



## Response of a Li-glass/multi-anode photomultiplier detector to collimated thermal-neutron beams<sup>☆</sup>

E. Rofors<sup>a</sup>, N. Mauritzson<sup>a</sup>, H. Perrey<sup>a,b</sup>, R. Al Jebali<sup>b,c</sup>, J.R.M. Annand<sup>c</sup>, L. Boyd<sup>c</sup>, M.J. Christensen<sup>d</sup>, U. Clemens<sup>f</sup>, S. Desert<sup>e</sup>, R. Engels<sup>f</sup>, K.G. Fissum<sup>a,b,\*</sup>, H. Frielinghaus<sup>f</sup>, C. Gheorghe<sup>g</sup>, R. Hall-Wilton<sup>b,c,h</sup>, S. Jaksch<sup>f</sup>, K. Kanaki<sup>b</sup>, S. Kazi<sup>i</sup>, G. Kemmerling<sup>f</sup>, I. Llamas Jansa<sup>i</sup>, V. Maulerova<sup>a,b,1</sup>, R. Montgomery<sup>c</sup>, T. Richter<sup>d</sup>, J. Scherzinger<sup>a,b,2</sup>, B. Seitz<sup>c</sup>, M. Shetty<sup>d</sup>

<sup>a</sup> Division of Nuclear Physics, Lund University, SE 221 00 Lund, Sweden

<sup>b</sup> Detector Group, European Spallation Source ERIC, SE 221 00 Lund, Sweden

<sup>c</sup> SUPA School of Physics and Astronomy, University of Glasgow, Glasgow G12 8QQ, Scotland, UK

<sup>d</sup> Data Management and Software Centre, European Spallation Source, Ole Maaloes Vej 3, 2200 Copenhagen, Denmark

<sup>e</sup> LLB, CEA, CNRS, Université Paris-Saclay, 91191 Gif-sur-Yvette, France

<sup>f</sup> Jülich Centre for Neutron Science JCNS, Forschungszentrum Jülich, D-52425 Jülich, Germany

<sup>g</sup> Integrated Detector Electronics AS, Gjerdrums Vei 19, N-0484 Oslo, Norway

<sup>h</sup> Dipartimento di Fisica "G. Occhialini", Università degli Studi di Milano-Bicocca, Piazza della Scienza 3, 20126 Milano, Italy

<sup>i</sup> Institute for Energy Technology, Instituttveien 18, 2007 Kjeller, Norway

### ARTICLE INFO

#### Keywords:

SoNDe thermal-neutron detector  
GS20 scintillator  
Li-glass  
H12700A multi-anode photomultiplier  
Position-dependent response  
Neutrons  
European spallation source

### ABSTRACT

The response of a position-sensitive Li-glass scintillator detector being developed for thermal-neutron detection with 6 mm position resolution has been investigated using collimated beams of thermal neutrons. The detector was moved perpendicularly through the neutron beams in 0.5 to 1.0 mm horizontal and vertical steps. Scintillation was detected in an  $8 \times 8$  pixel multi-anode photomultiplier tube on an event-by-event basis. In general, several pixels registered large signals at each neutron-beam location. The number of pixels registering signal above a set threshold was investigated, with the maximization of the single-hit efficiency over the largest possible area of the detector as the primary goal. At a threshold of  $\sim 50\%$  of the mean of the full-deposition peak,  $\sim 80\%$  of the events were registered in a single pixel, resulting in an effective position resolution of  $\sim 5$  mm in X and Y. Lower thresholds generally resulted in events demonstrating higher pixel multiplicities, but these events could also be localized with  $\sim 5$  mm position resolution.

### 1. Introduction

The scientific program to be performed at the European Spallation Source (ESS) [1–4] requires position-sensitive,  $^3\text{He}$ -free [5–8], thermal-neutron detectors with high counting-rate capability. Small-angle neutron-scattering experiments requiring two-dimensional position sensitivity [9–18] will be performed with Solid-state Neutron Detectors (SoNDe) [19–22]. The SoNDe concept employs an array of detector modules to instrument large areas with a reconstruction accuracy of  $\sim 6$  mm on the position of the detected neutron. A SoNDe "module" consists of a thin, thermal-neutron sensitive, Li-glass scintillator

sheet (GS20) attached to a 64-pixel multi-anode photomultiplier tube (MAPMT). Signals resulting from the scintillation light are processed using custom electronics [22]. In the envisioned operation mode at ESS, known as "Time of flight" mode (TOF), these electronics timestamp all pixels having signals above threshold if any single pixel amplitude exceeds threshold. Events involving the firing of a single pixel (multiplicity  $M = 1$  events) are thus straightforward to interpret. At the boundaries between pixels and in the corners, scintillation light sufficient to trigger several pixels (multiplicity  $M > 1$  events) is often registered. The behavior of clusters of bordering pixels in these regions is thus of interest. LEDs and laser light have been used extensively to

<sup>☆</sup> The data set [doi:10.5281/zenodo.4095210](https://doi.org/10.5281/zenodo.4095210) is available for download from <https://zenodo.org/record/4095210>.

\* Corresponding author at: Division of Nuclear Physics, Lund University, SE 221 00 Lund, Sweden.

E-mail address: [kevin.fissum@nuclear.lu.se](mailto:kevin.fissum@nuclear.lu.se) (K.G. Fissum).

<sup>1</sup> Present address: CERN, Route du Meyrin 285, CH-1211 Genève 23, Switzerland and Hamburg University, Mittelweg 177, 20148 Hamburg, Germany.

<sup>2</sup> Present address: Thermo Fisher Scientific Messtechnik GmbH, Frauenaucher Str. 96, 91056 Erlangen, Germany.

<sup>3</sup> The 2.0 and 2.4 Å neutron beams employed in this work had energies lying at the upper end of the cold-neutron energy window and are thus almost thermal. As the neutron interaction cross section at cold energies is larger than that at thermal energies, we anticipate a slightly higher detection efficiency at these marginally lower energies.

study the detailed responses of several different MAPMTs [23–32]. Previously, scans of a  $\sim 1$  mm collimated beam of  $\sim 4$  MeV  $\alpha$ -particles from an  $^{241}\text{Am}$  source [33] and  $\sim 100$   $\mu\text{m}$  diameter beams of 2.5 MeV protons and deuterons [34] have been used to study the position-dependent response of a SoNDe detector prototype in regions well-removed from the edges of the detector acceptance. Thermal neutrons have also been used to perform first tests both on similar detectors [35] and on SoNDe modules [22]. The thermal-neutron interaction with the  $^6\text{Li}$  of the Li-glass has a  $Q$ -value of 4.78 MeV and results in an  $\alpha$ -particle (2.05 MeV) and a triton (2.73 MeV). In this work, a SoNDe module has been systematically scanned through beams of thermal<sup>3</sup> neutrons. The goals were to:

1. complement the existing  $\alpha$ -particle, proton, and deuteron studies of the position-sensitive behavior of the detector, for events triggering multiple pixels to establish the response at the pixel boundaries and the corners where four pixels meet
2. provide thermal-neutron data with  $\sim 1$  mm precision (matching the existing  $\alpha$ -particle studies) on the position sensitivity of the detector for events triggering only one pixel, as a single-pixel mode-of-operation is anticipated as the ESS default
3. map the response of the detector as a function of both threshold and beam position for events which only trigger one pixel
4. determine the detector threshold that maximizes the number of single-pixel events
5. study regions within the detector where the position-reconstruction accuracy for an event better than  $\sim 6$  mm may be obtained for  $M > 1$  events, matching the resolution requirement of the envisioned scientific application of the detector at ESS [17]
6. provide a thermal-neutron dataset at the edge of the detector, for adjacent pixels with the highest gain contrast, to aid our understanding of the SoNDe module at its periphery
7. implement a SoNDe module into a prototype of the ESS data-acquisition architecture.

## 2. Apparatus

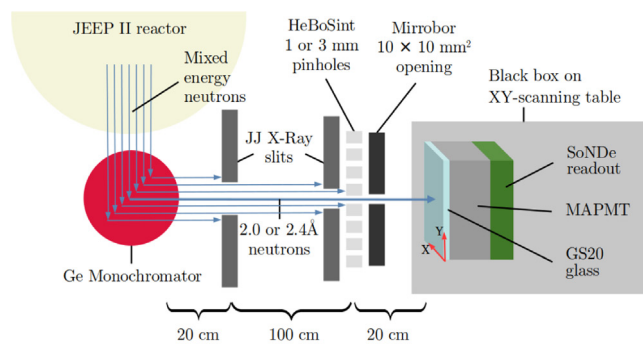
### 2.1. Neutron beams

The measurements were performed at the R2D2 beamline at the JEEP II reactor [36,37] at the Institute for Energy Technology (IFE) [38] in Norway. The setup is illustrated in Fig. 1. The nominal central-beam flux was  $10^5/\text{s}/\text{cm}^2$  with a  $\sim 0.6^\circ$  divergence. Thermal-neutron beams (2.0  $\text{\AA}$ ,  $\sim 18$  meV and 2.4  $\text{\AA}$ ,  $\sim 13$  meV) were defined using a composite Ge wafer monochromator [37]. The resulting thermal-neutron beams drifted  $\sim 20$  cm to the first of a pair of JJ X-ray IB-C80-AIR slits [39] which employed borated-aluminum blades to control the beam flux. The slit spacing was  $\sim 100$  cm, with the downstream slit located  $\sim 20$  cm upstream of the detector. A 5 mm thick HeBoSint mask [40] with pinholes was used to further collimate the beam to either  $\sim 1$  mm or  $\sim 3$  mm in diameter. A stack of three 2 mm thick Mirrobor sheets [41] with a square-shaped  $100\text{mm}^2$  aperture acted as a final barrier to any neutrons surviving the upstream collimation. The resulting beam was incident upon a black box containing the SoNDe module under investigation.

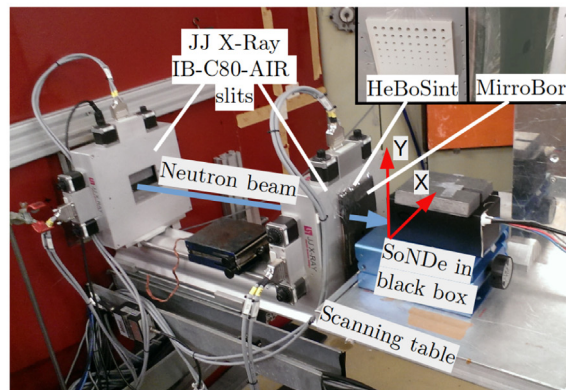
### 2.2. Sonde module

A SoNDe module (Fig. 2) described in the following sections consists of three basic components:

1. a 1 mm thick Li-glass scintillator sheet
2. a H12700 A MAPMT
3. purpose-built SoNDe readout electronics.



(a) Beamline schematic (not to scale).

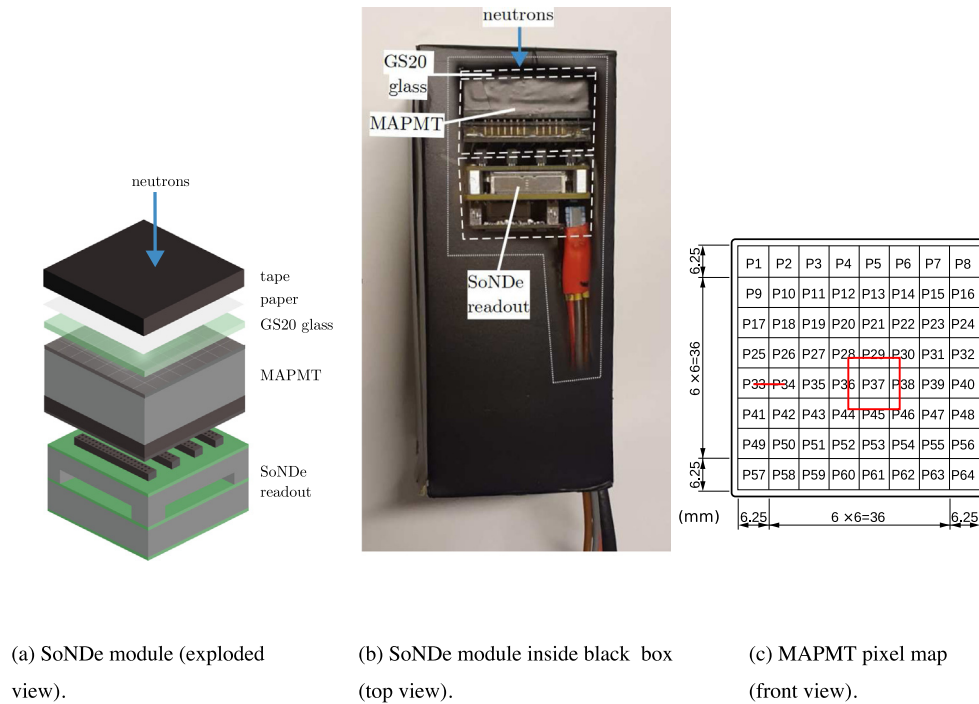


(b) Beamline photo.

Fig. 1. R2D2 thermal-neutron beamline at the JEEP II reactor at IFE. 1(a): The reactor (yellow half circle, top left) produced a continuous beam of neutrons (blue arrows) which was collimated and monochromated (red circle). The beam was then shaped with a pair of slits, a pinhole mask, and a shielding sheet before striking a black box containing the SoNDe module from the left. 1(b): A photograph of the experimental setup. The neutron beam is shown by the blue arrow from the left. The insets (top right) show the pinhole mask and the shielding sheet. The SoNDe module was located in the black box which was mounted on a motorized XY scanning table. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

#### 2.2.1. Li-glass scintillator

The scintillator employed was a cerium-activated lithium-silicate glass known as GS20 [43–46]. Provided by Scintacor [47], it was developed for the detection of thermal neutrons. The  $50\text{mm} \times 50\text{mm} \times 1$  mm sheet with polished faces and rough-cut 1 mm edges was dry-fitted to the MAPMT window and held in place with tape along the edges. The dry-fit approach was chosen to avoid the degradation of any optical-coupling medium. A piece of standard white copy paper ( $136\text{g}/\text{cm}^2$ ) placed over the upstream face of the GS20 diffusely reflected scintillation light back towards the MAPMT, increasing the amount of scintillation light reaching the MAPMT by  $\sim 40\%$ . The (assumed uniform) density of  $^6\text{Li}$  in GS20 is  $1.58 \times 10^{22}$  atoms/ $\text{cm}^3$ . The cross section for the  $n(25\text{meV}) + ^6\text{Li} \rightarrow ^3\text{H}(2.73\text{MeV}) + \alpha(2.05\text{MeV})$  capture reaction is 940 b, which yields a detection efficiency of  $\sim 75\%$  for the 1 mm sheet. The average ranges of the  $^3\text{H}$  and  $\alpha$ -particle in the GS20 are  $34.7\mu\text{m}$  and  $5.3\mu\text{m}$ , respectively [48]. The 4.78 MeV capture reaction results in a  $\sim 6600$  scintillation photon full-deposition peak [22] (roughly equivalent to 20%–30% of anthracene) peaked at  $\sim 390\text{nm}$  [49]. Scintillation-light transport from the GS20 (refractive index 1.55 at 395 nm) across a  $\sim 100\mu\text{m}$  air gap (refractive index 1) due to the concavity of the MAPMT borosilicate-glass window and then into the MAPMT window (refractive index 1.53) is generally inefficient.



**Fig. 2.** A SoNDe module. **2(a):** 3D view of a SoNDe module. From the top: light-tight tape, paper reflector, GS20 glass, MAPMT, and readout electronics. The beam of neutrons (blue arrow) is incident towards the front face of the detector. **2(b):** Photograph of the SoNDe module (dashed white box) in the black box. The GS20 glass, MAPMT, and readout electronics are labeled. The beam of neutrons (blue arrow) is incident towards the front face of the detector. **2(c):** Numbering scheme for the MAPMT pixels (front view) [42]. The region of irradiation with the 3 mm FWHM beam is indicated by the red box centered on P37. The trajectory of irradiation with the 1 mm FWHM beam is indicated by the red line connecting the center of edge pixel P33 to the center of P34. The beam of neutrons is incident into the page. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

### 2.2.2. Multi-anode photomultiplier tube

An  $8 \times 8$  pixel ( $\sim 6 \text{ mm} \times \sim 6 \text{ mm}$  per pixel) 10 stage Hamamatsu type H12700A MAPMT with a borosilicate glass window was employed. The sensitive bi-alkali photocathode area is  $48.5 \text{ mm} \times 48.5 \text{ mm}$ , while the outer dimensions are  $52 \text{ mm} \times 52 \text{ mm}$ , resulting in 87% of the surface being active. The peak quantum efficiency of the photocathode is  $\sim 33\%$  at  $\sim 380 \text{ nm}$ , nicely overlapping the  $\sim 390 \text{ nm}$  (peak) scintillation-light distribution produced by GS20. For the MAPMT used here, at a cathode-to-anode voltage of  $-1000 \text{ V}$ , the Hamamatsu data sheet specified a gain of  $2.09 \times 10^6$ , a dark current of  $2.67 \text{ nA}$ , a (worst-case) factor 1.7 gain difference between pixels, and  $\sim 2\%$  electronic crosstalk between pixels. Calvi et al. [30] report that the electronic crosstalk is actually dependent upon both the pixel and the position within this pixel and that it fluctuates differently near horizontal and vertical edges (see below). The operating voltage was  $-900 \text{ V}$ . Corrections for pixel-to-pixel gain variations were performed offline using a gain map measured with the  $3 \text{ mm}$ ,  $2.4 \text{ \AA}$  neutron beam used to irradiate each pixel center consecutively, as described in [34].

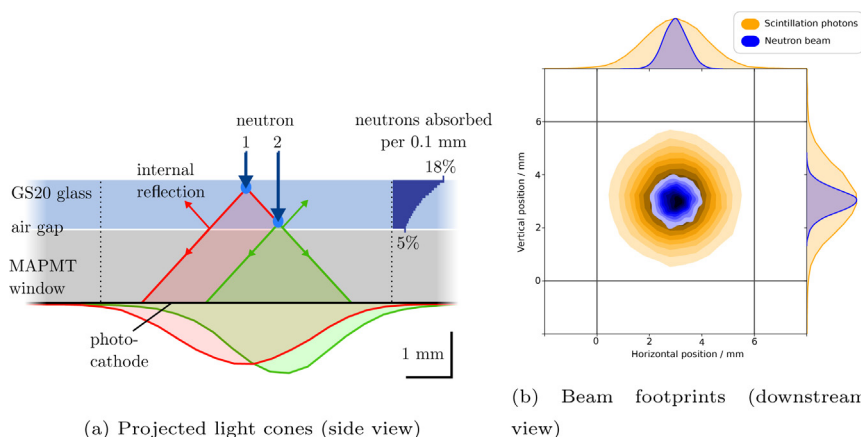
### 2.2.3. Readout electronics

A compact readout module designed for the H12700A MAPMT by IDEAS [50] was employed for data acquisition. The  $113 \text{ g}$  module is  $50 \text{ mm} \times 50 \text{ mm} \times 55 \text{ mm}$  (deep). It consists of two boards: front-end and controller [22]. The front-end board uses four 16-channel IDE3465 ASICs [51] to digitize the MAPMT signals with a precision of 14 bits. The controller board accommodates an FPGA and a MiniIO port for ethernet communication. The electronics can operate in two modes: the TOF mode previously discussed and "All-channel Spectroscopy" (ACS) that was used here. In ACS mode, when any pixel-amplitude threshold was exceeded, the digitized signal amplitudes from all 64 pixels were read out. The ACS rate limitation is  $\sim 10 \text{ kHz}$ , which corresponds to  $\sim 4 \text{ MHz/m}^2$ . A hardware threshold of 500 ADC channels was employed, which corresponds to  $\sim 5\%$  of the mean channel of the  $4.78 \text{ MeV}$  full-energy deposition peak measured by a single pixel for irradiation at its

center. Higher thresholds were applied offline. Control, visualization, and data logging were provided by the ESS Event Formation Unit (EFU) [52–54] and the ESS Daquiri visualization tool [55] both running on a Centos 7 PC. The SoNDe module and the EFU were connected via switched 1 Gbit/s Ethernet. The SoNDe module is configured with TCP/IP and transmits readout data over UDP/IP to the EFU [56] in a manner similar to that anticipated for operation at ESS. The EFU, designed for use by ESS instruments, employs an acquisition that also closely resembles the anticipated operation mode for ESS.

### 2.3. GEANT4 simulation

A detailed computer model of a SoNDe module is nearing completion [57]. This C++ model employs the GEANT4 Monte Carlo toolkit [58] version 4.10.6 [59]. It includes the GS20 sheet together with the glass window and photocathode of the MAPMT. The model may be configured to include a coupling medium such as optical grease between the GS20 and the MAPMT window, but in the present case was set to  $0.05 \text{ mm}$  air to replicate the dry fit employed in this work. The model simulates the interactions of ionizing radiation in the GS20 to the level of the emission of scintillation light and includes the transport of the scintillation photons to the MAPMT cathode. It also includes a model for electronic crosstalk. Electronic crosstalk results from voltage divider biasing, stray capacitances leading to AC coupling between pixel anodes, and charge sharing across neighboring dynode chains, all known to affect the performance of the H12700 MAPMT. It results in signal from the illuminated pixel leaking into neighboring pixels. Defined for each neighboring pixel as the ratio of the induced signal to the signal registered in the illuminated pixel, it has been reported to be up to  $\sim 3\%$  in vertically adjacent pixels and up to  $\sim 7\%$  in horizontally adjacent pixels [30]. The probability of signal leakage has been shown to be the lowest at a pixel center, while the highest probabilities are at the pixel edges. Based on these measurements, electronic crosstalk was modeled on an individual scintillation-photon basis with the crosstalk

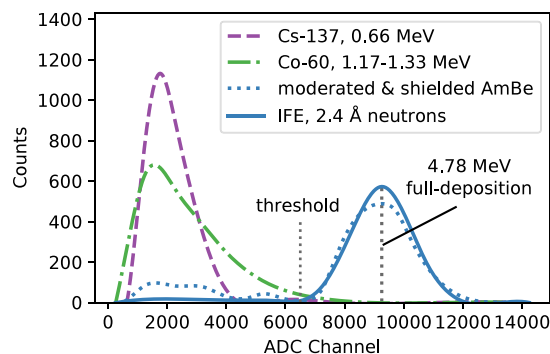


**Fig. 3.** Scintillation-light distributions. **3(a):** 2D projections of GEANT4 simulated scintillation-light cones (triangles) and yields (inverted normal distributions) resulting from interactions at the upstream GS20 surface (neutron 1, red) and the downstream GS20 surface (neutron 2, green). Internal reflection at the air gap between the scintillator and the MAPMT window is also illustrated. Vertical dotted lines indicate pixel edges. The number of interacting neutrons is shown as a function of penetration depth into the GS20 in the top-right corner. **3(b):** GEANT4 simulated 1 mm FWHM thermal neutron beam (blue) and resulting  $\sim 3.5$  mm FWHM scintillation-light distribution (yellow) detected at the MAPMT photocathode for a central irradiation of a pixel. XY projections of the neutron beam and the scintillation-photon distributions are also shown. The vertical and horizontal lines represent the pixel boundaries. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

probability increasing linearly as the pixel edge was approached. For uniform pixel illumination, the crosstalk model was configured so that once all scintillation photons were detected, adjacent pixels each registered 5% of the signal detected in the illuminated pixel. Fig. 3 presents some results from the GEANT4 simulation of the scintillation light. In Fig. 3(a), 2D projections of scintillation-light cones resulting from individual neutrons interacting at the upstream and downstream faces of the GS20 sheet are shown, together with the resulting simulated distributions of scintillation light arriving at the photocathode of the MAPMT. The projections represent the opening angle defined by total internal reflection, and have been drawn to guide the eye. The scintillation photons from the absorption of a single upstream neutron have a normal distribution with  $\sim 3.5$  mm FWHM at the photocathode, and can completely illuminate an entire pixel, even reaching into the adjacent pixels. The absorption of a single downstream neutron results in a  $\sim 3$  mm FWHM normal distribution of scintillation photons at the photocathode. In the first 0.1 mm of the GS20 sheet,  $\sim 18\%$  of the incoming neutrons are absorbed by the  ${}^6\text{Li}$  in the scintillator. This process continues exponentially so that  $\sim 5\%$  of the incident neutrons are absorbed in the last 0.1 mm of the 1 mm thick GS20. There is a 22% chance of a neutron passing through the GS20 without interacting. In Fig. 3(b), the relationship between the 1 mm FWHM extended  $2.0 \text{ \AA}$  neutron beam incident upon the middle of a pixel and the  $\sim 3.5$  mm FWHM distribution of scintillation photons at the photocathode of a pixel is shown. The 3 mm FWHM neutron-beam scintillation-light footprint covers essentially the entire MAPMT pixel. The paper shown in Fig. 2(a) diffusely scatters the scintillation light back towards the photocathode, resulting in a  $\sim 40\%$  increase in the detected yield, and increasing the width of the light cone by  $\sim 2\%$ .

### 3. Measurement

Collimated thermal-neutron beams were used to irradiate the SoNDe module at well-defined positions. After passing through the hole in the Mirrobor sheet, neutrons entered the black box which was positioned on an XY coordinate scanner instrumented with two translation stages (M-IMS600 and M-IMS300) and a motor controller (ESP301), all from MKS Newport Corporation [60]. The SoNDe module was located inside the black box and positioned so that its face was parallel to the upstream side of the black box, and both were perpendicular to the neutron beam. The beam struck the upstream face of the GS20 sheet after passing through a thin layer of tape and white paper. The SoNDe module was stepped through the neutron beams with a stepsize of 0.5–1 mm in the X and Y directions. The anode signals from each of the

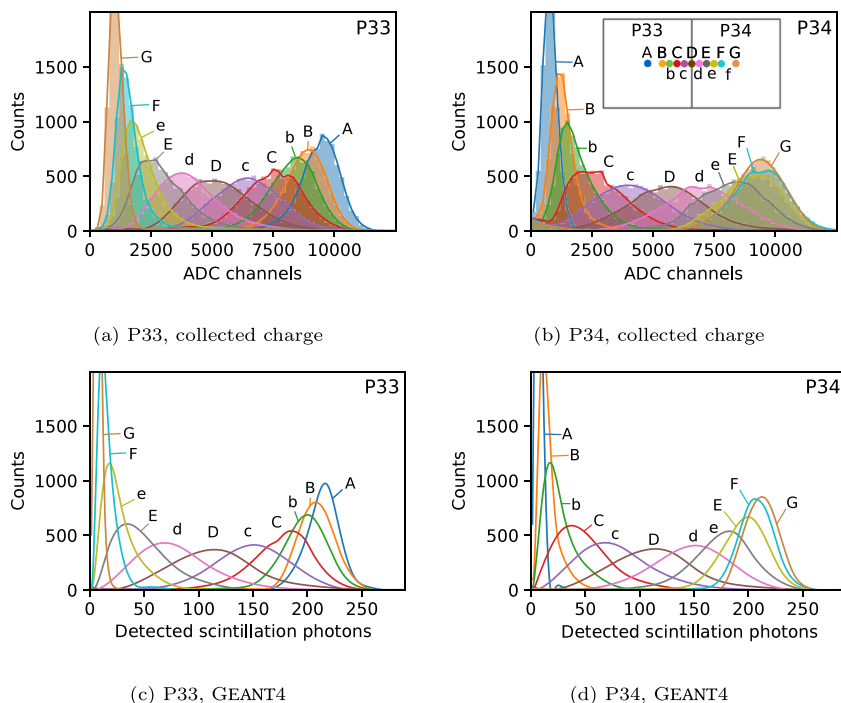


**Fig. 4.**  $\gamma$ -ray and neutron calibration spectra for P37. Events resulting in the largest signal in P37 are displayed. Spectra from sources of  $\gamma$ -rays (long purple dashes, green dash-dots) lie to the left (ADC channel 3000) while spectra from neutrons (solid blue, blue dots) lie to the right (ADC channel 9090). A typical threshold is indicated. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

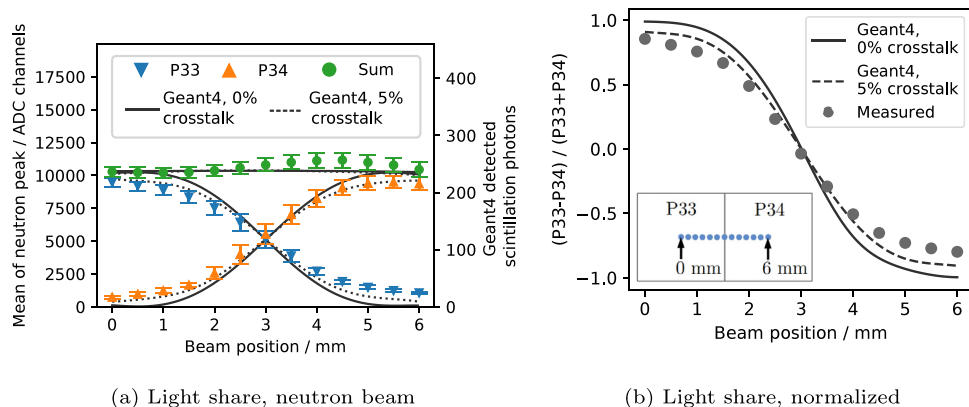
MAPMT pixels were processed using the dedicated SoNDe electronics. Negative polarity analog pulse heights for each event with at least one pixel producing a signal above the 500 ADC channel threshold were recorded. Data were recorded for  $\sim 15$  s (10000 events) at each point on a scan, followed by a motor translation, so that a complete scan of  $2 \times 2$  pixels with 0.5 mm spacing took several hours. Analysis of data was performed using Python-based [61] pandas [62] and SciPy [63] tools.

Fig. 4 shows calibration results from irradiations of the detector with uncollimated  ${}^{60}\text{Co}$  ( $E_\gamma = 1.17, 1.33 \text{ MeV}$ ) and  ${}^{137}\text{Cs}$  ( $E_\gamma = 662 \text{ keV}$ )  $\gamma$ -ray sources, a moderated and heavily  $\gamma$ -ray shielded (but uncollimated) Am/Be neutron source, and the 3 mm,  $2.4 \text{ \AA}$  neutron beam directed at the center of P37. Note that none of the source irradiations occurred in situ, but instead were performed subsequently with the same experimental equipment and setup parameters at the Source-Testing Facility at the University of Lund in Sweden [64]. For each irradiation, events corresponding to the largest signal in P37 were selected. The 4.78 MeV full-deposition peak resulting from neutron capture on  ${}^6\text{Li}$  ( $\sim 6600$  scintillation photons) is located at about ADC channel 9090 ( $\sim 0.53 \text{ keV/channel}$ ). The distribution from  $\sim 1 \text{ MeV}$   $\gamma$ -rays, which are typical backgrounds at accelerator facilities such as ESS, slightly overlap the neutron peak. A discriminator threshold of  $\sim 72\%$





**Fig. 5.** Scintillation-light sharing. Horizontal scan from P33 to P34 using the 1 mm diameter FWHM 2.0 Å neutron beam in 0.5 mm steps. The colors and beam locations are presented in the key (inset, top right panel) which applies to both the measured gain-corrected charge distributions (filled histograms, top panels) and the GEANT4 simulations of the scintillation-light yields (open histograms, bottom panels). Spectra for positions a and f are not shown as they are almost identical to the spectra obtained at the adjacent positions. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)



**Fig. 6.** Scintillation-light sharing. Horizontal scan from P33 to P34 through the 1 mm diameter FWHM 2.0 Å neutron beam in 0.5 mm steps. **6(a):** Data and GEANT4 simulations. Points are the means from fits to the distributions of scintillation light registered in the pixels shown in Fig. 5. The curves are spline fits to the corresponding GEANT4 simulations. The uncertainties in the means are smaller than the symbols. The error bars correspond to  $\pm\sigma/4$  of the fitted distributions. The curves result from the GEANT4 simulations of the scintillation light, and have been normalized to the measurements as previously discussed. The curves show the simulations including different levels of electronic crosstalk between pixels. **6(b):** Light-sharing ratio derived from **6(a)**. The uncertainties are smaller than the widths of the lines. The curves are as in Fig. 6(a).

of the mean of the full-deposition peak (ADC channel 6500) discriminates against ~93% of the detected ~1 MeV  $\gamma$ -rays while retaining the neutron peak.

#### 4. Results

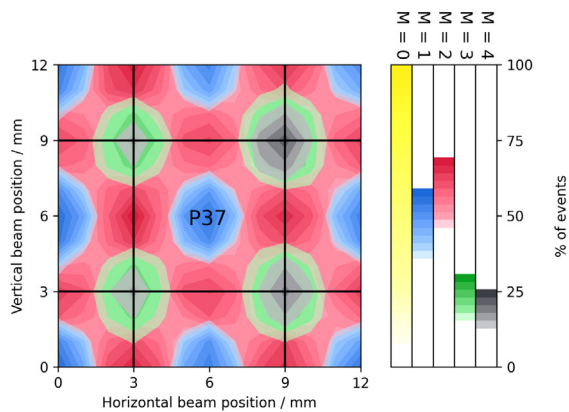
As previous work [23–28,31,33,34] has clearly demonstrated that MAPMT pixel-gain maps depend strongly upon the method of illumination, all of the results presented below have been pedestal subtracted and gain corrected with pixel-gain maps produced from 3 mm, 2.4 Å neutron-beam irradiations of the pixel centers. Fig. 5 shows data and GEANT4 simulations for a horizontal scan of the SoNDe module through the 1 mm FWHM, ~2.0 Å neutron beam. The module was moved from position A (center, edge pixel P33) to position G (center, P34)

in 0.5 mm steps. For 13 scan positions, neutron pulse-height spectra and GEANT4-simulated scintillation-light yields are shown for P33 (Figs. 5(a), 5(c)) and P34 (Figs. 5(b), 5(d)). The simulations include the nominal 5% electronic crosstalk contribution previously discussed (see also the discussion associated with Fig. 6 below). The location of the neutron beam determines the amount of scintillation light collected by a single pixel. The single-pixel signal amplitude is largest when the neutron beam strikes the pixel center, smaller when the neutron beam strikes the pixel edge, and smallest when the neutron beam strikes the pixel corner. Scintillation produced when the beam strikes the border between two pixels is equally shared, which results in equal signal amplitudes after pedestal and gain correction. The excellent agreement between the data and the simulation indicates that both the scintillation-light sharing and the resulting pixel response are well-understood.

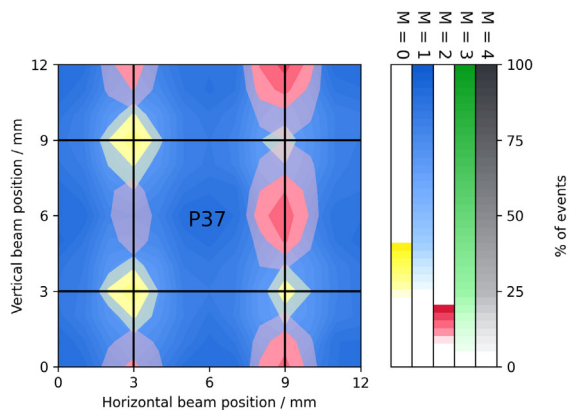
Fig. 6 illustrates scintillation-light sharing by edge pixel P33 and adjacent P34 as the SoNDe module was translated horizontally through the 1 mm FWHM,  $\sim 2.0\text{\AA}$  neutron beam. Fig. 6(a) shows the means of the pulse-height distributions (from Fig. 5) obtained for each beam position. The scan shows that signal leakage to adjacent pixels is  $\sim 7\%$ – $12\%$  when the neutron beam strikes the center of either pixel. This represents a larger spread of scintillation signal into the adjacent pixel than was the case for previous investigations of relatively central pixels with charged-particle beams [33,34] and may be related to the diffusely reflecting white paper placed at the front face of the GS20 sheet.  $\alpha$ -particle scan results for (non-edge) P36, P37, P44, and P45 [33] demonstrated summed gain-corrected charge distributions that were flat across the pixels and boundary regions. Proton- and deuteron-scan results for (non-edge) P37 and P38 [34] demonstrated summed gain-corrected charge distributions that were slightly convex and centered at the pixel edge. This was because the pixels together collected slightly more of the scintillation light produced from an event at the boundary between them than they collected from an event at the center of either pixel, with the missing light collected by the surrounding pixels. Here, the measured distributions may indicate a light-collection enhancement when P34 is irradiated. Simulations including the nominal 5% level of electronic crosstalk underestimate the amount of signal leaking into the adjacent pixel. Crosstalk measurements [30] suggest that 5% is already an overestimate. Fig. 6(b) shows the light-sharing ratio between P33 and P34 defined as  $(P33-P34)/(P33+P34)$ . For the nominal 5% level of electronic crosstalk, the simulation results in too much signal in the irradiated pixel relative to the adjacent pixel. Agreement at the border between pixels is very good. Very recent simulations [57] indicate that the matt white reflector (paper) upstream of the scintillator and the glass surface finish will likely contribute to the redistribution of scintillation light in a fashion similar to electronic crosstalk. Unfolding these effects will require a detailed study with a laser.

The pixel-hit multiplicity ( $M = 1$ ,  $M = 2$ , etc.) for events as a function of the beam-spot position has previously been studied with scans of  $\sim 1$  mm FWHM beams of  $\alpha$ -particles [33] and  $\sim 100\text{ }\mu\text{m}$  diameter beams of protons and deuterons [34]. A hit occurred if a pixel amplitude exceeded a threshold which was varied offline. Here, the 1 mm FWHM  $2.0\text{\AA}$  neutron beam was employed in a complementary study. Neutron-beam irradiations with a stepsize of 1 mm in X and Y were performed resulting in a  $13 \times 13$  matrix of data. Fig. 7 displays 2D position dependence of multiplicity distributions obtained near P37 for software thresholds of 1360 (Fig. 7(a)) and 4545 (Fig. 7(b)) ADC channels, which correspond to  $\sim 15\%$  and  $\sim 50\%$  of the mean of the P37 pixel-centered full-deposition neutron peak (4.78 MeV, 0.52 keV/channel), respectively. For the 1360 ADC channel threshold,  $\sim 2\%$  of events are lost, and  $M = 1$  events ( $\sim 22\%$ ) lie within  $\pm 1$  mm of the pixel center. The  $M = 2, 3$ , and 4 data are all localized to  $\sim 5$  mm, within the 6 mm position-resolution constraint for SoNDe operation at ESS. A threshold of 4545 ADC channels results in  $\sim 18\%$  event loss, maximizes both the number of  $M = 1$  events ( $\sim 78\%$ ) and the area of the detector where  $M = 1$  events may be detected. Fig. 7(c) shows event multiplicity as a function of applied threshold for  $M = 1-4$ . The curves all have well-defined maxima so that the multiplicity  $M$  for a dataset can be selected by enforcing the appropriate threshold. For example, for a dataset of  $\sim 78\% M = 1$ ,  $\sim 4\% M = 2$ , and a negligible number of  $M = 3, 4$  events, a threshold of 4545 ADC channels must be applied. A result of operating the SoNDe module with this relatively high threshold is that  $\sim 18\%$  of events have  $M = 0$  and are thus lost.

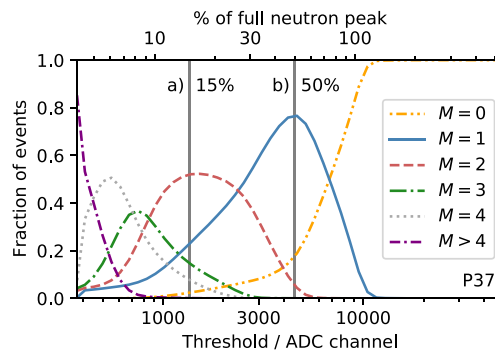
Fig. 8 shows the GEANT4-simulated position-reconstruction accuracy as a function of pixel-hit multiplicity for a 15% threshold corresponding to 33 scintillation photons (Fig. 8(a)) and a 50% threshold corresponding to 100 scintillation photons (Fig. 8(b)). The accuracy is defined as the distance between the reconstructed and simulated capture vertex. The area in the vicinity of P37 was uniformly illuminated with  $2.0\text{\AA}$  neutrons. For each event, the reconstructed hit position was given by the average of the locations of the pixel centers of all pixels registering



(a) Spatial, threshold 15% of full peak (1360 ADC channels)



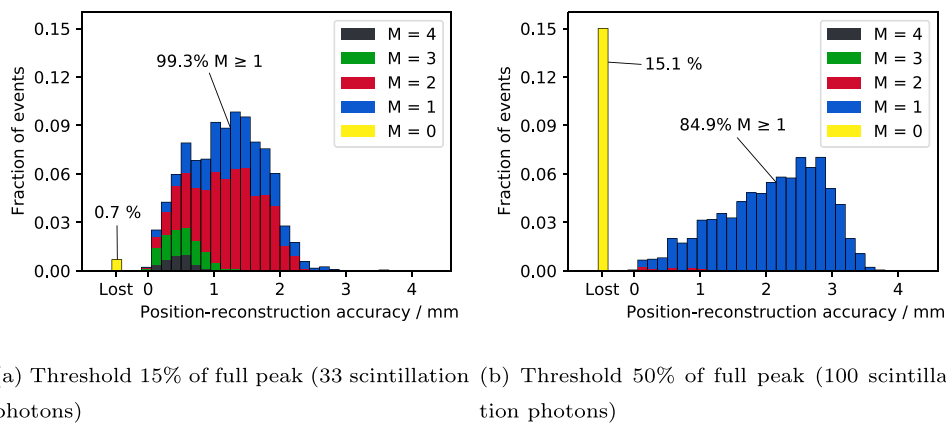
(b) Spatial, threshold 50% of full peak (4545 ADC channels)



(c) Fractional, various thresholds

Fig. 7. Multiplicity distributions for the 1 mm FWHM  $2.0\text{\AA}$  beam incident on P37 and the surrounding pixels. In 7(a) (threshold 1360 ADC channels, 15% of the full-deposition peak) and 7(b) (threshold 4545 ADC channels (50% of the full-deposition peak), the black lines denote the pixel boundaries. Yellow indicates  $M = 0$  events, blue indicates  $M = 1$  events, red indicates  $M = 2$  events, green indicates  $M = 3$  events, and gray indicates  $M = 4$  events. The lighter the shade of the color, the fewer the number of events. 7(c) presents the fraction of events registered in P37 for each multiplicity as a function of threshold common to all pixels, with the 1360 and 4545 ADC channel thresholds shown as vertical lines. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

a signal above threshold. This method results in  $M = 1$  events being “assigned” a position in the center of the pixel above threshold, while  $M = 2$  events are positioned at the edge between the two pixels above threshold.  $M = 3$  events are positioned within the triangle defined by the pixel centers, offset 1 mm in both X and Y from the common corner, and  $M = 4$  events are positioned at the common corner. For the 33 scintillation-photon threshold,  $\sim 1\%$  of events are lost. Event positions for multiplicities  $M = 1$  ( $\sim 31\%$ ) and  $M = 2$  ( $\sim 56\%$ ) events



**Fig. 8.** GEANT4 simulations of the position-reconstruction accuracy for 2.0 Å neutrons incident on P37 and the surrounding pixels. In 8(a) (threshold 33 scintillation photons, 15% of the full-deposition peak) and 8(b) (threshold 100 scintillation photons, 50% of the full-deposition peak), the color code is the same as in Fig. 7. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

are reconstructed to better than  $\sim 3$  mm, while  $M = 3$  ( $\sim 9\%$ ) and  $M = 4$  ( $\sim 4\%$ ) events are reconstructed to better than  $\sim 1$  mm. For the 100 scintillation-photon threshold,  $\sim 15\%$  of events are lost.  $M = 1$  ( $\sim 84\%$ ) events are reconstructed to better than  $\sim 4$  mm, and  $M = 2$  ( $\sim 1\%$ ) events are reconstructed to better than  $\sim 1$  mm. While agreement between the simulation and the data is presently not perfect, it is very encouraging, and development continues.

## 5. Summary and discussion

Collimated beams of 13 meV and 18 meV neutrons from the IFE reactor (Fig. 1) have been used to investigate the position-dependent response of a pixelated neutron detector known as a SoNDe module (Fig. 2). A SoNDe module consists of a 1 mm thick sheet of GS20 scintillating glass coupled to a 64 pixel H12700 A MAPMT with dedicated readout electronics. The software layer of the data acquisition (EFU) was a prototype under development for ESS. The amplitudes of the pixel signals were investigated for different irradiation positions by scanning the module through the beam in steps of 0.5–1 mm using a motor-driven XY table. A GEANT4 model of the SoNDe module greatly aided in the interpretation of the data (Fig. 3). The amount of scintillation light detected by the MAPMT was increased by  $\sim 40\%$  by placing a sheet of diffusely reflecting white paper at the front face of the GS20.  $\gamma$ -rays and neutrons could generally be discriminated with a simple threshold cut (Fig. 4). The amplitudes of the gain-corrected signals were highly dependent on where the neutron beam struck the detector (Fig. 5). When directed towards a central-pixel region,  $\sim 5\%$  of the signal was detected in an adjacent pixel. However, within  $\sim 1$  mm of the boundary,  $\sim 30\%$  of the signal was registered in the adjacent pixel. At the boundary, the signal was evenly split between pixels. Overall agreement between the data and GEANT4 simulations was good when a 5% level of interpixel electronic crosstalk was considered. The signal in a pixel adjacent to an edge pixel when the edge pixel was irradiated was underestimated (Fig. 6). The discrepancy may be due to divergent reflections of scintillation light at the unpolished, non-uniform edges of the GS20 wafer. This effect is currently not modeled in the simulation. For different beam positions, the effect of raising the pixel threshold on the hit multiplicity was studied (Fig. 7). When the threshold was set at  $\sim 50\%$  of the mean of the neutron full-deposition peak,  $\sim 78\%$  of the data had  $M = 1$ ,  $\sim 4\%$  were  $M = 2$ , and  $\sim 18\%$  were undetected. Increasing the threshold to higher values resulted in  $M = 1$  event loss and a reduction of the sensitive area of the detector. Decreasing the threshold to  $\sim 15\%$  of the mean of the neutron full-deposition peak resulted in  $\sim 2\%$  event loss,  $\sim 22\%$   $M = 1$  data, and  $\sim 66\%$   $M > 1$  data. The GEANT4 simulation was employed to investigate the position-reconstruction accuracy (Fig. 8) of the measured multiplicity regions shown in Fig. 7. For the threshold set at  $\sim 50\%$  of the mean of the

neutron full-deposition peak, the majority of the events registered with  $M = 1$  could be reconstructed to better than 3.5 mm. The majority of the events registered with  $M = 2$  could be reconstructed to better than 1 mm. For the threshold set at  $\sim 15\%$  of the mean of the neutron full-deposition peak, the majority of the  $M = 1$  and  $M = 2$  events could be reconstructed to better than 2.5 mm, and the majority of the  $M = 3$  and  $M = 4$  events could be reconstructed to better than 1 mm. Thus, all the IFE data could be reconstructed with an accuracy better than the 6 mm position resolution required for the operation of SoNDe at ESS.

## CRediT authorship contribution statement

**E. Rofors:** Conceptualization, Methodology, Software, Validation, Formal analysis, Investigation, Data Curation, Writing - original draft, Writing - review & editing, Visualization, Project administration. **N. Mauritzson:** Investigation. **H. Perrey:** Methodology, Software, Investigation, Data Curation, Visualization, Supervision. **R. Al Jebali:** Conceptualization, Methodology, Investigation, Resources, Supervision. **J.R.M. Annand:** Conceptualization, Methodology, Software, Validation, Formal Analysis, Writing - original draft, Writing - review & editing, Visualization, Supervision. **L. Boyd:** Supervision. **M.J. Christensen:** Software. **U. Clemens:** Resources. **S. Desert:** Supervision. **R. Engels:** Supervision. **K.G. Fissum:** Conceptualization, Methodology, Validation, Formal analysis, Resources, Writing - original draft, Writing - review & editing, Visualization, Supervision, Project administration, Funding acquisition. **H. Frielinghaus:** Resources. **C. Gheorghe:** Resources, Supervision. **R. Hall-Wilton:** Conceptualization, Resources, Supervision, Project administration, Funding acquisition. **S. Jaksch:** Conceptualization, Resources, Supervision, Project administration, Funding acquisition. **K. Kanaki:** Conceptualization, Resources, Supervision, Project administration, Funding acquisition. **S. Kazi:** Investigation. **G. Kemmerling:** Resources, Supervision. **I. Llamas Jansa:** Investigation. **V. Maulerova:** Investigation. **R. Montgomery:** Supervision. **T. Richter:** Software, Resources, Supervision. **J. Scherzinger:** Supervision. **B. Seitz:** Supervision. **M. Shetty:** Software.

## Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Acknowledgments

Support for this project was provided by the European Union via the Horizon 2020 Solid-State Neutron Detector Project (Proposal ID 654124) and the BrightNESS Project (Proposal ID 676548). Support

was also provided by the UK Science and Technology Facilities Council (Grant No. ST/P004458/1) and the UK Engineering and Physical Sciences Research Council Centre for Doctoral Training in Intelligent Sensing and Measurement (Grant No. EP/L016753/1). The authors thank the staff of IFE for providing the beams of neutrons.

## References

- [1] M. Lindroos, et al., Nucl. Instrum. Methods Phys. Res. B. 269 (2011) 3258, <http://dx.doi.org/10.1088/1402-4896/aa9bff>.
- [2] S. Peggs, et al., ESS Technical Design Report, 2013. ISBN 978-91-980173-2-8. [https://europenspallationsource.se/sites/default/files/downloads/2017/09/TDR\\_online\\_ver\\_all.pdf](https://europenspallationsource.se/sites/default/files/downloads/2017/09/TDR_online_ver_all.pdf).
- [3] R. Garoby, et al., Phys. Scr. 93 (2018) 014001, <http://dx.doi.org/10.1016/j.nimb.2011.04.012>.
- [4] The European Spallation Source (ESS), Lund, Sweden, European Spallation Source ERIC, <https://europenspallationsource.se/>. (Accessed 19 September 2020).
- [5] R.T. Kouzes, PNNL-18388 Pacific Northwest National Laboratory, Richland, WA, 2009. <http://dx.doi.org/10.2172/956899>.
- [6] D.A. Shea, D. Morgan, Technical Report R41419, Congressional Research Service, 2010.
- [7] K. Zeitelhack, Neutron News 23 (4) (2012) 10–13, <http://dx.doi.org/10.1080/10448632.2012.725325>.
- [8] O. Kirstein, et al., PoS Proc. Sci. (Vertex2014) 029 (2014) [arXiv:1411.6194](https://arxiv.org/abs/1411.6194) [physics.ins-det].
- [9] M. Heiderich, et al., Nucl. Instrum. Methods Phys. Res. A 305 (1991) 423, [http://dx.doi.org/10.1016/0168-9002\(91\)90562-5](http://dx.doi.org/10.1016/0168-9002(91)90562-5).
- [10] R. Engels, et al., IEEE Trans. Nucl. Sci. 44 (3) (1997) <http://dx.doi.org/10.1109/23.603701>.
- [11] R. Engels, et al., IEEE Trans. Nucl. Sci. 45 (3) (1998) 6666, <http://dx.doi.org/10.1109/NSSMIC.1997.672706>.
- [12] R. Engels, et al., IEEE Trans. Nucl. Sci. 46 (4) (1999) <http://dx.doi.org/10.1109/23.790694>.
- [13] G. Kemmerling, et al., IEEE Trans. Nucl. Sci. 48 (2001) 1114, <http://dx.doi.org/10.1109/23.958733>.
- [14] R. Engels, et al., IEEE Trans. Nucl. Sci. 49 (3) (2002) <http://dx.doi.org/10.1109/TNS.2002.1039590>.
- [15] G. Kemmerling, et al., IEEE Trans. Nucl. Sci. 51 (2004) 1098, <http://dx.doi.org/10.1109/TNS.2004.829576>.
- [16] G. Kemmerling, et al., IEEE Nucl. Sci. Symp. Conf. Rec 03CH 37515 (2003) 722, <http://dx.doi.org/10.1109/NSSMIC.2003.1351801>.
- [17] S. Jaksch, et al., Nucl. Instrum. Methods Phys. Res. A 762 (2014) 22, <http://dx.doi.org/10.1016/j.nima.2014.04.024>.
- [18] A.V. Feoktystov, et al., J. Appl. Crystallogr. 48 (2015) 61, <http://dx.doi.org/10.1107/S1600576714025977>.
- [19] Scintillation detector with a high count rate, Patent no. PCT/EP2015/074 200, 2015.
- [20] Solid-State Neutron Detector – A new Neutron Detector for High-Flux Applications. EU Publications Office. <https://cordis.europa.eu/project/id/654124>. (Accessed 19 August 2020).
- [21] S. Jaksch, et al., [arXiv:1707.08679](https://arxiv.org/abs/1707.08679) [physics.ins-det].
- [22] S. Jaksch, et al., Proc. Int. Conf. Neutron Optics (NOP2017), JPS Conf. Proc. 22 (2018) 011019. <http://dx.doi.org/10.7566/JPSCP.22.011019>.
- [23] S. Korpar, et al., Nucl. Instrum. Methods Phys. Res. A 442 (2000) 316, [http://dx.doi.org/10.1016/S0168-9002\(99\)01242-5](http://dx.doi.org/10.1016/S0168-9002(99)01242-5).
- [24] K. Rielage, et al., Nucl. Instrum. Methods Phys. Res. A 463 (2001) 149.
- [25] T. Matsumoto, et al., Nucl. Instrum. Methods Phys. Res. A 521 (2004) 367, <http://dx.doi.org/10.1016/j.nima.2003.11.384>.
- [26] K. Lang, et al., Nucl. Instrum. Methods Phys. Res. A 545 (2005) 852, <http://dx.doi.org/10.1016/j.nima.2005.02.041>.
- [27] P. Abbon, et al., Nucl. Instrum. Methods Phys. Res. A 595 (2008) 177, <http://dx.doi.org/10.1016/j.nima.2008.07.074>.
- [28] R.A. Montgomery, et al., Nucl. Instrum. Methods Phys. Res. A 695 (2012) 326, <http://dx.doi.org/10.1016/j.nima.2011.11.026>.
- [29] R.A. Montgomery, Nucl. Instrum. Methods Phys. Res. A 732 (2013) 732, <http://dx.doi.org/10.1016/j.nima.2013.08.012>.
- [30] M. Calvi, et al., J. Instrum. 10 (2015) P09021, <http://dx.doi.org/10.1088/1748-0221/10/09/p09021>.
- [31] R.A. Montgomery, et al., Nucl. Instrum. Methods Phys. Res. A 790 (2015) 28, <http://dx.doi.org/10.1016/j.nima.2015.03.068>.
- [32] X. Wang, et al., Chinese Phys. C 40 (2016) 086003, <http://dx.doi.org/10.1088/1674-1137/40/8/086003>.
- [33] E. Rofors, et al., Nucl. Instrum. Methods Phys. Res. A 929 (2019) 90, <http://dx.doi.org/10.1016/j.nima.2019.03.014>.
- [34] E. Rofors, et al., Nucl. Instrum. Methods Phys. Res. A 982 (2020) 164604, <http://dx.doi.org/10.1016/j.nima.2020.164604>.
- [35] F. Zai-Wei, et al., Chinese Phys. C 36 (2012) 1095, <http://dx.doi.org/10.1088/1674-1137/36/11/010>.
- [36] Jeep II Reactor at Kjeller, Norway. <https://ife.no/en/laboratory/jeep-ii-reactor-at-kjeller-norway/>. [Accessed 19 September 2020].
- [37] B.C. Hauback, et al., J. Neutron Res. 8 (2000) 215, <http://dx.doi.org/10.1080/10238160008200055>.
- [38] Institute for Energy Technology (IFE). Kjeller, Norway. <http://ife.no/en/>. (Accessed 20 September 2020).
- [39] JJ X-ray Danish Science Design JJ X-ray IB-C80-Air slits. <https://www.jjxray.dk/p/ib-c80-air/>. (Accessed 19 September 2020).
- [40] Henze Boron Nitride Products. HeBoSint sintered Boron Nitride Components. <https://www.henze-bnp.com/en/hebosint-boron-nitride-components.php>. (Accessed 19 September 2020).
- [41] Mirrotron Ltd. Radiation Shielding. <https://mirrotron.com/en/products/radiation-shielding>. (Accessed 19 September 2020).
- [42] Flat Panel Type Multianode PMT Assembly H12700 Series/H14220 Series, Iwata City Japan, Hamamatsu Photonics K.K. [https://www.hamamatsu.com/resources/pdf/etd/H12700\\_H14220\\_TPMH1379E.pdf](https://www.hamamatsu.com/resources/pdf/etd/H12700_H14220_TPMH1379E.pdf). (Accessed 19 August 2020).
- [43] F.W.K. Firk, et al., Nucl. Instrum. Methods 13 (1961) 313, [http://dx.doi.org/10.1016/0029-554X\(61\)90221-X](http://dx.doi.org/10.1016/0029-554X(61)90221-X).
- [44] A.R. Spowart, Nucl. Instrum. Methods 135 (1976) 441, [http://dx.doi.org/10.1016/0029-554X\(76\)90057-4](http://dx.doi.org/10.1016/0029-554X(76)90057-4).
- [45] A.R. Spowart, Nucl. Instrum. Methods 140 (1977) 19, [http://dx.doi.org/10.1016/0029-554X\(77\)90059-3](http://dx.doi.org/10.1016/0029-554X(77)90059-3).
- [46] E.J. Fairley, et al., Nucl. Instrum. Methods 150 (1978) 159, [http://dx.doi.org/10.1016/0029-554X\(78\)90360-9](http://dx.doi.org/10.1016/0029-554X(78)90360-9).
- [47] 6-Lithium Enriched Glass Scintillators, Cambridge, England, Scintacor. <https://scintacor.com/products/6-lithium-glass/>. (Accessed 19 September 2020).
- [48] B. Jamieson, et al., Nucl. Instrum. Methods Phys. Res. A 790 (2015) 6, <http://dx.doi.org/10.1016/j.nima.2015.04.022>.
- [49] C.W.E. Van Eijk, et al., Nucl. Instrum. Methods Phys. Res. A 529 (2004) 260, <http://dx.doi.org/10.1016/j.nima.2004.04.163>.
- [50] Integrated Detector Electronics AS. Oslo, Norway. <https://ideas.no/>. (Accessed 19 September 2020).
- [51] D. Meier, et al., <http://dx.doi.org/10.1109/NSSMIC.2013.6829764>.
- [52] A.H.C. Mukai, et al., J. Instrum. 13 (2018) T10001.
- [53] M.J. Christensen, et al., J. Instrum. 13 (2018) T11002, <http://dx.doi.org/10.1088/1748-0221/13/11/T11002>.
- [54] M.J. Christensen, et al., Brightness deliverable report 5.6, software neutron event data processing, 2018, <http://dx.doi.org/10.17199/brightness.d5.6>.
- [55] M. Shetty, A. Mukai, J. Nilsson, v.0.0.1 available at Zenodo. <http://dx.doi.org/10.5281/zenodo.1298512>.
- [56] M.J. Christensen, T. Richter, <http://dx.doi.org/10.1088/1748-0221/15/09/T09005>.
- [57] L. Boyd, et al., Nucl. Instrum. Methods Phys. Res. A. (2020) [arXiv:2010.15695](https://arxiv.org/abs/2010.15695) [physics.ins-det] (submitted for publication).
- [58] S. Agostinelli, et al., Nucl. Instrum. Methods Phys. Res. A. 506 (2003) 250, [http://dx.doi.org/10.1016/S0168-9002\(03\)01368-8](http://dx.doi.org/10.1016/S0168-9002(03)01368-8).
- [59] J. Allison, et al., IEEE Trans. Nucl. Sci. 53 (2006) 270, <http://dx.doi.org/10.1109/TNS.2006.869826>.
- [60] MKS Newport Corporation. Photonics Technologies. <https://www.newport.com/>. (Accessed 19 September 2020).
- [61] G. van Rossum, F.L. Drake (Eds.), Python reference manual, pythonlabs, 2001, Available at <https://www.python.org/>.
- [62] The Pandas Development Team, Pandas 1.1.0. Available at <https://pandas.pydata.org/>. <http://dx.doi.org/10.5281/zenodo.3964380>. (Accessed 19 August 2020).
- [63] W. McKinney, Proceedings of the 9th Python in Science Conference, 445 (2010) 56. <http://dx.doi.org/10.25080/Majora-92bf1922-00a>.
- [64] F. Messi, et al., IAEA Technical Meeting on Modern Neutron Detection (F1-TM-55243), [arXiv:1711.10286](https://arxiv.org/abs/1711.10286) [physics.ins-det].