative error) in the transverse plane can be parametrised with two Gaussian. A narrow Gaussian (G1) reflects the resolution of the tracker system and a wide Gaussian (G2) presents 'non-Gaussian' tails. The parametrisation is done with the



Figure 3: Flight path  $(L, L_{xy})$  and the errors of secondary vertex along  $(\Delta L_{xy}, \Delta L)$  and perpendicular  $(\Delta L_{xy}^T, \Delta L^T)$  to the flight path. On the left  $(a \div c)$ : distributions are in the transverse plane, on the right  $(d \div f)$ : in space.

following functions:

$$P_i = A_i exp(-\frac{(x-\mu_i)^2}{2\sigma_i^2}) \tag{5}$$

where i = 1, 2 corresponds to narrow and wide Gaussian distributions, respectively.

To parameterise the relative error the full range of the transverse momentum of  $B_s$  has been divided in five regions:  $\leq 10 GeV$ ,  $10 \div 12 GeV$ ,  $12 \div 16 GeV$ ,  $16 \div 24 GeV$  and  $24 \div 40 GeV$ . On Fig. 4 one can see the fit of the relative error with a function 'G1+G2' for the full region and the first three regions of transverse momentum of  $B_s$ . The regions have been chosen to have approximately the same number of events in each of them. The width of the narrow Gaussian depends on transverse momentum of  $B_s$  and varies from 6.5% to 4% while the width of wide Gaussian stays almost constant at the level of 32%-34%. Fig. 5 shows the parametrisation of the narrow Gaussian versus transverse momentum of  $B_s$ . The function used for the parametrisation is the following:

$$\frac{\sigma_t}{L_{xy}} = \sqrt{(\frac{P_1}{p_t})^2 + P_2^2}$$
(6)

Here,  $\sigma_t$  is the error of the flight path in the transverse plane,  $L_{xy}$  is the transverse flight path,  $p_t$  is the transverse momentum of  $B_s$ ,  $P_1$  and  $P_2$  are parameters of the fit. Parameter  $P_2$  has the meaning of an asymptotic resolution.



Figure 4: Fitted relative flight path error distributions for full (a) and the first three ranges of  $P_t$  of  $B_s$  ( $b \div d$ ).

For the CMS tracker in the low luminosity option for the vertex detector (the radius of the first pixel layer in the barrel is 3.8*cm*) the asymptotic resolution in the transverse plane is about 4%.

To simulate the reconstruction of the  $B_s$  decay vertex in the channel (1) one needs to count on the wide Gaussian. The relative fraction of narrow and wide Gaussian can be determined by calculating the ratio of the integrals:

$$I_i = \int P_i dx = \int A_i exp(-\frac{(x-\mu_i)^2}{2\sigma_i^2}) dx \tag{7}$$

which is equal to:

$$\frac{I_1}{I_2} = \sqrt{2\pi} \frac{A_1}{A_2} \frac{\sigma_1 \sigma_2}{\sqrt{\sigma_2^2 - \sigma_1^2}}$$
(8)

For all considered ranges of the  $B_s$  transverse momentum this ratio is almost constant and equal to ~ 1. The width of the wide Gaussian about 5 times larger than the width of the narrow one.

# 5 Conclusion

In this Note the parametrisation of the secondary vertex resolution for  $B_s$  decay channel (1) is done using the low luminosity option of the CMS vertex detector. Track reconstruction and vertex fitting are performed with CMSIM



Figure 5:  $B_s$  decay vertex resolution versus  $p_t$  of  $B_s$ .

v.112 package. Secondary vertex resolution can be parametrised as the sum of two Gaussian distributions with the relative fraction of 1:1. The width of the narrow Gaussian can be parametrised with the function (6), where P1=37% and P2=4%. The width of the wide Gaussian is about 5 times bigger than the width of the narrow one.

The ssymptotic secondary vertex resolution of 4% is comparable with the resolutions estimated by LHC-B ( $\sim 3\%$  [5]) and ATLAS ( $\sim 4.5\%$  [6]) Collaborations.

Such a good performance the CMS vertex detector and the whole tracker provides a possibility to measure  $x_s$  value up to upper limit (about 30) predicted by the Standard Model [7] in the decay channel (1).

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# References

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