

AA LONG TERM NOTE No. 16Summary of the meeting of August 17, 1982

Present : B. Autin, T. Dorenbos, S.X. Fang, F. Ferger, L. Henny, C. Johnson,
C. Metzger, W. Pirkel, G. Nassibian, L. Rinolfi, K.H. Schindl,
J.C. Schnuriger, R. Sherwood, C. Taylor, H.H. Umstaetter, E.J.N. Wilson

Topic : Pulsed Target and AC Acceptance.

The copies of the transparencies for this talk are attached. A discussion on the possibilities of the current carrying target, including the effect of proton beam defocusing was presented.

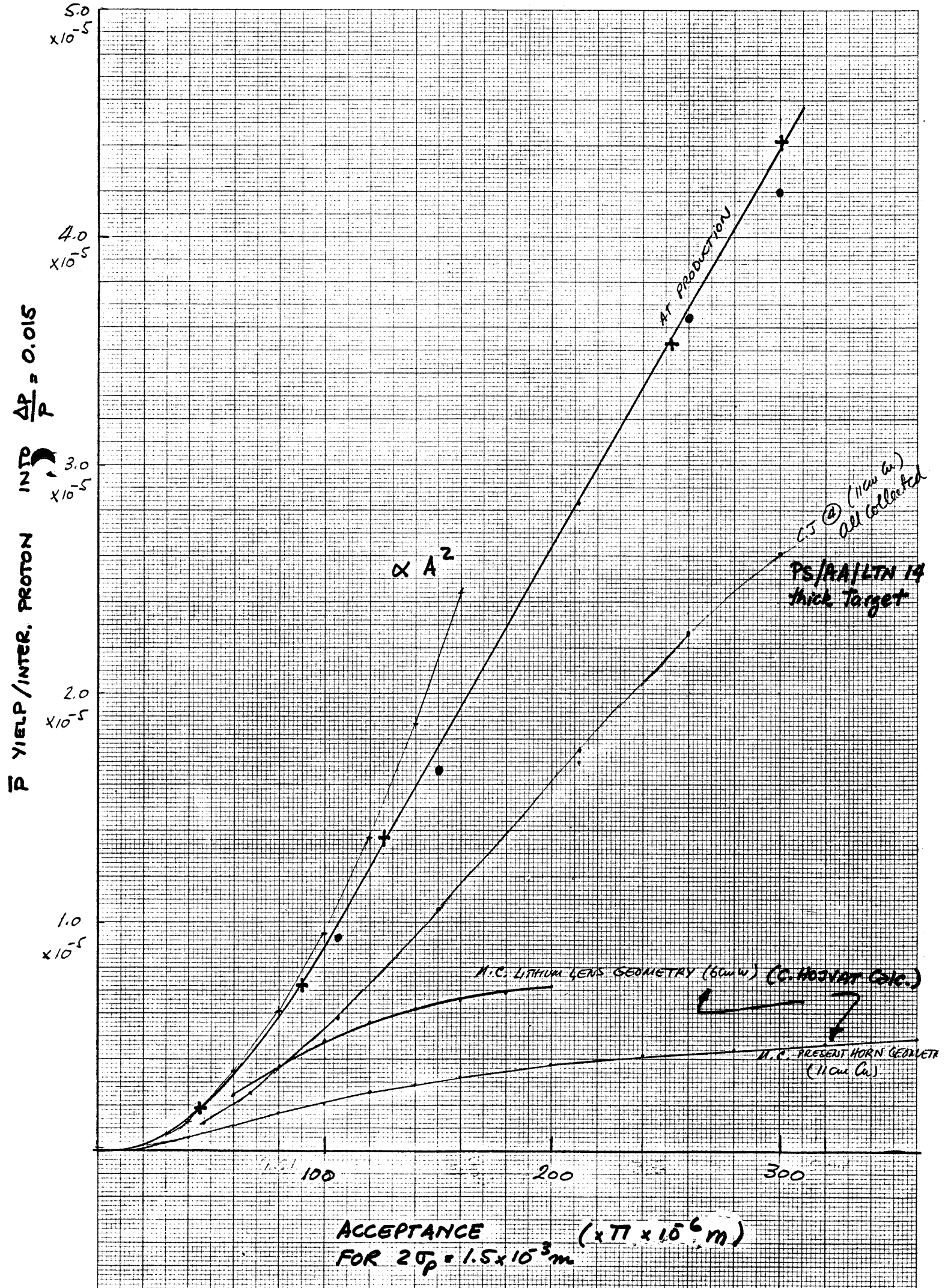
Although more detail calculations on antiproton yields are required it looks that the AC acceptance can be limited to 200π mm.mrad to obtain greater than 10^8 \bar{p} /pulse for a total $\Delta p/p$ of 0.06.

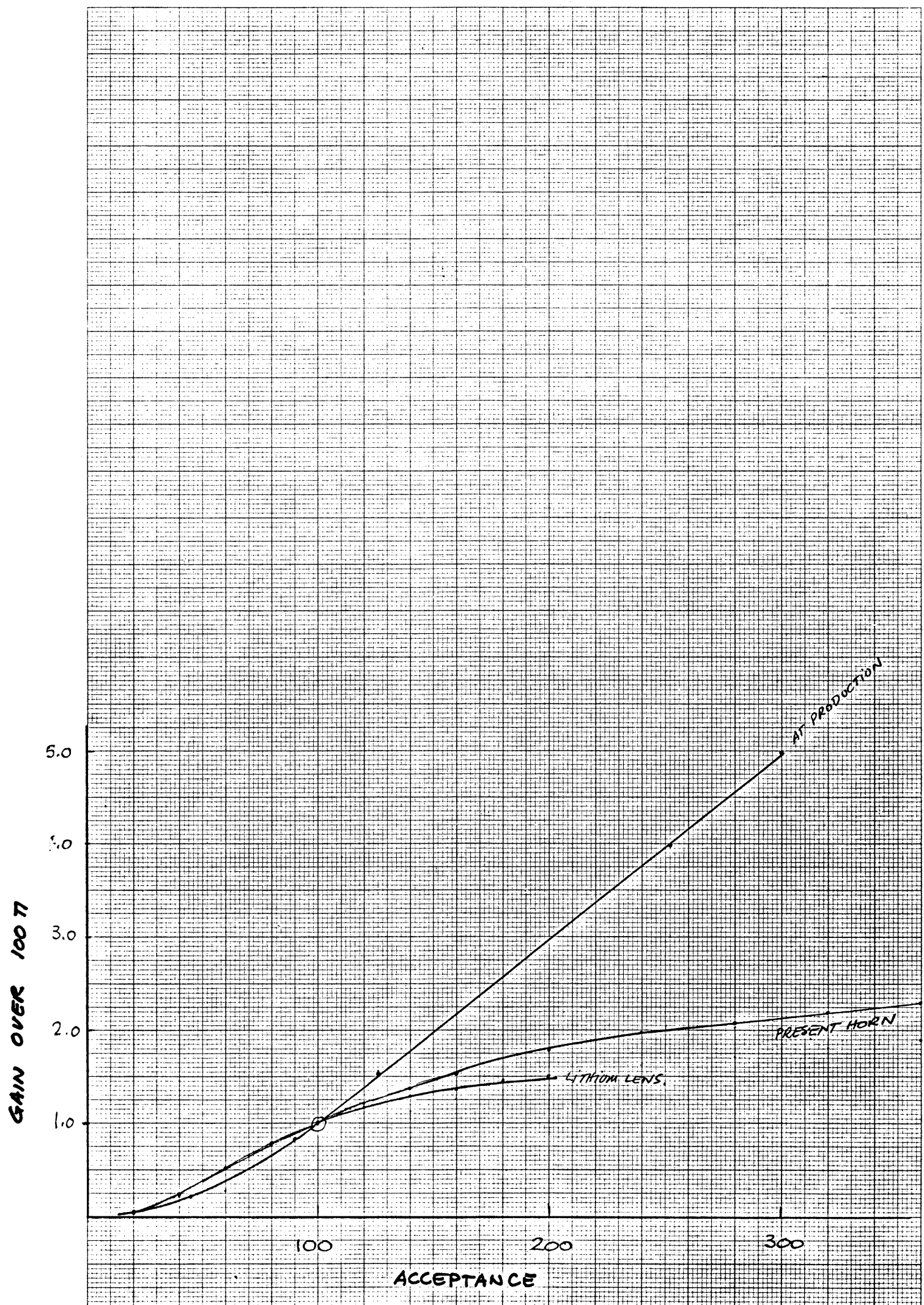
The geometry and magnetic fields for the first conducting target experiment in the AA were also presented.

C. Hojvat

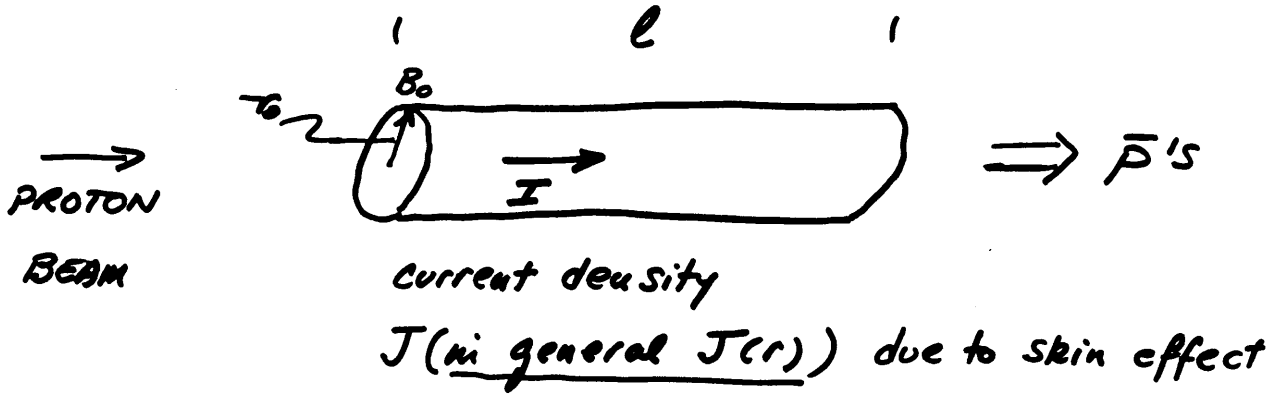
Distribution

PS/2 list





COMMENTS ON PULSED TARGET



For the case $J = \text{cte}$ $B(r) = \mu_0 \frac{r}{2} J$, $g = \frac{\mu_0}{2} J$
 Like a quadrupole for both planes (x, y)

① For \bar{p} 's (FOCUSING CASE)

$$M(A|B) : \begin{pmatrix} \cos L & \frac{1}{\sqrt{K}} \sin L \\ -\sqrt{K} \sin L & \cos L \end{pmatrix} \Rightarrow \begin{cases} \beta = \beta_0 \cos^2 L + \frac{1}{K \beta_0} \sin^2 L \\ \alpha = \left(\sqrt{K} \beta_0 - \frac{1}{\sqrt{K} \beta_0} \right) \sin L \cos L \end{cases}$$

where:

$$L = l \sqrt{K}$$

$$K = 0.31952 \frac{g}{\beta^2} \quad (\text{protons}/\bar{p})$$

FOR $\sqrt{K} = \frac{1}{\beta_0}$ then

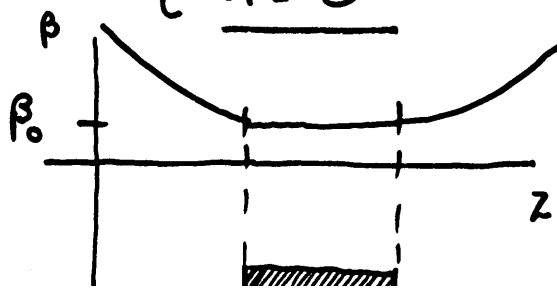
1 period:

$$L = 2\pi$$

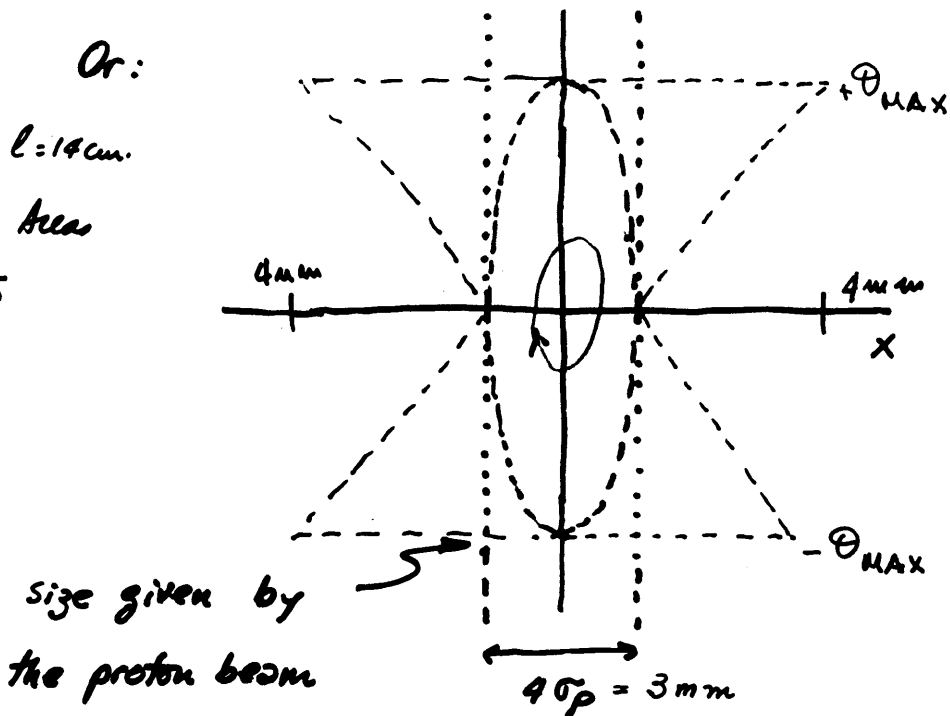
$$l \sqrt{K} = 2\pi$$

$$l = 2\pi / \sqrt{K}$$

$\beta = \beta_0$
 $\alpha = 0$ throughout the target



Or:
 @ 100 π $l = 14$ cm.
 Ratio of Area
 2.5



σ_p AS SMALL
 AS COMPATIBLE
 WITH
 BEAM
 INTENSITY

In order not to dilute, then:

$$2\sigma_p = \sqrt{\beta_0 \epsilon_p}$$

$$\Rightarrow \beta_0 = \frac{(2\sigma_p)^2}{\epsilon_p}$$

Assume $2 \times \sigma_p = 1.5 \times 10^{-3}$ m

3.575 GeV/c $r_{\text{Tar}} = 2\sigma_p = 1.5 \times 10^{-3}$ m

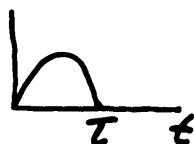
$$\epsilon_p = 2\sigma_p \theta_{\text{MAX}}$$

ϵ_p $\times \pi 10^{-6}$ m	θ_{MAX} mrad	β_0 m	\sqrt{k} m^{-1}	γ T/m	J KA/cm ²	I kA
50	33.3	0.045	22.22	5887	937	66
100	66.6	0.0225	44.44	23548	3.75×10^3	265
150	100	0.0150	66.67	53000	8.4×10^3	596
200	133	0.0113	88.89	94210	15×10^3	1060
250	166	9×10^{-3}	111.11	147.2 K	23×10^3	1657
300	200	7.5×10^{-3}	133.33	212 K	34×10^3	2385

14cm one period

e.g. ENTHALPY CONTENT of Cu ϕ 3mm 11 cm long \sim 1700 JOULES

FOR $\frac{\delta}{\tau} = 0.64$ (or J \sim cte) $\tau = 36 \mu\text{s}$



500°C $I_{\text{MAX}} = 300$ kA

BEAM \sim 1.25 KW \sim 3000 Joules PS/AA/ME/10

② FOR PROTONS (DEFOCUSING CASE)

③

$$M(A|B): \begin{pmatrix} \cosh L & \frac{1}{\sqrt{k}} \sinh L \\ \sqrt{k} \sinh L & \cosh L \end{pmatrix} \Rightarrow \begin{cases} \beta = \beta_0 \cosh^2 L + \frac{1}{k\beta_0} \sinh^2 L \\ \alpha = -\left(\sqrt{k}\beta_0 + \frac{1}{\sqrt{k}\beta_0}\right) \sinh L \cosh L \end{cases}$$

(around a waist β_0 $\alpha=0$)

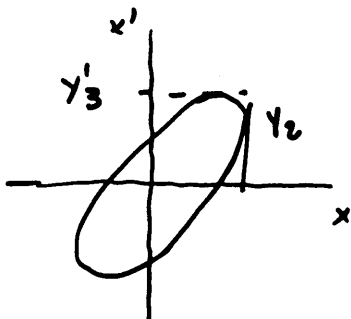
now \sqrt{k} is for the same gradient g as for \bar{p}

but for the proton energy:

eg $\sqrt{k} = 44.44 \quad \therefore \sqrt{k} = 16.21$

For a waist at the Target centre then at a distance d away: $\beta_0 = (1.5 \times 10^3)^2 / 2.5 \times 10^{-6} = 0.915 \text{ m}$

CASE OF $E_p / 100 \text{ MeV}$	d m	β m	α	y_2 $y_3' = \sqrt{\epsilon} \delta$ radians
-0.07		2.7	+35.7	2.6×10^{-3} 0.22
-0.06		2.1	+25.5	2.3×10^{-3} 0.18
-0.04		1.4	+12.6	0.108
-0.02		1.0	+5.2	0.052
+0.02		1.0	-5.2	
+0.04		1.4	-12.6	
+0.06		2.1	-25.5	
+0.07		2.7	-35.7	
0.0		0.92	0	0.0105



VERY STRONG DEFOCUSING EFFECT
 → LIMITS THE "EFFECTIVE" TGT LENGTH
 → REQUIRES MORE THAN CONVENT. FOCUSING

The production length (effective) due to the change in size of the proton beam is given by:

$$L_{\text{Prod}} = \int \frac{\beta_0}{\beta} dl$$

$$L_p = \int \frac{\beta_0}{\beta_0 \cosh^2 L + \frac{1}{k \beta_0} \sinh^2 L} dl$$

$$L_p = \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{1}{\cosh^2 L + \frac{1}{k \beta_0^2} \sinh^2 L} dl$$

from L.C. Teng note:

$$L_p = 2 \beta_0 \tan^{-1} \left(\frac{1}{\beta_0 \sqrt{k}} \tanh \frac{L}{2} \right)$$

For case above: (where β_0, k are for 26 GeV protons)

$$l = 0.14$$

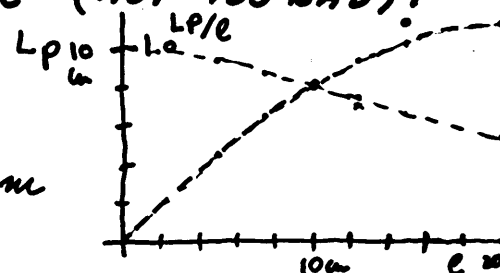
$$\beta_0 = 0.915 \text{ m}$$

$$\sqrt{k} = 16.21$$

$$L = 0.14 \text{ m} = \sqrt{k}$$

$$L_p = 0.10 \text{ m (NOT TOO BAD)!}$$

FOR L very large $\left(\int_{-\infty}^{+\infty} \right)$ $L_p = 0.123 \text{ m}$



CONCLUSIONS FOR 100% CASE

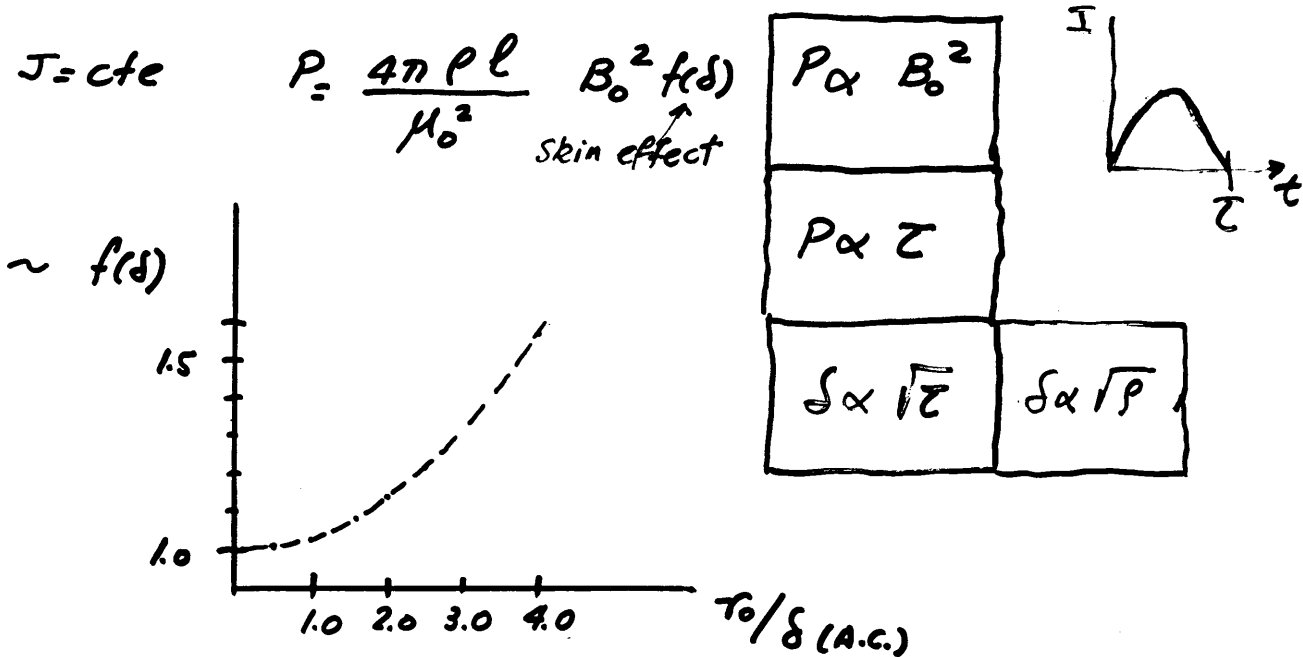
- ① IT LOOKS PROMISING
- ② GAIN: SAME # OF \bar{p} 'S CONTAINED IN A ELLIPSE AT/TGT EXIT
(NOW ALL SUFFER ABSORPTION IN TARGET)
- ③ PROTON FOCUSING NEEDS TO BE LOOKED AT

PULSED TARGET:

LOWER POWER DISSIPATED AND
TRADE P FOCUSING WITH P DEFOCUSING.

- ① Lower Current below $\sqrt{K} = \frac{1}{\beta_0}$ (L. Teng for FNAL)
- ② Lower skin depts by reducing pulse length τ
- ③ Utilize $\frac{1}{\tau}$ field outside target (R.S/E.J.)

POWER IN TARGET:



- ① and ③ Lower P
- ② Increases P (depending on choice of τ or P)

SAME IS TRUE FOR LITHIUM LENS

LIMITS: 1. MAGNETIC FORCES

2. POWER DISSIPATED

INCOMPLETE YIELD OF ANTI-PROTONS

IN $\frac{\Delta p}{p} = 0.015$

FOR 1×10^{13} incident protons.

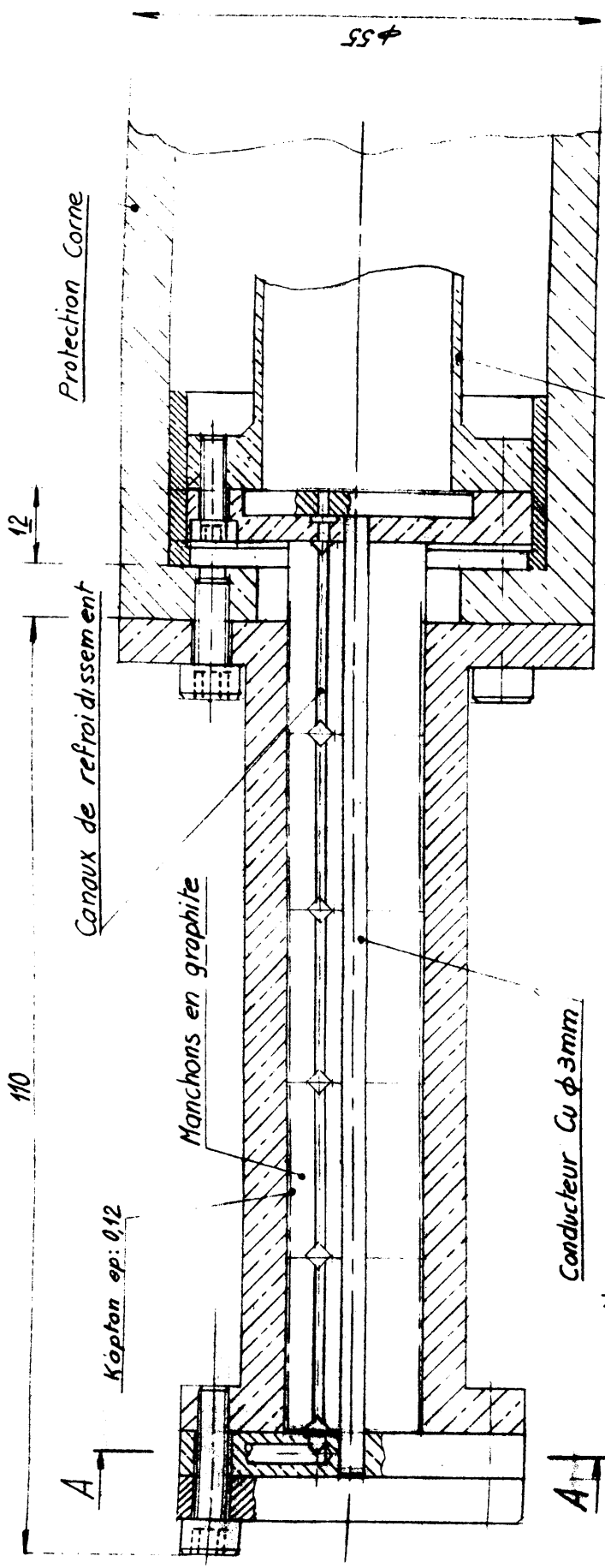
?

$E_{\bar{p}}$ $\pi 10^{-6} m$	Relative at Production	Present Horn CH/AVE	Proposed Li Lens 1cm 1000 μ m	Pulsed Target + Collector	
100	1.0	1.4×10^7 1.0	2.2×10^7 1.0	E.J.R.S. HORN 3.2×10^7 ~ Li Lens 5.0×10^7	
200	3.0	2.5×10^7 1.2	3.3×10^7 1.5	?	
300	5.0	3.0×10^7 2.1	4.0×10^7 1.8	?	

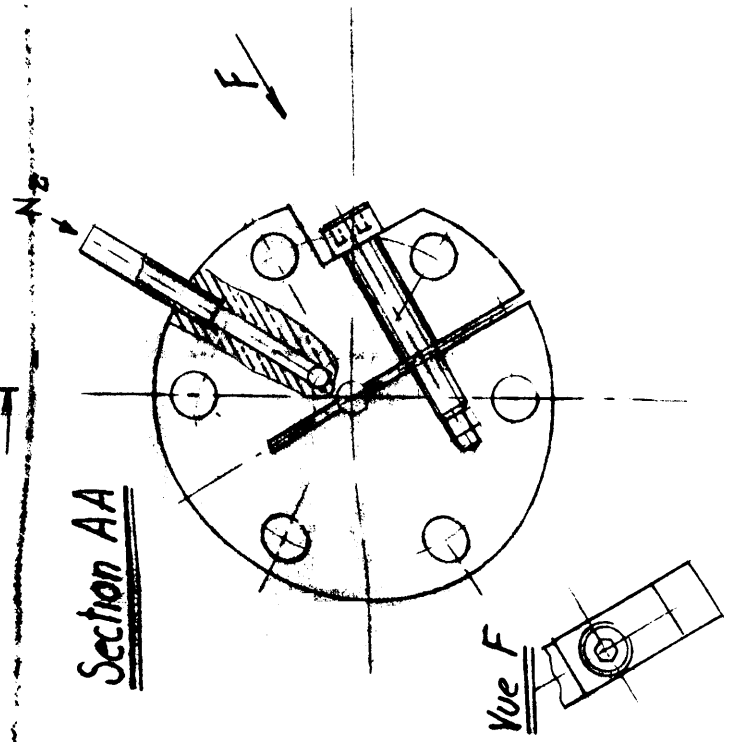
WITH PLANNED $\frac{\Delta p}{p} = 0.06$ (x 4 bigger than table)

$10^8 \bar{p}/\text{pulse}$ can be achieved

with $\frac{\Delta p}{p} = 0.06$ and $E_{\bar{p}} \leq 200 \pi$



Section AA



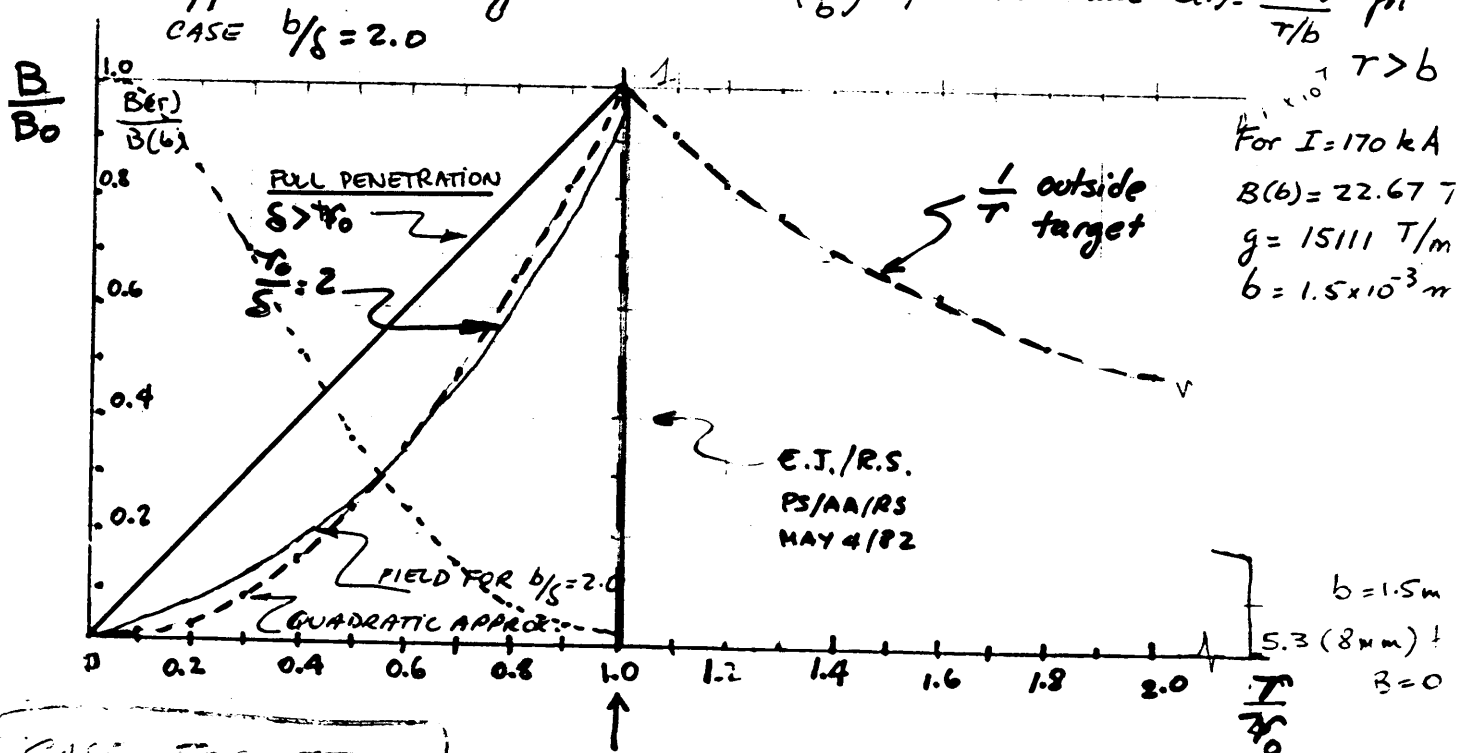
CIBLE CONDUCTRICE

— Ensemble —

— Projet —

sch: 2/4 R. Valfreva le 6/8/1982

approximated γ $N = 200 \times (\frac{1}{b})$ $\tau > b$ $\tau > b$



CASE FOR TEST
 END TARGET

TARGET
 RADIUS ($2.5\sigma_p$)

$b = 1.5 \text{ m}$
 $5.3 (8 \mu\text{m})!$
 $B = 0$