

Determination of the strong-coupling constant from the Z-boson transverse-momentum distribution with ATLAS

ICHEP, July 2024

Kristof Schmieden, on behalf of the ATLAS Collaboration

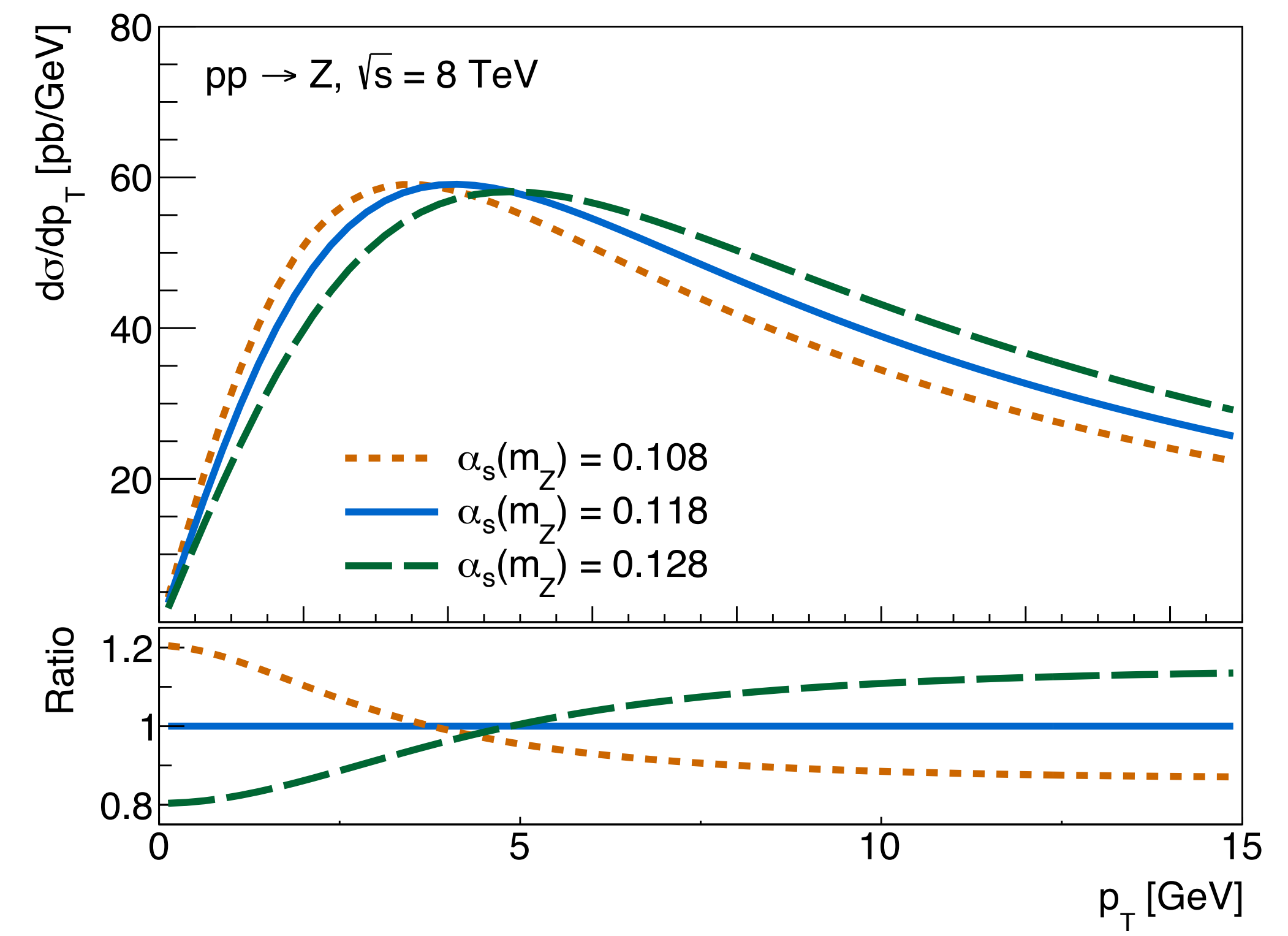
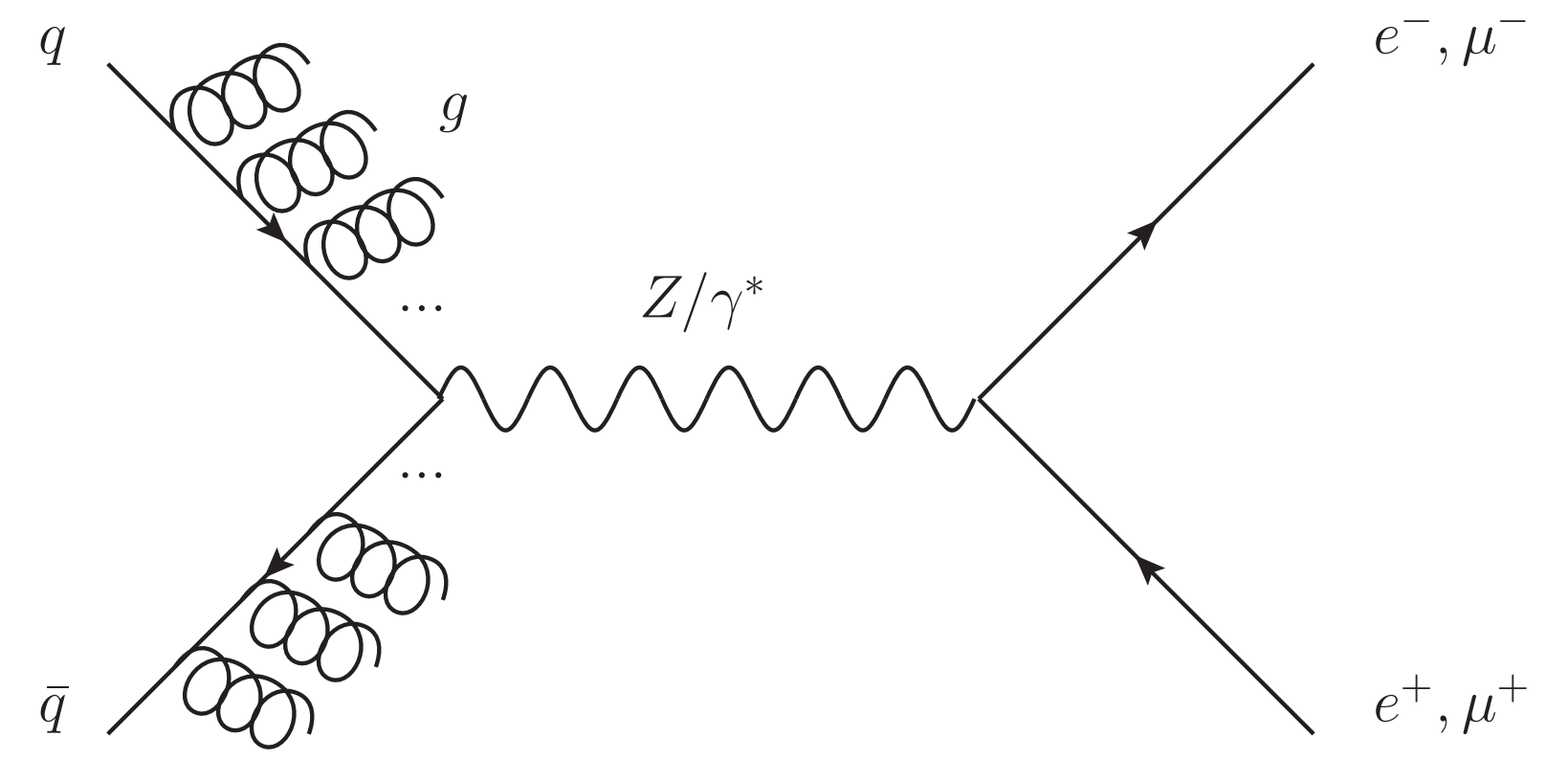


Based on
<https://arxiv.org/abs/2309.12986>
and
[Eur.Phys.J.C 84 \(2024\) 315](#)



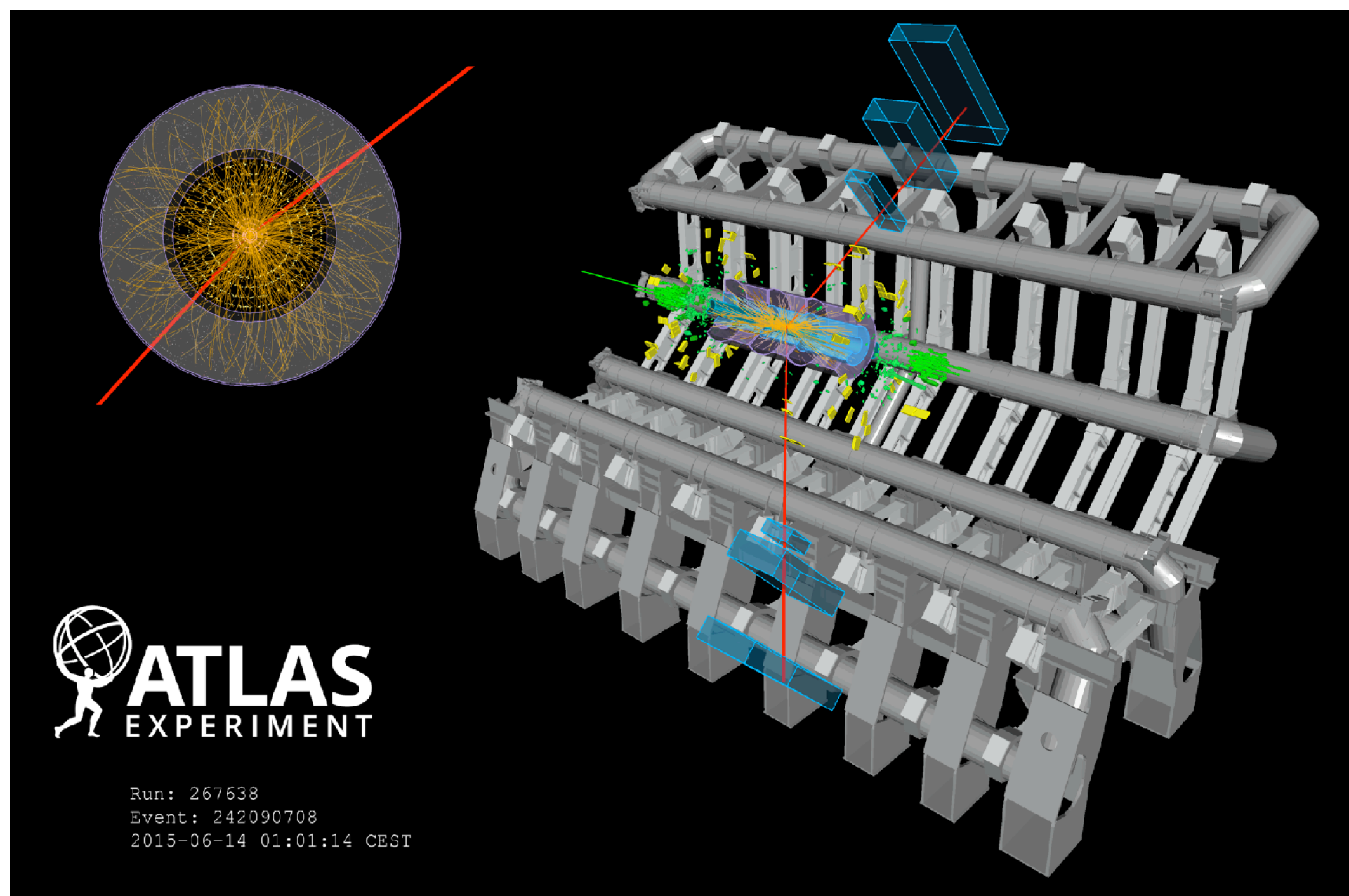
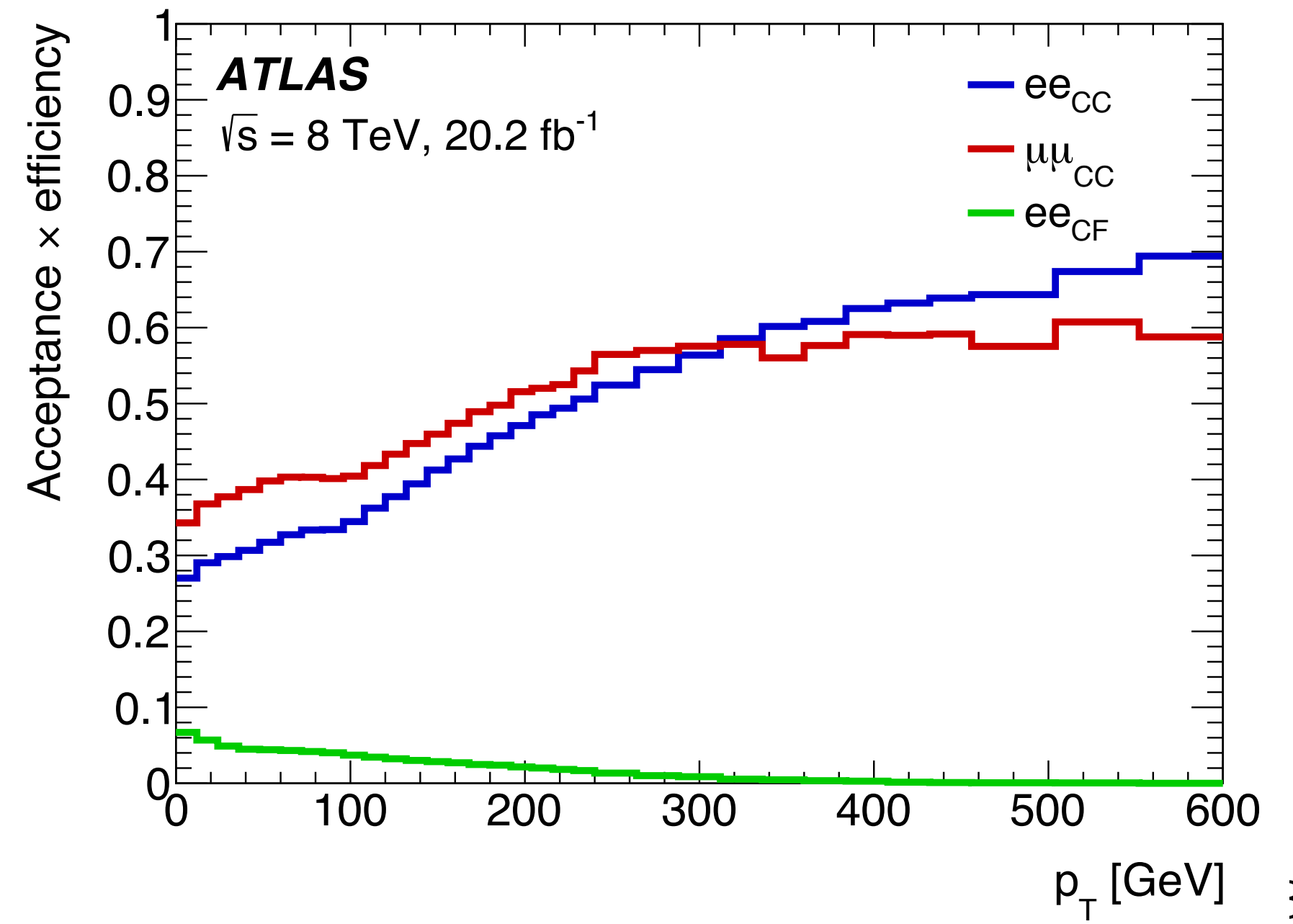
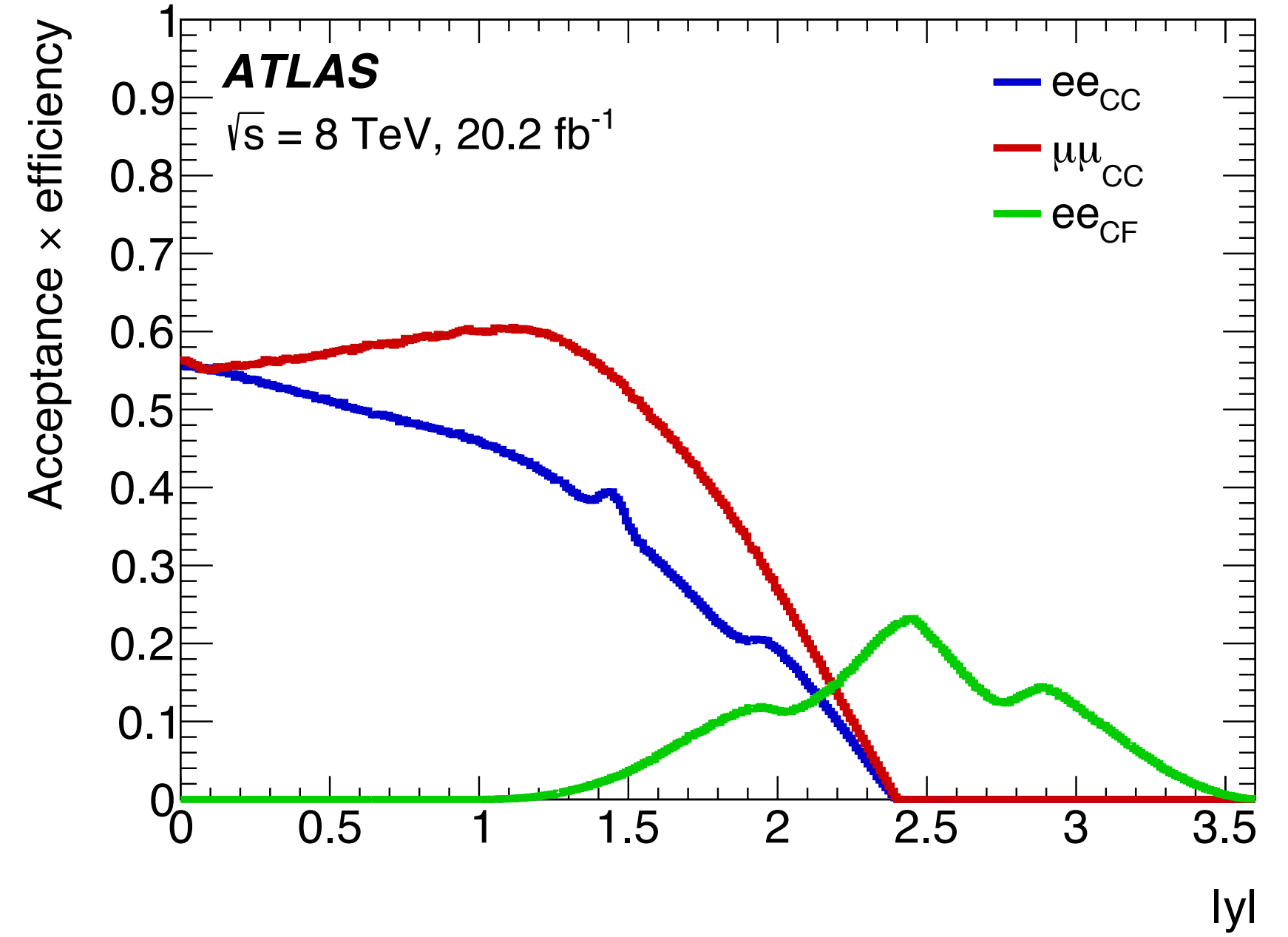
Measure $\alpha_s(m_Z)$ from the $p_T(Z)$ distribution

- Z bosons produced in hadron collisions recoil against QCD initial-state radiation:
 - By momentum conservation, ISR gluons will boost the Z in the transverse plane
- The Sudakov factor is responsible for the existence of a peak in the Z-boson p_T distribution, at values of approximately 4 GeV
- The position of the peak is sensitive to $\alpha_s(m_Z)$
- Semi-inclusive observable, which has advantages of
 - Exclusive obs. (higher exp. sensitivity)
 - Inclusive obs. (higher order theory, smaller non-pQCD effects)



Event Selection

- Three channels
 - **eeCC**: two electrons (6.2M events)
 - $p_T > 20 \text{ GeV}$, $|\eta| < 2.4$
 - **$\mu\mu$ CC**: two muons (7.8M events)
 - $p_T > 20 \text{ GeV}$, $|\eta| < 2.4$
 - **eeCF**: central + forward electron (1.2M events)
 - Forward: $p_T > 20 \text{ GeV}$, $2.5 < |\eta| < 4.9$
- $80 \text{ GeV} < m_{ll} < 100 \text{ GeV}$



$$\frac{d\sigma}{dpdq} = \frac{d^3\sigma^{U+L}}{dp_T dy dm} \left(1 + \cos^2 \theta + \sum_{i=0}^7 A_i(y, p_T, m) P_i(\cos \theta, \phi) \right)$$

- Differential cross section decomposed into basis of spherical harmonics

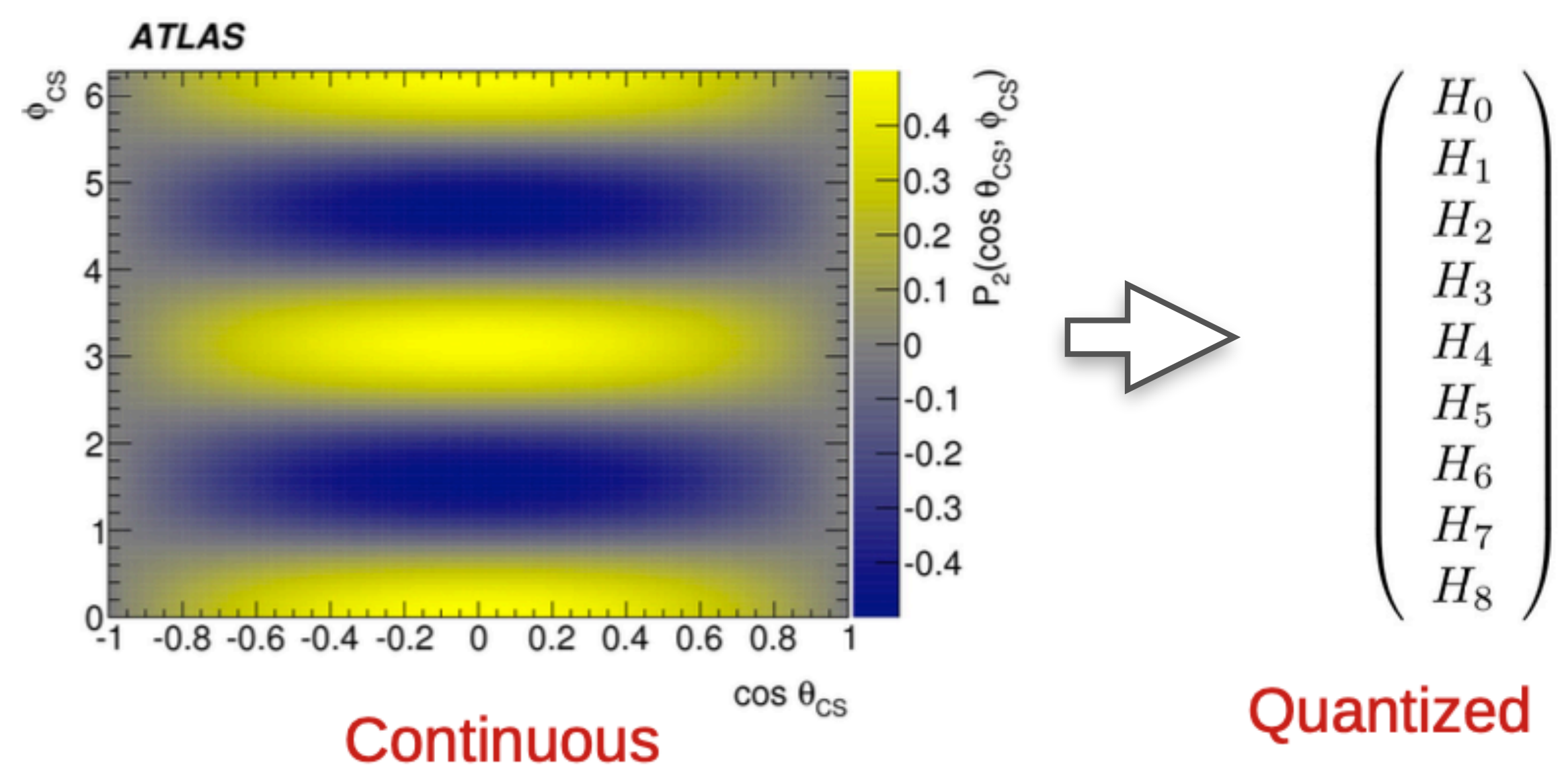
Expected Yield

Reco $(p_T^z, y^z, m^z, \cos\theta, \phi)$ bin

$$N_{\text{exp}}^n(A, \sigma, \theta) = \left\{ \sum_{j=1}^{N_{\text{bins}}^{\text{ana}}} \mathcal{L}_{\sigma_j} \left[t_{8j}^n(\beta) + \sum_{i=0}^7 A_{ij} t_{ij}^n(\beta) \right] \right\} \gamma^n + \sum_{B}^{bkg} T_B^n(\beta)$$

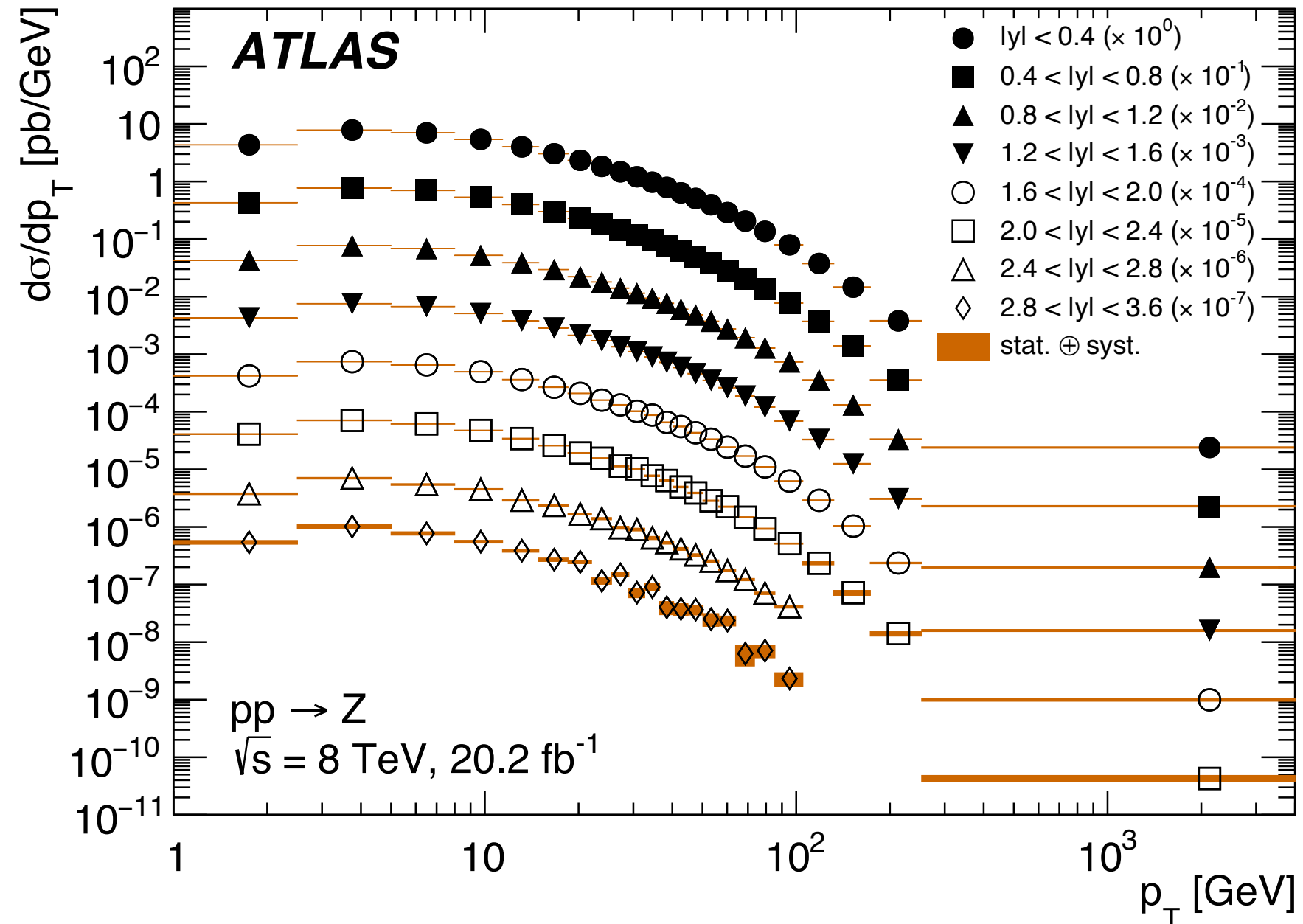
Truth (p_T^z, y^z, m^z) bin Cross section Angular coefficient Templated polynomial Background template

- Likelihood defined in 22528 $(\cos\theta, \phi, p_T, y)$ bins
- Parameters of interests are the 8 A_i + 1 cross section in p_T - y bins: 9 parameters in 176 bins

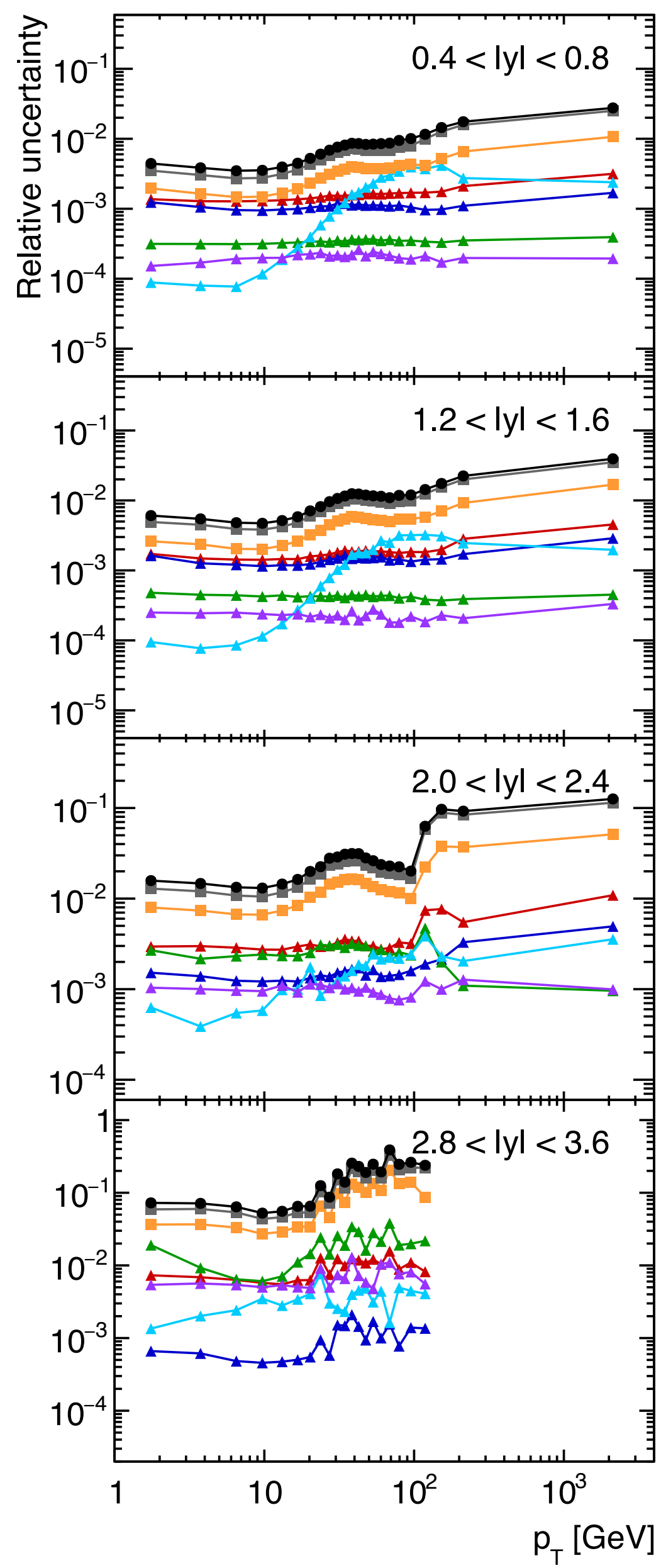
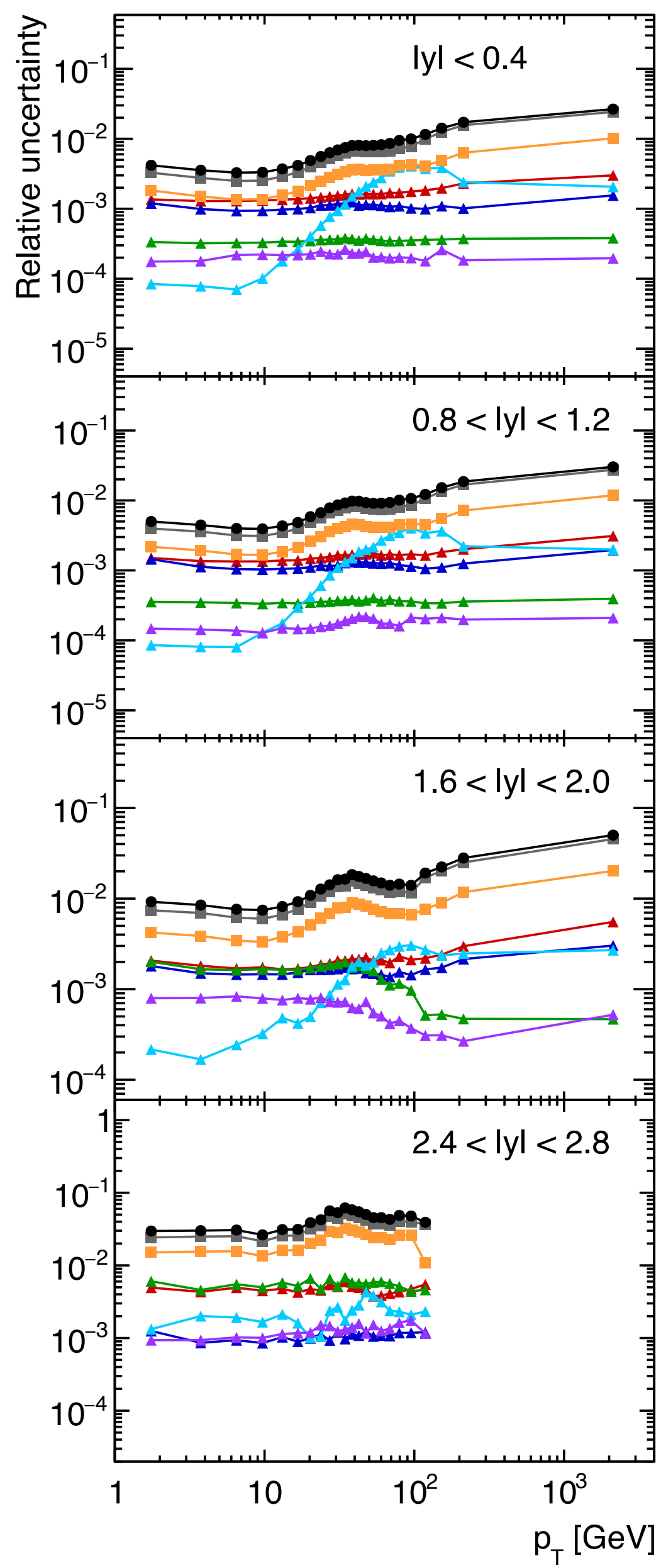


- Measuring the angular coefficients corresponds to building a synthetic “quantized” representation of the $(\cos\theta, \phi)$ kinematic space
- Trade systematics for statistics
- **Very powerful:**
 - Avoids theoretical extrapolation of fiducial lepton cuts to full phase space
 - Thereby opens the door to a rich field of precise interpretations

Double-Differential Drell Yan Measurement



- First measurement at the LHC of full lepton phase space cross section
- Statistically dominated measurement
- Negligible theory uncertainties:
 - Cross-sections are parameters of the fit, not the result of an extrapolation



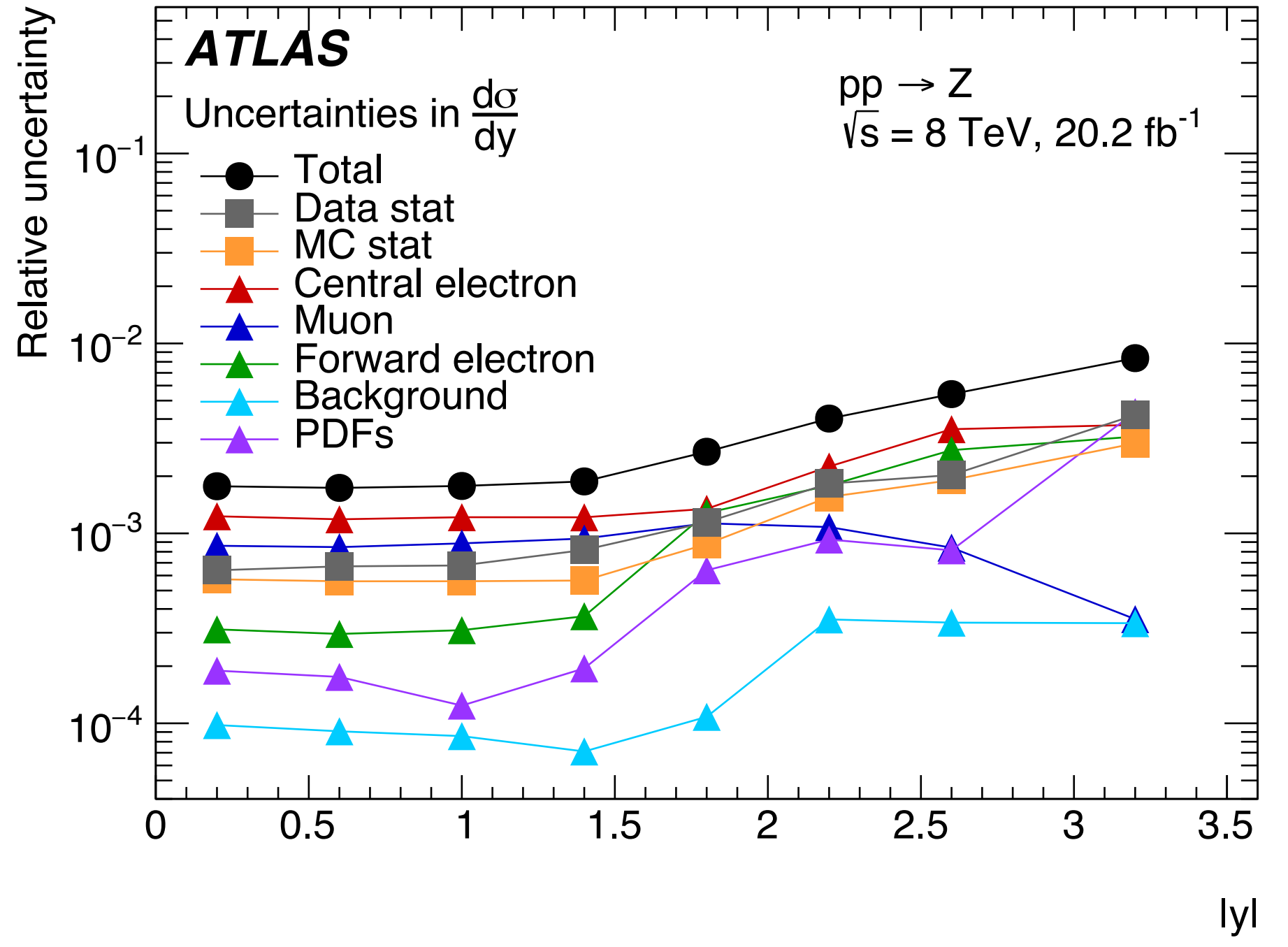
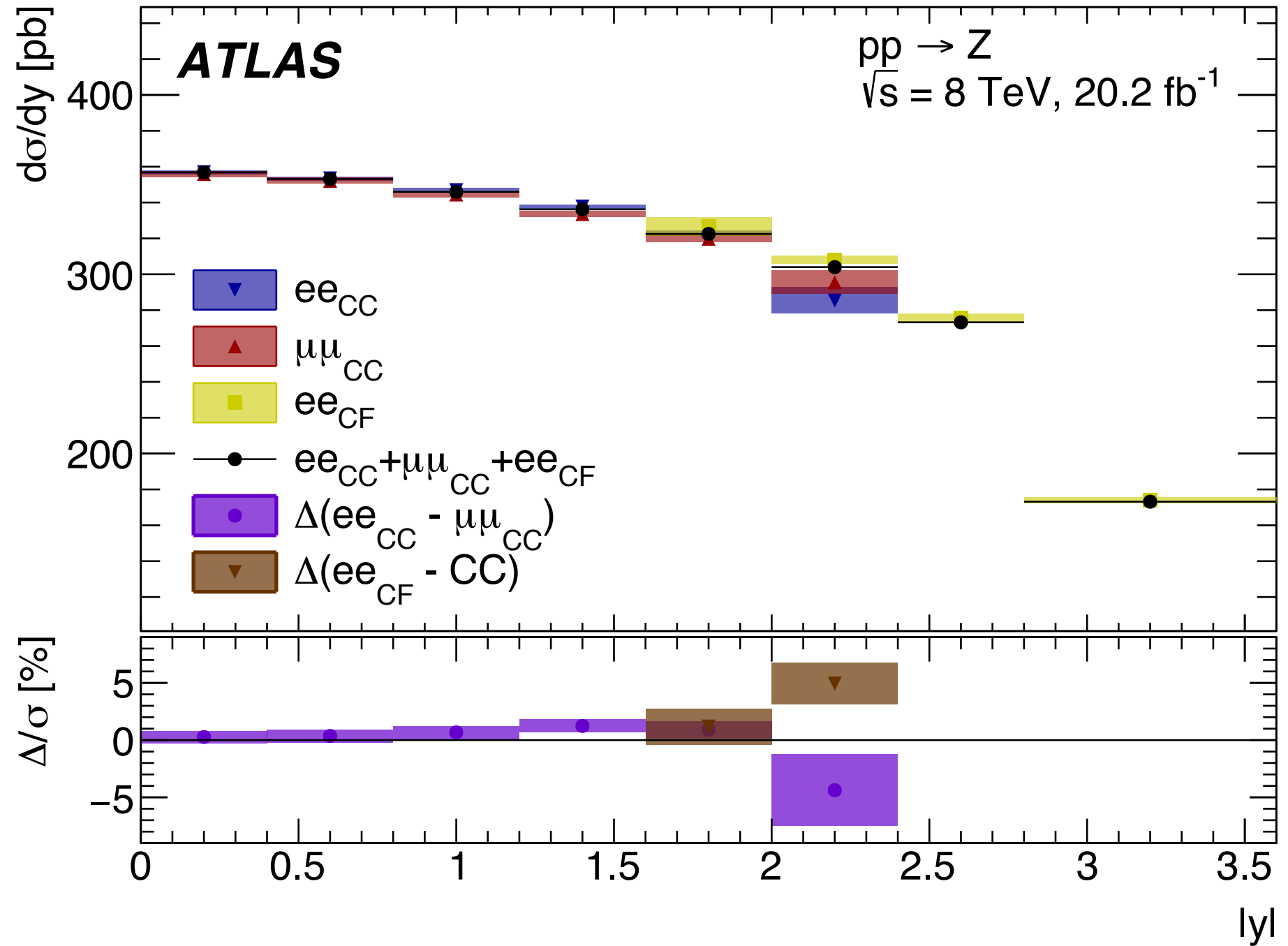
ATLAS

$pp \rightarrow Z$
 $\sqrt{s} = 8 \text{ TeV}, 20.2 \text{ fb}^{-1}$

Uncertainties in $\frac{d\sigma}{dp_T}$

- Total
- Data stat
- MC stat
- ▲ Central electron
- ▲ Muon
- ▲ Forward electron
- ▲ Background
- ▲ PDFs

Comparison of Analysis Channels

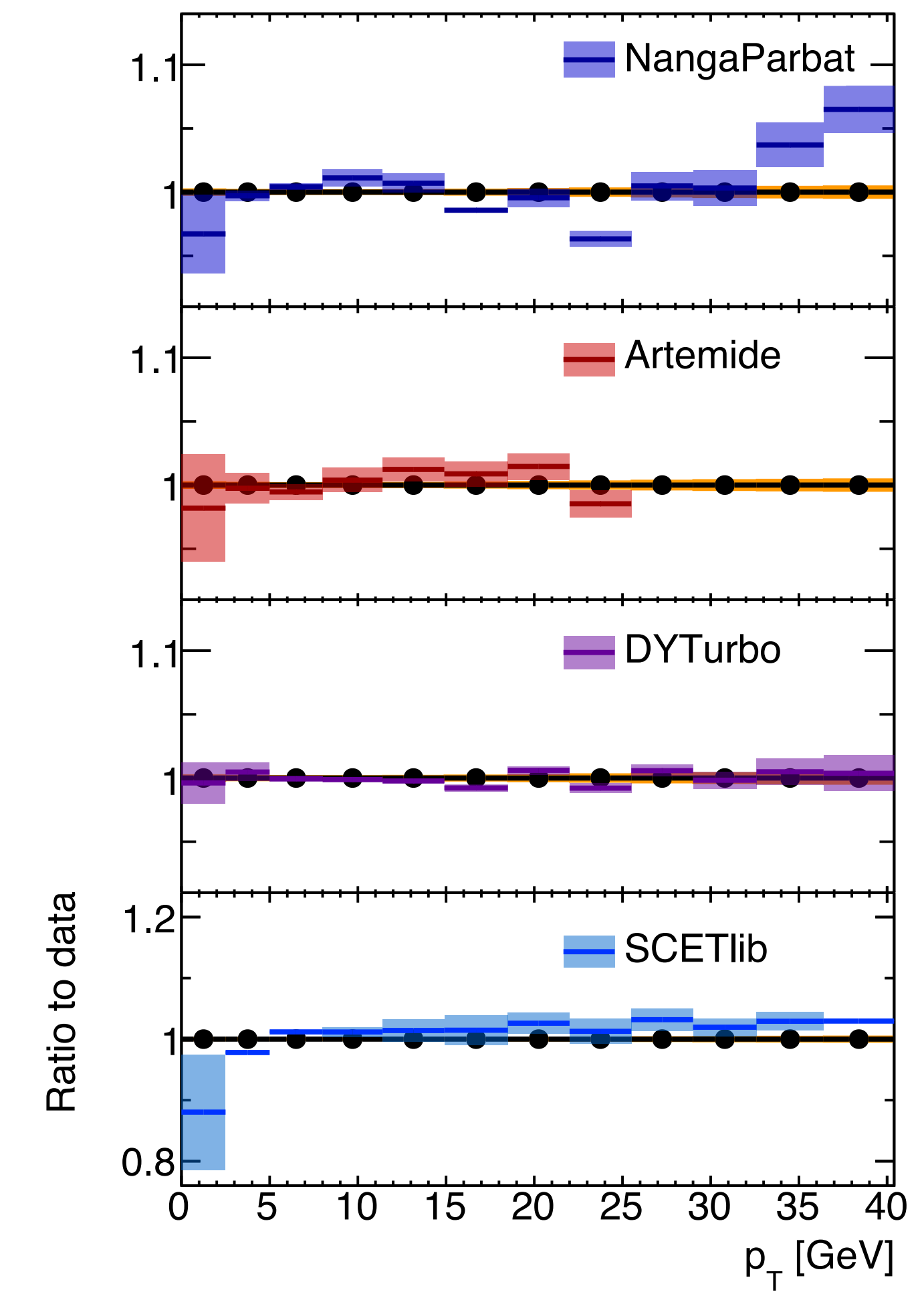
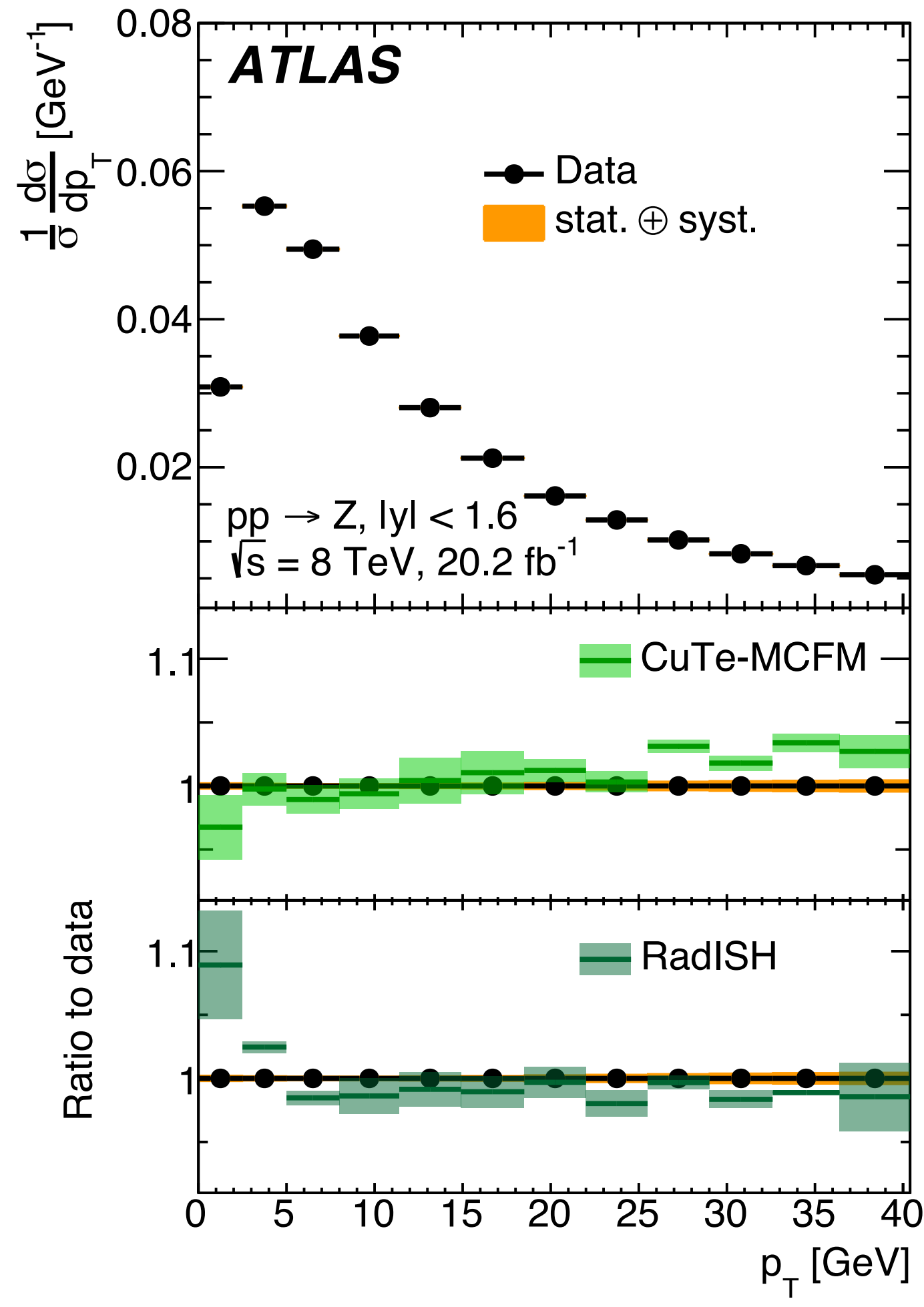


- All three channels (ee_{CC} , $\mu\mu_{CC}$, ee_{CF}) yield compatible results
 - Important cross-check of detector calibration
 - Forward electrons allow to minimize PDF dependence

- Exquisite per-mille level precision in the central region
- Sub-percent uncertainties up to $|y| < 3.6$ thanks to dedicated forward electron calibration

Differential p_T cross-section Measurement

- Measurement compared to predictions currently employing N³LL / N⁴LL logarithmic accuracy calculations
- Including $O(\alpha_s^3)$ matching from MCFM / NNLOJET
- Excellent agreement between data and predictions
- Impressive progress understanding of the boson p_T modelling from experimental and theoretical points of view



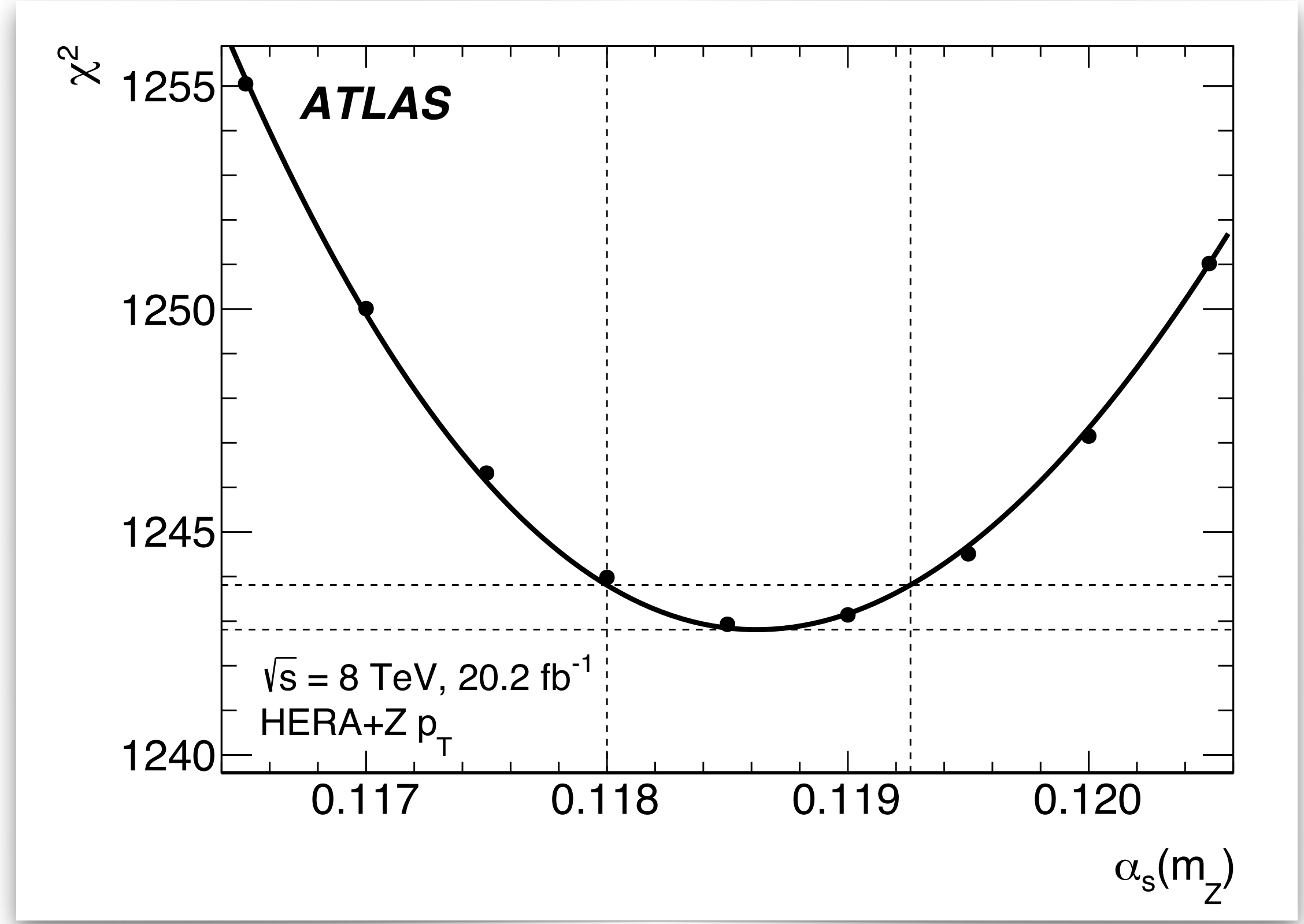
Extraction of the Strong Coupling Constant

Methodology

- Theory expectation: DYTurbo interfaced to xFitter arXiv:1410.4412
- Evaluate $\chi^2(\alpha_s)$ over variations provided in LHAPDF for each α_s point

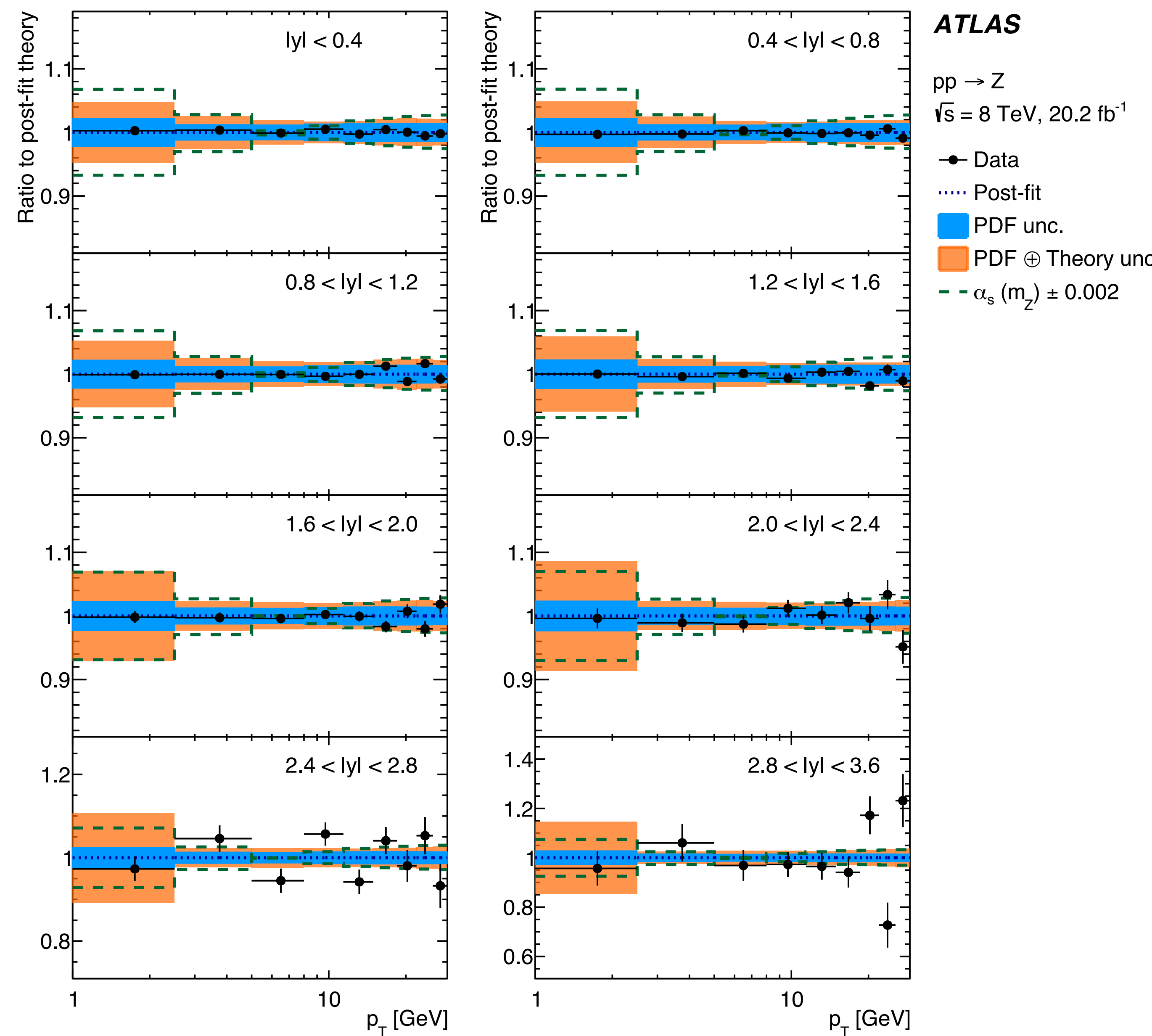
$$\chi^2(\beta_{\text{exp}}, \beta_{\text{th}}) = \sum_{i=1}^{N_{\text{data}}} \frac{\left(\sigma_i^{\text{exp}} + \sum_j \Gamma_{ij}^{\text{exp}} \beta_{j,\text{exp}} - \sigma_i^{\text{th}} - \sum_k \Gamma_{ik}^{\text{th}} \beta_{k,\text{th}} \right)^2}{\Delta_i^2} + \sum_j \beta_{j,\text{exp}}^2 + \sum_k \beta_{k,\text{th}}^2$$

- At each value of $\alpha_s(m_Z)$ the Γ_{ik}^{th} encodes PDF Hessian
- Experimental covariance matrix $\Gamma_{ij}^{\text{exp}} \beta_{j,\text{exp}}$
- Equivalent to including the new dataset in the PDF without refitting, using profiling / reweighting Eur.Phys.J.C 75(2015) 9, 458



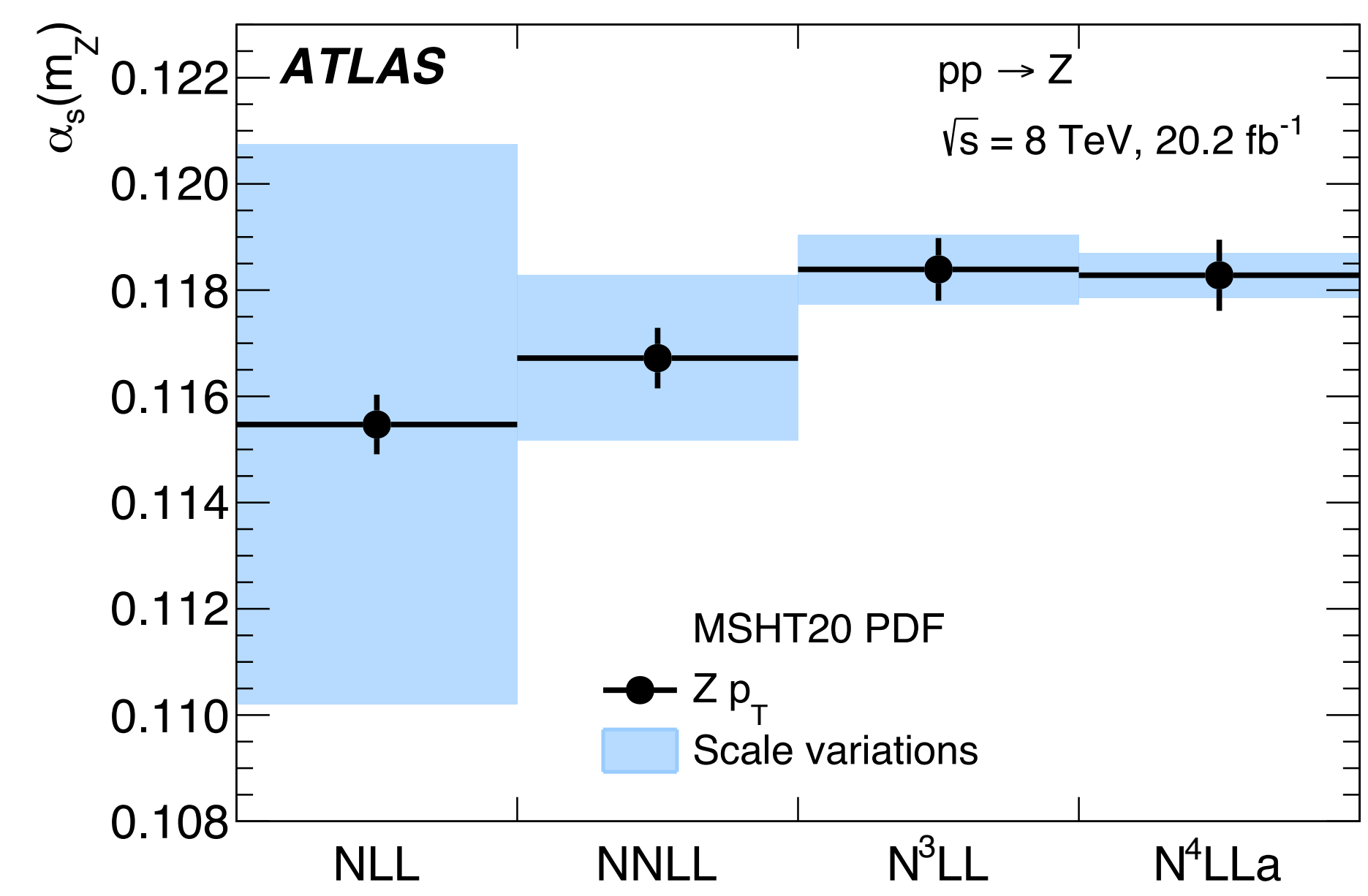
Determination of $\alpha_s(m_Z)$ from $p_T(Z)$ at 8 TeV

- $\alpha_s(m_Z)$ extracted from fit to double-differential p_T - yz cross section measured in the full lepton phase space
- Postfit $\chi^2/\text{dof} = 82/72$
- Determination performed at lower orders, demonstrating good convergence of the perturbative series



Determination of $\alpha_s(m_Z)$ from $p_T(Z)$ at 8 TeV - Uncertainties

- Use MSHT20 PDF (one of two PDFs available at N3LO)
- Repeat fit using lower orders (also with MSHT20)
 - α_s at higher orders always within uncertainties of lower orders
- Scale Variations:
 - Independent variations of renormalisation and factorisation scales and Q variations



| | | |
|---------------------------------|--------------|--------------|
| Experimental uncertainty | ± 0.44 | |
| PDF uncertainty | ± 0.51 | |
| Scale variation uncertainties | ± 0.42 | |
| Matching to fixed order | 0 | -0.08 |
| Non-perturbative model | +0.12 | -0.20 |
| Flavour model | +0.40 | -0.29 |
| QED ISR | ± 0.14 | |
| N ⁴ LL approximation | ± 0.04 | |
| Total | +0.91 | -0.88 |

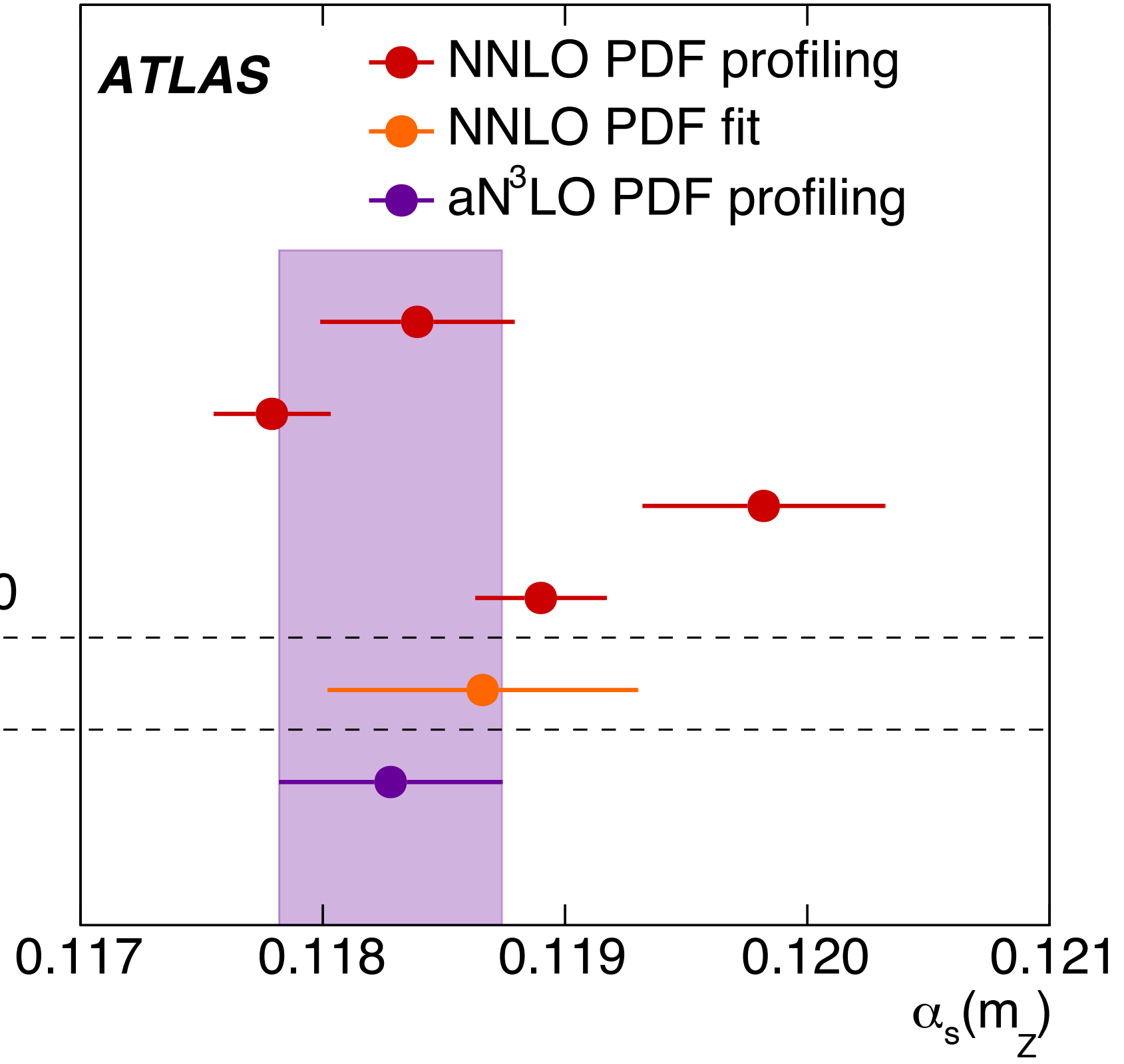
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NNLO MSHT20
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 NNLO HERAPDF2.0

 HERA+Z p_T PDF fit

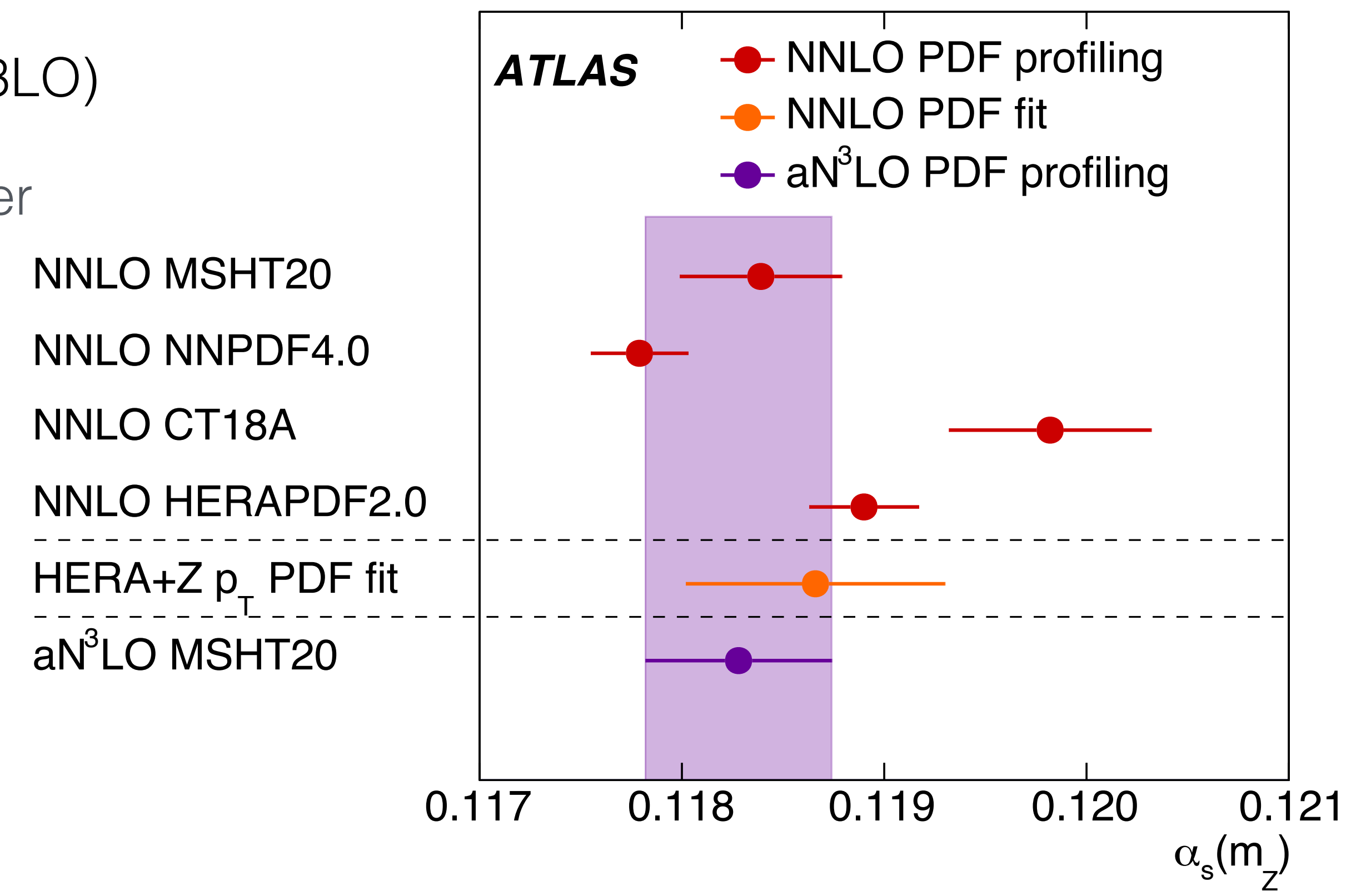
 aN³LO MSHT20



| PDF set | $\alpha_s(m_Z)$ | PDF uncertainty | $g [GeV^2]$ | $q [GeV^4]$ |
|-----------------|-----------------|-----------------|-------------|-------------|
| MSHT20 [37] | 0.11839 | 0.00040 | 0.44 | -0.07 |
| NNPDF4.0 [84] | 0.11779 | 0.00024 | 0.50 | -0.08 |
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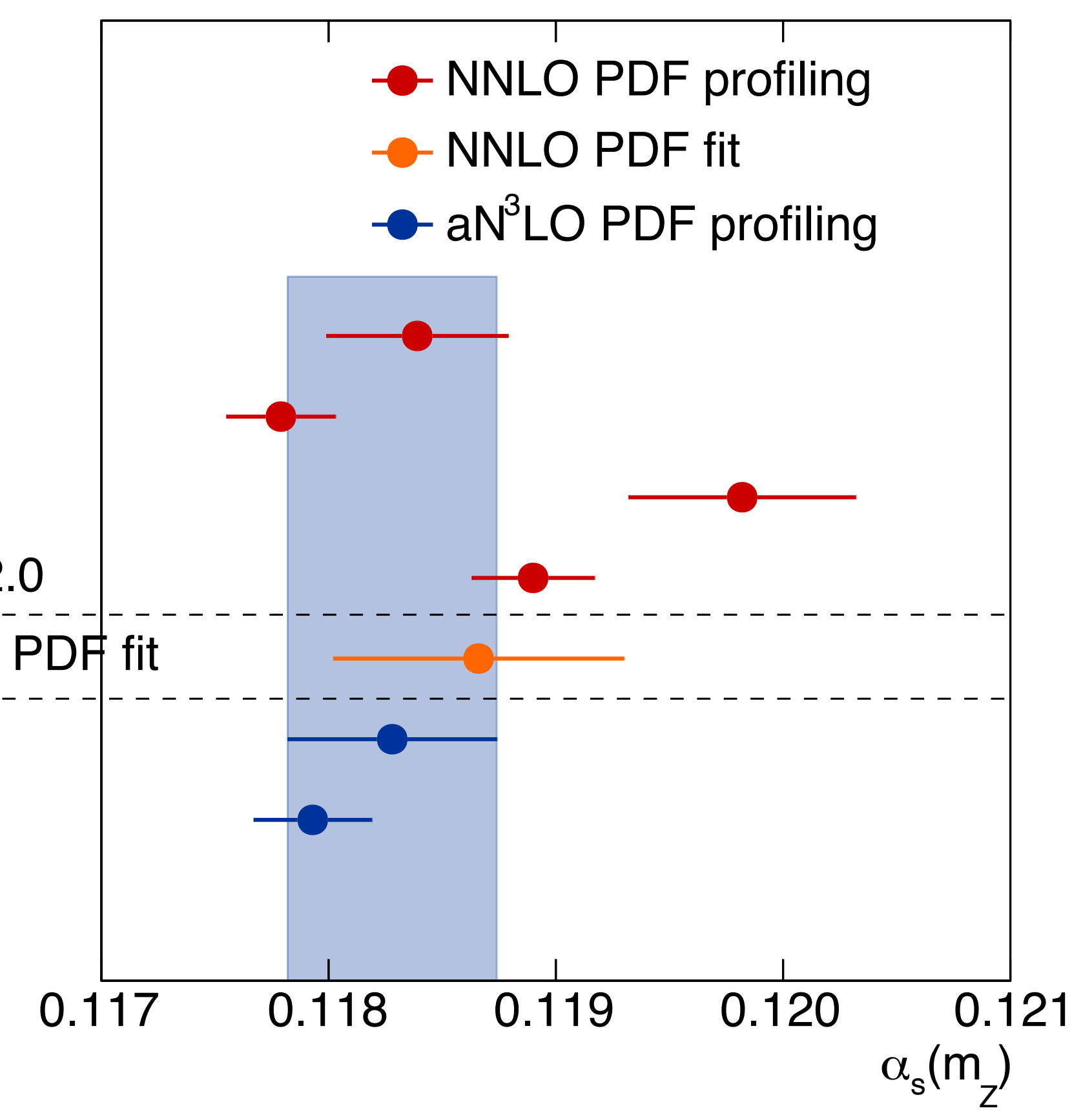


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aN³LO NNPDF4.0

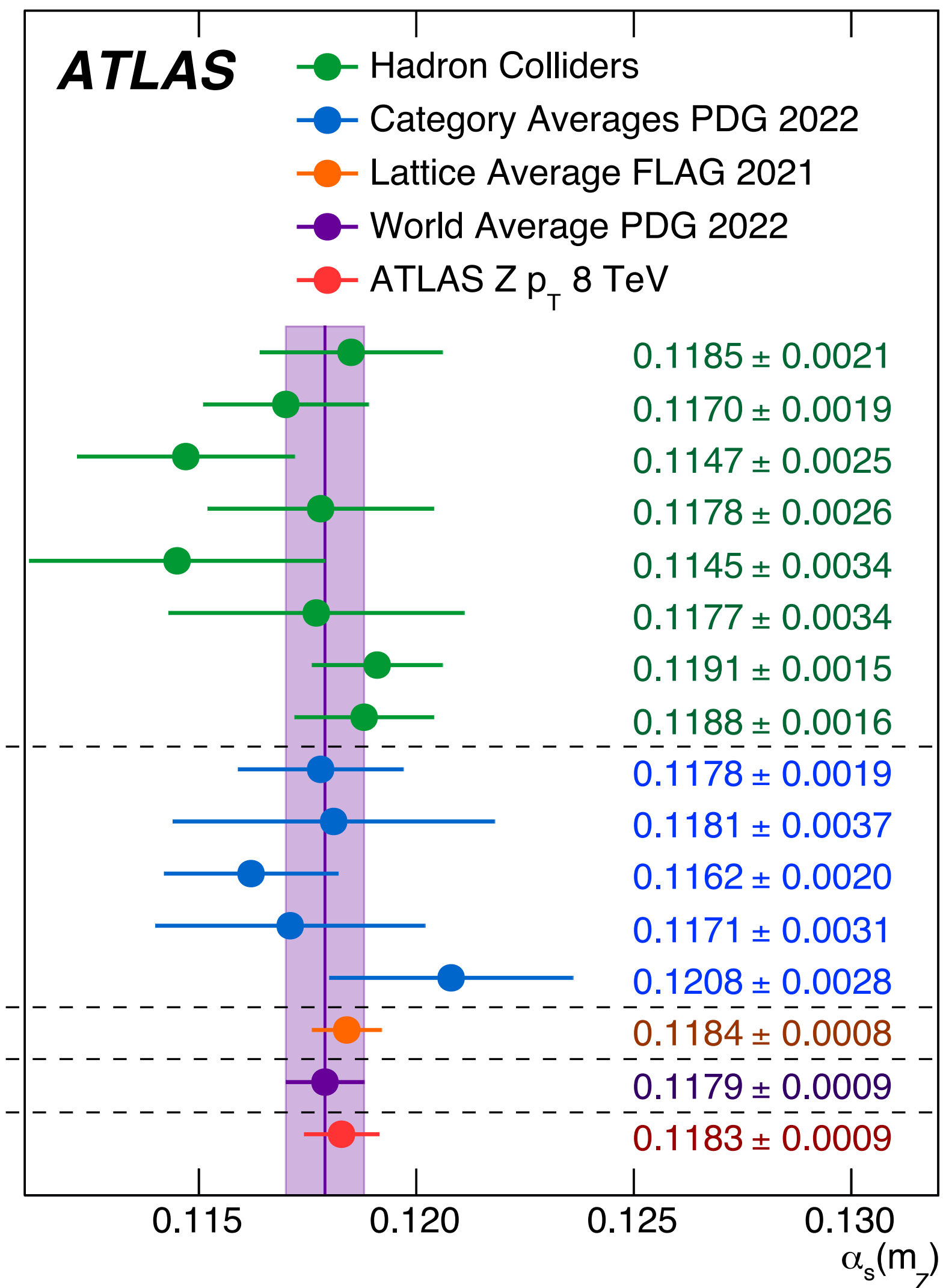


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$$\alpha_s(m_Z) = 0.11828^{+0.00084}_{-0.00088}$$

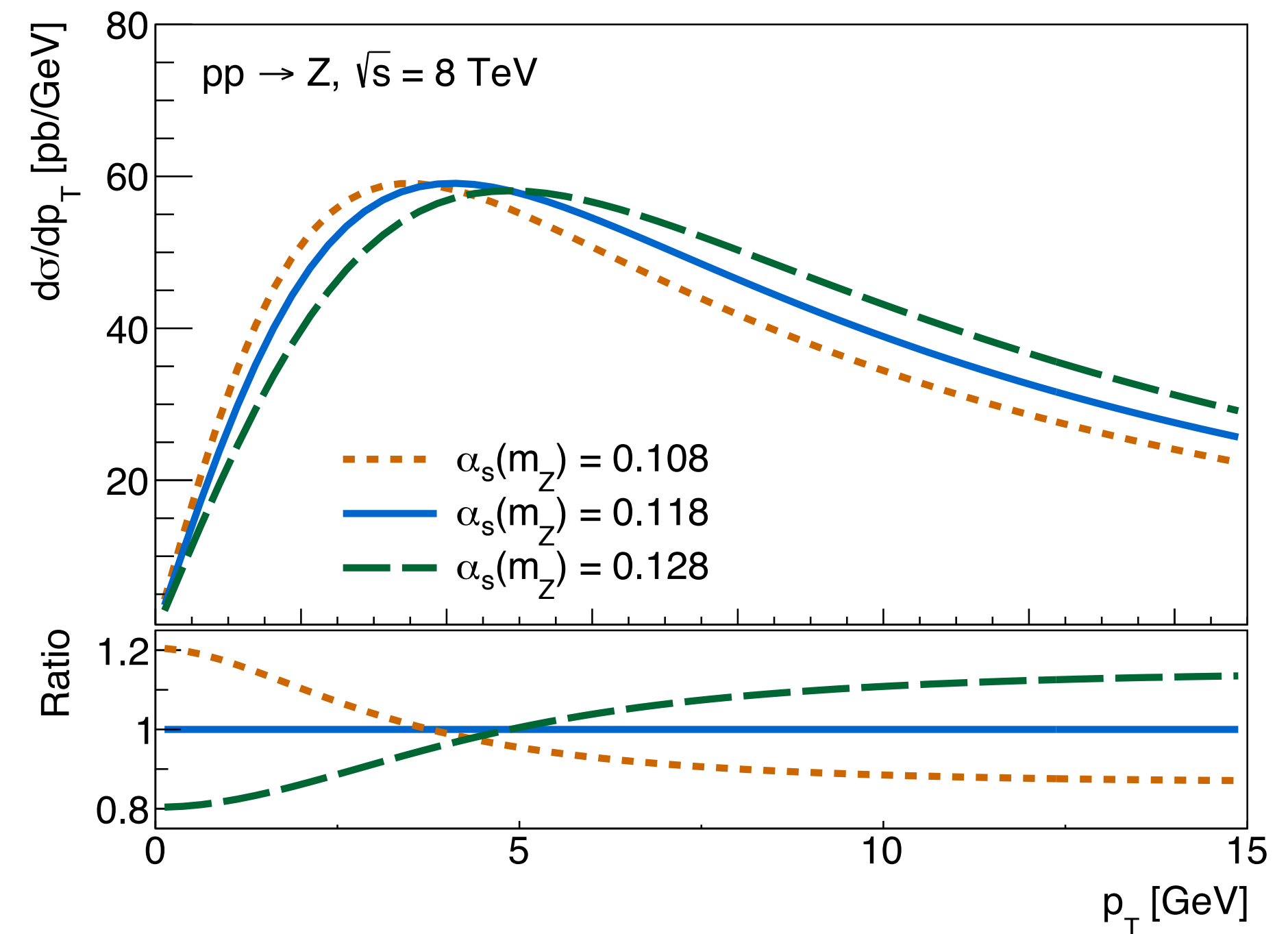
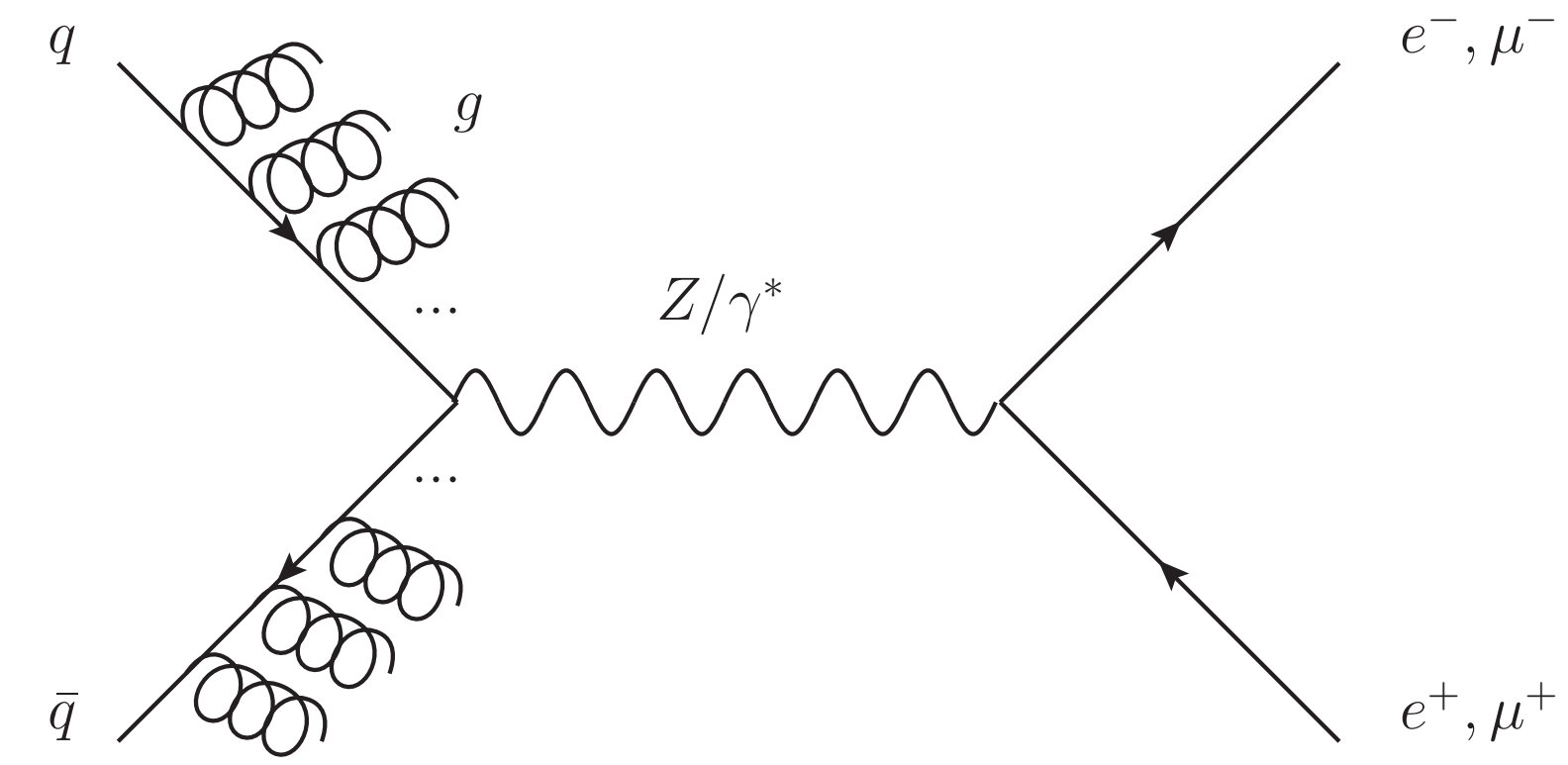
- Most precise experimental determination of $\alpha_s(m_Z)$
 - As precise as PDG and Lattice world averages!
- First $\alpha_s(m_Z)$ determination at N³LO + N⁴LL
- Determination focusing on Sudakov region (usually avoided to determine α_s)
- Observable not suitable for inclusion in PDF fits
 - No correlation with $\alpha_s(m_Z)$ determination from PDF fits

ATLAS ATEEC
 CMS jets
 H1 jets
 HERA jets
 CMS t \bar{t} inclusive
 Tevatron+LHC t \bar{t} inclusive
 CDF Z p $_T$
 Tevatron+LHC W, Z inclusive
 τ decays and low Q²
 Q \bar{Q} bound states
 PDF fits
 e⁺e⁻ jets and shapes
 Electroweak fit
 Lattice
 World average
 ATLAS Z p $_T$ 8 TeV



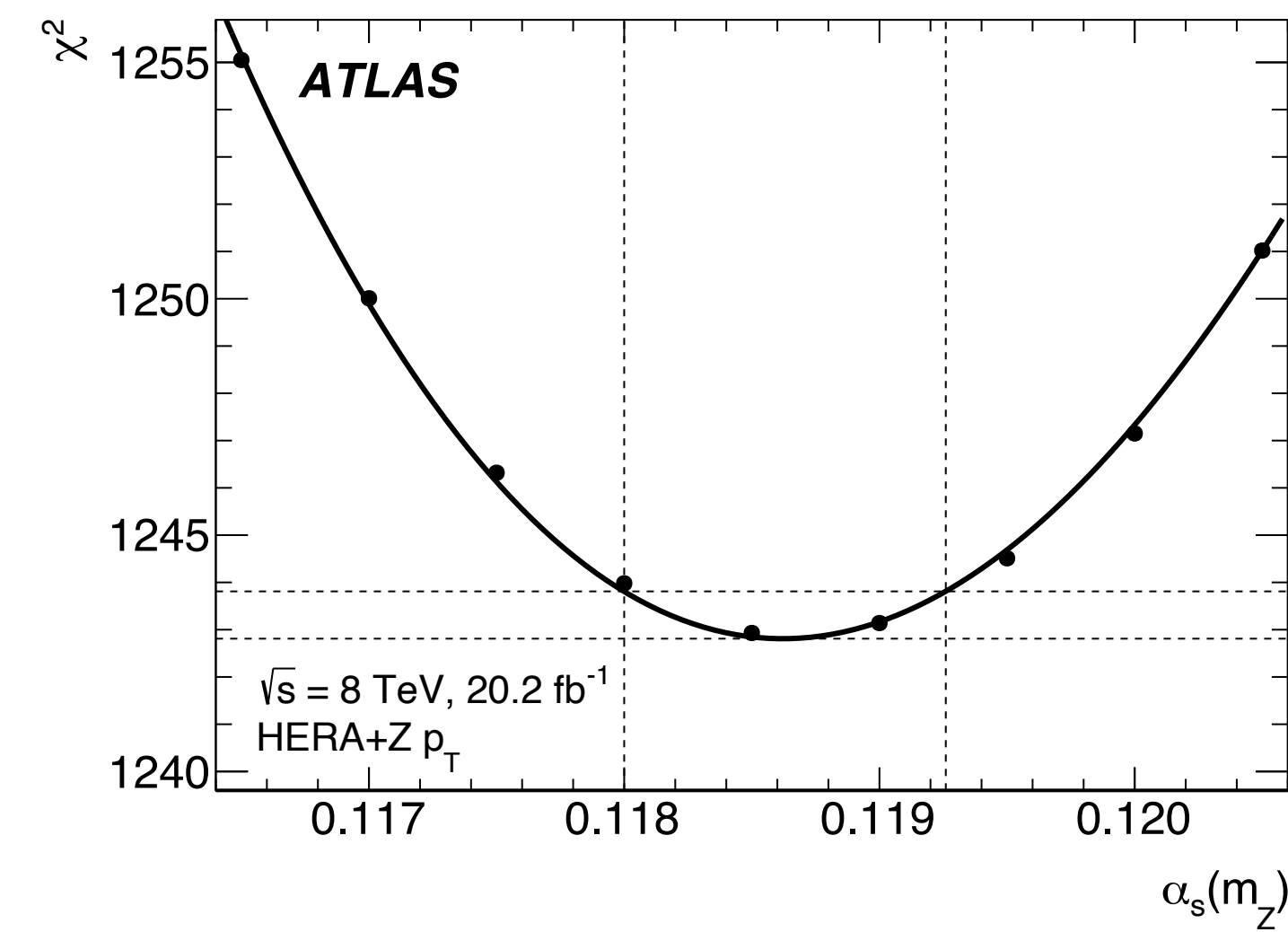
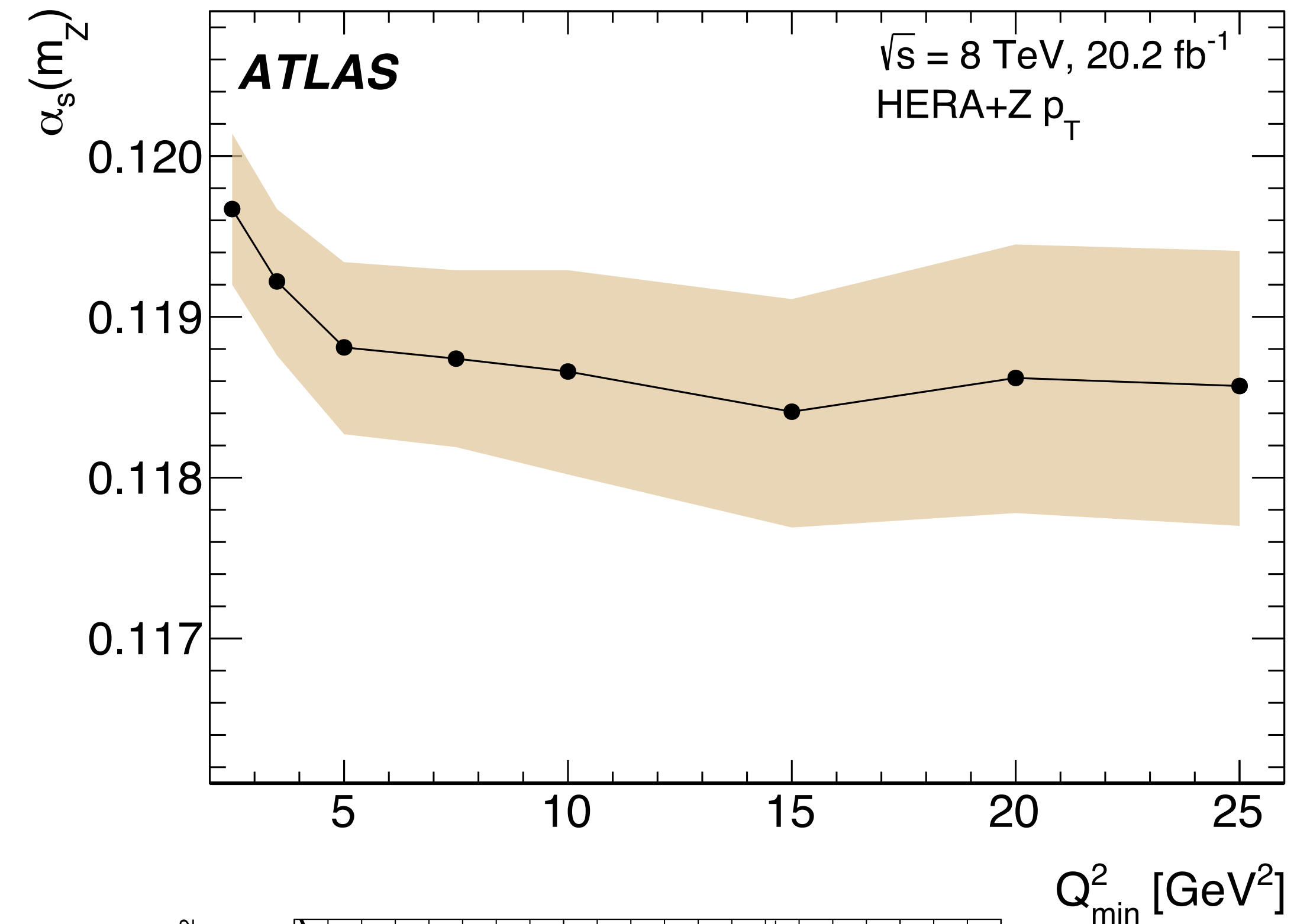
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- New window for the determination of $\alpha_s(m_Z)$, the strong coupling constant
 - Using the transverse momentum of Z bosons
- New measurements might **reduce PDF uncertainties** further
- New measurements required to **constrain further non-perturbative effects**

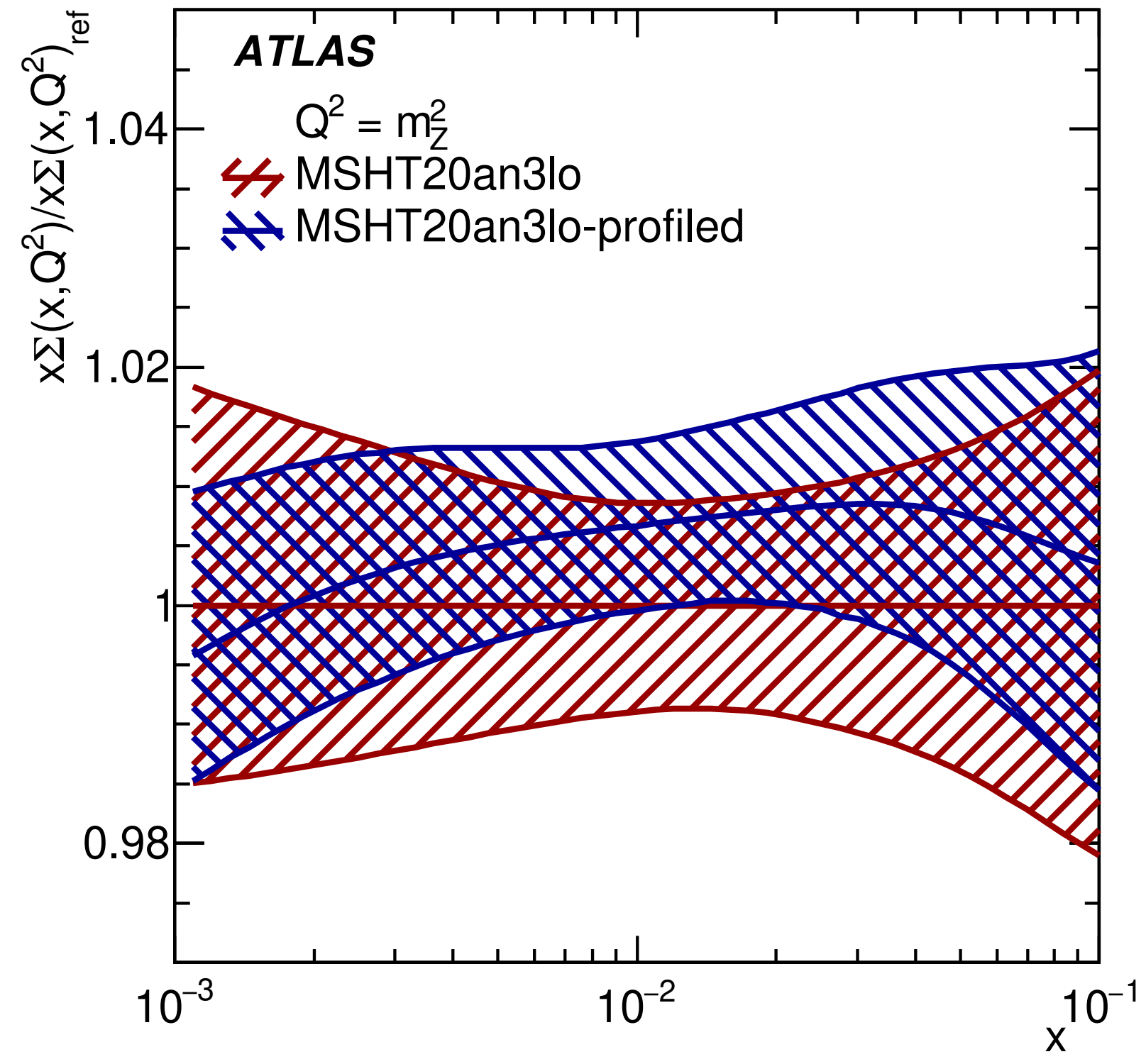
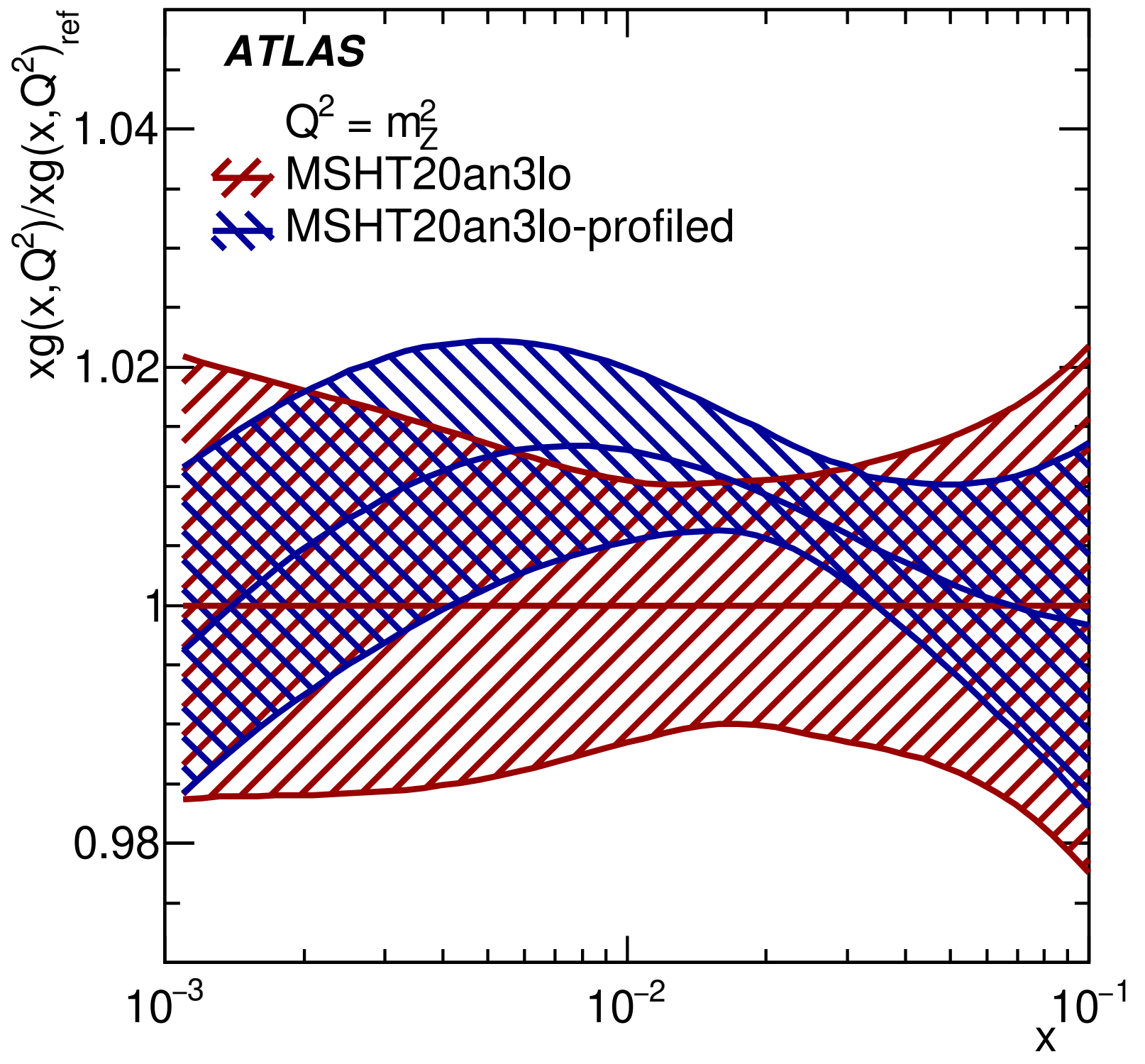


BACKUP

- Simultaneous determination of α_s , the PDFs, and the non-perturbative parameters
 - N3LL+N3LO, with PDFs evolved at NNLO.
- The light-quark coefficient functions of the DIS cross sections are calculated in the \overline{MS} scheme.
- The heavy quarks (c , b) generated dynamically
 - using general-mass variable-flavour-number scheme, with up to five active quark flavours.
- Fits performed at fixed values of α_s via a quadratic interpolation of the χ^2 function
 - 0.11866 ± 0.00064
- The dependence of α_s on the minimum squared four-momentum transfer Q^2 of the HERA data is studied in the range from 2.5 GeV to 25 GeV
 - No sign. dependence is observed for >5 GeV.



- PDF profiling at the best $\alpha_s(m_Z)$ shows reduction of gluon and sea quark PDF uncertainties



Non Perturbative QCD model

- NP model is generally determined from the data, parameters values depend on the chosen prescription to avoid the Landau pole in b-space

$$S_{NP}(b) = \exp \left[-g_j(b) - g_K(b) \log \frac{m_{\ell\ell}^2}{Q_0} \right]$$

$$g_j(b) = \frac{g b^2}{\sqrt{1 + \lambda b^2}} + \text{sign}(q) \left(1 - \exp \left[-|q| b^4 \right] \right)$$

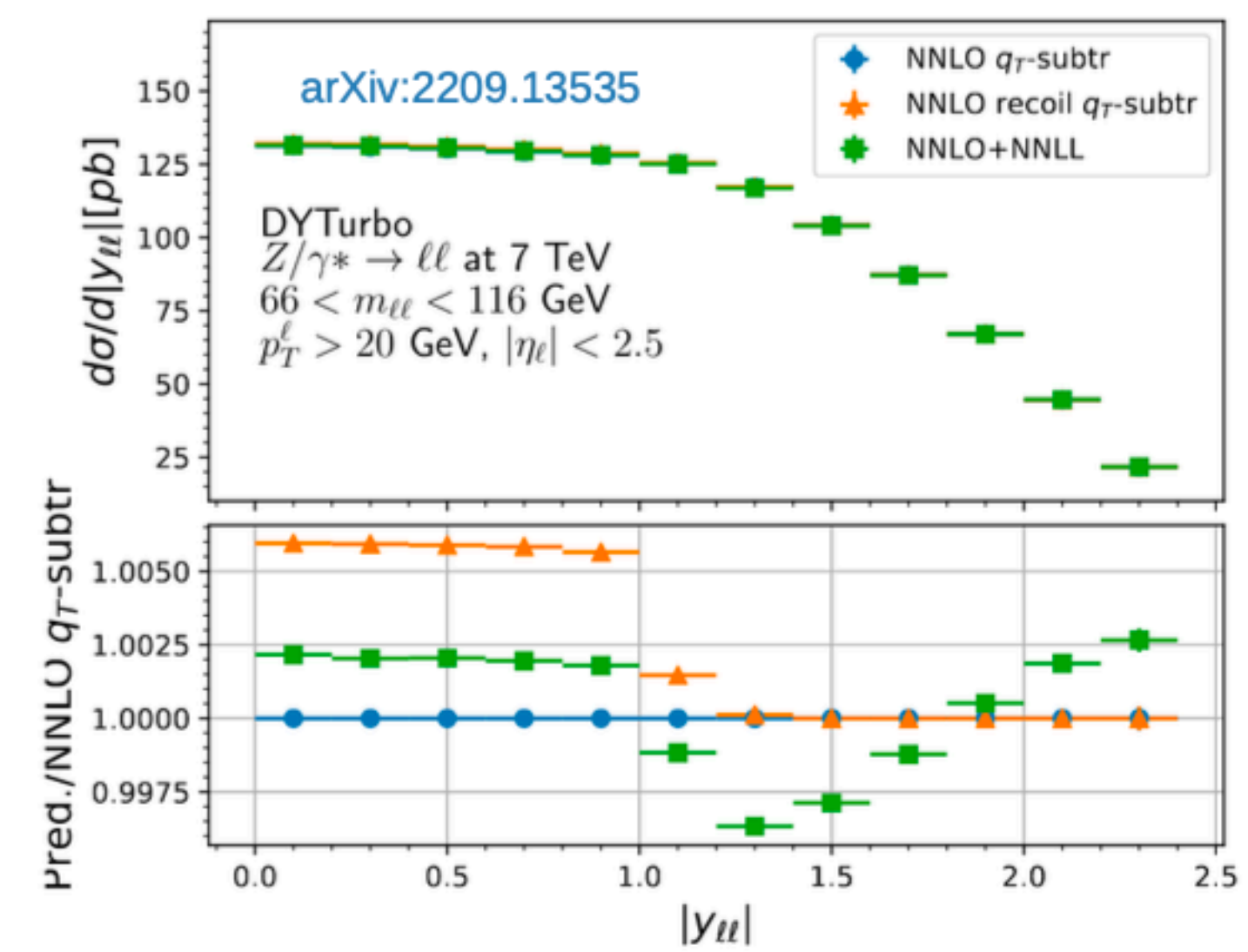
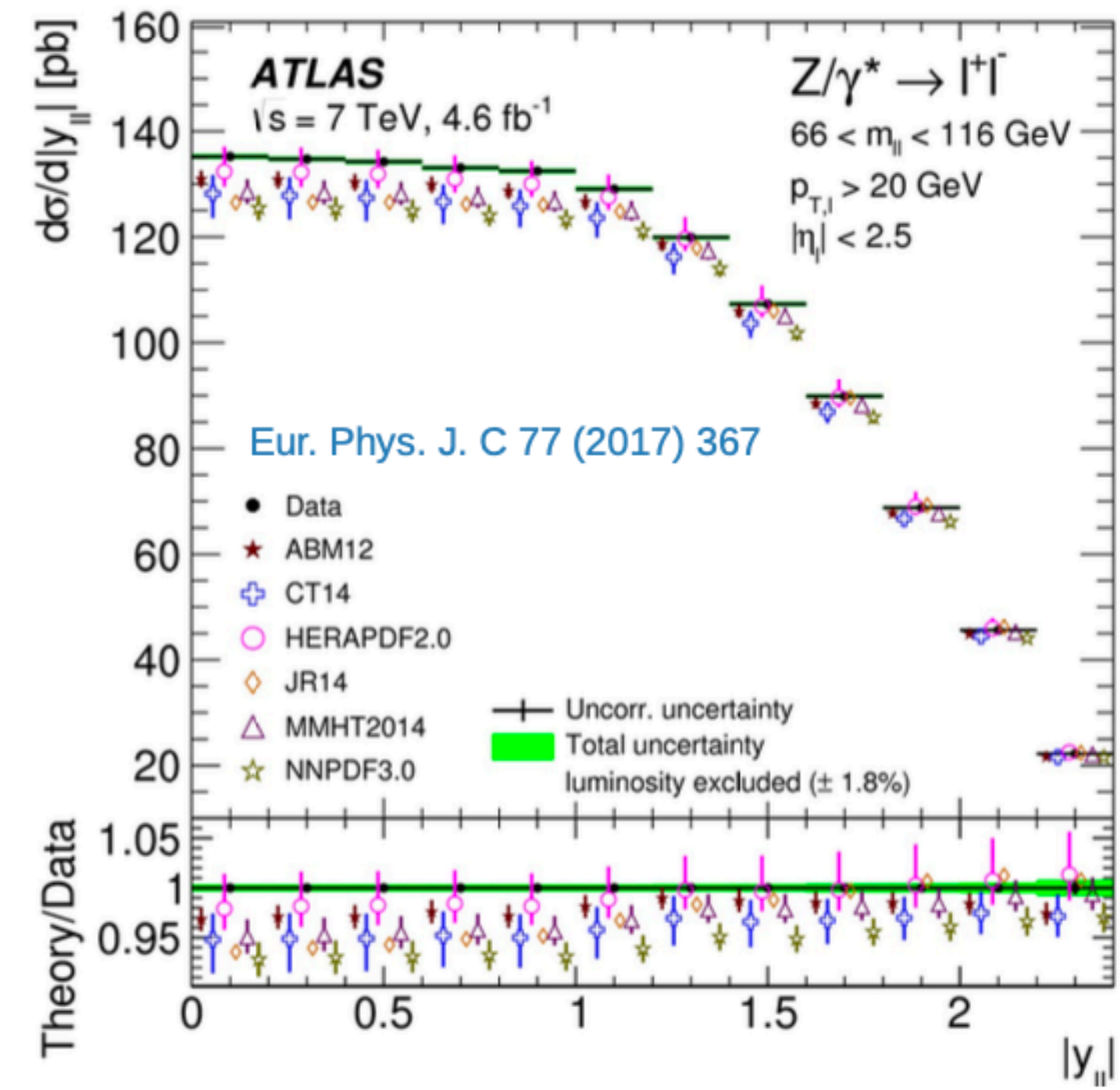
$$g_K(b) = g_0 \left(1 - \exp \left[-\frac{C_F \alpha_s(b_0/b_*) b^2}{\pi g_0 b_{lim}^2} \right] \right)$$

$$b_* = \frac{b}{1 + b^2/b_{lim}^2}$$

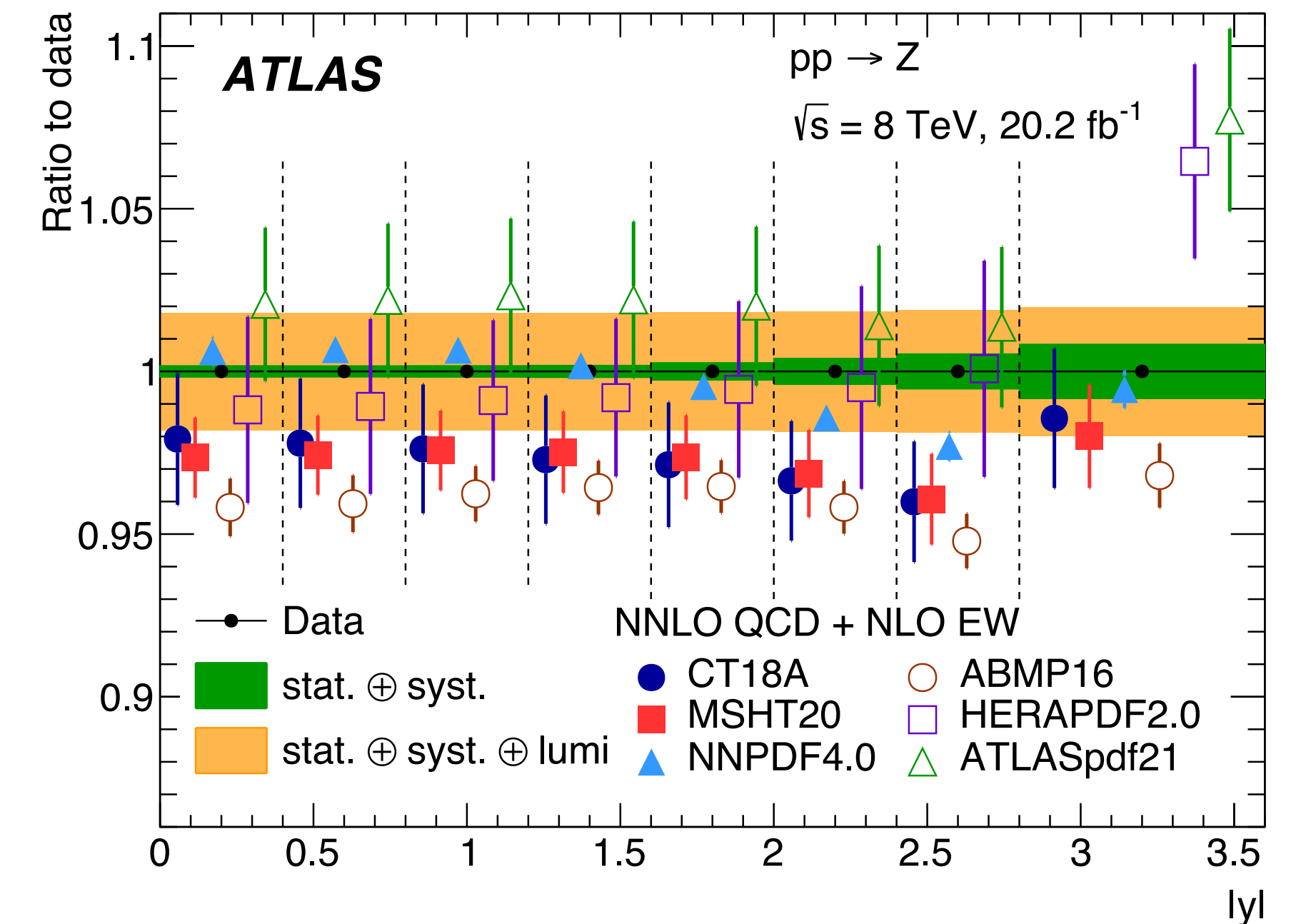
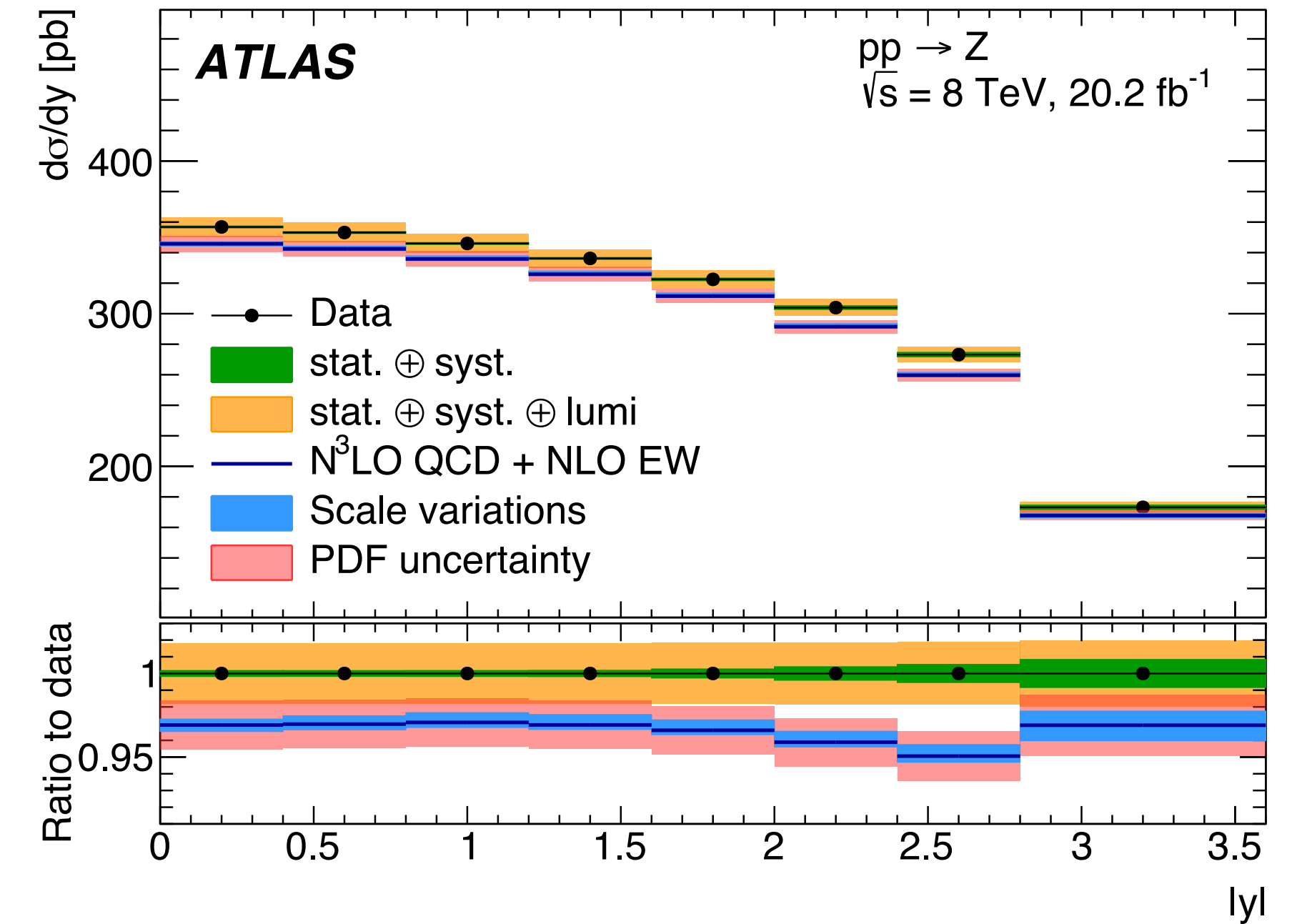
- g_j functions include a quadratic/quartic term: g and q free parameters of the fit
 - The theory should not depend on b_{lim} (freezing scale) and Q₀ (starting scale), provided SNP is flexible enough. Q₀ and b_{lim} estimated as parameterisation unc.
- g₀ controls the very high b (very small p_T) behaviour, should be fitted to data, but we have no sensitivity to it, so it is varied
- Lambda controls transition from Gaussian to exponential: varied between 0.5-2
 - Fits excluding 0-5 GeV yields α_s(m_Z) with a spread of + 0.00017 – 0.00010
 - Fit uncertainty increased from 0.00067 to 0.00071
 - Correlation between α_s(m_Z) and g largely reduced

Full Lepton Phase-Space Rapidity Cross-Section

- Interpretation of fiducial cross sections hampered by breakdown of fixed order perturbation theory
 - Problem: low $p_T(Z)$ spectrum impacts p_T lepton spectrum
- Proposed solutions:
 - Change the definition of fiducial cuts arXiv:2106.08329
 - Use Ai theory predictions to extrapolate the measured cross sections arXiv:2001.02933
 - Include resummation corrections into predictions arXiv:2209.13535 Amoroso et al.
 - All above solutions introduce either experimental or theoretical uncertainties/problems
- Ai-based elegant solution:
 - Fiducial cuts removed by analytic integration of $(\cos\theta, \Phi)$ in the full phase space of the decay leptons through the measured Ai coefficients
 - few permille total uncertainties for $d\sigma/dy$ and negligible theoretical uncertainties

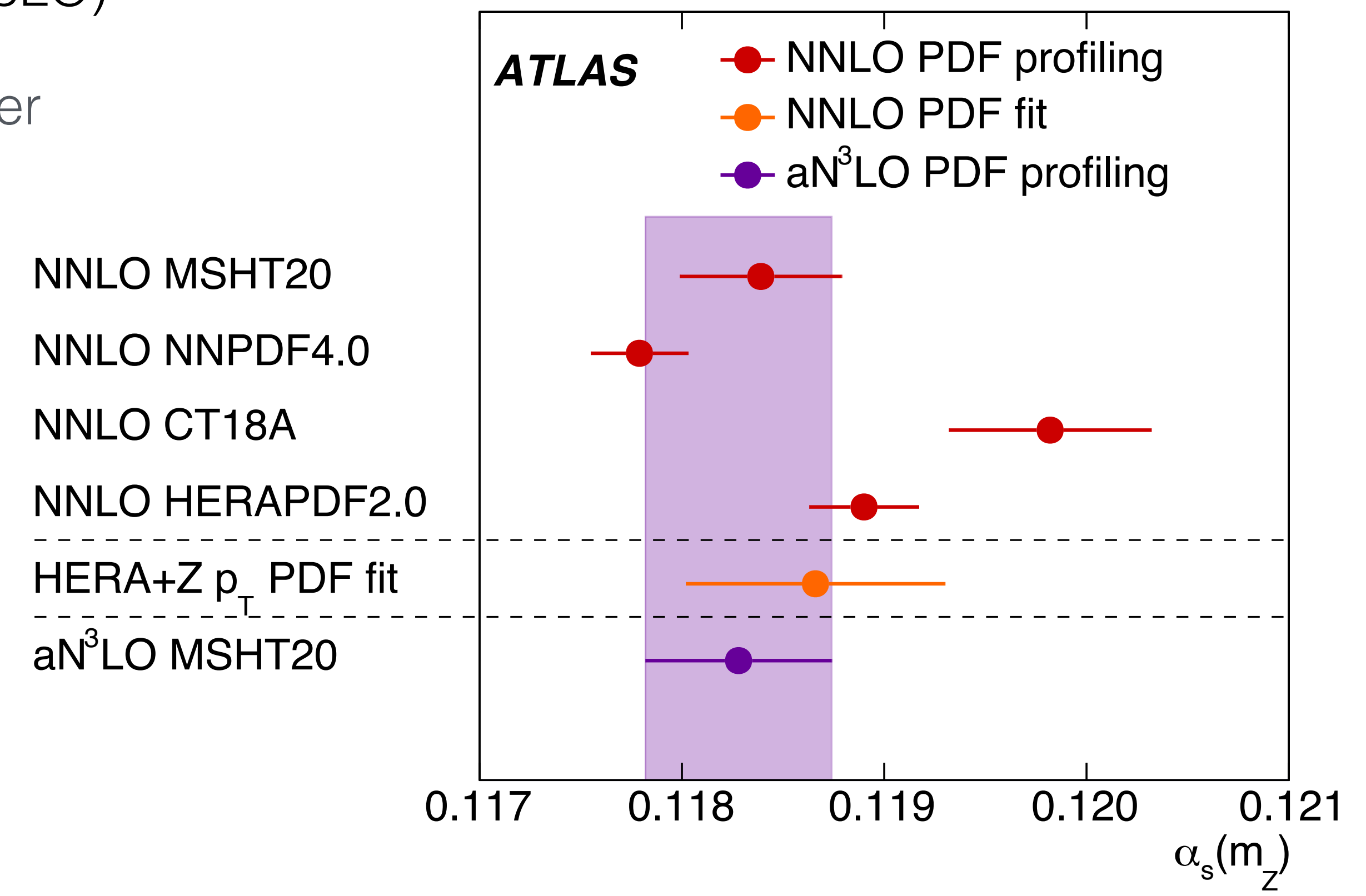


- First comparison to N3LO QCD predictions
- Enables precise and unambiguous PDF interpretation with QCD scale variations now smaller than PDF uncertainties



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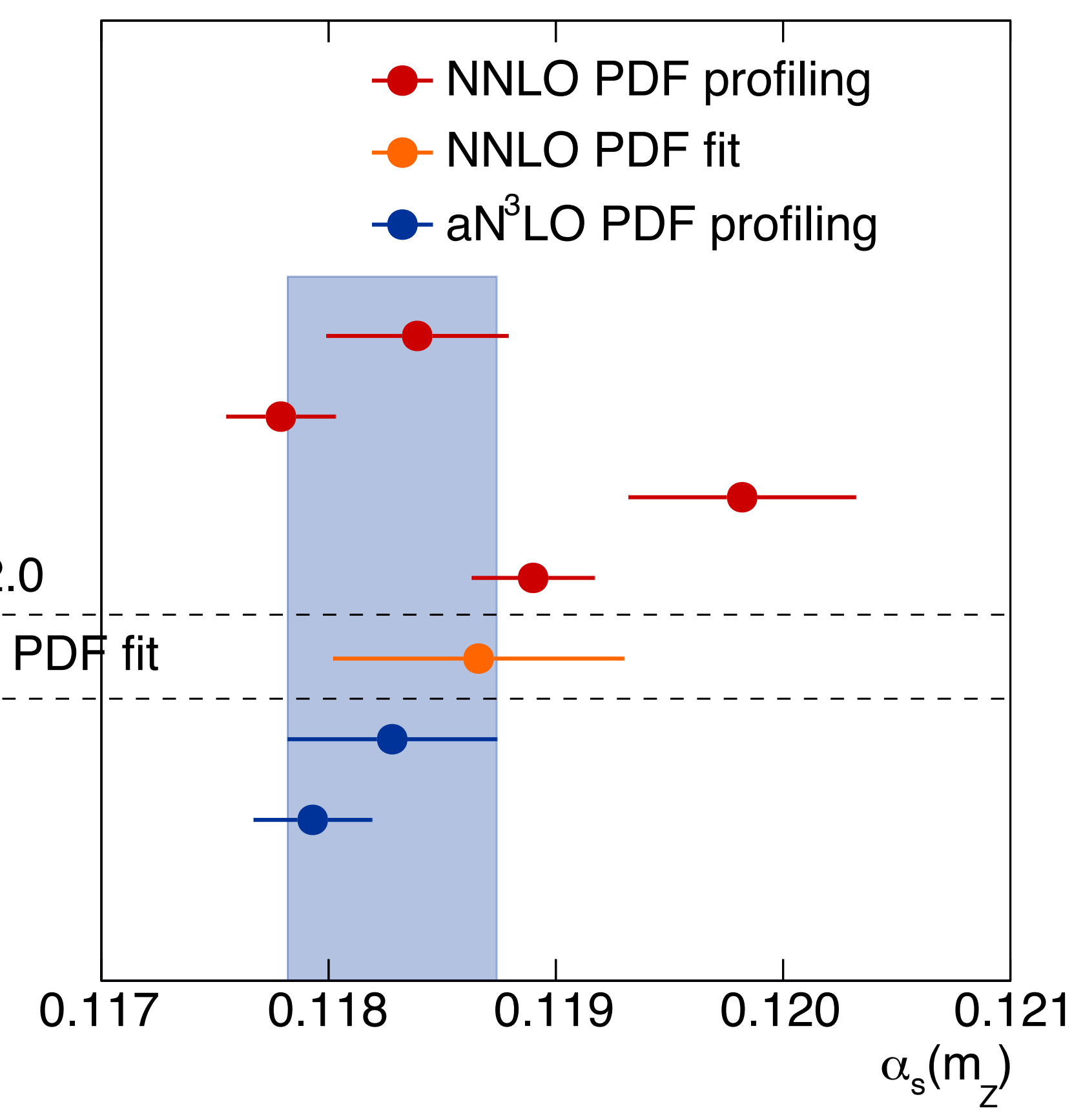


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