#### EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH ORGANISATION EUROPEENNE POUR LA RECHERCHE NUCLEAIRE

**CERN - PS DIVISION** 

PS/PA/Note96-27 (PPC)

## MINUTES OF THE PPC MEETING HELD ON 21ST JUNE, 1996

D. Manglunki

Geneva, Switzerland 28 June 1996

### Minutes of the PPC meeting held on June 21st, 1996

Present: G.Arduini, R.Cappi (Chairman), K.Cornelis, R.Jung, D.Manglunki (Secretary), M.Martini, J.P.Riunaud, G.Roy, K.Schindl

Agenda: PS-SPS transfer of LHC beams:

- Identification of the problems
- Status of PS beam studies
- Common strategy: Who will do what? How and when?

Introduction (R.Cappi):

- As a reminder, specifications of the PS beam for LHC:
  - 84 bunches
  - total intensity 1.5x10<sup>13</sup> protons per pulse
  - $\beta \gamma = 28.1$
  - $\varepsilon_x^* \approx \varepsilon_y^* \approx 3 \,\mu m$   $(\beta \gamma \, \sigma_{x,y}^2 / \beta_{x,y})$
  - $4 \sigma_t \approx 4 \text{ ns}$
  - $2 \sigma_{\rm p}/{\rm p} \approx 2.5 {\rm x} 10^{-3}$
  - bunch spacing = 25 ns
- Problem number one is the conservation of transverse emittances in the transfer from PS to SPS. The main causes for blow-up would be:
  - the stray field in the PS magnet downstream the extraction septum
  - mismatch
  - missteering

#### Status of the study of the stray field(M.Martini):

(see attached copy of poster and transparencies)

- The model that had been used up to now (localised multipole components, tracked with MAD) gave a computed geometric blow-up factor of more than 50%.
- The stray field has been measured on a laboratory magnet. The effect of shims (absent on the measured unit) has been added, and particles were tracked. The computed Twiss parameters are in better agreement with the measurements. (geometric blow-up less than 15%)
- Dispersion studies is the next step: reproduce on computer simulations the observed behaviour of the matching with respect to Δp/p

Discussion on common strategy

- The computed geometric blow-up is not to be neglected but is usually pessimistic.
- The measurements of the magnetic field on the PS magnets in the laboratory, and the computations have been performed with PFW currents that are different from the ones that are currently used in operation. They have to be refined with the proper currents.
- A correction scheme for matching with respect to  $\Delta p/p$  has to be studied, possibly by the addition of sextupoles inTT2.
- The SPS semgrids that are currently used in TT10 are not adapted to the LHC beam. A new device based on Mylar/Aluminum screens, using transition radiation (surface effect) is being studied by SL/BI.
- Machine developments to compare the emittances of the circulating beam in PS and SPS, using wire scanners, will take place during the following weeks.

#### TRAJECTORY AND OPTICAL PARAMETERS IN A NON-LINEAR STRAY FIELD

D. Manglunki, M. Martini, CERN

I. Kirsten, Heidelberg University

#### Abstract

A new optics for the main CERN Proton Synchrotron magnet is modelled to allow a precise description of the ejected beams. For that purpose, field maps of the magnet have been measured for the various operational current settings. They include the central field, the end stray field and the lateral stray field. In order to get a functional form which can be inserted in the equations of motion for a charged particle in a magnetic field, the discrete field maps are converted into bi-dimensional polynomials of degree up to fourteen. These equations of motion can be written as a set of four first order differential equations which are solved simultaneously. Two of them are non-linear and describe the centroid motion, the other two are linear and apply to the betatron motion. The method has been validated by producing extraction conditions which have been verified experimentally with the 26 GeV/c beam for the future LHC



The CERN PS lattice consists of 10 super-periods made of 10 combined function magnets, eight 1.0 m and two 2.4 m long straight sections. Each magnet is composed of two half-units with gradient of opposite sign, separated by a central junction. The half-units are made of five blocks with small wedge gaps in between lined up on the central orbit.

The necessity to extract 26 GeV/c protons in a 2.4 m straight section with little angle deflection (29 mrad) imposes the downstream half-unit adjacent to the ejection septum to be open to ease the fitting of the extraction pipe across the magnet aperture. The ejection trajectory in this region remains close to the central orbit and thus the aberrations in the magnetic fields are kept at a reasonable value. When traversing the subsequent F half-unit the ejection trajectory moves away from the central orbit and field aberrations become strongly non-linear: the beam experiences a field gradient with a reverse sign, yielding large horizontal betatron function values at the magnet end. Reduction of the non-linear aberrations was done by shimming the F half-unit. Straight parallel shims have been mounted at different radial positions on the five blocks to shape a constant magnetic field over the ejected beam width.

0.01 ſŤÌ .n or [m] ſm∮ х 0.3 **Fitted field map** Difference between fitted and measured field



New magnetic measurements on operational CERN PS magnet working points were performed in 1992, including measurements of the central field, the end and lateral stray fields, and the field in the junction between the two half-units. Polynomials up to degree 14 in x and z have been retained to reach a good agreement with the measured field (accuracy within ±0.01 T). The fitting has been performed using the standard Mathematica fit function. × [m]



Fitted field map after considering the shims

Field calculations have been carried out on the five blocks equipped with shims using the twodimensional Poisson program. Polynomial fittings up to degree 25 in x of Poisson output have been carried out to get a functional form of the computed field.

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	$\beta_{r}[m]$	a	$D_x[\mathbf{m}]$	D <sub>x</sub>	β <sub>y</sub> [m]	a,
Field map	31.25	-2.71	3.10	0.26	7.11	0.74
MAD	33.79	-3.37	3.25	0.32	6.13	0.85

Optics parameters (at stray field exit) derived from the transfer matrices



Election trajectory and transfer matrix computations have been performed using the built-in Mathematica numerical differential equation solver with initial conditions given by MAD: the beam centroid enters the field map with coordinates x=91.6 mm, #=62.6 mrad, and exits the field map at x=345.0 mm, #=36.4 mrad. For comparison the angle of the ejection pipe with respect to the z-axis in the F half-unit 16 is 43 mrad.

> The optical parameters have been derived from the transfer matrix components and compared with previous models which consider the MAD stray field description as given by dipole, quadrupole and sextupole coefficients distributed over the magnet length.

Magnetic measurements have been done on a laboratory magnet unit in the absence of shims, thus the measured field map has to be corrected to consider the shimming effect.



Transverse emittance matching in the TT2 channel are obtained from beam profiles measured at three SEM-grid detectors. Using the computed optics parameters, the mismatch derived from measurements was found to be less than 15% for the horizontal plane and less than 10% for the vertical plane. This is a fairly good result (the best achieved so far) considering that an error on optical parameters transforms into a large mismatch error.

The emittance measurement program shows a mismatch less than 15%

MAGNET PS





TRAJ model

Trajectory & transfer matrix computations through ejection magnet stray field (see EPAC 36, I. Kirsten, D. Manglun ki, M. Martini)

📼 Emi	ttance and mism	natch
FT16.MSG257 H	ORIZONTAL	
ε(2σ):	0.66 πμm	
$\frac{4\sigma^2}{\beta}$ :	0.78 πμm	
Blow up:	(17.27 3)	
β (G):	15.91	0.86
<u>α (B):</u>	1.30	-0.03
Match	ning vector	
		Close

 $\frac{\Delta \mathcal{E}}{\mathcal{E}} = k \left( \frac{k}{2} + \sqrt{1 + \frac{k^2}{4}} \right)$ 

 $k = \left| \frac{(G-1)^2 + B^2}{G} \right|$ 

k is the matching vector (k=0 i.e G=1 and B=0 for perfect matching)

	Dpx [1]	4 0 1 0 0 0 0 0 0 0 0 0 0 0 0 0
	Dx [m]	3.03 -0.999 -0.06 -0.06 -1.759 -0.06 -0.06 -1.759 -1.759 -1.759 -1.232 -1.759 -1.232 -1.259 -1.259 -1.259 -1.259 -1.259 -1.259 -1.2500 -1.250
	A L px(co) [.001]	0.000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.0000 0.00000 0.000000
: TT2 r: 1	Z O N T x(co) [mm]	0.000
line supe	H O R I mux [2pi]	13.568 14.709 14.979 15.117 15.117 16.345 16.345 16.345 16.345 16.345 16.325 16.325 16.325 16.325 16.325 16.325 17 15.117
	I alfax [1]	-3.358 2.007 -0.106 0.929 -0.056 mux dmus beta Dx(n
WISS Ymm: F	betax [m]	33.241 19.130 7.544 32.398 5.341 00 mm
ions. T 0000 s	ы dist г л	$\begin{array}{c} 0.000\\ 162.011\\ 180.471\\ 198.931\\ 343.099\\ 343.0994\\ 343.0994\\ 0.0000\\ 0.0000\\ \end{array}$
funct: 0.00	SEQUENCI occ. no.	
r lattice (p)/p:	ELEMENT S element name	TT2 MSG257 MSG267 MSG277 TT2 TT2 length = [s) =
Lineal Delta	l pos. no.	begin 98 114 end total delta(

TRAJ MODEL See EPAC96 (I. Kishu, D. Munglun ki, M. Marhin)

# MAD model



	ннн	
	Dpx [1]	0 0 0 1 4 4 2 M 8 4 M 0 1 7 1 4 4 2 M 8 4 M 0 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1
	Dx [m]	3.046 0.612 2.311 40653 40653 47786 50718 24966
	A L px(co) [.001]	0.000 0.000 0.000 0.000 0.000 116. 16.
: TT2 r: 1	Z O N T x(co) [mm]	
line supe:	H O R I mux [2pi]	13.569 14.774 14.993 15.131 16.407 16.407 x x x x max) max) r.m.s.)
	alfax [1]	-2.633 1.142 0.225 0.899 0.899 mux bet Dx(
WISS Ymm: F	betax [m]	30.572 13.609 10.850 22.408 7.761 00 mm
POINT, A ions. T 0000 s	E dist I [m] I	$\begin{array}{c} 0.000\\ 162.011\\ 180.471\\ 198.931\\ 343.099\\ 343.0994\\ 343.0994\\ 0.0000\end{array}$
<pre> Γ, R, TO,     funct     0.00 </pre>	SEQUENC occ. no.	
ROM, POINT r lattice (p)/p:	ELEMENT ( element name	TT2 MSG257 MSG267 MSG277 TT2 length = (s) =
1TT2,F1 Linea: Delta	pos.	begin 98 106 114 end total delta

MAD MODEL



#### **Distribution**

G.Arduini	SL
E.Brouzet	SL
R.Cappi	PS
K.Cornelis	SL
R.Jung	SL
K.H.Kissler	SL
D.Manglunki	PS
M.Martini	PS
J.P.Riunaud	PS
G.Roy	SL
K.Schindl	PS
D.J.Simon	PS