COMPUTED EFFECTS OF SHIMS AT THE END OF THE AC BHN DIPOLES

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INTRODUCTION

Magnet programmes TOSCA and PE2D have been used in an attempt to predict the effects of shims placed at the ends of the BHN dipoles in the AC. Because only comparisons were required (with and without shim) the pole shape was only approximately represented (see Fig. 1). Even with this simplified geometry, obtaining consistent results from TOSCA proved difficult.

MAGNET REPRESENTATION

The pole shape is shown in Fig. 1; the major difference from reality is that the pole corners are not rounded. A single shim 6 mm thick, 27 mm wide and 75 mm high was moved in steps of 27 mm so as to cover the whole of one half of the pole end. The field integrals through the magnet were computed for $x = -105$ to $x = 105$ mm in steps of 15 mm with the intention of comparing them with the "no shim" case.

Because magnetic saturation is appreciable at pole-ends, the different cases were all run in the non-linear mode, the B-H curve employed being the same as that used for Q54. The half-length of the yoke as represented in TOSCA was 10 mm less than the actual half-length.

RESULTS

For unknown reason the problem was very reluctant to converge, and eventually most of the cases were terminated after 30 iterations, still unconverged. In order that the jobs should run overnight, they were split into two runs of 15 iterations each. TOSCA has always had the peculiarity that on re-starting from an old solution, the initial iterations fluctuate wildly so that after the second series of 15 iterations one is not much nearer convergence than before. Cases BHN3-BHN6 (shim covering $x = 0$ to 108 mm) were run in this way, and as can be seen from the fields at z = 0 (Table 1) good consistency was found.

For cases BHN7 (shim at -135 to -108 mm) and BHN8 (no shim), however, it was not possible to re-start from the old solution without a divide zero error occurring. This was particularly unfortunate in the BHN8 case, since this was to form the base from which all comparisons were to be made. However, after only 15 iterations (16 for BHN7) the fields at $z = 0$ were found to be sufficiently close to those of cases BHN3-BHN6 that reasonable confidence could be placed in the results, which were obtained as followx:

- 1. Check that the fields at $z = 0$ are similar for all cases.
- 2. Calculate $L_0 = (\int B dz)/B_0$ at $x = 0$. These are seen to be consistent in that L_0 increases as the shim is moved towards $x = 0$.
- 3. For each case calculate ΔL/Lo versus x.
- 4. For BHN3-BHN7 calculate $(\Delta L/L_0)$ $(\Delta L/L_0)$ BHN8 and adjust the base line so that all differences are positive. These are plotted in Fig. 2.
- 5. For BHN7, and perhaps BHN6, where the shim is far from $x = 0$, a better result is probably obtained by calculating $(\Delta L/L_0)_{-X}-(\Delta L/L_0)_{+X}$. This too is shown in Fig. 2.

Finally, the effect of a 27 mm-wide bump on the pole profile was computed in 2 dimensions using PE2D. The bump was placed symmetrically about $x = 0$, and the results processed as follows:

- 1. Calculate ΔB/Bo versus x (no bump).
- 2. Calculate ΔB/Bo versus x (with bump).
- 3. Subtract (1) from (2) amd normalise so that $(\Delta B/B_0)_{\text{max}} \approx 9.5$ *E-4.

This result is included in Fig. 2 (BHN3); it can be seen that a vertically extended shim at the pole-end has a wider radial influence than that of a bump of the same width in the two-dimensional profile.

Table 1.

