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Combination of searches for singly and doubly charged Higgs bosons produced via vector-boson fusion in proton–proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

The ATLAS Collaboration

A combination of searches for singly and doubly charged Higgs bosons, $H^{\pm}$ and $H^{\pm\pm}$, produced via vector-boson fusion is performed using 140 fb^{-1} of proton–proton collisions at a centre-of-mass energy of 13 TeV, collected with the ATLAS detector during Run 2 of the Large Hadron Collider. Searches targeting decays to massive vector bosons in leptonic final states (electrons or muons) are considered. New constraints are reported on the production cross-section times branching fraction for charged Higgs boson masses between 200 GeV and 3000 GeV. The results are interpreted in the context of the Georgi-Machacek model for which the most stringent constraints to date are set for the masses considered in the combination.

Contents

1 Introduction

An important avenue of research for physics beyond the Standard Model is to understand whether the Higgs boson discovered in 2012 [1–4] at the Large Hadron Collider (LHC) is part of an extended Higgs sector. Charged Higgs bosons are predicted in extended Higgs sectors with additional complex doublets [5, 6] or with additional higher-isospin scalar fields [7–9]. In the generic two Higgs doublet model the charged Higgs boson $H^{\pm}$ does not decay into WZ bosons as a result of CP-invariance which forbids a tree-level $H^{\pm}W^{\pm}Z$ coupling. A tree-level coupling to massive vector bosons, is, however, present in models with additional isotriplet scalar fields [10].

The Georgi–Machacek (GM) model [11, 12] is used as a benchmark in this Letter. The GM model extends the Higgs sector of the SM by including one real and one complex triplet. This preserves a custodial symmetry at tree level whereby the GM model is not strongly constrained [13]. A parameter, $\sin\theta_H$, characterises the contribution of the isotriplet scalar fields to the masses of the W and Z bosons. The physical scalar states are organized into distinct custodial multiplets: a quintuplet ($H_5^{\pm\pm}, H_5^{\pm}, H_5^{0}$) that is fermiophobic but couples to W and Z bosons, a triplet, and two singlets, one of which is identified as the observed 125 GeV SM-like Higgs boson. The physical states in each multiplet are degenerate in mass (denoted as m_{H_5} for the quintuplet) at tree level. Mass splittlings of up to a few GeV in the quintuplet due to higher order effects [14] are not considered in this analysis.

The production of the $H_5^{\pm}$ and $H_5^{\pm\pm}$ scalars is via vector-boson fusion (VBF) in the GM model. Figure 1 shows representative Feynman diagrams of production and decay of the $H_5^{\pm}$ and $H_5^{\pm\pm}$ states. The `H5plane` benchmark is considered [15], where the triplet states are assumed to be heavier than the quintuplet states. Thus, in this benchmark, the branching fraction of $H_5^{\pm} \rightarrow W^{\pm}Z$ and $H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$ decays is 100% [16]. The production cross-sections and the widths of the $H_5^{\pm\pm}$ and $H_5^{\pm}$ states are proportional to $\sin^2\theta_H$. For the parameter space explored in this combination, the considered intrinsic width of the $H_5^{\pm\pm}$ and $H_5^{\pm}$ states is below 5%, which is below the experimental resolution.

This Letter reports the combination of the ATLAS Collaboration searches for $H_5^{\pm} \rightarrow W^{\pm}Z$ [17] and $H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$ [18] produced via VBF using proton–proton (pp) collisions at $\sqrt{s}$ = 13 TeV. The dataset corresponds to an integrated luminosity of 140.1 ± 1.2 fb^{-1} [19, 20], collected with the ATLAS detector [21] during Run 2 of the LHC (2015–2018). An extensive software suite [22] is used in data simulation, in the reconstruction and analysis of real and simulated data, in detector operations, and in the trigger and data acquisition systems of the experiment. The searches target events where the vector boson decays include electrons or muons. A simultaneous search for the $H_5^{\pm}$ and $H_5^{\pm\pm}$ states in the VBF topology using

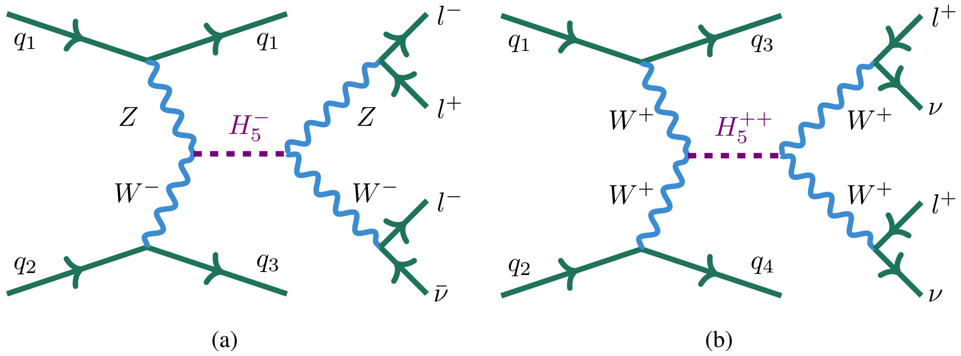


Figure 1: Representative Feynman diagrams of production and decay of the (a) $H_5^{\pm}$ and (b) $H_5^{\pm\pm}$ states.

the same leptonic decay modes was published by the CMS Collaboration [23], reporting upper bounds at 95% confidence level (CL) on the $\sin\theta_H$ parameter. The reported CMS upper bounds vary between ~ 0.2 and 0.55 in the mass range 200–2000 GeV.

Constraints on the GM model have also been reported by the ATLAS Collaboration considering the VBF production of the H_5^0 state and the $W^{\pm}W^{\mp}$ boson pair decay channel [24] at $\sqrt{s} = 13$ TeV. In addition, the ATLAS and CMS Collaborations have reported constraints on the GM model at $\sqrt{s} = 8$ TeV [25] and $\sqrt{s} = 13$ TeV [26], respectively, via searches for $H_5^{\pm} \to W^{\pm}Z$ and $H_5^{\pm\pm} \to W^{\pm}W^{\pm}$ using semileptonic final states. The results from Refs. [24, 25] are not included in this combination due to their limited sensitivity. Constraints from searches for heavy neutral Higgs bosons produced via VBF and decaying into WW or ZZ channels can also be interpreted in the GM model as discussed in Ref. [27], but a combination of all these searches is beyond the scope of this Letter.

2 Description of the nominal analyses

A brief overview of the $H_5^{\pm} \to W^{\pm}Z$ and $H_5^{\pm\pm} \to W^{\pm}W^{\pm}$ searches is given below. The detailed information about the reconstruction, identification and calibration of physics objects, as well as the simulation, triggers and event selection used in these results is given in the corresponding Refs. [17, 18]. The final states probed by these analyses consist in the leptonic decays (electrons or muons) of two massive vector bosons produced in association with two jets, denoted as $VVjj$. The event selection in the $H^{\pm}$ signal region (SR) requires three charged leptons, while the $H^{\pm\pm}$ SR requires a same-charge lepton pair. Requirements on the missing transverse momentum, with magnitude $E_{\mathrm{T}}^{\mathrm{miss}}$, in both channels exploit the presence of neutrinos in the final states. The VBF topology is characterised by requiring at least two jets with a large invariant mass, m_{jj}, and a large rapidity difference, $|\Delta y_{\mathrm{jj}}|$. The invariant mass of the two highest transverse momentum jets must satisfy $m_{\mathrm{jj}} > 100$ GeV in the $H^{\pm}$ SR and is further constrained as discussed below. The $H^{\pm\pm}$ SR is defined by requiring $m_{\mathrm{jj}} > 500$ GeV and $|\Delta y_{\mathrm{jj}}| > 2.0$.

The presence of only one neutrino in the $H_5^{\pm} \to W^{\pm}Z$ channel is exploited to estimate the longitudinal component of the neutrino momentum by constraining the invariant mass of the charged lepton and neutrino system to the pole mass of the W boson, where the charged lepton is the one assigned to the W boson candidate. The $E_{\mathrm{T}}^{\mathrm{miss}}$ is assumed to be due to the neutrino. The resulting quadratic equation leads to two solutions. If they are real, the one with the smaller magnitude of the neutrino momentum is

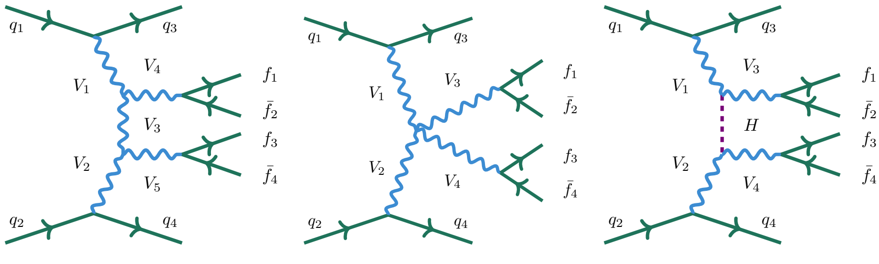


Figure 2: Representative Feynman diagrams for a VBS EW $VVjj$ production that either include (a) a triple-gauge-boson vertex, (b) a quartic gauge boson vertex, and (c) the exchange of a Higgs boson in the t-channel. The lines are labelled by quarks (q), vector bosons ($V = W/Z$), the Higgs boson (H) and fermions (f).

chosen, otherwise, the real part is chosen. The choice of the solution was optimised using generator-level information. The resulting reconstructed W boson four-momentum is used to calculate the invariant mass of the WZ system, m_{WZ}, which is used as a discriminating variable between the resonant $H^{\pm}$ signal and the SM backgrounds. The full kinematic reconstruction of the invariant mass of the $W^{\pm}W^{\pm}$ system is not attempted in the $H^{\pm\pm}$ SR due to the presence of two neutrinos. The transverse mass, m_{T}, defined as

$$m_{\mathrm{T}} = \sqrt{\left(E_{\mathrm{T}}^{\ell\ell} + E_{\mathrm{T}}^{\mathrm{miss}}\right)^2 - \left|\vec{p}_{\mathrm{T}}^{\,\ell\ell} + \vec{E}_{\mathrm{T}}^{\mathrm{miss}}\right|^2},$$

where $E_{\mathrm{T}}^{\ell\ell}$ is the transverse energy of the dilepton system, $\vec{p}_{\mathrm{T}}^{\,\ell\ell}$ is the vectorial sum of the lepton transverse momenta, and $\vec{E}_{\mathrm{T}}^{\mathrm{miss}}$ is the missing transverse momentum vector, is used as a discriminating variable in the $H^{\pm\pm}$ SR.

The SM production of the $VVjj$ final state constitutes the dominant background for these searches. The production of $VVjj$ at leading-order (LO) has contributions both from modes that involve only electroweak (EW) interaction vertices (EW $VVjj$) and from modes that involve strong interaction vertices (QCD $VVjj$). Representative Feynman diagrams for the vector-boson scattering (VBS) processes, which are part of a class of processes contributing to the EW $VVjj$ production, are shown in Figure 2. Representative Feynman diagrams for the QCD $VVjj$ processes are shown in Figure 3. The QCD and EW $W^{\pm}Zjj$, and EW $W^{\pm}W^{\pm}jj$ processes are the dominant backgrounds in the $H^{\pm}$ and $H^{\pm\pm}$ SRs, respectively. The QCD and EW $W^{\pm}Zjj$ processes are important backgrounds also in the $H^{\pm\pm}$ SR, contributing 22% of the overall expected event yield. It contributes when one of the leptons is not selected, typically because it is outside of the geometrical acceptance of the detector.

Leptons from hadron decays and jets misidentified as leptons are referred to as non-prompt leptons. The non-prompt lepton background is the third-largest background process in the $H^{\pm\pm}$ SR and arises mainly from W+jets and semileptonic $t\bar{t}$ processes. The non-prompt lepton, electron charge misidentification and photon conversion backgrounds are estimated using data-driven methods as described in Ref. [18]. The contribution of the ZZ process is a non-negligible background in the $H^{\pm}$ SR. Small backgrounds contributions from triboson VVV, $t\bar{t}V$, and tZq are also considered.

The dominant background processes are estimated with Monte Carlo (MC) simulated events and their modelling is constrained in dedicated signal-depleted control regions (CRs). An artificial neural network (ANN) with a binary classification task is used for the $H^{\pm}$ search to categorise the events as belonging to either the VBF process or to a background process. The $H^{\pm}$ SR is defined by requiring the ANN score to

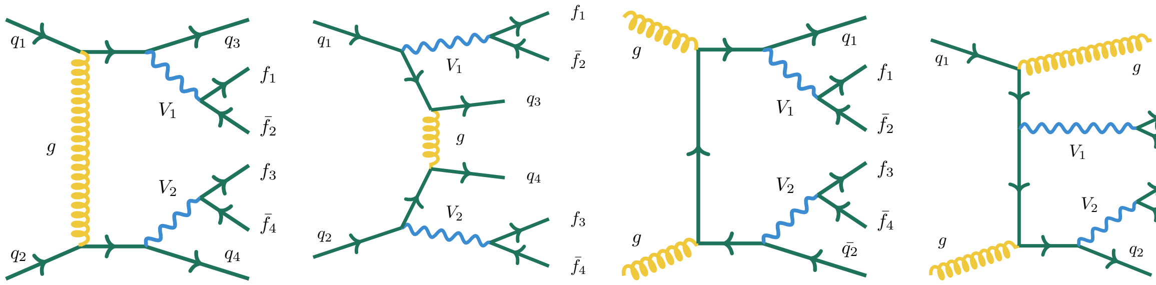


Figure 3: Representative Feynman diagrams for QCD $VVjj$ production with strong interaction vertices. The lines are labelled by quarks (q), vector bosons ($V = W/Z$), fermions (f), and gluons (g). The two diagrams on the right with gluons in the initial state are not present for $W^{\pm}W^{\pm}jj$ production.

be greater than 0.82, which maximizes the significance and effectively starts the SR at $m_{\mathrm{jj}} > 500$ GeV. Events with an ANN score of less than 0.82 and $m_{\mathrm{jj}} > 500$ GeV are used to define the QCD $W^{\pm}Zjj$ CR. To extract the ZZ background normalisation, a ZZ CR is defined by requiring four leptons in the final state and removing the $E_{\mathrm{T}}^{\mathrm{miss}}$ requirement.

Events with $200 < m_{\mathrm{jj}} < 500$ GeV are used in the $H^{\pm\pm}$ search to define the "low-m_{jj}" CR. This CR has a similar background composition to the SR and is used to control the uncertainties of major background contributions. A dedicated QCD $W^{\pm}Zjj$ CR is defined by requiring events with three charged leptons in the final state, $m_{\mathrm{jj}} > 200$ GeV, and a trilepton invariant mass greater than 106 GeV.

3 Combination strategy

A simultaneous binned maximum-likelihood fit is performed for the $H_5^{\pm} \rightarrow W^{\pm}Z$ and $H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$ signal extraction to discriminate between the resonant signal and non-resonant SM backgrounds. The m_{WZ} distribution is used in the fit for the $H^{\pm}$ SR. The m_{T} distribution in five regions of m_{jj} with boundaries at (500, 850, 1450, 2100, 2550, ∞) GeV, is used in the fit for the $H^{\pm\pm}$ SR. The SRs and the ZZ, low-m_{jj}, and QCD $W^{\pm}Zjj$ CRs are fitted simultaneously. The QCD $W^{\pm}Zjj$ CR defined in Ref. [18] overlaps considerably with the QCD $W^{\pm}Zjj$ CR defined in Ref. [17] and is not strictly orthogonal to the $H^{\pm}$ SR. Consequently, only the QCD $W^{\pm}Zjj$ CR defined in Ref. [17] is used in this simultaneous fit.

The signal, EW $W^{\pm}W^{\pm}jj$, QCD $W^{\pm}Zjj$, and ZZ background normalisations are kept as floating parameters in the fit and are constrained by the data in both the SRs and dedicated CRs. The relatively small contributions of the EW $W^{\pm}Zjj$ and QCD $W^{\pm}W^{\pm}jj$ processes are normalised to the SM predictions and allowed to vary within their uncertainties. The systematic uncertainties are included as nuisance parameters [28] with Gaussian priors. The nuisance parameters are profiled in the fit with the shape and normalisation of each distribution varying within the specified constraints. The results are driven by the statistical uncertainty of the data in the SRs and none of the considered systematic uncertainties have significant impact on the sensitivity of this search. The expected limits on $\sin\theta_H$ improve by up to 5%, depending on m_{H_5}, if the systematic uncertainties are not included in the simultaneous fit. The largest systematic uncertainties considered are briefly discussed in the following.

The dominant contributions to the systematic uncertainties stem from the theoretical uncertainties in the physics modelling of the GM signal, followed by the experimental uncertainty sources related to the jet energy calibration [29]. The $H^{\pm}$ ($H^{\pm\pm}$) signal samples are simulated with MadGraph5_aMC@NLO 2.7.2

(2.9.5) [30] interfaced to Pythia8.186 (8.245) [31, 32] for the modelling of the parton shower in the dipole recoil scheme, hadronisation and underlying event. The $H^{\pm}$ and $H^{\pm\pm}$ samples are simulated at next-to-LO (NLO) and LO in QCD, respectively. No significant differences in the shapes between the LO and NLO simulations for each signal distribution are seen. The next-to-NLO (NNLO) predictions [15, 16] and the uncertainties due to the renormalisation and factorisation scale variations, parton distribution function, and the strong coupling constant are used for the normalisation of the signal samples when deriving the constraints on $\sin\theta_H$. In addition, uncertainties from these sources affecting the shape of the $H^{\pm\pm}$ distributions at LO are considered. An uncertainty due to the missing NLO EW corrections is adopted, as recommended in Ref. [16].

The systematic uncertainties originating from common sources, such as those associated with the integrated luminosity, are treated as correlated between the $H_5^{\pm} \to W^{\pm}Z$ and $H_5^{\pm\pm} \to W^{\pm}W^{\pm}$ channels. Uncertainties related to signal modelling as discussed above are treated as correlated.

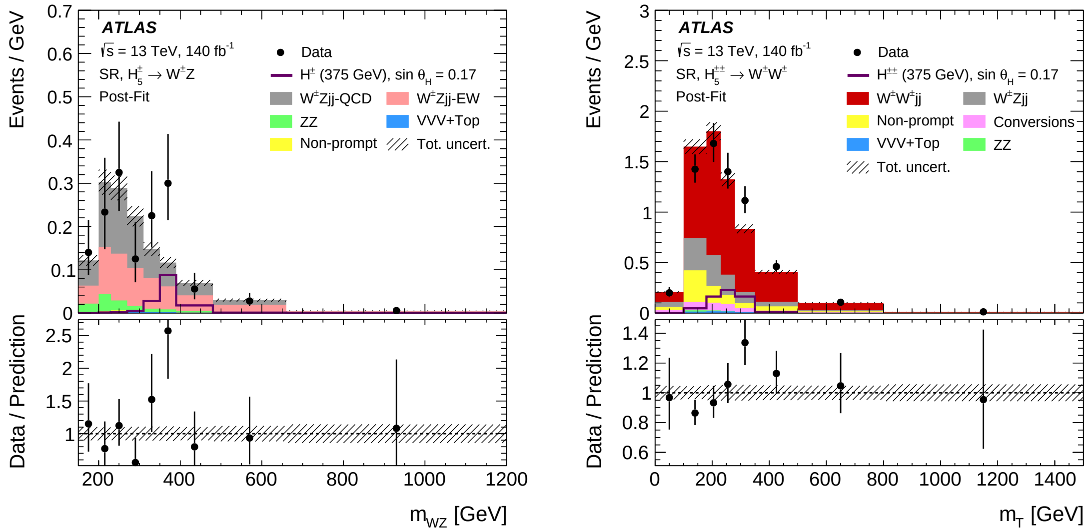
4 Results

The post-fit distribution of m_{WZ} in the $H_5^{\pm} \to W^{\pm}Z$ SR and the inclusive m_{T} distribution in the $H_5^{\pm\pm} \to W^{\pm}W^{\pm}$ SR are shown in Figure 4 for a fit performed under the SM background-only hypothesis. The expected contributions from a signal with $m_{H_5} = 375$ GeV and $\sin\theta_{\mathrm{H}} = 0.17$ are shown for illustration purposes. The corresponding post-fit background normalisation factors for the QCD $W^{\pm}Zjj$, EW $W^{\pm}W^{\pm}jj$ and ZZ processes are 0.73 ± 0.06, 1.16 ± 0.11, and 1.01 ± 0.15, respectively. These are consistent with the values reported in Refs. [17, 18]. No uncertainties are significantly constrained or pulled in the simultaneous fit.

The 95% CL upper limits on the production cross-section times branching fraction $\sigma_{\mathrm{VBF}}(H_5^{\pm}) \times \mathcal{B}(H_5^{\pm} \to W^{\pm}Z)$ and $\sigma_{\mathrm{VBF}}(H_5^{\pm\pm}) \times \mathcal{B}(H_5^{\pm\pm} \to W^{\pm}W^{\pm})$ for the VBF production of singly and doubly charged Higgs bosons as a function of m_{H_5} from 200 to 3000 GeV are shown in Figures 5(a) and 5(b). The 95% CL limits are derived using the $\mathrm{CL_s}$ method [33, 34]. The asymptotic approximation [35], whose validity was confirmed through studies with pseudo-experiments, is used to derive the upper limits.

The simultaneous fit results are interpreted in the context of the GM model, setting the most stringent constraints to date on the $\sin\theta_H$ parameter as a function of m_{H_5}. The limits are shown in Figure 5(c). The black hatched region represents the parameter space for which the total width of the quintuplet states exceeds 10% of m_{H_5}, where the model is not applicable due to considerations of perturbativity and vacuum stability, and indirect experimental constraints [16]. The combined expected limits on $\sin\theta_H$ are 10% to 26% more stringent, depending on m_{H_5}, than the respective limits obtained separately in the $H_5^{\pm} \to W^{\pm}Z$ [17] and $H_5^{\pm\pm} \to W^{\pm}W^{\pm}$ [18] channels. The expected limits are 10% to 50% stronger, depending upon m_{H_5}, compared to the respective limits obtained by the CMS Collaboration [23]. The observed 95% CL limits exclude $\sin\theta_H$ parameter values greater than 0.10–0.36 for the m_{H_5} between 200 and 1500 GeV.

The results show a local excess of events over the SM prediction at a resonance mass of around 400 GeV as can be seen in Figure 5. The significance of the excess has been evaluated for different m_{H_5} in terms of the local p–value. The largest excess is for $m_{H_5} = 375$ GeV, with a p–value of 5.7×10^{-4}, corresponding to a local significance of 3.3 standard deviations. The global significance of the excess was also evaluated [36], and yields a global p–value of 5.6×10^{-3}, corresponding to a global significance of 2.5 standard deviations. The largest local (global) significances obtained separately from the fits performed in the $H_5^{\pm} \to W^{\pm}Z$ and



(a) (b)

Figure 4: Post-fit (a) m_{WZ} and (b) inclusive m_{T} distributions in the signal regions for the SM background-only hypothesis. Data are shown as black markers with vertical error bars representing the statistical uncertainty. Filled histograms show contributions of various SM processes, with the hatched band representing the total uncertainty. The last bin of each distribution includes overflow events. The lower panel shows the ratio of the data to the SM prediction. The line shows the prediction of the GM model for $\sin\theta_{\mathrm{H}} = 0.17$ and $m_{H_5} = 375$ GeV, where the $\sin\theta_{\mathrm{H}}$ value corresponds to the expected 95% CL limit for that H_5 mass.

$H_5^{\pm\pm} \to W^\pm W^\pm$ channels used in this analysis are 2.8 (1.6) and 3.2 (2.5) for m_{H_5} values of 375 GeV and 450 GeV, respectively.

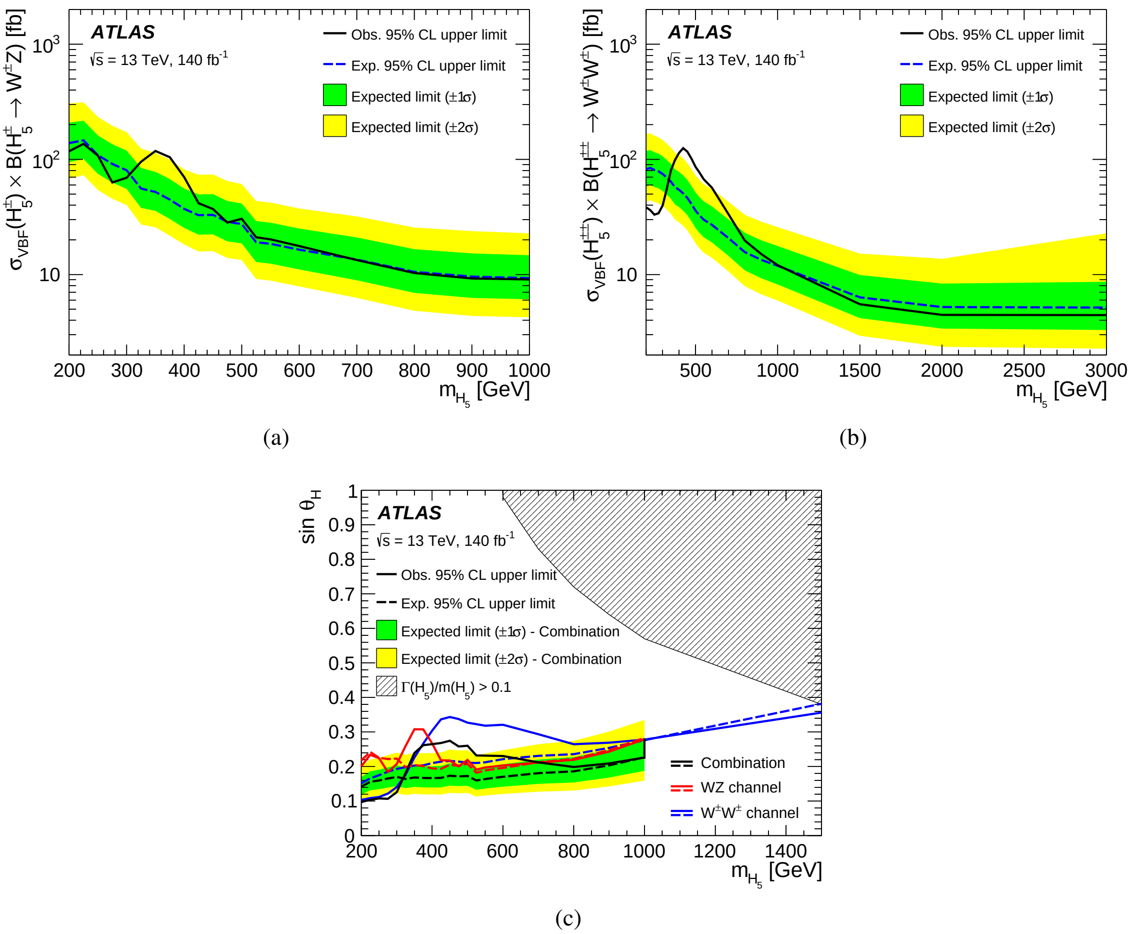


Figure 5: Expected and observed exclusion limits at 95% CL for (a) $\sigma_{\text{VBF}}(H_5^{\pm}) \times \mathcal{B}(H_5^{\pm} \to W^{\pm}Z)$ and (b) $\sigma_{\text{VBF}}(H_5^{\pm\pm}) \times \mathcal{B}(H_5^{\pm\pm} \to W^{\pm}W^{\pm})$ as a function of m_{H_5}. The inner (outer) band represents the 68% (95%) confidence interval around the median expected limit. The exclusion limits for (c) $\sin\theta_H$ are shown up to $m_{H_5} = 1500$ GeV. The limits on $\sin\theta_H$ obtained separately in the $H_5^{\pm} \to W^{\pm}Z$ and $H_5^{\pm\pm} \to W^{\pm}W^{\pm}$ channels are also shown for comparison. The hatched region covers the parameter space where the intrinsic widths of the $H_5^{\pm}$ and $H_5^{\pm\pm}$ bosons would be larger than 10% of m_{H_5} and is disfavoured in the GM model [16].

5 Conclusion

A combination of searches for singly and doubly charged Higgs bosons, $H_5^{\pm} \rightarrow W^{\pm}Z$ and $H_5^{\pm\pm} \rightarrow W^{\pm}W^{\pm}$, produced via vector-boson fusion is reported. The dataset corresponds to 140 fb^{-1} of proton–proton collision data at $\sqrt{s}$ = 13 TeV collected with the ATLAS detector during Run 2 of the LHC (2015-2018). Constraints are reported on the production cross-section times branching fraction for singly and doubly charged Higgs bosons. The simultaneous fit results are interpreted in the context of the Georgi-Machacek model, for which the most stringent constraints to date are set. The observed 95% CL limits exclude $\sin\theta_H$ parameter values greater than 0.10–0.36 for m_{H_5} between 200 and 1500 GeV. The largest deviation from the Standard Model occurs for a resonant mass near 375 GeV, with a global significance of 2.5 standard deviations.

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