An isochronous optics for EPA

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1 Introduction

An experiment using very short electron bunches in the EPA ring requires the orbit length to be independent of momentum and betatron amplitudes. The modification of the ring optics should be obtained preferably using the existing machine components, as the machine will return to normal lepton production immediately after the test. Equally, beam envelopes and trajectories in the injection area should change as little as possible.

This note presents the results of the beam optics calculations carried out to prepare the experiment.

2 Modification of α

The derivative of the orbit length w.r.t. $\Delta p/p$ is equal to the ring circumference C multiplied by the momentum compaction factor α . The latter depends on the value of the dispersion function in the bending magnets. If the dispersion is modified using a set of N quadrupoles the change of α is given [1] by

$$C \Delta \alpha = - \sum_{i=1}^{N} \Delta K_i L_i D_i D_i^*$$
(1)

with :

 ΔK_i the normalized gradient increment of quadrupole *i*

 L_i the length of quadrupole number i

 D_i the *unperturbed* dispersion at quadrupole i

 D_i^* the modified dispersion at quadrupole i



Figure 1: Layout of the quadrupole families and the dispersion function with the nominal optics in one half of the EPA ring

For small gradient increments the higher order terms in ΔK may be neglected and a crude estimation of the α change can be made using the dispersion values of the *nominal* machine only:

$$C \Delta \alpha \approx - \sum_{i=1}^{N} \Delta K_i L_i D_i^2$$
 (2)

This shows why the quadrupole families QFN and QDN have little effect, as at these locations D is zero in the nominal optics (see fig. 1). The QFI quadrupoles are not shown in this figure, they are assumed to belong to the QFN family and to have the same strength. The main contribution is made by QFL ($D \approx 2 m$). QTR and QFW may be used to control the tunes, leaving the QFN and QDN families unchanged in order to perturb the injection conditions as little as possible.

The nominal and modified dispersion functions have been calculated with the MAD program (figure 2). The gradient settings for both cases are listed in table 1. A summary of the MAD results is printed in figure 9. The maximum β values (26 and 37 m) are larger than in the nominal machine ($\beta_x \approx \beta_y \approx 15$ m).

In the machine experiment it will be useful to be able to fine tune the value of α by *manual* adjustment of one single family (QFL, see table 2).



Figure 2: Dispersion functions with nominal ($\alpha = 0.03$) and modified ($\alpha = 0$) optics

3 The sensitivity of α to $\Delta p/p$ and betatron amplitudes

By adjusting the chromaticities to values close to zero the variations of α with $\Delta p/p$ could be minimized in the MAD calculation. To this end both chromaticities were first set to zero. Then the orbit length variations across the $\Delta p/p$ range between -.005 and +.005 were minimized by fine tuning the F sextupole current (XNH), yielding $Q'_x = -0.8$ and $Q'_y = +0.7$. The optimum sextupole settings are included (XNH and XNV) in table 1. The maximum length variation in this momentum range could thus be kept below 30 μm (fig. 3).

The residual length variation is of third order in $\Delta p/p$. Simulations show that by adding octupoles in EPA the variation might be further reduced to less than 1 μm .

An additional advantage of cancelling the length variations was discovered when tracking was done for particles with *finite* betatron emittances. Provisional tracking results show that this procedure made the trajectory lengths equally insensitive to the *betatron* amplitudes. Confirmation and explanation of this observation will require further study.



Figure 3: orbit length vs. momentum

4 Evolution of the bunch length

Particles were tracked without r.f. voltage, with initial conditions $\sigma_e = 2.5 \ 10^{-3}$ and $\sigma_t = 7.5 \ ps$. After 1000 turns the bunch length increases to $\sigma_t = 141 \ ps$ (fig. 4) for particles with zero initial betatron amplitude.

The calculation of the exact contribution of the betatron amplitudes, although probably small as mentioned above, requires a more sophisticated description of the EPA bending magnet. The bunch length simulation was done with the combined function bending magnet model currently used for EPA in MAD, which has a *finite* magnetic length and may thus be used to calculate path length differences. On the other hand, this model is not sufficiently symplectic to allow tracking of transverse motion over 1000 turns and a *zero length* bend had to be used for this case. More work is required to build a model which allows simultaneous tracking in both transverse and longitudinal planes.

5 Adjustment of the revolution frequency

A +0.2 mrad kick in both bumpers BSW12 and BSW91 produces an orbit length increase of 1 mm ($\approx 8 ppm$) at the expense of an orbit distortion of only 3.4 mm. This may be useful in future tests where an adjustment of the revolution frequency is required.

Family	Normalized $[m^{-2}]$ o	d Gradient r $[m^{-3}]$		$\begin{array}{c} \text{Current} \\ [A] \end{array}$	
	Nominal	$\alpha = 0$	Nominal	$\alpha = 0$	I ^{max}
QFN QFI QDN QFW QTR QFL XNH XNV	$\begin{array}{c} 0.5659220\\ 0.5335500\\ -0.5689250\\ 1.1104300\\ -0.0521899\\ 1.3803500\\ 8.4380000\\ -6.1390000\\ \end{array}$	$\begin{array}{c} 0.565922\\ 0.533550\\ -0.568925\\ 1.180000\\ 0.280000\\ 1.610630\\ 5.511000\\ -12.866000\end{array}$	$\begin{array}{r} 33.5360\\ 28.7832\\ -33.7140\\ 51.5717\\ -3.0927\\ 74.4651\\ 28.4171\\ -20.6747\end{array}$	33.5360 28.7832 -33.7140 54.8027 16.5926 86.8879 18.5597 -43.3296	55 46 55 95 20 118 91 91
HI.QFD1 HI.QFD2	2.0250000 2.3700000	2.127000 2.381000	94.0471 110.0699	98.7843 110.5808	130 130

Table 1: Quadrupole and sextupole settings with the nominal and with the $\alpha = 0$ optics

6 Injection and envelopes

The optics in the injection area is left *locally* unchanged, as the QFN and QFD settings are not modified. This ensures that the injection bump remains closed.

The optics of the injection line is not sufficiently flexible with the present configuration of the quadrupole families. Therefore a larger betatron mismatch than usual will have to be accepted in the test.

The first-turn beam envelopes at 1 σ are shown in figures 5 and 6. Multi-turn beam envelopes are presented in figures 7 and 8. The first-turn envelopes are globally larger with the modified optics, which may result in a smaller injection efficiency. The multi-turn envelopes with the modified optics are larger only in the long straight sections, but not in the injection kickers and septa.

QFL Current [A]	$\begin{matrix} \alpha \\ [10^{-3}] \end{matrix}$
89.9 88.9 87.9	-12.0 -7.8 -3.8
86.9	0.0
85.9 84.9 83.9	$3.5 \\ 6.8 \\ 10.0$

Table 2: Fine tuning of α using one single quadrupole family

7 Conclusions

The above presented calculations show that EPA can be made isochronous across the momentum range of the incoming bunches. The 0th, 1st, 2d and 3d derivatives of the orbit length w.r.t. $\Delta p/p$ can be cancelled using respectively:

- bumpers BSW12 + 91
- quadrupoles QFL + QTR + QFW,
- sextupole XNH
- an octupole family, if this were available

The results obtained in simulation may be summarized as follows:

- $\alpha = 0$ can be obtained by adjusting 3 existing quadrupole families which do not change the injection conditions
- manual fine tuning of α is possible by adjusting the setting of one single family (QFL)
- tuning the HOR chromaticity (≈ -0.8) using the F sextupole family allows to keep α small in the entire $\Delta p/p$ range -.005 to +.005
- the residual orbit length variations are smaller than $0.030 \ mm$
- the resulting bunch length increase is of the order of 140 ps after 1000 turns



Figure 4: bunch shape at injection (left) and after 1000 turns (right)

Provisional simulation results suggest that with these settings the path length becomes equally independent of the *betatron* amplitudes. Further studies are required to confirm and explain this observation.

References

 T. Risselada, in Proceedings of the Fifth General Accelerator Physics Course, CERN 94-01, 1994, Vol. I, p.313



XENV

XENV

Figure 5: H and V first-turn envelopes at 1 σ with the nominal optics



Figure 6: H and V first-turn envelopes at 1 σ with the modified optics



Figure 7: H and V multi-turn envelopes at 1 σ with the nominal optics



Figure 8: H and V multi-turn envelopes at 1 σ with the modified optics

XENV

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Figure 9: MAD output for EPA with $\alpha = 0$