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# A Combination of Preliminary Results on Triple Gauge Boson Couplings Measured by the LEP Experiments

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#### Abstract

This note presents a combination of published and preliminary measurements of triple gauge boson couplings (TGCs) from the four LEP experiments. We combine the charged TGCs  $g_1^Z$ ,  $\kappa_{\gamma}$  and  $\lambda_{\gamma}$  for the first time since the inclusion of  $O(\alpha_{em})$  corrections to the WW production process in Monte Carlo simulations. The combinations of neutral TGCs  $h_i^V$  and  $f_i^V$  are also presented.

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The collaborations each take responsibility for the preliminary data of their own experiment. WWW access at http://www.cern.ch/LEPEWWG/lepww/tgc/

## 1 Introduction

During its second phase of operation, from 1996 until November 2000, the  $e^+e^-$  collider LEP at CERN steadily increased its centre-of-mass energy from 161 GeV reaching up to 209 GeV in the final year. Until the end of LEP-II operation, a total integrated luminosity of approximately 700 pb<sup>-1</sup> per experiment has been recorded.

The measurement of gauge boson couplings and the search for possible anomalous contributions due to the effects of new physics beyond the Standard Model are among the principal physics aims at LEP-II [1]. Combined preliminary measurements of triple gauge boson couplings are presented here. Results from W-pair production are combined for the first time including  $O(\alpha_{em})$  corrections. The combinations for neutral TGCs remain unchanged from our last report [2]. A new combination of quartic gauge coupling (QGC) results, including the sign convention as reported in [3,4] and the reweighting based on [3] is foreseen for our next report.

The W-pair production process,  $e^+e^- \rightarrow W^+W^-$ , involves charged triple gauge boson vertices between the W<sup>+</sup>W<sup>-</sup> and the Z or photon. During LEP-II operation, about 10,000 W-pair events were collected by each experiment. Single W ( $e\nu$ W) and single photon ( $\nu\bar{\nu}\gamma$ ) production at LEP are also sensitive to the WW $\gamma$  vertex. Results from these channels are also included in the combination for some experiments; the individual references should be consulted for details.

For the charged TGCs, Monte Carlo calculations (RacoonWW [5] and YFSWW [6]) incorporating an improved treatment of  $O(\alpha_{em})$  corrections to the WW production process have recently become available. These corrections affect the measurements of the charged TGCs in W-pair production. Preliminary results including these  $O(\alpha_{em})$  corrections have been submitted from all four LEP collaborations ALEPH [7], DELPHI [8], L3 [9] and OPAL [10]. Therefore, updated combinations are made for the charged TGC measurements.

At centre-of-mass energies exceeding twice the Z boson mass, pair production of Z bosons is kinematically allowed. Here, one searches for the possible existence of triple vertices involving only neutral electroweak gauge bosons. Such vertices could also contribute to  $Z\gamma$  production. In contrast to triple gauge boson vertices with two charged gauge bosons, purely neutral gauge boson vertices do not occur in the Standard Model of electroweak interactions.

#### 1.1 Charged Triple Gauge Boson Couplings

The parametrisation of the charged triple gauge boson vertices is described in References [1, 11–16]. The most general Lorentz invariant Lagrangian which describes the triple gauge boson interaction has fourteen independent complex couplings, seven describing the WW $\gamma$  vertex and seven describing the WWZ vertex. Assuming electromagnetic gauge invariance as well as C and P conservation, the number of independent TGCs reduces to five. A common set is  $\{g_1^Z, \kappa_Z, \kappa_\gamma, \lambda_Z, \lambda_\gamma\}$  where  $g_1^Z = \kappa_Z = \kappa_\gamma = 1$  and  $\lambda_Z = \lambda_\gamma = 0$  in the Standard Model. The parameters proposed in [1] and used by the LEP experiments are  $g_1^Z, \lambda_\gamma$  and  $\kappa_\gamma$  with the gauge constraints:

$$\kappa_{\rm Z} = g_1^{\rm Z} - (\kappa_{\gamma} - 1) \tan^2 \theta_{\rm W} , \qquad (1)$$

$$\lambda_{\rm Z} = \lambda_{\gamma} \,, \tag{2}$$

where  $\theta_W$  is the weak mixing angle. The couplings are considered as real, with the imaginary parts fixed to zero. In contrast to previous LEP combinations [2, 17], we are now quoting the measured coupling values themselves and not their deviation from the Standard Model.

The photonic couplings  $\lambda_{\gamma}$  and  $\kappa_{\gamma}$  are related to the magnetic and electric properties of the W-boson. One can write the lowest order terms for a multipole expansion describing the W- $\gamma$ 

interaction as a function of  $\lambda_{\gamma}$  and  $\kappa_{\gamma}$ . For the magnetic dipole moment  $\mu_W$  and the electric quadrupole moment  $q_W$  one obtains:

$$\mu_W = \frac{e}{2m_W} \left( 1 + \kappa_\gamma + \lambda_\gamma \right), \tag{3}$$

$$q_W = -\frac{e}{m_W^2} \left(\kappa_\gamma - \lambda_\gamma\right). \tag{4}$$

The inclusion of  $O(\alpha_{em})$  corrections in the Monte Carlo calculations has a considerable effect on the charged TGC measurement. Both the total cross-section and the differential distributions are affected. The cross-section is reduced by 1-2% (depending on the energy). Amongst the differential distributions, the effects are naturally more complex. The polar W<sup>-</sup> production angle carries most of the information on the TGC parameters; its shape is modified to be more forwardly peaked. In a fit to data, the effect manifests itself as a negative shift of the obtained TGC values with a magnitude of typically -0.015 for  $\lambda_{\gamma}$  and  $g_1^{\rm Z}$  and -0.04 for  $\kappa_{\gamma}$ .

#### 1.2Neutral Triple Gauge Boson Couplings

There are two classes of Lorentz invariant structures associated with neutral TGC vertices which preserve  $U(1)_{em}$  and Bose symmetry, as described in [12, 18].

The first class refers to anomalous  $Z\gamma\gamma^*$  and  $Z\gamma Z^*$  couplings which are accessible at LEP in the process  $e^+e^- \rightarrow Z\gamma$ . The parametrisation contains eight couplings:  $h_i^V$  with i = 1, ..., 4 and  $V = \gamma$ ,Z. The superscript  $\gamma$  refers to  $Z\gamma\gamma^*$  couplings and superscript Z refers to  $Z\gamma Z^*$  couplings. The photon and the Z boson in the final state are considered as on-shell particles, while the third boson at the vertex, the s-channel internal propagator, is off shell. The couplings  $h_1^V$  and  $h_2^V$  are CP-odd while  $h_3^V$  and  $h_4^V$  are CP-even.

The second class refers to anomalous  $ZZ\gamma^*$  and  $ZZZ^*$  couplings which are accessible at LEP-II in the process  $e^+e^- \rightarrow ZZ$ . This anomalous vertex is parametrised in terms of four couplings:  $f_i^V$  with i = 4, 5 and  $V = \gamma, Z$ . The superscript  $\gamma$  refers to  $ZZ\gamma^*$  couplings and the superscript Z refers to ZZZ<sup>\*</sup> couplings, respectively. Both Z bosons in the final state are assumed to be on-shell, while the third boson at the triple vertex, the s-channel internal propagator, is off-shell. The couplings  $f_4^V$  are CP-odd whereas  $f_5^V$  are CP-even. The  $h_i^V$  and  $f_i^V$  couplings are assumed to be real and they vanish at tree level in the Standard

Model.

#### $\mathbf{2}$ Measurements

The combined results presented here are obtained from charged and neutral electroweak gauge boson coupling measurements as discussed above. The individual references should be consulted for details about the data samples used.

The charged TGC analyses of ALEPH, DELPHI, L3 and OPAL use data collected at LEP-II up to centre-of-mass energies of 209 GeV. These analyses are using different channels, typically the semileptonic and fully hadronic W-pair decays [7–10]. The full data set is analysed by ALEPH, whereas DELPHI presently uses all data at and above 189 GeV, and L3 and OPAL at this stage use data from the semileptonic channel only. Anomalous TGCs affect both the total production cross-section and the shape of the differential cross-section as a function of the polar W<sup>-</sup> production angle. The relative contributions of each helicity state of the W bosons are also changed, which in turn affects the distributions of their decay products. The analyses presented by each experiment make use of different combinations of each of these quantities. In general, however, all analyses use at least the expected variations of the total production cross-section and the W<sup>-</sup> production angle. Results from  $e\nu W$  and  $\nu \bar{\nu} \gamma$  production are included by some experiments. Single W production is particularly sensitive to  $\kappa_{\gamma}$ , thus providing information complementary to that from W-pair production.

The *h*-coupling analyses of ALEPH, DELPHI and L3 use data collected up to centre-ofmass energies of 209 GeV. The OPAL measurements so far use the data at 189 GeV. The results of the *f*-couplings are now obtained from the whole data set above the ZZ-production threshold by all of the experiments. The experiments already pre-combine different processes and final states for each of the couplings. For the neutral TGCs, the analyses use measurements of the total cross sections of  $Z\gamma$  and ZZ production and the differential distributions: the  $h_i^V$ couplings [19–22] and the  $f_i^V$  couplings [19, 20, 23, 24] are determined.

# 3 Combination Procedure

The combination is based on the individual likelihood functions from the four LEP experiments. Each experiment provides the negative log likelihood,  $\log \mathcal{L}$ , as a function of the coupling parameters to be combined. The single-parameter analyses are performed fixing all other parameters to their Standard Model values. The two-parameter analyses are performed setting the remaining parameters to their Standard Model values. For the charged TGCs, the gauge constraints listed in Section 1.1 are always enforced.

The  $\log \mathcal{L}$  functions from each experiment include statistical as well as those systematic uncertainties which are considered as uncorrelated between experiments. For both single- and multi-parameter combinations, the individual  $\log \mathcal{L}$  functions are combined. It is necessary to use the  $\log \mathcal{L}$  functions directly in the combination, since in some cases they are not parabolic, and hence it is not possible to properly combine the results by simply taking weighted averages of the measurements.

The main contributions to the systematic uncertainties that are uncorrelated between experiments arise from detector effects, background in the selected signal samples, limited Monte Carlo statistics and the fitting method. Their importance varies for each experiment and the individual references should be consulted for details.

In the neutral TGC sector, the systematic uncertainties arising from the theoretical cross section prediction in Z $\gamma$ -production ( $\simeq 1\%$  in the  $q\bar{q}\gamma$ - and  $\simeq 2\%$  in the  $\nu\bar{\nu}\gamma$  channel) are treated as correlated. For ZZ production, the uncertainty on the theoretical cross section prediction is small compared to the statistical accuracy and therefore is neglected. Smaller sources of correlated systematic uncertainties, such as those arising from the LEP beam energy, are for simplicity treated as uncorrelated.

The combination procedure for neutral TGCs, where the relative systematic uncertainties are small, is unchanged with respect to the previous LEP combinations of electroweak gauge boson couplings [2,17]. The correlated systematic uncertainties in the h-coupling analyses are taken into account by scaling the combined log-likelihood functions by the squared ratio of the sum of statistical and uncorrelated systematic uncertainty over the total uncertainty including all correlated uncertainties. For the general case of non-Gaussian probability density functions, this treatment of the correlated errors is only an approximation; it also neglects correlations in the systematic uncertainties between the parameters in multi-parameter analyses.

In the charged TGC sector, systematic uncertainties considered correlated between the experiments are the theoretical cross section prediction (0.5%) for W-pair production and 5\% for

single W production), hadronisation effects, the final state interactions, namely Bose-Einstein correlations and colour reconnection, and the uncertainty in the radiative corrections themselves. The latter is by convention presently taken to be the difference between Monte Carlo samples with and without the  $O(\alpha_{em})$  corrections respectively, as a reliable and applicable estimate of this uncertainty is at this time unavailable. This conservative approach makes the radiative correction uncertainty by far the dominant contribution.

In case of the charged TGCs, the systematic uncertainties considered correlated between the experiments amount to roughly 70% of the combined statistical and uncorrelated uncertainties for  $\lambda_{\gamma}$  and  $g_1^Z$ , whilst for  $\kappa_{\gamma}$  their contribution is equal. Hence an improved combination procedure [25] is used for the charged TGCs. This procedure allows the combination of statistical and correlated systematic uncertainties, independently of the analysis method chosen by the individual experiments. At present, only single-parameter combinations of the charged TGCs are performed. Tests on the feasibility and effects of applying the same combination method in multidimensional fits are currently underway [26].

The combination of charged TGCs uses the likelihood curves and correlated systematic errors submitted by each of the four experiments. The procedure is based on the introduction of an additional free parameter to take into account the systematic uncertainties, which are treated as shifts on the fitted TGC value, and are assumed to have a Gaussian distribution. A simultaneous minimisation of both parameters (TGC and systematic error) is performed to the log-likelihood function.

In detail, the combination proceeds in the following way: the set of measurements from the LEP experiments ALEPH, DELPHI, OPAL and L3 is given with statistical plus uncorrelated systematic uncertainties in terms of likelihood curves:  $-\log \mathcal{L}_{stat}^A(x)$ ,  $-\log \mathcal{L}_{stat}^D(x)$  $-\log \mathcal{L}_{stat}^L(x)$  and  $-\log \mathcal{L}_{stat}^O(x)$ , respectively, where x is the coupling parameter in question. Also given are the shifts for each of the five totally correlated sources of uncertainty mentioned above; each source S is leading to systematic errors  $\sigma_A^S$ ,  $\sigma_D^S$ ,  $\sigma_L^S$  and  $\sigma_O^S$ .

Additional parameters  $\Delta^S$  are included in order to take into account a Gaussian distribution for each of the systematic uncertainties. The procedure then consists in minimising the function:

$$-\log \mathcal{L}_{total} = \sum_{E=A,D,L,O} \log \mathcal{L}_{stat}^{E} (x - \sum_{S=DPA,\sigma_{WW},HAD,BE,CR} (\sigma_{E}^{S} \Delta^{S})) + \sum_{S} \frac{(\Delta^{S})^{2}}{2}$$
(5)

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where x and  $\Delta_S$  are the free parameters, and the sums run over the four experiments and the five systematic errors. The resulting uncertainty on x will take into account all sources of uncertainty, yielding a measurement of the coupling with the error representing statistical and systematic sources. The projection of the minima of the log-likelihood as a function of x gives the combined log-likelihood curve including statistical and systematic uncertainties. The advantage over the previously used scaling method is that it treats systematic uncertainties that are correlated between the experiments correctly, whilst not forcing the averaging of these systematic uncertainties into one global LEP systematics scaling factor. In other words, the (statistical) precision of each experiment now gets reduced by its own correlated systematic errors, instead of an averaged LEP systematic error. The method has been cross-checked against the scaling method, and was found to give comparable results. The inclusion of the systematic uncertainties lead to small differences as expected by the improved treatment of correlated systematic errors, a similar behaviour as seen in Monte Carlo comparisons of these two combinations methods [26]. Furthermore, it was shown that the minimisation-based combination method used for the charged TGCs agrees with the OO-based method for any realistic ratio of statistical and systematic uncertainties. Further details on the improved combination method can be found in [25].

In all combinations, the one standard deviation uncertainties (68% confidence level) are obtained by taking the coupling values for which  $\Delta \log \mathcal{L} = +0.5$  above the minimum. The 95% confidence level (C.L.) limits are given by the coupling values for which  $\Delta \log \mathcal{L} = +1.92$ above the minimum. These cut-off values are used for obtaining the results of both single- and multi-parameter analyses reported here. Note that in the case of the neutral TGCs, double minima structures appear in the negative log-likelihood curves. For multi-parameter analyses, the two dimensional 68% C.L. contour curves for any pair of couplings are obtained by requiring  $\Delta \log \mathcal{L} = +1.15$ , while for the 95% C.L. contour curves  $\Delta \log \mathcal{L} = +3.0$  is required.

### 4 Results

We present results from the four LEP experiments on the various electroweak gauge boson couplings, and their combination. The combined charged TGC results including  $O(\alpha_{em})$  corrections are available for the first time, whilst the neutral TGC results remain unchanged since our last note [2]. The results quoted for each individual experiment are calculated using the methods described in Section 3. Therefore they may differ slightly from those reported in the individual references, as the experiments in general use other methods to combine the data from different channels, and to include systematics uncertainties. In particular for the charged couplings, experiments using an OO-based combination method (ALEPH, OPAL) obtain results with small differences compared to the values given by our combination technique. These small differences have been studied in Monte Carlo tests and are well understood [26]. For the *h*-coupling result from OPAL and DELPHI, a slightly modified estimate of the systematic uncertainty due to the theoretical cross section prediction is responsible for slightly different limits compared to the published results.

### 4.1 Charged Triple Gauge Boson Couplings

The individual analyses and results of the experiments for the charged couplings are described in [7-10].

The results of single-parameter fits from each experiment are shown in Table 1, where the errors include both statistical and systematic effects. The individual  $\log \mathcal{L}$  curves and their sum are shown in Figure 1. The results of the combination are given in Table 2. A list of the systematic errors treated as fully correlated between the LEP experiments, and their shift on the combined fit result are given in Table 3.

#### 4.2 Neutral Triple Gauge Boson Couplings in $Z\gamma$ Production

The individual analyses and results of the experiments for the h-couplings are described in [19–22].

#### Single-Parameter Analyses

The results for each experiment are shown in Table 7, where the errors include both statistical and systematic uncertainties. The individual  $\log \mathcal{L}$  curves and their sum are shown in Figures 2 and 3. The results of the combination are given in Table 4. From Figures 2 and 3 it is clear

that the sensitivity of the L3 analysis [21] is the highest amongst the LEP experiments. This is partially due to the use of a larger phase space region, which increases the statistics by about a factor two, and partially due to additional information from using an Optimal Observable technique.

#### Two-Parameter Analyses

The results for each experiment are shown in Table 5, where the errors include both statistical and systematic uncertainties. The 68% C.L. and 95% C.L. contour curves resulting from the combinations of the two-dimensional likelihood curves are shown in Figure 4. The LEP average values are given in Table 6.

## 4.3 Neutral Triple Gauge Boson Couplings in ZZ Production

The individual analyses and results of the experiments for the f-couplings are described in [19, 20, 23, 24].

### Single-Parameter Analyses

The results for each experiment are shown in Table 8, where the errors include both statistical and systematic uncertainties. The individual  $\log \mathcal{L}$  curves and their sum are shown in Figure 5. The results of the combination are given in Table 9.

#### Two-Parameter Analyses

The results from each experiment are shown in Table 10, where the errors include both statistical and systematic uncertainties. The 68% C.L. and 95% C.L. contour curves resulting from the combinations of the two-dimensional likelihood curves are shown in Figure 6. The LEP average values are given in Table 11.

# Conclusions

With the LEP-combined charged TGC results, the existence of triple gauge boson couplings among the electroweak gauge bosons is experimentally verified. No significant deviation from the Standard Model prediction is seen for any of the electroweak gauge boson couplings studied. As an example, these data allow the Kaluza-Klein theory [27], in which  $\kappa_{\gamma} = -2$ , to be excluded completely [28].

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Parameter	ALEPH	DELPHI	L3	OPAL
$g_1^{\mathrm{Z}}$	$1.008^{+0.034}_{-0.033}$	$1.002^{+0.041}_{-0.043}$	$0.952^{+0.053}_{-0.048}$	$0.976^{+0.039}_{-0.039}$
$\kappa_\gamma$	$0.926\substack{+0.095\\-0.085}$	$0.966\substack{+0.106\\-0.106}$	$0.892\substack{+0.099\\-0.095}$	$0.819\substack{+0.086\\-0.079}$
$\lambda_\gamma$	$-0.006\substack{+0.035\\-0.034}$	$0.013\substack{+0.048\\-0.045}$	$-0.030\substack{+0.057\\-0.054}$	$-0.060\substack{+0.040\\-0.039}$

Table 1: The measured central values and one standard deviation errors obtained by the four LEP experiments. In each case the parameter listed is varied while the remaining two are fixed to their Standard Model values. Both statistical and systematic errors are included. The values given here differ slightly from the ones quoted in the individual contributions from the four LEP experiments, as a different combination method is used. See text in section 3 for details.

Parameter	68% C.L.	95% C.L.
$g_1^{\mathrm{Z}}$	$0.990^{+0.023}_{-0.024}$	[0.944, 1.035]
$\kappa_\gamma$	$0.896\substack{+0.058\\-0.056}$	[0.786, 1.009]
$\lambda_{\gamma}$	$-0.023^{+0.025}_{-0.023}$	[-0.069, 0.026]

Table 2: The combined 68% C.L. errors and 95% C.L. intervals obtained combining the results from the four LEP experiments. In each case the parameter listed is varied while the other two are fixed to their Standard Model values. Both statistical and systematic errors are included.

Source	$g_1^{\mathrm{Z}}$	$\lambda_\gamma$	$\kappa_\gamma$
$O(\alpha_{em})$ correction	0.015	0.015	0.039
$\sigma_{WW}$ prediction	0.003	0.005	0.014
Hadronisation	0.004	0.002	0.004
Bose-Einstein Correlation	0.005	0.004	0.009
Colour Reconnection	0.005	0.004	0.010
$\sigma_{singleW}$ prediction	-	-	0.011

Table 3: The systematic uncertainties considered correlated between the LEP experiments in the charged TGC combination and their effect on the combined fit results.

Parameter	95% C.L.
$h_1^\gamma$	[-0.056, +0.055]
$h_2^\gamma$	[-0.045, +0.025]
$h_3^\gamma$	[-0.049, -0.008]
$h_4^\gamma$	[-0.002, +0.034]
$h_1^Z$	[-0.13, +0.13]
$h_2^Z$	[-0.078, +0.071]
$h_3^Z$	[-0.20, +0.07]
$h_4^Z$	[-0.05, +0.12]

Table 4: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) obtained combining the results from the four experiments. In each case the parameter listed is varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included.

Parameter		ALEPH	]	DELPHI		L3
$h_1^{\gamma}$	[-0.32,	+ 0.32]	[-0.28,	+0.28]	[-0.17,	+0.04]
$h_2^{\gamma}$	[-0.18,	+0.18]	[-0.17,	+0.18]	[-0.12,	+0.02]
$h_3^\gamma$	[-0.17,	+0.38]	[-0.48,	+ 0.20]	[-0.09,	+0.13]
$h_4^\gamma$	[-0.08,	+0.29]	[-0.08,	+0.15]	[-0.04,	+ 0.11]
$h_1^Z$	[-0.54,	+0.54]	[-0.45,	+0.46]	[-0.48,	+0.33]
$h_2^Z$	[-0.29,	+0.30]	[-0.29,	+0.29]	[-0.30,	+0.22]
$h_3^Z$	[-0.58,	+0.52]	[-0.57,	+0.38]	[-0.43,	+0.39]
$h_4^Z$	[-0.29,	+0.31]	[-0.31,	+0.28]	[-0.23,	+0.28]

Table 5: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) measured by ALEPH, DELPHI and L3. In each case the two parameters listed are varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included.

Parameter	95% C.L.	Correlations
$h_1^{\gamma}$	[-0.16, +0.05]	1.00 + 0.79
$h_2^{\gamma}$	[-0.11, +0.02]	+0.79 1.00
$h_3^\gamma$	[-0.08, +0.14]	1.00 + 0.97
$h_4^\gamma$	[-0.04, +0.11]	+0.97 1.00
$h_1^Z$	[-0.35, +0.28]	1.00 + 0.77
$h_2^Z$	[-0.21, +0.17]	+0.77 1.00
$h_3^Z$	[-0.37, +0.29]	1.00 + 0.76
$h_4^Z$	[-0.19, +0.21]	+0.76 1.00

Table 6: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) obtained combining the results from ALEPH, DELPHI and L3. In each case the two parameters listed are varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included. Since the shape of the log-likelihood is not parabolic, there is some ambiguity in the definition of the correlation coefficients and the values quoted here are approximate.

Parameter	ALEPH	DELPHI	L3	OPAL
$h_1^\gamma$	[-0.14, +0.14]	[-0.15, +0.15]	[-0.06, +0.06]	[-0.13, +0.13]
$h_2^\gamma$	[-0.07, +0.07]	[-0.09, +0.09]	[-0.053, +0.024]	[-0.089, +0.089]
$h_3^\gamma$	[-0.069, +0.037]	[-0.047, +0.047]	[-0.062, -0.014]	[-0.16, +0.00]
$h_4^\gamma$	[-0.020, +0.045]	[-0.032, +0.030]	[-0.004, +0.045]	[+0.01, +0.13]
$h_1^Z$	[-0.23, +0.23]	[-0.24, +0.25]	[-0.17, +0.16]	[-0.22, +0.22]
$h_2^Z$	[-0.12, +0.12]	[-0.14, +0.14]	[-0.10, +0.09]	[-0.15, +0.15]
$h_3^Z$	[-0.28, +0.19]	[-0.32, +0.18]	[-0.23, +0.11]	[-0.29, +0.14]
$h_4^Z$	[-0.10, +0.15]	[-0.12, +0.18]	[-0.08, +0.16]	[-0.09, +0.19]

Table 7: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) measured by the ALEPH, DELPHI, L3 and OPAL. In each case the parameter listed is varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included.

Parameter		ALEPH	I	DELPHI		L3		OPAL
$f_4^{\gamma}$	[-0.26,	+ 0.26]	[-0.26,	+0.28]	[-0.24,	+0.26]	[-0.36,	+0.36]
$f_4^Z$	[-0.44,	+0.43]	[-0.49,	+0.42]	[-0.43,	+0.41]	[-0.55,	+0.64]
$f_5^{\gamma}$	[-0.54,	+0.56]	[-0.48,	+ 0.61]	[-0.48,	+0.56]	[-0.82,	+0.72]
$f_5^Z$	[-0.73,	+0.83]	[-0.42,	+0.69]	[-0.46,	+1.2]	[-0.96,	+0.31]

Table 8: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) measured by ALEPH, DELPHI, L3 and OPAL. In each case the parameter listed is varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included.

Parameter	95% C.L.
$f_4^{\gamma}$	[-0.17, +0.19]
$f_4^Z$	[-0.31, +0.28]
$f_5^\gamma$	[-0.36, +0.40]
$f_5^Z$	[-0.36, +0.39]

Table 9: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) obtained combining the results from all four experiments. In each case the parameter listed is varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included.

Parameter		ALEPH	]	DELPHI		L3		OPAL
$f_4^{\gamma}$	[-0.26,	+0.26]	[-0.26,	+0.28]	[-0.24,	+0.26]	[-0.36,	+0.36]
$f_4^Z$	[-0.44,	+0.43]	[-0.49,	+0.42]	[-0.43,	+0.41]	[-0.54,	+0.63]
$f_5^{\gamma}$	[-0.52,	+0.53]	[-0.52,	+ 0.61]	[-0.48,	+0.56]	[-0.77,	+0.73]
$f_5^Z$	[-0.77,	+0.86]	[-0.44,	+ 0.69]	[-0.46,	+ 1.2]	[-0.96,	+0.44]

Table 10: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) measured by ALEPH, DELPHI, L3 and OPAL. In each case the two parameters listed are varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included.

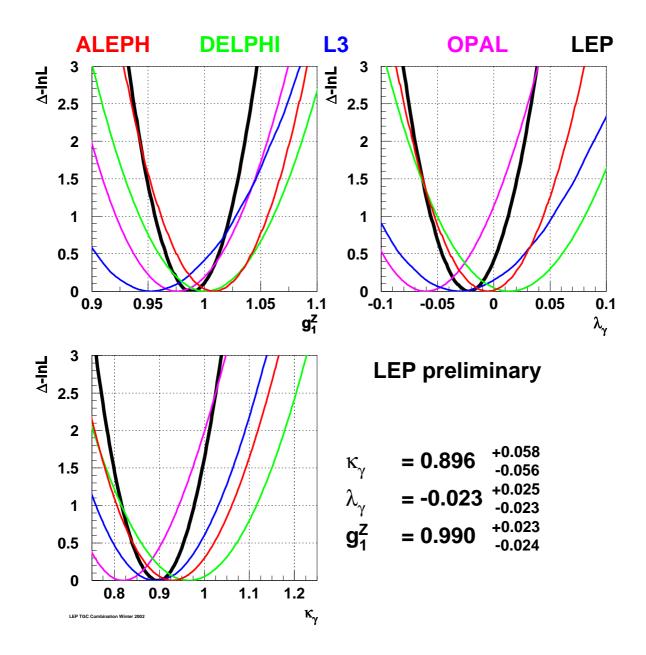


Figure 1: The log  $\mathcal{L}$  curves of the four experiments (thin lines) and the LEP combined curve (thick line) for the three charged TGCs  $g_1^Z$ ,  $\kappa_{\gamma}$  and  $\lambda_{\gamma}$ . In each case, the minimal value is subtracted.

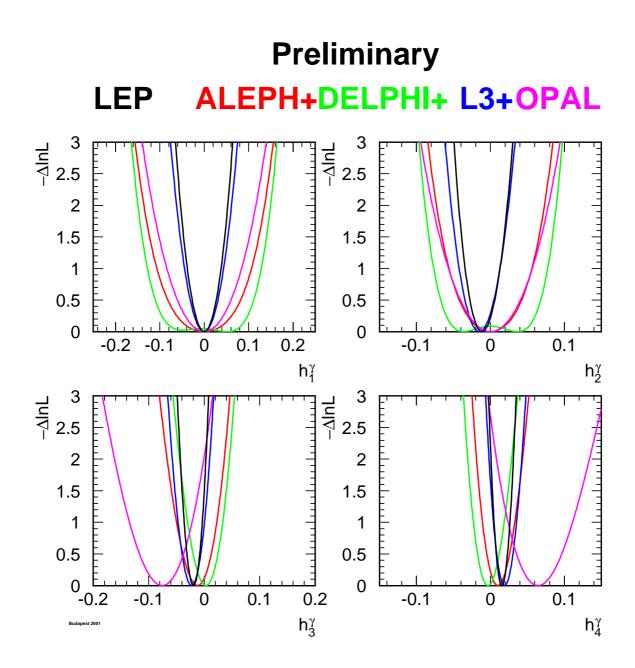


Figure 2: The log  $\mathcal{L}$  curves of the four experiments, and the LEP combined curve for the four neutral TGCs  $h_i^{\gamma}$ , i = 1, 2, 3, 4. In each case, the minimal value is subtracted.

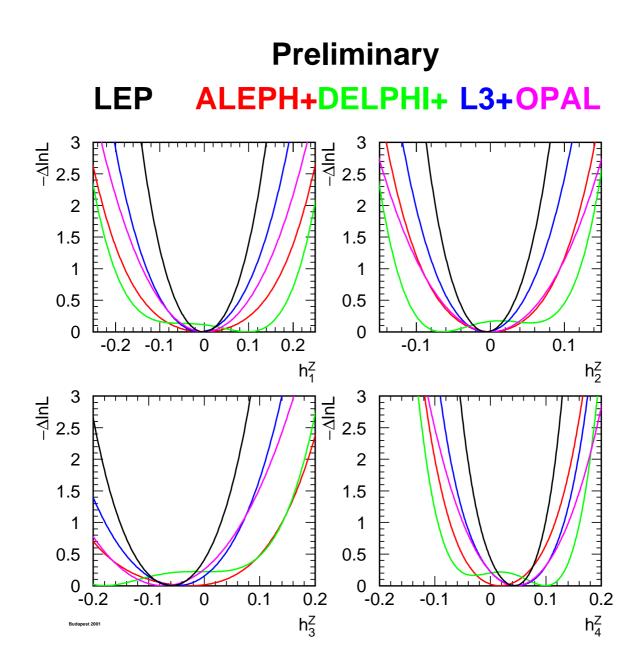


Figure 3: The log  $\mathcal{L}$  curves of the four experiments, and the LEP combined curve for the four neutral TGCs  $h_i^Z$ , i = 1, 2, 3, 4. In each case, the minimal value is subtracted.

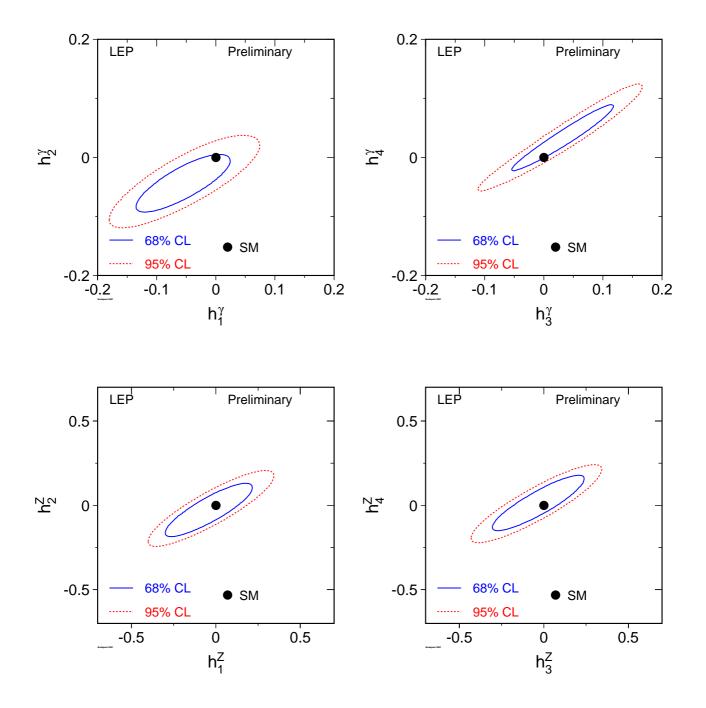


Figure 4: Contour curves of 68% C.L. and 95% C.L. in the planes  $(h_1^{\gamma}, h_2^{\gamma})$ ,  $(h_3^{\gamma}, h_4^{\gamma})$ ,  $(h_1^Z, h_2^Z)$  and  $(h_3^Z, h_4^Z)$  showing the LEP combined result.

Parameter	95%	C.L.	Correlations		
$f_4^{\gamma}$	[-0.17,				
$f_4^Z$	[-0.30,				
$f_5^\gamma$	[-0.34,	+0.38]	1.00	-0.18	
$f_5^Z$	[-0.36,	+ 0.38]	-0.18	1.00	

Table 11: The 95% C.L. intervals ( $\Delta \log \mathcal{L} = 1.92$ ) obtained combining the results from all four experiments. In each case the two parameters listed are varied while the remaining ones are fixed to their Standard Model values. Both statistical and systematic uncertainties are included. Since the shape of the log-likelihood is not parabolic, there is some ambiguity in the definition of the correlation coefficients and the values quoted here are approximate.

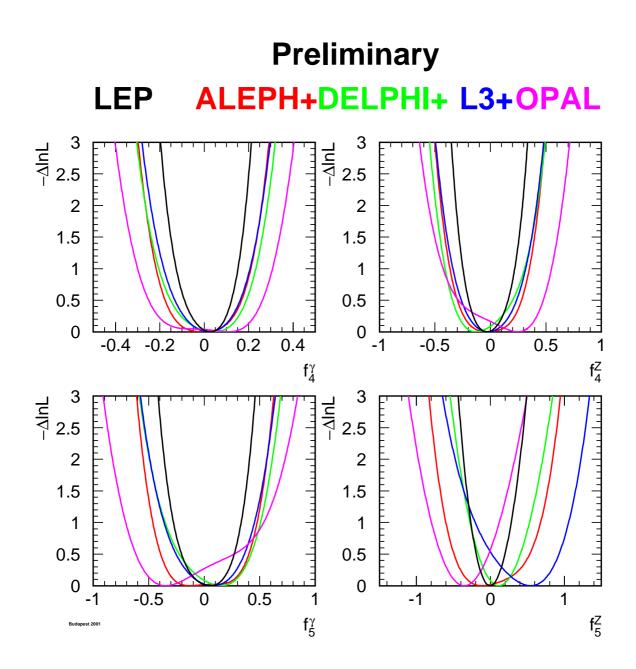


Figure 5: The log  $\mathcal{L}$  curves of the four experiments, and the LEP combined curve for the four neutral TGCs  $f_i^V$ ,  $V = \gamma, Z$ , i = 4, 5. In each case, the minimal value is subtracted.

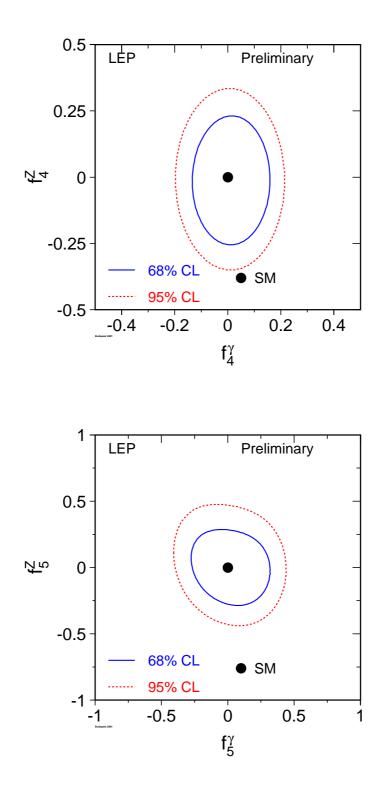


Figure 6: Contour curves of 68% C.L. and 95% C.L. in the plane  $(f_4^{\gamma}, f_4^Z)$  and  $(f_5^{\gamma}, f_5^Z)$  showing the LEP combined result.

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The physics notes describing the preliminary results of the LEP experiments submitted to the winter conferences 2002 are available at:

ALEPH:	http://alephwww.cern.ch
DELPHI:	http://delphiwww.cern.ch/~pubxx/delsec/conferences/moriond02/

- L3: http://13.web.cern.ch/13/conferences/Moriond2002/
- OPAL: http://opal.web.cern.ch/Opal/PPwelcome.html