

Rare and semileptonic heavy flavour decays at LHCb

 B^0

 B_s^0

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On behalf of the LHCb collaboration

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Why search indirectly for new interactions?





- Precise predictions in the SM
- Rare phenomena \rightarrow New interactions can be major contribution
- New interactions can have different symmetries from the SM
- Charm and beauty probe complementary couplings

Example Scalar interaction Higgs-like boson C_S, C_P Vector interaction Z' C_V, C_A Over-constraining new interaction couplings is crucial to understand their origin

LHCb experiment





- pp collisions at $\sqrt{s} = 7, 8, 13$ TeV
- 3 (6) fb^{-1} in Run 1 (Run 2)

$$B^0_{d,s} \to \mu^+ \mu^-$$
 decays

Very rare decays

 $B(B_s^0 \to \mu^+ \mu^-) = (3.66 \pm 0.14) \times 10^{-9}$ $B(B^0 \to \mu^+ \mu^-) = (1.03 \pm 0.05) \times 10^{-10}$

[Beneke, Bobeth, Szafron, JHEP10(2019) 232] Similar predictions using the correlation with B mixing,

[Buras, Venturini -2109.11032]

- Summer 2020: 3 LHC experiments combined
- 2.1 σ from SM in the 2D plane

New LHCb analysis

- Full statistics: 9 fb⁻¹, two-fold increase in statistics w.r.t previous analysis
- Branching fractions and effective lifetime measurement
- Added $B_s^0 \to \mu^+ \mu^- \gamma$ search



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$B^0_{d,s} \to \mu^+ \mu^-$ analysis with full statistics

- Normalised to two channels: $B^+ \to J/\psi K^+$ and $B^0 \to K^+ \pi^-$
- Multivariate operator against combinatorial background
- Tight PID calibrated on data against misID
- Significant improvement in hadronisation fraction $\frac{f_s}{f_d}(13 \text{ TeV}) = 0.2539 \pm 0.0079$ from combined measurement [LHCb-PAPER-2020-046 - PRD 104, 032005 (2021)]



 $B^0_{d,s} \to \mu^+ \mu^-$ analysis with full statistics Results



$$\begin{split} \mathcal{B}(B^0_s \to \mu^+ \mu^-) &= \left(3.09^{+0.46+0.15}_{-0.43-0.11}\right) \times 10^{-9} \\ \mathcal{B}(B^0 \to \mu^+ \mu^-) &= \left(1.2^{+0.8}_{-0.7} \pm 0.1\right) \times 10^{-10} < 2.6 \times 10^{-10} \\ \mathcal{B}(B^0_s \to \mu^+ \mu^- \gamma)_{m\mu\mu > 4.9 \, \text{GeV}} &= (-2.5 \pm 1.4 \pm 0.8) \times 10^{-9} < 2.0 \times 10^{-9} \\ \text{No significant signal for } B^0 \to \mu^+ \mu^- \text{ and } B^0_s \to \mu^+ \mu^- \gamma, \text{ upper limits at } 95\% \\ \text{First world limit on } B^0_s \to \mu^+ \mu^- \gamma \text{ decay} \\ \text{Measured effective lifetime } \tau_{\text{eff}}(B^0_s \to \mu^+ \mu^-) = 2.07 \pm 0.29 \pm 0.03 \, \text{ps} \\ \text{Consistent at } 1.5\sigma \text{ and } 2.2\sigma \text{ with the heavy and light } B^0_s \text{ eigenstates lifetimes} \end{split}$$



- Prior to LHC(b) orders of magnitude enhancements of the $B^0_{d,s} \to \mu^+ \mu^$ branching fractions were allowed
- Now closed to about 20% distance
- This tightens the phase-space for possible new physics that would cause (pseudo)-scalar or axial-vector $bs\mu\mu$ couplings

- Very rare baryonic modes
- Run 2 analysis, normalised to $B^0 \to K^+ \pi^-$ decays
- $B^0 \rightarrow p\bar{p}$ (re)-observed with large significance and branching fraction $\mathcal{B}(B^0 \rightarrow p\bar{p}) = (1.27 \pm 0.15 \pm 0.05 \pm 0.04) \times 10^{-8}$
- Rarer $B_s^0 \rightarrow p\bar{p}$ not observed, with upper limit $\mathcal{B}(B_s^0 \rightarrow p\bar{p}) < 4.37(5.03) \times 10^{-9}$ at 90% (95%)
- Combined Run 1-2 branching fraction $\mathcal{B}(B^0 \to p\bar{p}) = (1.27 \pm 0.13 \pm 0.05 \pm 0.03) \times 10^{-8}$





Search for B to four muons

- $B^0_{(s)} \rightarrow \mu^+ \mu^- \mu^+ \mu^-$ very rare in the SM (non-resonant) $\mathcal{B}(B^0_s) \sim 10^{-10}, \ \mathcal{B}(B^0) \sim 10^{-12}$
- Many extension of the SM can give contributions orders of magnitude larger, such as MSSM [Demidov, Gorbunov] *
- In particular light axions that could explain the g 2 anomaly
 [Bauer, Neubert, Thamm PRL119,031802(2017)]
 [Liu, Wagner, Wang JHEP 03 (2019) 008]
 [Chala, Egede, Spannowsky Eur.Phys.J.C 79 (2019) 5, 431]
- Use full Run1-2 statistics $(9 \, \text{fb}^{-1})$, supersedes previous results
- Search for non-resonant $B_{(s)}^0 \to \mu^+ \mu^- \mu^+ \mu^-$, axion mediated $B \to aa$ with $m_a = 1$ GeV, and $B_{(s)}^0 \to J/\psi (\mu^+ \mu^-) \mu^+ \mu^-$
- Normalisation to $B_s^0 \to J/\psi(\mu^+\mu^-)\phi(\mu^+\mu^-),$ $\mathcal{B} = (1.74 \pm 0.14)10^{-8}$
- Search in bins of a BDT trained against combinatorial background
- Misidentified background found to be negligible







Search for B to four muons

- No excess above background expectation found
- Limit with CLs method in GAMMACOMBO The limits at 95% confidence are

$$\begin{split} & \mathcal{B} \left(B^0_s \to \mu^+ \mu^- \mu^+ \mu^- \right) < 8.6 \times 10^{-10} \,, \\ & \mathcal{B} \left(B^0 \to \mu^+ \mu^- \mu^+ \mu^- \right) < 1.8 \times 10^{-10} \,, \\ & \mathcal{B} \left(B^0_s \to a \left(\mu^+ \mu^- \right) a \left(\mu^+ \mu^- \right) \right) < 5.8 \times 10^{-10} \,, \\ & \mathcal{B} \left(B^0 \to a \left(\mu^+ \mu^- \right) a \left(\mu^+ \mu^- \right) \right) < 2.3 \times 10^{-10} \,, \\ & \mathcal{B} \left(B^0_s \to J/\psi \left(\mu^+ \mu^- \right) \mu^+ \mu^- \right) < 2.6 \times 10^{-9} \,, \\ & \mathcal{B} \left(B^0 \to J/\psi \left(\mu^+ \mu^- \right) \mu^+ \mu^- \right) < 1.0 \times 10^{-9} \,. \end{split}$$

First search for $B \to aa$ with $m_a = 1 \text{ GeV}$ First limit on $B^0_{(s)} \to J/\psi (\mu^+ \mu^-) \mu^+ \mu^-$ decays Factor 2 improvement on the non resonant channels.



F. Det See also first search for $B^0 \to \phi(K^+K^-)\mu^+\mu^-$ [hep-ex/2201.10167] in Mbaickup CD 22 10/18

Measurement of the photon polarization in $\Lambda_b^0 \to \Lambda \gamma$

- + $b\to s\gamma$ photon polarisation in SM predominantly left-handed (deviations $\propto m_s^2/m_b^2)$
- Photon recoils against Λ that is in turn polarised
- Measure photon polarisation α_{γ} from proton angular distribution
- Sensitive to $C_7^{(\prime)}$ couplings
- Analysis based on $6 \, \text{fb}^{-1}$ data triggered by the hadrons or photon
- BDT based selection against combinatorial background, small $\Lambda^0_b\to\Lambda\eta$ residual background



Measurement of the photon polarization in $\Lambda_b^0 \to \Lambda \gamma$



• Signal described by $\frac{\mathrm{d}\Gamma}{\mathrm{d}\cos(\theta_p)} \propto 1 - \alpha_\Lambda \alpha_\gamma \cos(\theta_p)$ times acceptance function

- Use Λ decay parameter $\alpha_{\Lambda} = 0.754 \pm 0.004$ from BESIII [Nature Phys. 15 (2019) 631]
- Background by a 4^{th} order polynomial with coefficients determined from mass sidebands
- The photon polarization, with all systematic uncertainties $\alpha_{\gamma} = 0.82 \pm 0.23 \pm 0.13$
- Charge separated as measure of CP violation $\alpha_{\gamma}^- = 1.26 \pm 0.42 \pm 0.20$ $\alpha_{\gamma}^+ = -0.55 \pm 0.32 \pm 0.16$

Measurement of the photon polarization in $\Lambda_b^0 \to \Lambda \gamma$

• Confidence intervals using Feldman-Cousins method imposing physical limits

$$\begin{split} &\alpha_{\gamma}^{-} > 0.56(0.44) \text{ at } 90\% \ (95\%) \text{ CL} \\ &\alpha_{\gamma}^{+} = -0.56^{+0.36}_{-0.33}(stat.)^{+0.16}_{-0.09}(syst.), \end{split}$$

 $\alpha_{\gamma} = 0.82^{+0.17}_{-0.26}(stat.)^{+0.04}_{-0.13}(syst.)$



- This represents the first measurement of the photon polarisation in $b\to s\gamma$ decays
- Consistent with SM predictions and CP symmetry
- Constraints on Wilson coefficients with **FLAVIO** software
- Tightly constraining new C_7 currents: solving two fold ambiguity with C'_7 currents

Angular analysis of $D^0 \to h^+ h^- \mu^+ \mu^-$ decays

- First full angular analysis of very rare $D^0 \rightarrow h^+ h^- \mu^+ \mu^-$ decays
- Full $9 \, \text{fb}^{-1}$ statistics
- Regions dominated by resonances $D^0 \rightarrow h^+ h^- R$ used as SM null tests
- Differential rate expressed in terms of 9 angular coefficients I_i
- Both flavour averages and CP asymmetries measured



Angular analysis of $D^0 \to h^+ h^- \mu^+ \mu^-$ decays



- First angular analysis of these very rare D decays
- The null test observables in agreement with SM at 0.3σ and 2.7σ for the $D^0 \rightarrow \pi^+\pi^-\mu^+\mu^$ and $D^0 \rightarrow K^+K^-\mu^+\mu^-$ decays
- All measurements are consistent with SM predictions (where present) and CP symmetry



 $\Lambda_b^0 - \overline{\Lambda}_b^0$ production asymmetry at $\sqrt{s} = 7$ and 8 TeV

- Production asymmetry fundamental for CP violation measurements
- Use semi-leptonic $\Lambda_b^0 \to \Lambda_c^+ \mu \bar{\nu}_\mu X$ decays
- Measure asymmetry in 3 fb^{-1} of pp data at $\sqrt{s} = 7$ and 8 TeV, in y and p_T bins

 $A_{\rm raw} = A_{\rm Prod} + A_{\rm det} + A_{\rm int} + A_{\rm PID} + A_{\rm trig}$

120

100

80

60

40

20

LHCh

 $\sqrt{s} = 8 \text{ TeV}$

- Fit

+ Data 2fb⁻¹

Signal

Comb. bkg.

2250

2300

 $m(pK^{-}\pi^{+})$ [MeV/c²]

Candidates per (2 MeV/c

- A_{int} from external measurements and $\Lambda \to p\pi$ control channel
- A_{det} canceled by swapping magnetic field
- $A_{\rm PID}$ and $A_{\rm trig}$ calibrated from control channels in data





2350

 $A_b^0-\bar{A}_b^0$ production asymmetry at $\sqrt{s}=7$ and $8\,{\rm TeV}$

- Asymmetry is observed at 5.8σ significance, at $\sim 1\%$ level on average
- Evidence for a dependence on the rapidity
- Different Pythia tunings mostly overestimate the asymmetry, only one Colour Recombination model predicts correct low- $p_{\rm T}$ behaviours.
- Good comparison with Heavy-quark recombination model [PRD91(2015)054022]



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Summary and conclusion



Legacy analyses from LHCb Run1-2 are being produced

- $B^0_{(s)} \to \mu^+ \mu^-(\gamma)$ with world best single experiments results:
 - * first limit on $B_s^0 \to \mu^+ \mu^- \gamma$ decays
 - \star Closing the phase space of (pseudo-)scalar or axial-vector new interactions
 - \star looking forward to the full Run 1-2 analyses from ATLAS and CMS
- New $B^0_{(s)} \to p\bar{p}$ measurement
- Updated $B^0_{(s)} \to \mu^+ \mu^- \mu^+ \mu^-$ search
 - * First search for $B \rightarrow aa$ with mass also around 1 GeV
 - \star Strong constraints on all branching fractions
- Observation of photon polarisation in $\Lambda_b^0 \to \Lambda \gamma$ decays
 - * Constraining new C_7 couplings
- Observation of Λ_b^0 production asymmetry fundamental for future CP violation studies
- Measurements with charm decays tighten the space for models not constrained by the ${\cal B}$

All of the very rare decays are statistically limited, and will be for some time Looking forward to the collected data in Run 3 with the upgraded LHCb detector!



Backup

F. Dettori

Moriond QCD 22 18/18

Combined measurement of hadronisation fraction ...and B_s^0 branching fractions

Breaking the recursive problem: combine information of different measurements Measure production ratios from ratio of decays with known rate (semileptonic) or known rate ratios $(B \rightarrow Dh)$, and cross-check dependencies with decays of high rate $(B \to J/\psi X).$



 $\rightarrow D\pi$

7 TeV

 p_{T}^{30} [GeV/c]

LHC

2 fb⁻¹ 8 TeV

 p_{T}^{30} [GeV/c]

Recent LHCb combination $\frac{f_s}{f_s}(13 \text{ TeV}) = 0.2539 \pm 0.0079$



LHCb-PAPER-2020-046

PRD

about a factor 2

13 TeV

 p_{T}^{s0} [GeV/c]

LHCb

 p_{π}^{30} [GeV/c]

2 fb⁻¹ 13 TeV

 $B \rightarrow D \pi$

20

0.22 0.2





ž,

0.24 0.22 $B \rightarrow D \pi$

7 TeV

 p_{τ}^{30} [GeV/c]

LHCb

1 fb⁻¹ 7 TeV

³⁰ p_m [GeV/c]

0.28 0.26

0.22

\$

0.22

Combined measurement of hadronisation fraction ... and B_s^0 branching fractions

More than 50 B^0_s meson branching fractions updated, reducing significantly their uncertainties.

				Decay mode	Opdated branching fraction	Frevious result		
Decay mode	Updated branching fraction (2.75 + 0.18 + 0.12 + 0.13 + 0.24) - 10 ⁻⁵	Previous result (2.52 + 0.17 + 0.11 + 0.20 + 0.12) 10-5	[5.6]	$B_s^0 \rightarrow \pi^+\pi^-$ $B_s^0 \rightarrow K^-\pi^+$ $B_s^0 \rightarrow K^+K^-$	$\begin{array}{l} (7.60\pm0.58\pm0.69\pm0.25\pm0.25)\times10^{-7} \\ (6.15\pm0.49\pm0.49\pm0.20\pm0.20)\times10^{-6} \\ (2.63\pm0.08\pm0.16\pm0.09\pm0.09)\times10^{-5} \end{array}$	$\begin{array}{c} (6.91\pm0.54\pm0.63\pm0.40\pm0.19)\times10^{-7} \\ (5.4\pm0.4\pm0.4\pm0.4\pm0.2)\times10^{-6} \\ (2.30\pm0.07\pm0.14\pm0.17\pm0.07)\times10^{-5} \end{array}$	[70] [71] [71]	*
$B_s \rightarrow \psi^{\uparrow}$ $B_s^0 \rightarrow \mu^+ \mu^-$ $B_s^0 \rightarrow \overline{K}^{*0} \mu^+ \mu^-$ $B_s^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$	$\begin{array}{c} (3.73 \pm 0.18 \pm 0.12 \pm 0.12 \pm 0.12 \pm 0.24) \times 10^{-9} \\ (3.26 \pm 0.65^{-10.22}_{-0.11} \pm 0.10) \times 10^{-9} \\ (3.09 \pm 1.07 \pm 0.21 \pm 0.10 \pm 0.22) \times 10^{-8} \\ (8.66 \pm 1.50 \pm 0.47 \pm 0.28 \pm 0.60) \times 10^{-8} \end{array}$	$\begin{array}{c} (3.0\pm0.6^{+0.2}_{-0.1}\pm0.2)\times10^{-9}\\ (3.0\pm0.6^{+0.2}_{-0.1}\pm0.2)\times10^{-9}\\ (2.9\pm1.0\pm0.2\pm0.2\pm0.2)\times10^{-8}\\ (8.6\pm1.5\pm0.5\pm0.5\pm0.7)\times10^{-8} \end{array}$	[50] * [57] [58] [59] *	$\begin{array}{l} B^0_s \to K^0_{\rm S} K^0_{\rm S} \\ B^0_s \to K^0_{\rm S} \pi^+ \pi^- \\ B^0_s \to K^0_{\rm S} K^\pm \pi^\mp \end{array}$	$\begin{array}{l} (8.28\pm1.60\pm0.90\pm0.26\pm0.81)\times10^{-6} \\ (5.21\pm0.74\pm0.85\pm0.17\pm0.23)\times10^{-6} \\ (4.64\pm0.19\pm0.30\pm0.15\pm0.21)\times10^{-5} \end{array}$	$\begin{array}{c} (8.3\pm1.6\pm0.9\pm0.3\pm0.8)\times10^{-6} \\ (4.7\pm0.7\pm0.8\pm0.3\pm0.2)\times10^{-6} \\ (4.22\pm0.18\pm0.28\pm0.25\pm0.17)\times10^{-5} \end{array}$	[72] [73] [73]	
$B_s^0 \rightarrow \phi \mu^+ \mu^-$ $q^2 \in [1.0, 6.0]$	$(7.54^{+0.43}_{-0.41} \pm 0.30 \pm 0.36) \times 10^{-7}$ $(2.44^{+0.31}_{-0.30} \pm 0.07 \pm 0.12) \times 10^{-8}$	$(7.97^{+0.45}_{-0.43} \pm 0.32 \pm 0.60) \times 10^{-7}$ $(2.58^{+0.33}_{-0.31} \pm 0.08 \pm 0.19) \times 10^{-8}$	[14] * [14] *	$B_s^0 \rightarrow K^{*0}\overline{K}^{*0}$ $B_s^0 \rightarrow K^{*\pm}K^{\mp}$ $B_s^0 \rightarrow K^{*-}\pi^+$	$(2.70 \pm 0.44 \pm 0.43 \pm 0.09 \pm 0.19) \times 10^{-5}$ $(1.23 \pm 0.18 \pm 0.13 \pm 0.04 \pm 0.07) \times 10^{-5}$ $(3.21 \pm 1.07 \pm 0.41 \pm 0.10 \pm 0.18) \times 10^{-6}$	$(2.81 \pm 0.46 \pm 0.43 \pm 0.34 \pm 0.13) \times 10^{-5}$ $(1.27 \pm 0.19 \pm 0.13 \pm 0.07 \pm 0.10) \times 10^{-5}$ $(3.3 \pm 1.1 \pm 0.4 \pm 0.2 \pm 0.3) \times 10^{-6}$	[74] [75] [75]	*
$q^2 \in [15.0, 19.0]$ $q^2 \in [0.1, 2.0]$ $q^2 \in [2.0, 5.0]$	$(3.82^{+0.38}_{-0.36} \pm 0.12 \pm 0.18) \times 10^{-8}$ $(5.54^{+0.69}_{-0.65} \pm 0.13 \pm 0.27) \times 10^{-8}$ $(2.42^{+0.40} + 0.06 \pm 0.12) \times 10^{-8}$	$(4.04^{+0.30}_{-0.38} \pm 0.13 \pm 0.30) \times 10^{-8}$ $(5.85^{+0.73}_{-0.29} \pm 0.14 \pm 0.44) \times 10^{-8}$ $(2.56^{+0.72}_{-0.29} \pm 0.06 \pm 0.19) \times 10^{-8}$	[14] * [14] * [14] *	$B_s^0 \rightarrow p\bar{p}K^{\pm}\pi^{\mp}$ $B_s^0 \rightarrow p\Lambda K^{\mp}$	$\begin{array}{c} (1.41\pm0.23\pm0.12\pm0.05\pm0.11)\times10^{-6} \\ (6.01\pm0.66\pm0.62\pm0.20\pm0.57)\times10^{-6} \end{array}$	$\begin{array}{l} (1.30\pm 0.21\pm 0.11\pm 0.09\pm 0.08)\times 10^{-6} \\ (5.46\pm 0.61\pm 0.57\pm 0.32\pm 0.50)\times 10^{-6} \end{array}$	[76] [77]	-
$q^2 \in [5.0, 8.0]$ $q^2 \in [11.0, 12.5]$	$(3.03^{+0.42}_{-0.38} \pm 0.00 \pm 0.12) \times 10^{-8}$ $(3.03^{+0.42}_{-0.42} \pm 0.07 \pm 0.15) \times 10^{-8}$ $(4.45^{+0.42}_{-0.42} \pm 0.14 \pm 0.21) \times 10^{-8}$	$(3.21^{+0.44}_{-0.45} \pm 0.08 \pm 0.24) \times 10^{-8}$ $(4.71^{+0.45}_{-0.45} \pm 0.15 \pm 0.36) \times 10^{-8}$	[14] * [14] *	$B_s^0 \rightarrow \phi \overline{K}^{*0}$ $B_s^0 \rightarrow \phi$	$(1.27 \pm 0.28 \pm 0.16 \pm 0.04 \pm 0.07) \times 10^{-6}$ $(2.02 \pm 0.05 \pm 0.08 \pm 0.07 \pm 0.11) \times 10^{-5}$	$(1.10 \pm 0.24 \pm 0.13 \pm 0.08 \pm 0.06) \times 10^{-6}$ $(1.84 \pm 0.05 \pm 0.07 \pm 0.11 \pm 0.12) \times 10^{-5}$	[78] [79]	*
$q^2 \in [15.0, 17.0]$ $q^2 \in [17.0, 19.0]$	$(4.28^{+0.34}_{-0.24} \pm 0.11 \pm 0.21) \times 10^{-8}$ $(3.75^{+0.34}_{-0.51} \pm 0.13 \pm 0.18) \times 10^{-8}$	$(4.52^{+0.34}_{-0.54} \pm 0.12 \pm 0.34) \times 10^{-8}$ $(3.96^{+0.37}_{-0.54} \pm 0.14 \pm 0.30) \times 10^{-8}$	[14] *	οφο 	$(3.82 \pm 0.25 \pm 0.19 \pm 0.30) \times 10^{-6}$ $(2.36 \pm 0.61 \pm 0.30 \pm 0.19) \times 10^{-6}$	$(3.48 \pm 0.23 \pm 0.17 \pm 0.35) \times 10^{-6}$ $(2.15 \pm 0.54 \pm 0.28 \pm 0.21) \times 10^{-6}$	[80] [81]	*
			-È	Decay mode	Updated branching fraction	Previous result		_
			$\mathcal{A}^{\mathcal{Y}}$	$\begin{array}{c} B^0_s ightarrow D^{*-}_s \mu^+ u_\mu \ B^0_s ightarrow D^s \mu^+ u_\mu \end{array}$	$\begin{array}{c} (5.19\pm 0.24\pm 0.47\pm 0.13\pm 0.14)\times 10^{-2} \\ (2.40\pm 0.12\pm 0.15\pm 0.06\pm 0.10)\times 10^{-2} \end{array}$	$\begin{array}{c} (5.38\pm 0.25\pm 0.48\pm 0.20\pm 0.15)\times 10^{-2} \\ (2.49\pm 0.12\pm 0.16\pm 0.09\pm 0.11)\times 10^{-2} \end{array}$	[53] [53]	
Decay mode	Updated branching fraction	Previous result	<u> </u>	$B_s^0 \rightarrow D^+ D_s^-$ $B_s^0 \rightarrow D^+ D^-$	$(3.01 \pm 0.32 \pm 0.10 \pm 0.08 \pm 0.34) \times 10^{-4}$ $(2.47 \pm 0.46 \pm 0.23 \pm 0.08 \pm 0.22) \times 10^{-4}$	$(2.7 \pm 0.3 \pm 0.1 \pm 0.2 \pm 0.3) \times 10^{-4}$ $(2.2 \pm 0.4 \pm 0.1 \pm 0.1 \pm 0.3) \times 10^{-4}$	[82] [83]	
$B_s^0 \rightarrow J/\psi K_S^0$ $B_s^0 \rightarrow J/\psi K_S^0 K^{\pm} \pi^{\mp}$ $R^0 \rightarrow \psi(2S) \overline{K}^{*0}$	$(2.06 \pm 0.08 \pm 0.06 \pm 0.07 \pm 0.08) \times 10^{-5}$ $(5.01 \pm 0.35 \pm 0.33 \pm 0.16 \pm 0.44) \times 10^{-4}$ $(3.62 \pm 0.37 \pm 0.26 \pm 0.12 \pm 0.25) \times 10^{-5}$	$(1.93 \pm 0.08 \pm 0.05 \pm 0.11 \pm 0.07) \times 10^{-5}$ $(4.6 \pm 0.3 \pm 0.3 \pm 0.3 \pm 0.4) \times 10^{-4}$ $(3.25 \pm 0.24 \pm 0.24 \pm 0.19 \pm 0.22) \times 10^{-5}$	[60] [61] * [62]	$B_s^0 \rightarrow D^0 \overline{D}{}^0$ $B_s^0 \rightarrow D_s^+ D_s^-$ $B_s^0 \rightarrow D^{*\pm} D^{*\mp}$	$\begin{array}{l}(1.83\pm 0.29\pm 0.29\pm 0.05\pm 0.18)\times 10^{-4}\\(4.38\pm 0.23\pm 0.31\pm 0.11\pm 0.49)\times 10^{-3}\\(8.38\pm 1.02\pm 0.12\pm 0.26\pm 0.81)\times 10^{-5}\end{array}$	$\begin{array}{l} (1.9 \pm 0.3 \pm 0.2 \pm 0.2 \pm 0.3) \times 10^{-4} \\ (4.0 \pm 0.2 \pm 0.2 \pm 0.2 \pm 0.4) \times 10^{-3} \\ (8.41 \pm 1.02 \pm 0.12 \pm 0.39 \pm 0.79) \times 10^{-5} \end{array}$	[83] [83] [84]	
$B_s^0 \rightarrow \psi(2S)K^+\pi^-$ $B_s^0 \rightarrow J/\psi\eta$ $B_s^0 \rightarrow J/\psi\eta'$	$(3.43 \pm 0.23 \pm 0.14 \pm 0.11 \pm 0.24) \times 10^{-5}$ $(4.04 \pm 0.35^{+0.32}_{-0.34} \pm 0.13 \pm 0.28) \times 10^{-4}$ $(3.67 \pm 0.32^{+0.34}_{-0.38} \pm 0.12 \pm 0.25) \times 10^{-4}$	$\begin{array}{c} (3.12\pm0.21\pm0.13\pm0.18\pm0.22)\times10^{-5} \\ (3.79\pm0.31^{+0.20}_{-0.20}\pm0.28\pm0.56)\times10^{-4} \\ (3.42\pm0.30^{+0.14}_{-0.35}\pm0.26\pm0.51)\times10^{-4} \end{array}$	[62] [63] * [63] *	$\begin{array}{l} B^0_s \to D^{+(*)}_s D^{-(*)}_s \\ B^0_s \to D^{*\pm}_s D^{\mp}_s \\ B^0_s \to D^{*+}_s D^{*-}_s \end{array}$	$\begin{array}{l} (3.36\pm 0.11\pm 0.14\pm 0.09\pm 0.38)\times 10^{-2}\\ (1.49\pm 0.06\pm 0.07\pm 0.04\pm 0.17)\times 10^{-2}\\ (1.39\pm 0.09\pm 0.10\pm 0.04\pm 0.16)\times 10^{-2} \end{array}$	$\begin{array}{l} (3.05\pm0.10\pm0.13\pm0.14\pm0.34)\times10^{-2} \\ (1.35\pm0.06\pm0.06\pm0.06\pm0.15)\times10^{-2} \\ (1.27\pm0.08\pm0.09\pm0.06\pm0.14)\times10^{-2} \end{array}$	[85] [85] [85]	
$B_s^0 \rightarrow \psi(2S)\phi$ $B_s^0 \rightarrow \chi_{c1}\phi$	$(4.98 \pm 0.26 \pm 0.24 \pm 0.24) \times 10^{-4}$ $(1.92 \pm 0.18 \pm 0.14 \pm 0.09) \times 10^{-5}$	$(5.33 \pm 0.28 \pm 0.26^{+1.37}_{-1.12}) \times 10^{-4}$ $(1.98 \pm 0.19 \pm 0.15 \pm 0.20) \times 10^{-5}$	[12] * [64] *	$B_s^0 \rightarrow \overline{D}{}^0 K_S^0$ $B_s^0 \rightarrow \overline{D}{}^{*0} K_S^0$	$\begin{array}{c} (4.69\pm 0.51\pm 0.28\pm 0.15\pm 0.64)\times 10^{-4} \\ (3.05\pm 1.13\pm 0.40\pm 0.10\pm 0.41)\times 10^{-4} \end{array}$	$\begin{array}{l}(4.3\pm0.5\pm0.3\pm0.3\pm0.6)\times10^{-4}\\(2.8\pm1.0\pm0.3\pm0.2\pm0.4)\times10^{-4}\end{array}$	[86] [86]	
$B_s^0 \rightarrow J/\psi \pi^+ \pi^-$ $B_s^0 \rightarrow J/\psi \phi \phi$ $B_s^0 \rightarrow J/\psi \overline{K}^{*0}$ $B_s^0 \rightarrow J/\psi p \overline{\rho}$ $B^0 \rightarrow J/\psi p \overline{\rho}$	$\begin{array}{l} (2.01\pm 0.05\pm 0.05\pm 0.10)\times 10^{-4} \\ (1.17\pm 0.12^{+0.02}_{-0.00}\pm 0.06)\times 10^{-5} \\ (4.12\pm 0.19\pm 0.13\pm 0.20)\times 10^{-5} \\ (3.54\pm 0.19\pm 0.24\pm 0.16)\times 10^{-6} \\ (3.94\pm 0.35\pm 0.26\pm 0.13)\times 10^{-7} \end{array}$	$\begin{array}{l} (2.16\pm 0.05\pm 0.06^{+0.31}_{-0.42})\times 10^{-4}\\ (1.19\pm 0.12^{+0.05}_{-0.42})\pm 0.10)\times 10^{-5}\\ (4.20\pm 0.20\pm 0.13\pm 0.36)\times 10^{-5}\\ (3.58\pm 0.19\pm 0.24\pm 0.30)\times 10^{-6}\\ (4.51\pm 0.40\pm 0.30\pm 0.32)\times 10^{-7} \end{array}$	$[11] \\ \star \\ [15] \\ \star \\ [65] \\ \star \\ [66] \\ \star \\ [66] \\ \star \\]$	$B_a^0 \rightarrow \overline{D}^0 \overline{K}^{*0}$ $B_a^0 \rightarrow \overline{D}^0 K^- \pi^+$ $B_a^0 \rightarrow \overline{D}^0 \phi$ $B_a^0 \rightarrow \overline{D}^{*0} \phi$ $B_a^0 \rightarrow \overline{D}^0 K^+ K^-$	$\begin{array}{c} (5.31\pm1.22\pm0.54\pm0.17\pm0.35)\times10^{-4}\\ (1.11\pm0.05\pm0.07\pm0.04\pm0.06)\times10^{-3}\\ (3.25\pm0.38\pm0.19\pm0.11\pm0.18)\times10^{-5}\\ (4.01\pm0.48\pm0.27\pm0.13\pm0.23)\times10^{-5}\\ (6.13\pm0.59\pm0.28\pm0.20\pm0.56)\times10^{-5} \end{array}$	$\begin{array}{l} (4.72\pm1.07\pm0.48\pm0.37\pm0.74)\times10^{-4} \\ (1.00\pm0.04\pm0.06\pm0.08\pm0.10)\times10^{-3} \\ (3.0\pm0.3\pm0.2\pm0.2\pm0.2)\times10^{-5} \\ (3.7\pm0.5\pm0.2\pm0.2\pm0.2)\times10^{-5} \\ (5.7\pm0.5\pm0.2\pm0.3\pm0.5)\times10^{-5} \end{array}$	[87] 5 [88] 5 [89] 5 [89] 5 [89] 5 [90] 5	* * *
$\begin{array}{l} B^0_s \rightarrow \psi(2S)\eta \\ B^0_s \rightarrow \psi(2S)\eta' \end{array}$	$\begin{array}{c} (3.35\pm0.57\pm0.48\pm0.50)\times10^{-4} \\ (1.42\pm0.33\pm0.06\pm0.20)\times10^{-4} \end{array}$	$(3.15 \pm 0.53 \pm 0.45^{+0.61}_{-0.67}) \times 10^{-4}$ $(1.32 \pm 0.31 \pm 0.05^{+0.67}_{-0.28}) \times 10^{-4}$	[67] * [68] *	$\begin{array}{c} B^0_s \rightarrow D^{\mp}_s K^{\pm} \\ B^0_s \rightarrow D^{-}_s \pi^+ \pi^- \pi^+ \end{array}$	$\begin{array}{c} (2.41\pm 0.05\pm 0.06\pm 0.14)\times 10^{-4} \\ (6.43\pm 1.18\pm 0.64\pm 0.38)\times 10^{-3} \end{array}$	$\begin{array}{l}(2.29\pm 0.05\pm 0.06\pm 0.17)\times 10^{-4}\\(6.01\pm 1.11\pm 0.60\pm 0.48)\times 10^{-3}\end{array}$	[91] s [92] s	2
$B_s^0 \rightarrow J/\psi \pi^+ \pi^- \pi^+ \pi$ $B_s^0 \rightarrow \psi(2S)\pi^+ \pi^-$	$\begin{array}{l} - (7.49 \pm 0.30 \pm 0.44 \pm 0.42) \times 10^{-5} \\ (6.87 \pm 0.81 \pm 0.65 \pm 0.39) \times 10^{-5} \end{array}$	$(7.62 \pm 0.36 \pm 0.64 \pm 0.42) \times 10^{-5}$ $(7.3 \pm 0.9 \pm 0.6^{+1.9}_{-1.6}) \times 10^{-5}$	[69] * [67] *	$B_s^0 \rightarrow D_s^- K^+ \pi^- \pi^+$ $B_s^0 \rightarrow D_{s1}(2536)^- \pi^-$	$(3.34 \pm 0.32 \pm 0.19 \pm 0.73) \times 10^{-4}$ + $(2.57 \pm 0.64 \pm 0.26 \pm 0.56) \times 10^{-5}$	$\begin{array}{c} (3.13\pm 0.30\pm 0.18\pm 0.76)\times 10^{-4} \\ (2.41\pm 0.60\pm 0.24\pm 0.58)\times 10^{-5} \end{array}$	[93] [93]	

[LHCb-PAPER-2020-046

PRD 104, 032005 (2021)]

INFN

Search for $B^0 \rightarrow \phi \mu^+ \mu^-$ decays

- Rare decay in the SM (penguin CKM / OZI suppressed) $b \to d \mu^+ \mu^-$ FCNC
- Short distance $\mathcal{B} \sim 10^{-12}$
- Including $\omega \phi$ mixing could raise at $10^{-11} 10^{-10}$ level
- New physics contributions such as Z' could enhance this



Search for $B^0 \to \phi \mu^+ \mu^-$ decays

- Normalised to $B_s^0 \to \phi \mu^+ \mu^-$ decays
- $B_s^0 \to J/\psi \phi$ decays as control channel
- Main background: $B^0 \to K^* \mu^+ \mu^-, \Lambda^0_b \to p K^- \mu^+ \mu^-$



No excess over background expectation, upper limit

$$\begin{split} \mathcal{R} &= \frac{\mathcal{B}(B^0 \to \phi \mu^+ \mu^-)}{\mathcal{B}(B_s^0 \to \phi \mu^+ \mu^-)} < 4.4 \times 10^{-3} \text{ at } 90\% \text{ CL} \; . \\ \mathcal{B}(B^0 \to \phi \mu^+ \mu^-) < 2.3(3.2) \times 10^{-9} \text{ at } 90\% \text{ CL} \; . \\ \text{excluding } \phi \text{ and charmonia dimuon regions (extrapolating to full } q^2) \end{split}$$

Distributions of $B_s^0 \to \mu^+ \mu^-$ decays kinematics in data



LHCb

Decay time acceptance for $B_s^0 \to \mu^+ \mu^-$ decays



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Lifetime measurement cross-check with $B_s^0 \to K^+ K^-$ decay



$$\tau_{B^0_s \to K^+ K^-} = 1.435 \pm 0.026 \,\mathrm{ps}$$

In agreement with published

$$\tau_{B^0_s \to K^+ K^-} = 1.407 \pm 0.016 \,\mathrm{ps}$$

Lifetime measurement cross-check with $B^0 \to K^+\pi^-$ decays \longrightarrow



$$\tau_{B^0 \to K^+ \pi^-} = 1.510 \pm 0.015 \,\mathrm{ps}$$

In agreement with published

$$\tau_{B^0 \to K^+\pi^-} = 1.524 \pm 0.011 \, \mathrm{ps}$$

Correlation between branching fraction and effective lifetim 🖗 🛲 🏧

The branching fraction measurement is affected by the effective lifetime, through the efficiency *

 $\rightarrow\,$ Hence there is a correlation between the two measurements

Both are thus sensitive to $A_{\Delta\Gamma}$



^{*}See e.g. [F.D., Guadagnoli, Phys.Lett.B 784 (2018) 96-100]

Moriond QCD 22 27/18

Calibration

Search in mass distribution in bins of multivariate discriminant (BDT)

- BDT shape calibrated from simulation and $B \to h^+ h^-$ in data
- Mass shape calibrated from quarkonia and $B \to h^+ h^-$ in data





$B_{d,s}^0 \to \mu^+ \mu^-$ observables

1. Branching fraction

$$\mathcal{B}^{t=0}(B_s^0 \to \mu^+ \mu^-) = \frac{G_F^4 M_W^4}{\pi^2} \tau_{B_s^0} f_{B_s}^2 m_{B_s}^3 \sqrt{1 - \frac{4m_\mu^2}{m_{B_s}^2}} |V_{ib} V_{ts}^*|^2 \left(|2\frac{m_\mu}{m_{B_s}} (C_{10} - C_{10}') + C_P - C_P'|^2 + |C_S - C_S'|^2 \right)$$

2. Ratio of branching fractions

$$\mathcal{R} = \frac{\mathcal{B}(B^0 \to \mu^+ \mu^-)}{\mathcal{B}(B_s^0 \to \mu^+ \mu^-)} = \frac{\tau_{B_d}}{\tau_{B_s}} \left(\frac{f_{B_d}}{f_{B_s}}\right)^2 \left|\frac{V_{td}}{V_{ts}}\right|^2 \frac{m_{B_d}}{m_{B_s}} \sqrt{1 - \frac{4m_{\mu}^2}{m_{B_d}^2}}$$
3. Effective lifetime
 B_s^0 mesons oscillate and mix into
their mass eigenstates, the effective
lifetime depends on which eigenstate
decays to $\mu^+ \mu^-$

$$\tau_{\mu\mu} = \frac{\tau_{B_s}}{(1 - y_s^2)} \frac{1 + 2y_s \mathcal{A}_{\Delta\Gamma} + y_s^2}{1 + y_s \mathcal{A}_{\Delta\Gamma}}$$

$$\frac{10}{\sqrt{p}} = \frac{1}{1/2} \frac{1}{\sqrt{p}} \frac{1}{1/2} \frac{1}{\sqrt{p}} \frac{1}{$$

Moriond QCD 22

29/18

F. Dettori

 $\tau_{\mu\mu} =$

 $B_s^0 \to \mu^+ \mu^- \gamma$ as an additional observable

🎯 🕼 🕊

The radiative $B_s^0 \to \mu^+ \mu^- \gamma$ decay is very interesting:

- Not helicity suppressed as rare as $B_s^0 \to \mu^+ \mu^-$
- Sensitive to vector couplings (C_9) (not just scalar or axial-vector)
- Can be split in initial (ISR) and final state radiation (FSR bremsstrahlung)

New method: measure the $B_s^0 \to \mu^+ \mu^- \gamma$ rate without photon reconstruction from the left sideband of the $B_s^0 \to \mu^+ \mu^-$ analysis.



Backgrounds



[PRL128(2022)041801] [PRD105(2022)012010]



Normalisation



Convert yields to branching fractions by normalising to channels of known rate



Use two channels

• $B^+ \to J/\psi K^+$ - same trigger & PID as signal

•
$$B^0 \to K^+ \pi^-$$
- same topology of signal



Measurement of the effective lifetime



Consistent at 1.5σ and 2.2σ with the heavy and light B_s^0 eigenstates lifetimes $(\tau_L = 1.423 \pm 0.005 \text{ ps} \text{ and } \tau_H = 1.620 \pm 0.007 \text{ ps})$





The
$$B^0_{d,s} \to \mu^+ \mu^-$$
 decays

Extremely rare decays

- Flavour changing neutral currents
- Helicity suppressed

Most recent Standard Model predictions

$$B(B_s^0 \to \mu^+ \mu^-) = (3.66 \pm 0.14) \times 10^{-9}$$
$$B(B^0 \to \mu^+ \mu^-) = (1.03 \pm 0.05) \times 10^{-10}$$

[Beneke, Bobeth, Szafron, JHEP10(2019) 232]

- Impressively precise predictions
- Any significant deviations from these values is sign of new interactions beyond the SM

• Dominated by parametric uncertainties Using the correlation of $\Delta F = 1$ rare decays with $\Delta F = 2$ B mixing,

using experimental ΔM values can also be predicted to be:

$$B(B_s^0 \to \mu^+ \mu^-) = (3.62^{+0.15}_{-0.10}) \times 10^{-9}$$

$$B(B^0 \to \mu^+ \mu^-) = (0.99^{+0.05}_{-0.03}) \times 10^{-10}$$

[Buras, Venturini -2109.11032]



 $B^0_{d,s} \to \mu^+ \mu^-$ analysis with full statistics Final invariant mass fit

Run 2







- Simultaneous fit in 10 bins 2 datasets (Run 1, 2) \times 5 BDT bins
- External constraints on yield and shape of misidentified backgrounds
- Combinatorial background free
- Signal shapes calibrated and constrained
 - All systematic uncertainties directly propagated

Run 1