



An experiment for electron-hadron scattering at the LHC

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Received: 23 June 2021 / Accepted: 21 December 2021

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Abstract Novel considerations are presented on the physics, apparatus and accelerator designs for a future, luminous, energy frontier electron-hadron (eh) scattering experiment at the LHC in the thirties for which key physics topics and their relation to the hadron-hadron HL-LHC physics programme are discussed. Demands are derived set by these physics topics on the design of the LHeC detector, a corresponding update of which is described. Optimisations on the accelerator design, especially the interaction region (IR), are presented. Initial accelerator considerations indicate that a common IR is possible to be built which alternately could serve eh and hh collisions while other experiments would stay on hh in either condition. A forward-backward symmetrised option of the LHeC detector is sketched which would permit extending the LHeC physics programme to also include aspects of hadron-hadron physics. The vision of a joint eh and hh physics experiment is shown to open new prospects for solving fundamental problems of high energy heavy-ion physics including the partonic structure of nuclei and the emergence of hydrodynamics in quantum field theory while the genuine TeV scale DIS physics is of unprecedented rank.

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1 Introduction

The Standard Model (SM) of particle physics is based on a non-Abelian gauge theory with a symmetry group $SU(2)_L \times U(1) \times SU_c(3)$. The SM has and continues to enjoy great success in describing a wide span of phenomena emerging from interactions of particles at a range of energies that is accessible experimentally. That said, the SM is not a satisfactory theory of fundamental interactions nor does it explain a number of phenomena in nature. It is of paramount importance to the field of particle physics to establish whether and how the SM breaks down in laboratory conditions. This is expected to be achieved by pushing the boundaries of energy and precision frontiers, and through various sensitive experiments at low energy. Theory is currently less predictive than ever after the birth of the SM, such that experimentation based on novel designs acquires a particular eminence for the decades ahead.

Deep inelastic scattering (DIS) of electrons and photoproduction off high energy protons and ions with high instantaneous luminosity offers a unique opportunity to enhance the precision frontier in particle physics, for which examples are provided in this paper. The intense, unique hadron beams of the LHC represent a salient opportunity to create a

new laboratory for particle physics and energy frontier DIS, the Large Hadron electron Collider (LHeC), at affordable cost: a larger than TeV centre of mass system (cms) energy new collider is in sight by adding an energy recovery linac to the LHC, in possibly staged phases. The present paper is mainly devoted to an update of the detector, describing relevant physics, apparatus and accelerator design considerations and new results. The LHeC would be the fifth large collider experiment at the LHC facility, sustaining its future and exploiting the hitherto biggest investment in particle physics.

A first comprehensive design concept for the LHeC was published in 2012 [1], just weeks prior to the Higgs boson (H) discovery and incorporating the findings of a review pursued by twenty experts in experimental, theoretical or accelerator physics. Following nearly ten years of LHC operation and analysis, incorporating technology progress, accounting for the advent of experimental Higgs physics and relying on the brilliant LHC performance, a further detailed report appeared recently, written again by representatives of more than a hundred institutions [2]. That paper presented the energy recovery, linac ring electron-hadron collider as the default configuration with luminosity parameters an order of magnitude enhanced compared to before. It suggested to downscope the electron beam energy from originally 60 to 50 GeV in an attempt to economise both investments and efforts as the racetrack electron accelerator circumference then became comparable to that of the SPS.¹ This concept, its main ingredients, had been submitted to the European Strategy Update and published a year before the extended paper [3].

The LHeC has been maintained as an option and complement of HL-LHC in strategic considerations of the future, especially the forthcoming European road-maps on detector and accelerator research and development. Its main role, naturally, is that of complementing the TeV scale exploration with the LHC and a possible future e^+e^- collider, much like HERA was coupled to the Tevatron and LEP before.

The present paper summarises key physics topics, on hadron structure, parton dynamics, Higgs, top and BSM physics and it enlists the detector requirements derived from each of these topics leading to an update of the LHeC apparatus concept which is here detailed.

For long, the LHeC development followed the common understanding that with the long shutdown (LS) 4, in the early thirties, the operation of the LHC as a heavy-ion (HI) collider would be terminated in order to maximise pp luminosity, which would have enabled using the Interaction Point (IP) 2 for a new experiment. Meanwhile one considers operating LHC with heavy ions further hence. New considera-

tions have appeared for a follow-up heavy-ion experiment at IP2, configured to study soft heavy-ion interactions [4], while heavy-ion physics is newly discussed at LHCb and ATLAS and CMS continue their HI programmes too.

In order to conceptually resolve a possible clash of the LHeC and ALICE3 plans in a productive manner, it had been suggested by the International Advisory Committee on the LHeC, see Appendix in [2] and [5], to investigate whether an adjusted LHeC detector, or a commonly designed apparatus,² and the interaction region (IR) could be configured to register both ep/eA and, alternately, AA scattering events. This paper comprises an initial study of the accelerator and IR modifications which suggest that a common IR may indeed be feasible, with features that are further to be evaluated. In this context, a symmetrised first version of the LHeC detector is here presented together with a discussion of a future joint eh and hh heavy ion physics programme.

This article is devoted to experimenting at the LHeC with the option of alternately registering electron-hadron (p, A) and hadron-hadron interactions in a common detector. Section 2 presents key physics topics in DIS scattering with emphasis on constraints on the detector layout and performance. Section 3 presents an updated LHeC detector layout, with an enlarged tracker radius and choosing liquid argon technology for the electromagnetic calorimeter. A detailed evaluation of a variety of aspects of the accelerator design, especially the IR and its possible configuration to alternately serve eh and hh interactions, is given in Sect. 4. Section 5 is dedicated to a comprehensive description of the heavy-ion physics potential for a joint eh/hh experiment and a brief presentation of a symmetrised version of the LHeC detector to enable its use for hh physics. The paper concludes with a summary in Sect. 6.

2 Physics

The physics programme at the LHC and the DIS programme at the LHeC are both extremely rich and stand on their own. However, they also have much in common: with the necessity of understanding hadron structure and parton dynamics for searches for new physics and precision measurements at the LHC, with the opportunity to explore the Higgs mechanism at per cent level, novel top quark physics, precision SM measurements and further, in the search for new physics and in the understanding of nuclear parton structure and the phenomenon of the Quark Gluon Plasma and heavy ion physics

¹ Pending on the finally achievable operational gradients in the SRF linac, an upgrade to 60 GeV might still be feasible at a second phase of the LHeC operation.

² While this paper, including the concept of a combined eh and hh detector at IP2, was in the peer-review process, the concept of the ALICE3 detector is evolving, possibly extending the initial tracking device concept [4] to also include an electromagnetic calorimeter, a higher field solenoid and a muon spectrometer following an absorber [6].

in general. With a view on the resulting detector constraints and for illustrating the exciting physics programme that the LHeC entails, we have chosen these key topics for a brief description of the potential of the “Experiment for Electron-Hadron Scattering at the LHC” we here describe. Some special emphasis is given to heavy ion physics in view of the idea, mentioned above, of possibly realising this experiment in a configuration that may jointly be used by DIS oriented and more heavy ion interested communities. Similar illustrations of the physics potential and experimental requirements could be provided for the physics at small Bjorken x , for diffraction, photo-production physics and other areas, see ref. [2]. One should be aware that the opening of an unexplored kinematic range, accessed with so high luminosity, may lead to surprises and should not pretend to be able to predict everything. This also regards, technically, the development of analysis tools, for which the past decade on LHC physics brought many examples of results exceeding in their depth and precision the initial expectations by far. Finally, new theoretical insight or surprises from other particle physics experiments, may indeed shift the focus.

2.1 Partons and proton structure

One may distinguish four phases, including the LHeC, of the experimental development of the physics of parton structure of the proton which was opened with the SLAC-MIT lepton-hadron scattering experiment at Stanford in 1968: fixed target experiments, HERA, pp Drell–Yan scattering and the LHeC. The role of a next, luminous energy frontier ep scattering experiment becomes obvious when one revisits the past and realises the unique potential of the LHeC, recently presented in much detail [2].

Partons, quarks and gluons, are confined inside the proton; still a major puzzle for modern physics, they cannot be observed directly. Quarks of any type q (and \bar{q}) have a certain probability of carrying a fraction x of the proton’s momentum, described by a momentum density function $xq(x)$, called a parton distribution function (PDF). The characteristics of the proton are given by the valence quark content of two up and one down quarks. The relative distribution of the proton’s momentum among the quarks varies with x . It changes as we resolve the proton more deeply in lepton-hadron deep inelastic scattering, i.e. through a virtual photon or a Z or W^\pm boson, of virtuality Q^2 , interacting with a quark. The strong interaction between quarks is mediated by gluons, discovered in 3-jet events in e^+e^- , which carry a half of the proton’s momentum. The quark-gluon interactions are described within QCD, with a coupling constant $\alpha_s(Q^2)$. With rising Q^2 , the coupling decreases logarithmically such that asymptotically quarks become free and the strong interaction at the parton level can be described as a perturbation theory. These and further fundamental properties have been

established³ in a first era of PDF physics enabled by a series of neutrino, electron and muon scattering experiments on stationary hadronic targets.

HERA was the first ep collider. It extended the kinematic range of DIS experiments, given as $s = Q_{max}^2 = 4E_e E_p$, by two orders of magnitude but fell short against those by another about two orders of magnitude in luminosity. Its contributions to the understanding of parton structure and dynamics, nevertheless, can not be underestimated. Of special importance has been, firstly, the extension of DIS into the very high Q^2 region with a) the validation of the linear DGLAP evolution law to Q^2 values beyond the weak boson masses, $\gtrsim 10^4 \text{ GeV}^2$, and b) the simultaneous use within one experiment of the charged (CC) and neutral current (NC) weak interaction, besides the electromagnetic NC photon exchange, to determine PDFs, including first determinations of the charm and bottom quark densities through impact parameter measurements. Since the accessible x range towards small x is extended $\propto 1/s$, HERA was able to resolve, in addition, the gluon, sea and valence quark behaviour at small x . It established the dominance of the gluon density xg at small x but could not convincingly answer the question of whether non-linear gluon-gluon interactions occur, which would damp the rise of xg towards small x and lead outside the validity range of the DGLAP equations. The HERA NC and CC collider data did permit a first and far reaching set of PDFs to be determined, without using additional data, with their canonical uncertainties understood at the $\Delta\chi^2 = 1$ level [8], and they are the inevitable part of any modern PDF determination.

Following HERA, with the advent of the LHC and its Drell–Yan measurements, the art of extracting PDFs from so-called global data has become an active field of particle physics, to test QCD and to understand LHC measurements using maximum suitable data, and novel analysis and mathematical methods. Such analyses carry a number of severe theoretical and practical difficulties which, despite impressive successes by the various PDF analysis groups, lead to a principally unsatisfactory situation due to the nature of hadron-hadron scattering with respect to DIS. This comprises effects of initial state quark radiation (leading to resummation), of hadronisation and reconstruction arbitrariness in jet data, the incompatibility of many data sets leading to the rather ad-hoc inflation of uncertainty bands or even exclusion of the most precise data, such as the ATLAS inclusive W,Z data from CT18 [9], for example. A reflection of these effects is the observed difference between PDF sets of different groups

³ Despite their phenomenological success there continue to exist certain doubts about the whole parton picture based on principles for the structure of nature going back to Newton [7] with testable hypotheses at the LHeC, which therefore has been termed the “Newtonian Telescope of CERN”.

which is often larger than the claimed precision of fits. A conceptual difficulty is the uncertainty at high mass, corresponding to large $x \gtrsim 0.5$, where the occurrence of new physics is possible, such that the LHC data should be excluded from PDF fits. The current status of the determination of α_s to about 2% uncertainty limits the PDF determination, and the precision of predictions such as the $gg \rightarrow H$ production cross section, being $\propto \alpha_s^2$. Simulation studies on future PDF determinations from the LHC assume that the data compatibility problems may disappear, while the principal problems will in fact remain.

A precision physics era at the HL-LHC will be maximally precise if it was accompanied by the LHeC PDF programme, the possible fourth phase of PDF physics. As described in detail in [2]: (i) the increased energy will make the CC DIS data for the first time a useful base extending over 4 orders of magnitude in x and Q^2 ; (ii) all PDFs, $xq(x, Q^2)$ and $xg(x, Q^2)$,⁴ can be determined in a single DIS experiment over many orders of magnitude, with $q = u_v, d_v, u, \bar{u}, d, \bar{d}, s, c, b$ and also t ; (iii) the kinematic range, unlike at HERA or lower energy fixed target or ep collider experiments, extends to such low values of x in the DIS region that one will be able to settle the question of non-linear gluon-gluon interactions, etc. An unprecedented precision on these distributions is in reach, as has been simulated more than once but conclusively in [2], including per mille accuracy of α_s . This will comprehensively test pQCD and the underlying parton dynamics view at the highest level; will enable new physics, possibly occurring in the high mass tails from interference contact interaction effects, to be discovered at LHC; lead to possible discoveries in QCD such as the breaking of factorisation not only in diffraction; and enable precision electroweak and Higgs physics at the joint ep/pp LHC facility to a stunning level of precision.

Such an ambitious programme, including precision measurements of the strange, charm and bottom quark distributions and of the longitudinal structure function $F_L(x, Q^2)$, sets important constraints for the experiment here presented: (i) it is very desirable that such data exist while HL-LHC operates. Therefore a dedicated study [2] has been made of the LHeC PDF prospects for an initial data set of 50 fb^{-1} , see Fig. 1.

Such a luminosity is a factor of 100 larger than that which H1 and ZEUS collected in their 15 year lifetime, while being expected in the first LHeC running period [10]; (ii) the detector acceptance should extend maximally to small hadron final state angles to cover larger x and to low elec-

⁴ PDFs depend on a factorisation scale μ_F^2 which in DIS is taken to be equal to Q^2 . Also note that we make no distinction between Bjorken x as the kinematic variable in DIS and x as the momentum fraction of the proton carried by the probed parton as they coincide in the parton model.

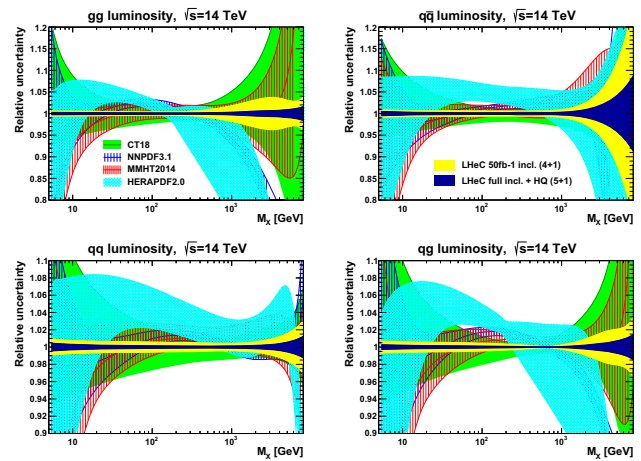


Fig. 1 Expected precision for the determination of parton-parton luminosities as function of M_X in Drell-Yan scattering at the 14 TeV LHC. Light blue: HERA, yellow: initial LHeC run, dark blue: full LHeC data set, overlaid with three, recent global fit results. For more information see [2]

tron scattering angles to cover low $Q \sim 1 \text{ GeV}^2$, even when one can extend the region of acceptance considerably with lower beam energy runs; (iii) hermiticity of the apparatus is required to apply an $E - p_z$ balance criterion which diminishes the radiative corrections substantially; (iv) cross calibration of the hadronic and electromagnetic calorimeter as well as polar angle measurements should ensure a below per cent level accuracy of the energy scales keeping the experimental scale uncertainties small; (v) high resolution hadron energy measurements are important for controlling hadronic backgrounds and reconstructing missing energy; (vi) heavy flavour reconstruction requires impact parameter resolutions of order $10 \mu\text{m}$ resulting from novel tracking technology and the small beam size of about $7 \mu\text{m}$ transversally, twenty times better than at HERA; (vii) the large photo-production background shall be tagged, for its own physics study and for subtracting it in DIS measurements. A major demand in ep scattering is the control of halo and synchrotron radiation backgrounds through a carefully designed interaction region, see Sect. 4.3.

2.2 Higgs boson physics

The Higgs boson is of fundamental importance as it is related to the spontaneous breaking of a locally symmetric gauge theory, to a mechanism in which the intermediate vector bosons are explained to be massive while the photon is kept massless. Fermions acquire a mass through the Higgs field. This spectacular theoretical prediction was confirmed in 2012 at the LHC, which boosted Higgs physics to the top of particle physics investigations. The nearest task is to verify the theory in more detail, especially to reconstruct the complete

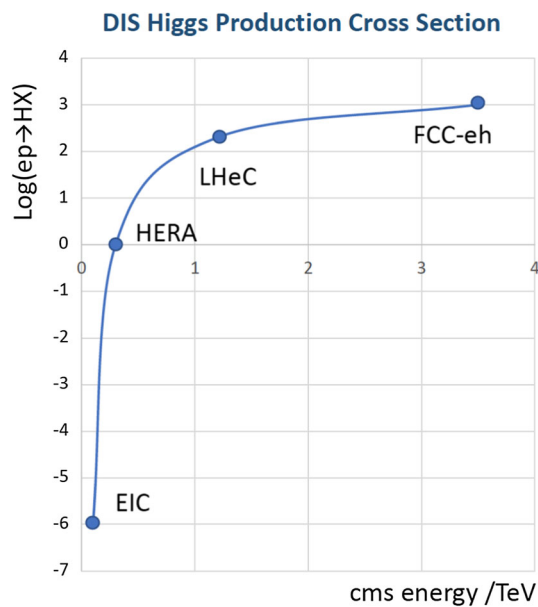


Fig. 2 Inclusive Higgs cross section in charged current $e^- p$ DIS, plotted as $\text{log}(\sigma/fb)$, as a function of the cms ep scattering energy, $\sqrt{s} = 2\sqrt{E_e E_p}$ in TeV

decay spectrum of the Higgs boson, besides searching for possible extensions of the Higgs mechanism and relations of the Higgs boson to exotic particles.

Prior to the discovery of the Higgs boson, initial work on the Higgs physics potential in DIS was centered around the dominant decay channel into $b\bar{b}$ [11], and a first comprehensive study was published with the LHeC design report in 2012 [1]. It had been realised that the theoretical understanding of Higgs production in ep was particularly suited for precision measurements with small QED and QCD corrections [12, 13].

In deep inelastic scattering, the Higgs boson is produced predominantly in charged current reactions, through the t -channel reaction $WW \rightarrow H$, with the neutral current cross section being smaller but still measurable. The production cross section depends strongly on the cms energy as is shown in Fig. 2.

The cross section at the LHeC is about 200 fb, which is close to the $e^+e^- \rightarrow HZ$ cross section in electron-positron scattering at 250 GeV energy. With thousand times its luminosity, HERA would have had a chance to observe the Higgs boson in DIS while that is beyond the reach of lower energy ep colliders such as the EIC. At the FCC-eh the SM H boson is calculated to have a pb cross section which implies it has an outstanding potential for Higgs physics as has been demonstrated in [14].

A comprehensive analysis of SM Higgs boson physics at the LHeC and the FCC-eh has recently been presented in [2]. The LHeC is sensitive to the six most frequent decay channels, $b\bar{b}$, W^+W^- , gg , $\tau\tau$, $c\bar{c}$ and ZZ , which repre-

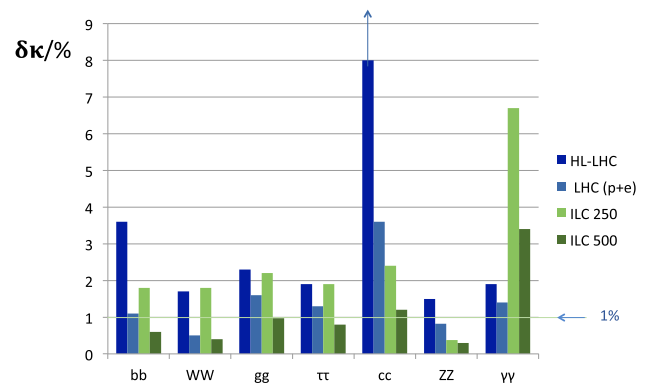


Fig. 3 Results of prospect evaluations of the determination of Higgs couplings in the SM kappa framework for HL-LHC (dark blue), LHC with LHeC combined (p+e, light blue), ILC 250 (light green) and ILC-500 (dark green) [15]

sent about 99.6% of the total SM Higgs boson decay width. Owing to the high precision with which the couplings to the b -quark and the W and Z bosons can be measured, the latter cleanly in $WW \rightarrow H \rightarrow WW$ in CC and $ZZ \rightarrow H \rightarrow ZZ$ in NC, one finds that the total of the (six) SM decays can be reconstructed at per cent level accuracy.

A dedicated analysis has also been presented in [2] of the potential which the LHC facility at large had for precision Higgs physics if one was eventually able to combine the HL-LHC results with the ones from LHeC. This analysis is illustrated in Fig. 3 as a comparison of the LHC facility, singly pp and ep & pp , with the International Linear Collider (ILC) in its initial and an upgraded configuration.

Two observations may be emphasised: (i) the addition of the LHeC to the LHC improves the proton-proton result very significantly for the dominant decay channels, $b\bar{b}$ and WW ,⁵ also for ZZ . The improvements on the gluon and tau channels arise from measurements which are roughly as precise in ep as in pp . A striking result is the expectation to measure the charm coupling accurately in ep , to $\delta\kappa_c \simeq 4\%$, which is considered to be inaccessible in pp scattering for large combinatorial background reasons; (ii) the combination of LHC and LHeC promises to deliver results some which are better than those at the ILC, and vice versa too. One observes an intriguing complementarity between ep and e^+e^- . The dominant production mechanisms are CC ep scattering and Z Higgsstrahlung, which explains why κ_W is estimated to be most precise in DIS and κ_Z in e^+e^- . Similar observations have been made in the comparison of the FCC-eh and FCC-ee prospects at an even higher level of precision [14].

Such comparisons have to be taken with care, they are the end product of advanced but prospect analyses, with many

⁵ This underlines the observation that the ep configuration has a unique sensitivity to the HWW vertex, which may lead to the observation of anomalous effects when measured with high precision [16].

inherent accelerator and experiment assumptions. One yet may conclude that the one percent accuracy level for SM Higgs physics is indeed in reach for the hadron (pp and ep) colliders, LHC and LHeC taken together, not worse than a next e^+e^- collider with a $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ luminosity horizon, having the important additional feature to provide a model independent Higgs width measurement.

It has been emphasised [17] that a triangle of precise pp , ep and e^+e^- measurements is essential for resolving inherent correlations in the Higgs sector. i.e. there is more involved than a competition for highest luminosity prospects.⁶

As discussed below, the LHeC is a single top facility. The cross section for the production of the Higgs boson in association with a single top is sufficiently large for measurements to be affected. In the SM, the production of the Higgs production in association with a single top is heavily suppressed due to negative interference. As such, it is very difficult to access this production mechanism at the LHC.

The LHeC provides a unique opportunity to study the CP structure of the Higgs boson Yukawa coupling [20]. One can introduce CP-phase ζ_t of the $t\bar{t}h$ coupling, where $\zeta_t = 0$ corresponds to the SM. As a result of the strong enhancement of $pe^- \rightarrow \bar{t}h\nu_e$ for $\zeta_t > 0$, strong limits can be set on deviations from the SM.

Assuming the Yukawa coupling to have the same structure as in the SM, the coupling size could be measured at the LHC with an accuracy of 17% with 1 ab^{-1} of integrated luminosity [20]. The use of multivariate techniques and additional channels, not employed so far, will further improve the precision of the measurement.

For the LHeC accelerator, especially the luminosity targets, and the detector, the striking Higgs potential has had a strong influence on their designs. The following specific demands have been derived in shaping the LHeC detector design: (i) b tagging is essential down to small angles corresponding to a pseudorapidity η of up to about 3; (ii) charm tagging requires an impact parameter resolution of $10 \mu\text{m}$ or better, which raises interest in an instrumented beam pipe; (iii) the hadronic calorimeter resolution should be as good as is feasible for reducing background; (iv) efficient photo-production tagging of backward low Q^2 processes is required for background rejection and normalisation; (v) scattered electrons shall be identified up to $\eta \simeq 4$, similarly for muons which may emerge from abundant, exotic $H \rightarrow \mu\mu$ decays and (vi) the reconstruction of missing energy poses a desir-

able η reach of up to 5. These considerations have been taken into account for the detector design.

The reconstruction, especially of rarer channels, sets a goal of $\mathcal{O}(1) \text{ ab}^{-1}$ of integrated luminosity, which is a desirable target too for BSM and very high x and high Q^2 DIS physics. The Higgs potential of the LHeC and FCC-eh is striking and worth every attention already in the design phase, setting, for example, an about 20 mA current goal for the energy recovery facility PERLE [21] which is being built at IJCLab Orsay.

2.3 Top quark physics

Electron-proton colliders at high energy are ideal to study the electroweak interactions of the top quark. The LHeC is an outstanding single top facility in its own right. The charged current cross section is 1.9 pb , compared to 0.05 pb of the photo-production of $t\bar{t}$. This provides an opportunity to measure the Wtb coupling with high precision and to search for anomalous contributions in the Wtb vertex [22]. With 100 fb^{-1} of integrated ep luminosity relative errors of order of 1% can be achieved in the measurement of the Wtb coupling. The Next-to-Leading Order corrections to the total and fiducial cross-sections are known [23] and do not significantly affect the ability of the LHeC to achieve precision. These may reduce the expected fiducial cross-section of single top production by 14%, while providing stability against scale variations. By contrast, measurements of single top production at the LHC are hampered by the large $t\bar{t}$ production cross-section. This is a further epitome of the complementarity of the LHeC and the LHC.

Given the high level of precision characteristic to the LHeC, other elements of the CKM matrix are also accessible with a precision superior to that of the LHC [24, 25]. Competitive measurements of V_{td} and V_{ts} could be performed at the LHC with $\mathcal{O}(1) \text{ ab}^{-1}$ of integrated luminosity.

In addition, important measurements of other top-quark properties, such as of the top-quark spin and polarisation, can be performed with competitive precision compared to the LHC, but with the advantage of theoretical simplicity and experimental clearness [26].

The photo-production of $t\bar{t}$ provides a window of opportunity to measure the $t\bar{t}\gamma$ magnetic and electric dipole moments [27]. Here an energetic photon couples only with the top quark so the cross-section depends directly on the $t\bar{t}\gamma$ coupling. The sensitivity of the LHeC here is superior to that achieved with measurements of the $b \rightarrow s\gamma$ transition and with the production of $t\bar{t}\gamma$ at the LHC.

The LHeC also provides access to Flavor Changing Neutral Current (FCNC) processes driven by the γtq and Ztq vertexes, where $q = u, c$ [28, 29]. This is achieved by measuring the process $e^- p \rightarrow e^- W^\pm q + X$. The expected sensitivity improves on current limits from the LHC by up to

⁶ This is an example for LHC and ILC, there are more electron-positron collider options under discussion. A comparative analysis of Higgs boson measurements at future colliders has been presented in [18]. The field, i.e. theory, HL-LHC prospects, e^+e^- collider prospects, see for example the ERL version of the ILC [19], and other aspects, such as the EFT view on these measurements, is very dynamic and cannot be comprehensively described here.

one order of magnitude in case of the γtu coupling, and is competitive with the expected accuracy from the HL-LHC.

2.4 Precision standard model physics

Through its Lagrangian, the Standard Model of particle physics defines all fundamental interactions of the elementary particles. While the SM is a highly successful theory, a main limitation of its exploration is that relevant free parameters, especially the coupling constant of the strong interaction $\alpha_s(M_Z)$ and the weak mixing angle, $\sin^2 \theta_W$, have only been measured with moderate precision. Any interference with new physics, from e.g. the dark sector, is expected to modify loop corrections of the SM, altering the SM predictions with increasing scales. Hence, the precise measurement of the fundamental parameters of the SM including their scale dependence is of crucial importance.

Strong coupling constant

The value of $\alpha_s(M_Z)$ is determined at the LHeC through precision measurements of scaling violations and as well from jet measurements in the Breit frame. Highest experimental precision is achieved since a color-neutral probe (the lepton) scatters off a colored parton inside the proton, which provides direct access to QCD phenomena and highest experimental precision through the charged lepton in the final state. The inclusive measurement over a large kinematic range in Q^2 and x requires precise vertexing and the measurement of the energy-momentum balance along the beam axis [30,31], aligned with the requirements to measure PDFs as listed above.

These measurements are improved with the challenging Silicon tracking detectors, here extended, and highly granular calorimeters with complete azimuthal angular acceptance. When determined together with the PDFs, the value of $\alpha_s(M_Z)$ will be determined at the LHeC with an uncertainty of ± 0.00022 [2]. This represents an impressive improvement by a factor of five as compared to the present world average value [32].

The measurement of inclusive jet cross sections is known to provide direct access to $\alpha_s(M_Z)$. These jet measurements require as well precise tracking, vertexing and a high resolution calorimeter with efficient discrimination between electromagnetic and hadronic clusters [33,34]. With its highly granular calorimeter, precise tracking, and the in-situ calibration (see below), the jet-energy scale uncertainty at the LHeC will be as small as about 0.5%. This will result in a determination of $\alpha_s(M_Z)$ with an experimental uncertainty of ± 0.00016 [2], which is about a factor 10 more precise than the best measurements from HERA jet data [35,36].

A per mille level measurement of α_s eventually challenges the so far dominating lattice result. It will pose new demands

for theoreticians requiring to match that small experimental uncertainty with improved higher-order predictions in QCD. The equally precise determination in inclusive and jet production in DIS is a crucial cross check for the theory, and so will be a comparison to future e^+e^- precision measurements. It is remarkable that this result will be free of higher twist and nuclear corrections. The LHC is unable to reach such precision but needs it, as for example for the interpretation of the gluon-gluon fusion production of the Higgs Boson being $\propto \alpha_s^2$.

At the LHeC the scale dependence of the strong coupling can be determined in a large range in (renormalization) scale, from a few GeV up to several hundreds of GeV, with small uncertainties. Such a unique measurement provides a direct test of the $SU_c(3)$ gauge-structure of the strong sector of the SM.

Electroweak measurements

The space-like momentum transfer in deep-inelastic scattering is mediated by electroweak interactions. With increasing Q^2 the contributions to NC scattering from γZ interference and pure Z exchange become more important, while CC DIS is mediated exclusively by a weak W boson. This provides the opportunity for direct measurements of fundamental properties of the electroweak sector of the SM. A highly granular electromagnetic calorimeter (EMC) improves the measurement through identification of final-state QED photons. At the LHeC the weak mixing angle will be measured through polarisation asymmetry and NC/CC ratio measurements in e^-p scattering as has been discussed in the initial CDR of the LHeC [1]. With an uncertainty of $\Delta \sin^2 \theta_W = \pm 0.00022$ [37] this is per se not an improvement over the precise Z -pole measurements at LEP+SLC [38]. DIS, however, is mediated in the t -channel and thus the scale dependence of $\sin^2 \theta_W$ can be measured for the first time with percent precision in the range of $\sqrt{Q^2}$ between 20 GeV up to nearly a TeV [2]. This could represent the first precision measurement of the weak mixing angle at high scales, away from the Z -pole, an inevitable measurement for particle physics, and a direct test of the gauge structure of the weak sector of the SM.

The Z -exchange directly probes the weak NC couplings of the light quarks inside the proton (u, d) to the Z -boson, and uncertainties of the vector and axial-vector couplings of the light-quarks of about 1% are achieved, which is an improvement by a factor of ten and more as compared to present best measurements [32,39].

The unique precision of the PDF and α_s measurements in ep are a requirement to measure the mass of the W boson to a few MeV precision, and similarly to extract a world leading measurement of the mixing angle from the LHC pp measurements as is discussed in [2].

2.5 Beyond the standard model searches

The clean environment of high-energy electron-hadron collisions provides an excellent framework for studying many extensions of the Standard Model. The excellent detector performance, the absence of pileup, and the large luminosity allows testing of entire classes of models that are difficult to study at the LHC. Many studies from recent years have been summarised succinctly in chapter 8 of ref. [2].

Prominent examples among these studies are searches for sterile neutrinos, for instance via lepton-trijets and displaced vertex signatures [40], heavy scalar particles with masses around the electroweak scale [41], and in general models with final states that look like ‘hadronic noise’ in proton-proton collisions [42,43].

Recent studies demonstrate that the LHeC could be a world-leading laboratory to study flavor-changing neutral currents in the charged lepton sector, in particular for processes that lead to electron-to-tau transitions, where the projected sensitivity could be an order of magnitude better than current and planned experiments in tau factories [44].

Scalar and fermion $SU(2)_L$ triplets can explain the observation of neutrino masses via the so-called type-II and type-III mechanisms, respectively. Both types of particles can be produced via their gauge interactions in vector boson fusion, but studying them at the LHC is very challenging due to the towering backgrounds. The prospects of finding triplet fermions via fat jet final states were shown to be feasible at the LHeC [45]. Triplet scalar searches at the LHeC were discussed in ref. [46].

Certain classes of leptoquarks can be studied at the LHeC if they interact with first generation fermions and have decay channels that are difficult to reconstruct at the LHC. It is possible to test explanations of the flavor anomaly $R_{D^{(*)}}$ via the R_2 leptoquark at the LHeC via its decays into τb final states [47].

Less minimal models with a \tilde{R}_2 leptoquark that has a dominant branching ratio into right-handed neutrinos may escape the LHC searches, but can be studied at the LHeC [48]. The specific signature of a displaced fat jet, stemming from the decay of a long lived heavy neutrino, would be a very promising sign of this model at the LHeC [49] and could already be observable within the first few months of operation.

Dark photons with masses below 10 GeV can be tested in a decay-agnostic approach via distinct non-dglap scaling violations, which may be the smoking gun for LHeC searches [50]. In the event that the dark photon in this mass range is long lived and decays dominantly to lepton pairs, LHeC searches for displaced dark photon decays would be sensitive to an otherwise challenging region of the parameter space [51].

BSM physics poses requirements on the detector design which are important to consider for the tracker, calorimeter

and muon detectors, such as an impact parameter resolution of $10\ \mu\text{m}$, very good momentum resolution, low mass photon identification, high resolution missing energy, detection of long lived particles outside the inner tracker, full hermiticity and maximum pseudo-rapidity coverage to $\eta \sim 5$.

3 Detector design

The LHeC detector, initially designed in [1] and recently updated [2] as a modern general purpose ep detector, is a composite system made of several subcomponents: beampipe, tracking, calorimetry, magnets and a muon system, each optimised for its purpose and adapted to the interaction region, which has the peculiarity of hosting 3 beams, the 2 proton or ion beams, of which one is a spectator while the other one is interacting with the counter-rotating electron beam, or, in the joint eh/hh configuration possibly the hadron beam.

In the following section the LHeC detector baseline design and some of its subcomponents are discussed illustrating few aspects and recent developments. Some consideration to adapt the detector for higher energy running (HE-LHC or FCC) are briefly addressed. A separate Sect. 5.2 below is devoted to modifications that would be required should the LHeC detector be used to also register hadron-hadron collisions.

3.1 Requirements

Requirements on the LHeC detector have been presented above concluding the discussions of the key physics topics. The detector should be highly hermetic in order to maximise coverage, in both the forward and backward directions,⁷ to provide a precise measurement of the hadronic final states and of scattered electrons towards very low- Q^2 . For charged current processes, the reconstruction of kinematic variables is only possible through the hadronic final state measurement where an excellent performance on calorimetry for hadrons is required to reconstruct the missing energies and distinguish CC from NC events at low Q^2 . The good hermeticity is also important for calibration of the detector through transverse momentum balance using NC DIS and photo-produced di-jet events.

Fine segmentation and good resolution for the electromagnetic calorimeter is required all over the angular coverage to tag both low- Q^2 and high- Q^2 neutral current events. Good resolution in the hadronic section is also important to measure the missing energies for CC DIS as well as for QCD studies using jets.

⁷ Positive z is defined as the incoming hadron beam direction for ep/eA collisions.

Excellent flavour tagging performance is desirable, especially in the forward direction, for flavour decomposition of jets and for tagging the SM Higgs decaying to $b\bar{b}$ and $c\bar{c}$, which are predominantly produced at large positive η .

There are also various constraints and considerations to take into account from the accelerator and technical aspect of detectors:

- The detector is required to have a magnet system consisting of a central solenoid along with a dipole system to steer the electron beam allowing for head-on eh collisions at the interaction point;
- The non-interacting proton/ion beam has to bypass the ep interaction yet to be guided through the same beam pipe housing the interacting electron and proton/ion beams;
- The shape of the beam pipe has to allow for the electron beam generated synchrotron fan to leave the interaction region unaffected and with minimal back-scattering;
- Good vertex resolution implies a small radius and thin beam pipe optimised in view of synchrotron radiation and background effects;
- The tracking and calorimetry in the forward and backward directions are set up to take into account the extreme asymmetry of the DIS production kinematics, see [1], with multi-TeV energies emitted in the forward, proton beam direction while the electromagnetic and hadron energies emitted backwards are limited by the electron beam energy.
- Very forward and backward detectors have to be set up to access diffractive produced events and to tag photo-production processes besides measuring the luminosity with high precision in Bether–Heitler scattering, respectively.

These and further specific requirements from inclusive DIS, see Sect. 2.1, are basically known from the H1 and ZEUS experiments at HERA. However, at the LHeC they are posed with extra severity because of the much enlarged beam energies, broad physics programme and more ambitious precision demands driven by new physics opportunities and owing to the huge increase in luminosity – as a few days of LHeC operation correspond to 15 years of HERA operation. New requirements are worth emphasising that arise from a much extended physics programme for which Higgs measurements are a prime example.

3.2 A detector for DIS at the LHC

The present LHeC detector is illustrated in Fig. 4. The detector is asymmetric in design, reflecting the beam energy asymmetry. The design is largely based on established technologies from the LHC general purpose detectors, ATLAS and CMS. More advanced technologies are introduced, and will

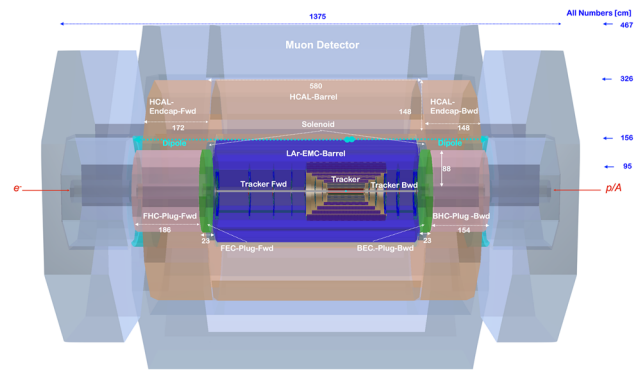


Fig. 4 Side view of the updated baseline LHeC detector concept, providing an overview of the main detector components and their locations. The detector dimensions are about 13m length and 9m diameter. The tracker is setup using pixel, macropixel and strip detectors. The barrel electromagnetic LAr-calorimeter EMC (in blue) surrounding the tracking region. The solenoid magnet is placed at radii immediately outside the EMC-Barrel, and is housed in a cryostat, which it shares with the weak dipole magnet that ensures head-on collisions. The hadronic calorimeter HCAL in the barrel part (colored orange); it uses steel and scintillating tiles) is located outside of the solenoid. The forward/backward electromagnetic calorimeters FEC/BEC (in green) and hadronic calorimeters FHC/BHC (in bright orange) are using Si-based sensitive and readout technology and as absorbers W/Pb and W/Cu, respectively [1, 2]. The muon detector (in grey) forms the outer shell of the detector. The detector description has been setup using **DD4hep** [52]

evolve with time, to fulfill the above described requirements and to adapt to different running conditions. The detector covers the angular range from 1° to 179° by the calorimeters to achieve the required hermeticity. Compared to pp running, the expected ep collision and background rates are about three orders of magnitude smaller relaxing the requirements on radiation hardness and also data acquisition. The pile-up rate is less than 0.1 per crossing at the LHeC for $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. The neutron field is expected to be a few orders of magnitude smaller than the LHC environment, estimated to be about 10^{13} neq/cm^2 .

As illustrated below in Sect. 4.5 (cf. Fig. 14) a dipole field is needed to steer the electron beam in the interaction region and allow for head-on collisions with the proton beam. The required dipole field (0.17 T over the range $z = [-8\text{m}, +8\text{m}]$) is combined in the central region with the central solenoid providing a field of 3 T, for a technical design see [1]. The synchrotron radiation generated by the electron beam in the dipole field is leaving the interaction region not affecting the detector performance thanks to the asymmetric design of the lightweight Beryllium beam pipe.

The main detector, from the interaction point to the outer radius, consists of: the silicon tracker with a central barrel part, forward and backward wheels; the electromagnetic calorimeter housed inside solenoid and dipole magnet; the hadronic calorimeter and the muon system. Not shown in the figure are backward (electron-side) detectors for low-

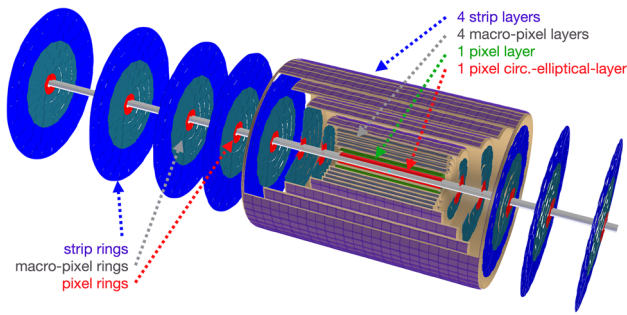


Fig. 5 Structure of the Silicon tracker for the LHeC detector

angle scattered electron to tag γp and γA collisions and forward detectors for neutrals ($n, \pi^0 \dots$) from the p/A remnant and protons spectrometer to measure proton momentum from elastic and quasi-elastic scattering.

This baseline design serves also a scalable configuration for HE-LHC and FCC-he where the main changes to be made for higher collision energy are the extensions for rapidity coverage in the tracking system and the depth (X_0, λ_I) in the calorimetry, both affecting mainly the size of the detector in the beam direction but only logarithmically. With respect to the earlier versions presented in the CDR and in the update, some optimisation has been done in particular to the silicon tracking, extending the radius from 60 to 80 cm, and to the calorimetry, choosing LAr for the EMC, which are described below in more detail. The larger tracking volume, with a longer lever arm measurement using more track points, allows for better resolution even at a slightly reduced B field. We expect that this configuration will deliver good and stable performance also in different experiment and accelerator configurations (*eh* and *hh* running) with certain modifications for *hh* as are described in Sect. 5.2.

3.2.1 Silicon tracking system

As described previously, excellent flavour tagging ability, including charm decays, is required across wide angular range, in particular towards forward rapidities. The decay particles from the SM Higgs may go beyond $|\eta| > 2.5$, the usual tracking coverage for the LHC pp detectors. The silicon tracker is shown in Fig. 5. It covers up to $|\eta| < 3.6$ with at least six hits and two hits for $-4.3 < \eta < 4.8$ using seven forward and five backward disks. In comparison to the previous LHeC tracker, the outer radius was extended from 60 cm to 80 cm and the number of layers in the barrel region from 7 to 10 layers while the magnetic field of the solenoid was reduced from 3.5 T to 3 T. Using the **tkLayout** tool [53] for optimising the tracker arrangement and minimising the material impact over a large region of η , the calculated radiation length figure shows tolerable levels as shown in Fig. 6. Some properties of the tracker setup are summarised in Table 1.

Radiation Length by Category

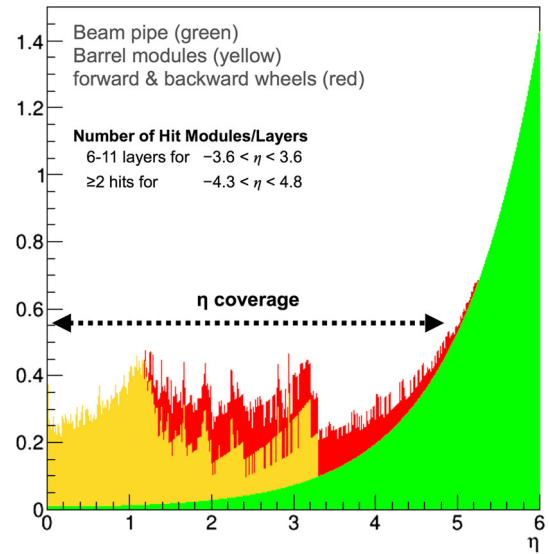


Fig. 6 Calculation of the radiation length budget of the LHeC tracker using **tkLayout** [53]. Support structures and services are not included

Table 1 Summary of the main properties of the tracker modules in the revised LHeC detector configuration based on calculations performed using **tkLayout** [53]. $\eta_{\max/\min}$ denotes the pseudo-rapidity range. #Layers_{Barrel} are the number of layers in the barrel and #Rings_{Wheels} the number of wheels in the End Caps, Fwd and Bwd tracker parts, respectively

LHeC Tracker Part	η_{\max}	η_{\min}	#Layers _{Barrel}
Inner Barrel			
pix	3.3	- 3.3	2
pix _{macro}	2.	- 2.	4
strip	1.3	- 1.3	4
#Rings _{Wheels}			
End Caps			
pix	4.1/- 1.1	1.1/- 4.1	2
pix _{macro}	2.3/- 1.4	1.4/- 2.3	1
strip	2./- 0.7	0.7/- 2.	1-4
pix	5.2	2.6	2
Fwd Tracker			
pix _{macro}	3.4	2.2	1
strip	3.1	1.4	4
pix	- 2.6	- 4.6	2
Bwd Tracker			
pix _{macro}	- 2.2	- 2.9	1
strip	- 1.4	- 2.5	4
Total $\eta_{\max/\min}$	5.2	- 4.6	

The relatively small radiation level allows to employ CMOS-based technology for the inner silicon tracker. Depleted CMOS sensors, also known as Depleted Monolithic

Active Pixel Sensors (DMAPS), are position sensitive detectors in industry standard CMOS or High Voltage-CMOS (HV-CMOS) processes [54]. These sensors are extremely attractive for experiments in particle physics as they integrate the sensing element and the readout electronics in a single layer of silicon, which removes the need for interconnection with complex and expensive solder bump technology. Depleted CMOS sensors also benefit from faster turnaround times and lower production costs when compared to hybrid silicon sensors. The final choice will depend on the region of exploration. Low-fill factor DMAPS have been or are being prototyped and produced for several experiments in particle physics, such as Mu3e [55], ATLAS [56], LHCb [57], CLIC [58] and ALICE [4,59] in a few different processes. Today’s most performant DMAPS detectors are 50 μm thin and have 50 μm × 50 μm cell size with integrated mixed analogue and digital readout electronics, 6ns time resolution and 2 × 10¹⁵ 1 MeV neq/cm² radiation tolerance. The development is ongoing and extends towards radiation hard technologies. Interesting for our purpose are the possibilities of features offered by CMOS imaging sensor technologies, called stitching, which allows developing a new generation of larger size MAPS using wafers that are 300mm in diameter. Moreover, the reduction of the sensor thickness to values as small as about 20–40 μm shall allow for exploiting the flexible nature of silicon to possibly implement large-area curved sensors. In this way, it looks feasible to build cylindrical or in general curved layers of silicon-only sensors, with a significant reduction of the material thickness by avoiding overlap between sensors [60–62].

The challenge in vertexing at the LHeC is that the beampipe has to accommodate the synchrotron radiation fan from the electron beam. To minimise the impact, the innermost barrel pixel layer is designed to follow the optimized circular-elliptic shape of the beampipe as shown in Fig. 7. Thanks to the integrated read-out electronics of the DMAPS sensors, the layout of the innermost layer can be flexible. The current design uses a scheme with many narrow sensors in the x–y coordinate plane, see Fig. 7, following the shape of beampipe as closely as possible. A possibility to use the bent sensors as described above is being studied.

3.2.2 Calorimetry

As described in Sect. 3.1, the LHeC requires well developed electromagnetic and hadronic calorimeter sections. The electromagnetic calorimeter (EMC) surrounds completely the silicon tracker and can be subdivided into a barrel, a forward and a rear system.

For the barrel region two options had been considered: a cold option using Liquid Argon, copper electrodes and lead absorbers, and a warm one based on lead absorbers and scintillator tiles. Liquid Argon is known for its high

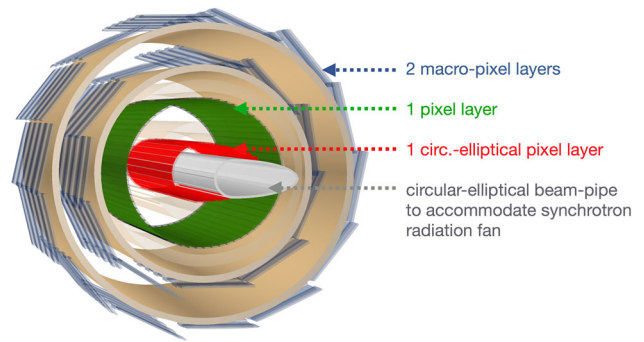


Fig. 7 A view of inner 4 layers of central barrel tracker with innermost circular-elliptical silicon pixel layer following the shape of beampipe

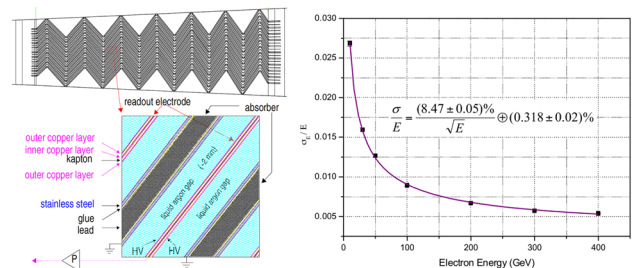


Fig. 8 Longitudinal view of one cell of the ATLAS LAr Calorimeter, showing the accordion structure (left). The LHeC LAr energy resolution for electrons between 10 and 400 GeV as simulated using GEANT4 [73] (right) [1]

resolution, linearity, long term stability and radiation tolerance confirmed over many years in ATLAS and H1 experiments [33,63–72]. The cryogenic system required for the LAr option can be combined in the LHeC detector with the one from the magnet system, which is directly surrounding the calorimeter. The flexibility in the longitudinal and transverse segmentation, and the possibility of implementing a section with narrow strips to measure the shower shape in its initial development, represent additional advantages [71].

Figure 8(left) shows a detail of the accordion-electrode structure. A basic cell consists of an absorber plate, a liquid argon gap, a readout electrode and a second liquid argon gap. The mean thickness of the liquid argon gap is constant along the whole barrel and along the calorimeter depth. The LHeC LAr electromagnetic calorimeter would also provide the required energy resolution and detector granularity (Fig. 8(right)). More details on the LAr calorimeters proposed can be found in the CDR [1]. As an option, in [2], alternatively a warm lead-scintillator electromagnetic calorimeter has been simulated for comparison. Its advantage compared to the LAr calorimeter are no cryostat walls in front of the barrel EMC and a convenient modularity. Given, however, its significantly better resolution performance, preference was decided to give to a barrel LAr calorimeter.

The hadronic calorimeter in the barrel part is a sampling calorimeter using steel and scintillating tiles as absorber and

active material, respectively, for good resolution. This also provides mechanical stability for the magnet/dipole cryostat and the tracking system.

Calorimetry in the forward and backward direction at the LHeC needs very fine granularity for position resolution, good e/π separation through shower shape and also good resolution, especially for the scattered electron. The very forward and to a much lesser extent – in eh – the backward parts of the calorimeter are exposed to high levels of particle radiation and must therefore be radiation hard by design. Tungsten (W) is considered as the absorber material, in particular for the forward inserts (electromagnetic and hadronic inserts), because of its very short radiation length. Since the backward inserts have looser requirements, the materials for the absorbers are chosen as lead (Pb) for the electromagnetic part and copper (Cu) for the hadronic one.

The active signal sensors have been chosen to be silicon-strip for the electromagnetic forward/backward calorimeters and silicon-pad for the hadronic fwd/bwd calorimeters. The demanding requirements of very forward /backward angle resolution favors fine segmentations of calorimeter cells interconnecting the tracking and calorimeter information for best particle-tracking and -identification. Si based tracking/imaging calorimeters appear to be appropriate to withstand the higher radiation load near the beam-pipe. They also open the opportunity to measure the neutral component of particle flow as already demonstrated by developments of the CALICE collaboration for the Linear Collider [74–76].

The hadronic calorimeter compensation algorithm would profit as well from knowing the neutral part of shower development best. The steel structures are in the central and plug calorimetry close the outer field of the central solenoid. The total depth of the electromagnetic section is about 30 radiation lengths on average in the barrel and backward regions. In the forward direction, where particle and energy densities are highest, the segmentation/granularity has to be finer, and it varies with radius and depth. The hadronic calorimeter has a depth of between 7.1 and 9.6 interaction lengths, with the largest values in the forward plug region. For each of the calorimeter modules, the pseudorapidity coverage, the types of the absorber and sensitive materials used, the number of layers, radiation or interaction lengths, and the energy resolutions obtained from GEANT4 simulations can be found in Ref. [2].

3.2.3 Muon system

Muon identification is an important aspect for any general purpose HEP experiment. In the baseline LHeC detector design, the muon system provides a reliable muon tag signature which is used in conjunction with the central detector for muon identification, triggering and precision measurements. Detection of long lived particle decays suggests to foresee

not too short track segment lengths in the muon detector. The detector elements are organised in a nearly hermetic envelope surrounding the hadronic calorimetry. In terms of technology choices, the options in use in the LHC general purpose experiments [77, 78] and their planned upgrades are adequate for LHeC since muon background rates are lower. A solution composed of layers of last generation Resistive Plate Chambers (RPC), providing the Level 1 trigger and a two coordinate (η, ϕ) measurement and possibly aided Monitored Drift Tubes for additional precision measurements appears as appropriate [2]. In the baseline design, the muon chambers have a compact multi-layer structure, providing a pointing trigger and a precise timing measurement which is used to separate muons coming from the interaction point from cosmics, beam halo and non prompt particles. This tagging feature does not include the muon momentum measurement, but is performed only in conjunction with the central detector. The muon detector will cover the region $|\eta| < 4.1$ and the transverse momentum coverage can be roughly estimated to be $p_T > 2-3$ GeV (at central pseudorapidities).

4 Accelerator considerations

The design of the machine is described in detail in the updated version of the LHeC design report [2]. It combines an optimised proton optics of the LHC interaction region (IR 2) with a dedicated new lattice for strong focusing of a 50 GeV electron beam that will collide with one of the counter rotating proton beams.

The electrons are accelerated in a new Energy recovery Linac (ERL) structure that allows high beam currents and at the same time a moderate energy consumption. The lattice of the combined interaction region and the beam optics of the three beams has to be flexible enough to provide in alternative manner standard LHC beam operation of pp (or AA) collisions for heavy ion physics and electron-proton (or electron-ion) collisions at the same LHC interaction point “IP2”.

The main challenges of this new layout are:

- the design of a high intensity high energy ERL that will provide a 50 GeV electron beam;
- an Interaction Region that is flexible enough for the alternative operation of ep and pp (or AA) collisions in IP2;
- and a beam combination and separation scheme to bring the electrons into head-on collisions with the 7 TeV proton beam (or 2.76 TeV/nucleon ion beam) and separate them after the IP.

Beyond these beam dynamics aspects, the aperture need of the mini-beta concept of electron as well as proton beam are defined by the fact that the ep mode in IR2 will have to

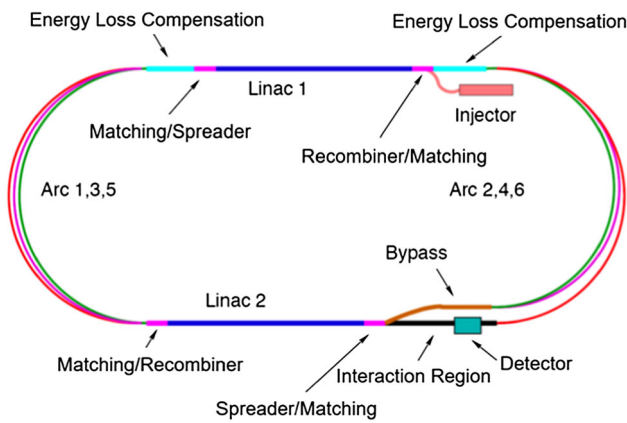


Fig. 9 ERL geometry, using two superconducting linear accelerators, connected by three-pass return arcs

be combined with standard *pp* collisions in the other LHC interaction points, IP1, IP5 and IP8. As a consequence, the new interaction region is optimised to house three beams: the colliding electron and proton beams as well as the second proton beam that will pass unaffected the same vacuum system. A special relaxed beam optics is foreseen for this non-colliding protons to limit the additional aperture need. For the *pp* operation mode – the standard scenario in present LHC operation – the so-called ATS optics will be applicable, that has been established in the context of the HL-LHC luminosity upgrade project, for both heavy ion and proton operation. The final focus system of the electron beam is embedded in this existing LHC proton structure. As a consequence its mini beta focusing system has to be extremely compact to be nested in the proton mini beta lattice. It will provide transverse beam dimensions of the electrons to match in both planes the dimensions of the LHC hadron beam. At the same time the electron lattice has to provide a sufficient beam separation and guide the electrons after the collision point back to the ERL return arc for the deceleration part of the energy recovery process.

The ERL is based on two super-conducting linacs of about 900 m length, which are placed opposite to each other and connected by three return arcs on both sides (Fig. 9). A final electron beam energy of 50 GeV is reached in this 3-turn racetrack ERL design. The concept allows to keep the overall energy consumption on a modest level for up to 20 mA electron current. The main parameter list is shown in Table 2.

4.1 Linac and RF system

The option to design a particle collider as Energy Recovery Linac, provides the opportunity to overcome or avoid a number of limitations of circular machines. In order to reach the luminosity of $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ with an electron energy of 50 GeV, the concept of an ERL offers the advantage of a

Table 2 ERL main parameters

Parameter	Unit	Value
Beam energy	GeV	50
Bunch charge	pC	499
Bunch spacing	ns	24.95
Electron current	mA	20
trans. norm. emittance	μm	30
RF frequency	MHz	801.58
Acceleration gradient	MV/m	20.06
Total length	m	6665

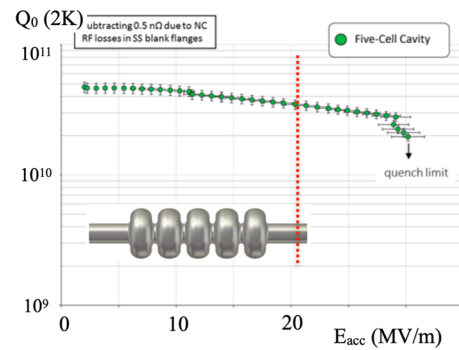


Fig. 10 Q-parameter of the 5 cell cavity prototype

high brightness beam, high beam currents with limited synchrotron radiation losses and it avoids limitations due to the beam–beam effect – a major performance limitation in many circular lepton colliders (e.g. LEP). On the other side, the current of the ERL as well as the emittance are determined by its source. An operational goal of $I_e = 20 \text{ mA}$ for the LHeC has been set, corresponding to a bunch charge of 500 pC at a bunch frequency of 40 MHz. Given three turns for the acceleration and deceleration, an overall current of 120 mA will be circulating in the ERL with impacts on the RF design, facing a virtual beam power of 1 GW. In order to limit RF losses, a super conducting (s.c.) RF system is foreseen with a required quality factor above $Q = 10^{10}$. In collaboration with JLab [79] prototypes have been developed: Fig. 10 shows the Q-value of a five cell s.c. resonator which lies comfortably above this value up to the required acceleration gradient, which is indicated by the red cursor line in the plot, and much beyond.

The further SRF technology development, the validation of these design concepts and the investigation of the ERL performance in terms of source brightness and stable and efficient operation are important tasks of the PERLE facility [21, 80].

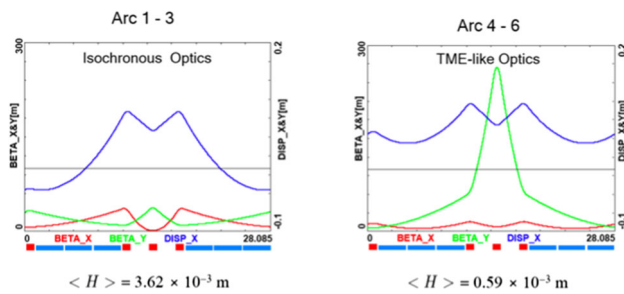


Fig. 11 Basic FMC cells of the ERL arcs: Isochronous (left) for arc 1, 2, 3 and TME lattice (right) for arc 4, 5, 6

4.2 Return arcs and spreaders

Special care has to be taken in the design of the ERL lattice: The optics of the three return arcs has to be optimised for the different challenges, that come along with the increasing beam energy [81]. At low energy, a flexible momentum compaction lattice will allow optimisation of the bunch length: An isochronous beam optics has been chosen for arc 1, 2, 3 to allow short bunches. At higher energies, in arc 4, 5, 6 an efficient emittance control is needed, as the effects of the emitted synchrotron light will take over. These arcs therefore are equipped with a theoretical minimum emittance optics (TME) to mitigate the emittance blow up (see Fig. 11).

The magnet structure of the linacs has to provide focusing for the complete energy range of the accelerating/decelerating beams. Here a FoDo structure has been chosen with a phase advance of 130° per cell. Different cell lengths have been investigated and simulation studies showed – not unexpectedly – an increasing performance for a shorter cell length. At the end of the Linac, the beam has to be guided into the return arc that corresponds to the beam rigidity at the given acceleration step. A combination of dipoles and quadrupole magnets provides the vertical bending and adapts the beam optics to the arc structure. This “spreader” (in front) and “re-combiner” (after the arc) represent a non-dispersive deflecting system to provide the necessary vertical off-set between the three arc modules and limit at the same time the detrimental effect on the vertical beam emittance.

4.3 Interaction region

The Interaction Region (IR) of the LHeC is a most challenging part of the machine: While seeking for highest luminosity in $ep(A)$ collisions, which includes mini-beta insertions for strong focusing of both beams, the colliding electron and hadron beams have to be separated after their collisions and guided to their lattice structures, to avoid parasitic bunch encounters. In addition, collisions and beam–beam effects with the second non-colliding proton beam have to be avoided and backgrounds such as beam halo be mitigated.

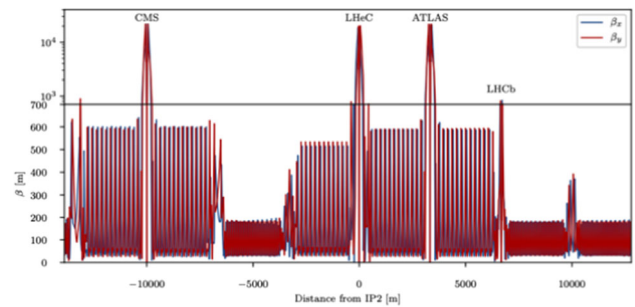


Fig. 12 LHC proton beam optics, optimised for the LHeC design values at the LHeC IP

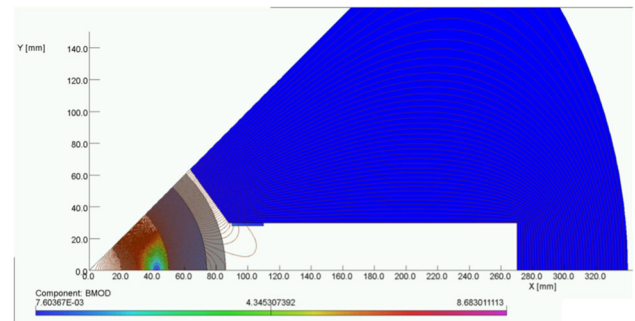


Fig. 13 Layout of the first proton quadrupole after beam separation. Special emphasis is put on minimising the remaining field in the electron aperture at a distance of 106 mm from the p design orbit

4.4 Proton beam optics

The optics of the colliding proton beam follows the standard settings of the HL-LHC. Figure 12 shows the proton optics at the interaction point of the LHeC. The long-ranging beta-beat which is an essential feature of the HL-LHC optics [82] is clearly visible on both sides of the IP and will be used for both, h-h and e-p collisions in IP2.

Special design effort is needed in the layout of the superconducting quadrupole “QA1” : Positioned right after the electron mini beta quadrupoles, it has to provide sufficient aperture and gradient to re-match the proton optics towards the arc structure. At the same time a field free region inside the cryostat is needed for the outgoing electron beam. Figure 13 shows a first layout of the magnet. The field calculations for both apertures are determined using the magnet design code ROXIE [83] with special emphasis on minimizing the remaining quadrupole field in the electron aperture: located at a distance of 106 mm from the proton design orbit – it has to be low enough not to distort the electron beam. Following the first layout and field calculations described, further R&D will be needed leading to a detailed design and construction of a prototype magnet in order to investigate the feasibility of the technical concept.

4.5 Electron beam optics and separation scheme

The design orbit of the electron beam – accelerated by the ERL and brought into collision at IP2 – will be merged with the proton orbit only in a short part of the lattice: Due to the different beam rigidities,

$$(B * \rho)_p = 23333 \text{ Tm} \quad (B * \rho)_e = 167 \text{ Tm}$$

a common focusing structure is not possible. The design of the IR therefore has to take a manifold of conditions into account: Focus the electron beam to the required β values in both planes, establish sufficient beam separation, optimise for smallest critical energy and synchrotron light power, and leave sufficient space for the detector hardware. A separation scheme has been established [84] that combines these requirements in one lattice structure (see Fig. 14). Due to the different rigidity of the beams, a separation is possible through the common effect of several magnetic fields: The spectrometer dipole of the LHeC detector, named $B0$ in the figure, is used to establish a first separation. Right after and as close as possible to the IP, the mini-beta quadrupoles of the electron beam are located. They provide focusing in both planes for matched beam sizes of protons and electrons at the IP:

$$\beta_x(p) = \beta_x(e), \quad \beta_y(p) = \beta_y(e)$$

Moreover, they are positioned off-center with respect to the electron beam, thus acting as combined function magnets to provide the same bending field as the separator dipole: A quasi-constant, soft bending of the electron beam is achieved throughout the magnet structure: $1/\rho_{B0} = 1/\rho_{quad_f} = 1/\rho_{quad_d}$. Additional conditions were put for a reduced beam size of the electron beam at the location of the first proton quadrupole. At this position, $L^*=15$ m, the reduced electron beam size leads automatically to a minimum of the required beam separation and as direct consequence to smallest synchrotron radiation effects. The optical functions of the electron beam in this optimised interaction region are shown in Fig. 15.

4.6 Joint $ehlhh$ operation options

The interaction region layout described above has been optimised for highest luminosity, matched beam sizes between electrons and the colliding proton beam and a smooth but efficient beam separation scheme. Still, an additional boundary condition arises from the second, “non-colliding” proton beam: A concurrent operation of the LHeC as electron-proton collider means that the operation as ep collider will be possible in parallel to the standard LHC proton-proton operation. During ep operation in IP2, with electrons provided by the ERL, the standard pp collisions in the LHC interaction points IP1, (ATLAS), IP5 (CMS) and IP8 (LHCb) will continue and

thus the second proton beam has to be guided through the new interaction region IR2, in parallel to the electron and proton beams. At IP2 therefore, in ep operation mode, the second non-colliding proton beam will be separated by a symmetric orbit bump to avoid direct collisions between the two proton beams as well as with the electron beam. Parasitic encounters with the subsequent bunches are suppressed by a vertical crossing angle. This scenario follows the LHC standard operation, where similar orbit bumps are applied during injection and acceleration phase of the two beams. Additional constraints arise from the need to preserve the overall LHC geometry: The two LHC proton beams will have to cross over from the inner ring to the outer, see Fig. 16.

All in all, two basic operation modes have to be established:

- Standard hh collisions in IP 1,2,5,8, no electron beam.
- Concurrent operation of eh collisions in IP2 and hh collisions in IP 1,5,8.

Concerning the first operation mode, the set up will be equivalent to the HL-LHC upgrade lattice and optics, with the two hadron beams colliding in all LHC interaction points. The magnets of the electron mini beta structure and beam separation scheme, shown schematically in Fig. 14, will be switched off.

For the second operation mode the colliding hadron beam (protons or ions) will be focused to match the size and position of the ERL electron beam at the IP. Electron beam focusing and beam separation between electrons and hadrons will follow the above mentioned scheme and due to the flexibility of the ERL optics a wide range of parameters (i.e. beam sizes) at the IP can be achieved, fulfilling the requests of both, e-p and e-A collisions. The second non-colliding LHC beam, however, will pass untouched through IR2, but still being used for collisions and data taking in IP 1,5 and 8. For this purpose, a sufficient beam separation between this non-colliding proton and the colliding beams in IP 2 is needed. Schematically the situation is shown in Fig. 17. The beam separation is established via the LHC standard separation bumps, that are used during beam injection and throughout the complete acceleration phase.

At the interaction point IP2, direct collisions are avoided by a horizontal offset of the non-colliding beam. In addition a vertical crossing angle is applied to prevent effects from the so-called parasitic encounters, that otherwise would occur at a distance of half a bunch spacing (25/2 ns). While this scheme is used during LHC standard operation it requires special attention for the concurrent ep/pp operation. A special beam optics for the non-colliding proton beam has to be established, to provide sufficient aperture for this new type of beam operation. The colliding proton beam will be focused strongly to achieve a β -function of 10 cm at the IP.

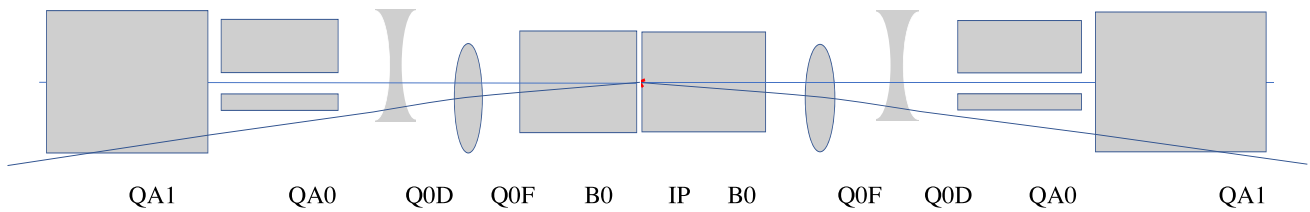


Fig. 14 Schematic view of the combined focusing and beam separation scheme

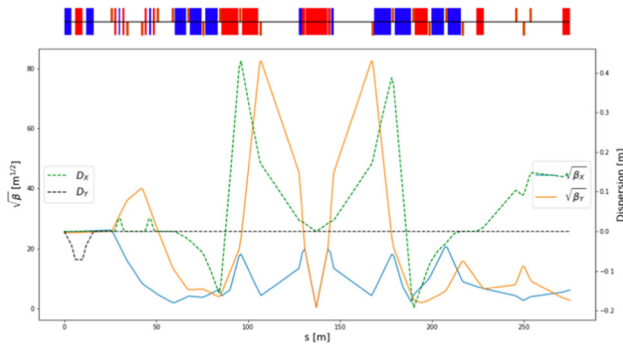


Fig. 15 Optical functions of the electron beam in the IR

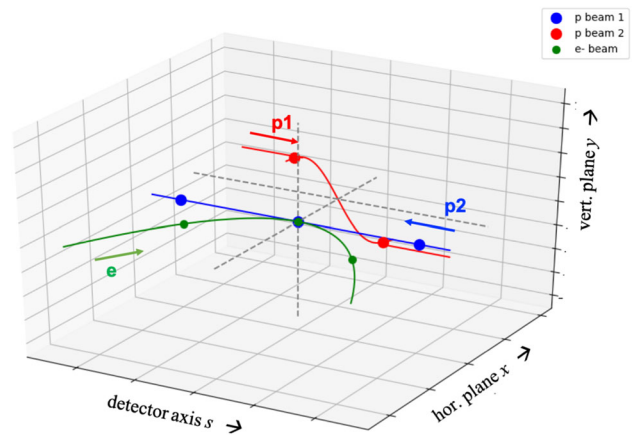


Fig. 17 Schematic view of the three beams in the interaction region. Collisions between electrons and proton beam 1 and a well separated proton beam 2

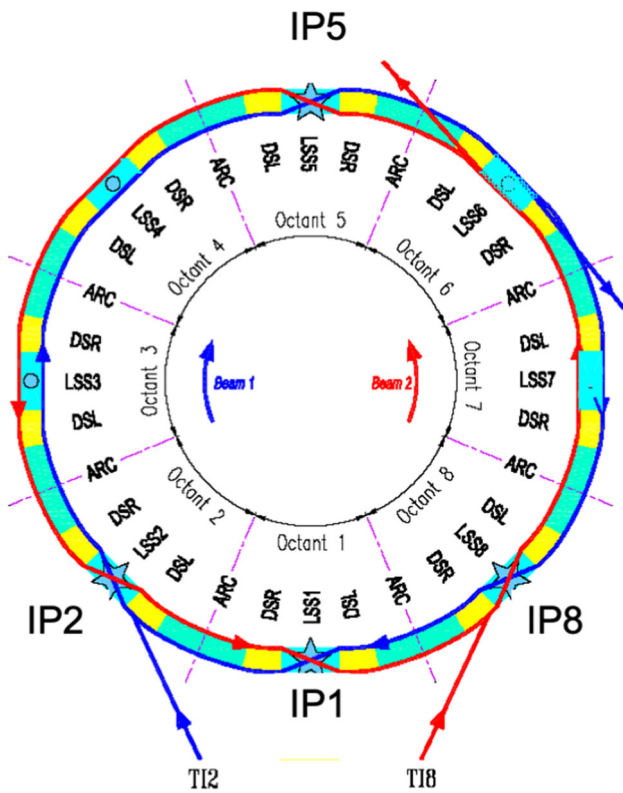


Fig. 16 Geometry of the two LHC beams, crossing from inner to outer ring in the four interaction points IP1, 2, 5 and 8

At the same time the non-colliding beam will see a relaxed optics with smallest achievable beam size in the proton mini beta quadrupoles. First estimates, based on an “injection type optics” with β^* of 15 m lead to an additional aperture request

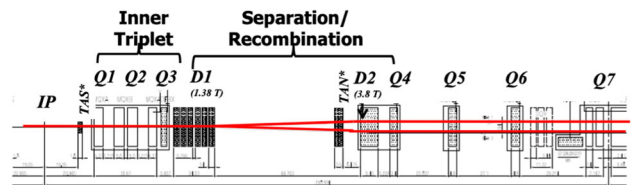


Fig. 18 Schematic view of the LHC proton beam separation scheme. The two separator dipoles D1 and D2 provide the hor. separation needed, before the beams enter their distinct magnet lattices in the arcs

of about 10% in the first proton quadrupole Q1A. Further downstream the two proton beams will follow the usual beam separation defined by the separator dipoles D1 and D2 (see Fig. 18). Further studies will concentrate on the level of flexibility of the different LHC magnet lattices which is a prerequisite for the proposed scenario. Beyond that, the beam-beam effect between the electron and the non-colliding proton beam – traveling for a considerable distance in parallel to each other – will be studied in detail.

4.7 Synchrotron light

The synchrotron light parameters, i.e. critical energy, radiation power and the geometry of the emitted light cone were determined with the simulation code BDSIM [85]. As expected, the synchrotron light conditions in the arcs become

Table 3 Critical energy and power of the emitted synchrotron light in the return arcs of the ERL

Arc	Energy (GeV)	Crit. energy (keV)	Power (MW)
1	8.75	3.2	0.01
2	17.00	23.9	0.21
3	25.25	78.5	0.75
4	33.5	183.3	2.45
5	41.75	354.8	5.87
6	50.0	609.3	12.17

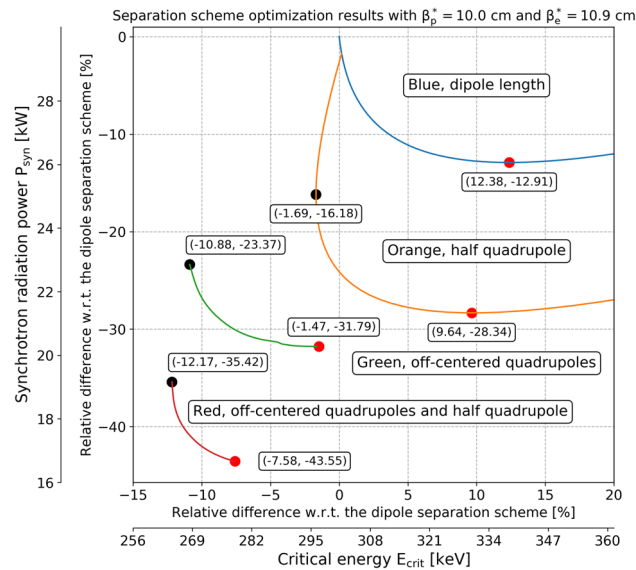


Fig. 19 Optimising the synchrotron light for lowest critical energy and power in the IR, details in the text

more serious turn by turn, reaching the highest level in the return arc 6, after the collision point. The values are summarised in Table 3. Special care is needed in the vicinity of the particle detector. The properties of the focusing elements, the separation scheme and the geometry of the interaction region have been optimised for smallest critical energies and power of the emitted light.

Figure 19 summarizes the results. The graph shows the reduction of the critical energy and power in the interaction region, due to the different steps of the optimisation procedure. Starting from a pure separator dipole design to establish the required beam separation, the concept of a half-quadrupole as first focusing element in the proton lattice is introduced as well as an improved beam separation of the electrons by off-centre quadrupoles.

The actual distribution of the detector dipole field and the off-centre quadrupoles has a considerable effect: The red and black points in the graph correspond to the minimum achievable critical energy and emitted power, respectively. Dedicated calculation of the synchrotron light cone and a

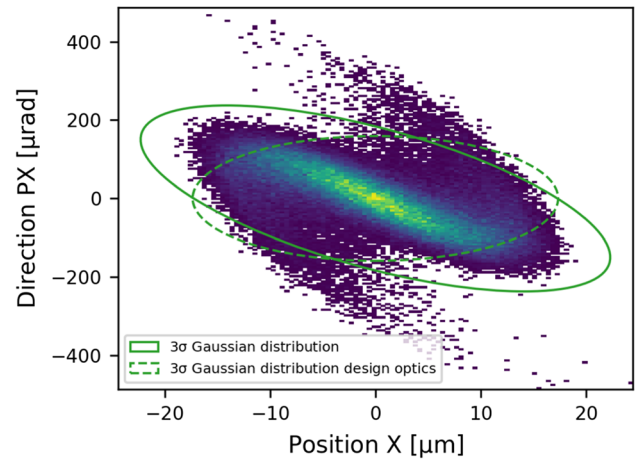


Fig. 20 Phase space of the electron distribution after beam collision, backtracked to the IP for matched optics conditions of electrons and the HL-LHC proton beam

Table 4 Luminosity and transverse emittance growth for the optical matching

Parameters	Optical matching
Luminosity	$8.2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
$\Delta\gamma\varepsilon$	15 mm mrad

sophisticated machine detector interface including absorbers will be needed to shield the detector parts and accelerator magnets.

4.8 Beam–beam effects

The beam–beam effect will always be the final limitation of a particle collider and care has to be taken, to preserve the beam quality and to limit detrimental effects on the emittance for assuring a successful energy recovery process in the ERL.

The beam–beam interaction has been simulated with a weak strong tracking simulation for a matched transverse beam size of the electron and proton beam at the IP. In Fig. 20 the situation post collision is represented in the (x, x') phase space. While tails in the transverse beam distribution as consequence of the beam–beam effect are clearly visible, the core of the beam still remains in a quasi ellipse like boundary. The coordinates obtained are used as starting conditions for the deceleration part of the ERL for a full front-to-end simulation.

The resulting emittance increase and luminosity, taking into account the beam–beam force are summarised in Table 4.

The beam–beam effect on the proton bunch remains in the shadow of the other effects and is considered as not critical. A careful alignment of the electron bunch at the IP, however, will be necessary as it could lead to undesirable proton emittance growth build up [2].

Table 5 Comparison of the electron distribution after non linear beam–beam interaction. The values represent the density of electrons for several rms emittance areas for the design optics as well as for the fitted ellipse of the post-collision distribution, in parenthesis

	Gaussian distribution (%)	Optical matching
1 σ	68.3	46.3% (70.7%)
2 σ	95.4	78.4% (95.4%)
3 σ	99.7	95.8% (98.4%)
4 σ	99.9	99.2% (99.5%)

The phase space distribution of the electrons after beam collision does not follow a Gaussian distribution. The non-linearity of the interaction distorts the electrons on the edges as well as modifies the Twiss parameters from the original design as shown in Fig. 20. The distortion of the phase space impacts the particle density and makes the core and tail of the distribution more populated than a Gaussian distribution. Nevertheless, the ellipse fitted to the post-collision distribution, that takes into account the modification of the Twiss parameters at the interaction point – including a so called capture optics – has a higher central density and the tails are less populated.

Further studies are needed regarding the impact of a smaller beam size of the electrons at the IP e.g. following the quest for a luminosity optimum. In fact, the optimal separation scheme may need to be adapted, the beam stay-clear aperture in the mini-beta quadrupoles would decrease and could be a showstopper for this luminosity optimisation scheme and finally the use of not matched lepton/hadron beam sizes could lead to instability for the proton bunch.

4.9 Front-to-end tracking studies

The tracking simulations of the ERL have been performed with the tracking code PLACET2 [86] and include, beyond the properties of the magnetic fields, the Incoherent Synchrotron Radiation (ISR) and the weak-strong beam–beam interaction at the interaction point. The studies focused on the achieved transmission and the beam quality along the ERL passages, i.e. the emittance budget required, for different machine circumferences that are considered for the basic machine layout. The beam parameters used for the tracking simulations correspond to the main parameter list, see Table 2.

The optics design of the multi turn ERL is shown in Fig. 21 and presents the sequence of linacs and arcs leading to the interaction region with a strong focusing and accordingly large vertical beta function in the mini beta quadrupoles. The other peaks are located in the matching sections between the linac optics and the periodic arc structure. The tracking takes place over three acceleration turns until the IP. Three

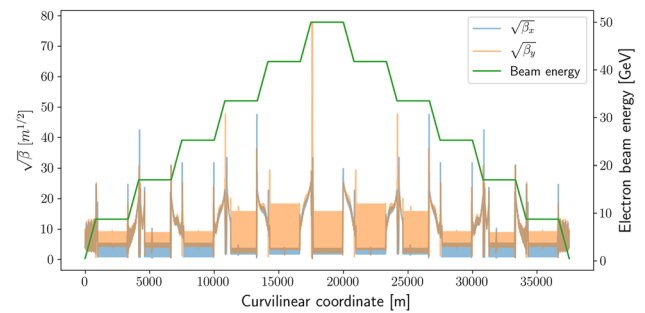


Fig. 21 Representation of the beta functions and the beam energy along the multi-turn ERL operation

Table 6 Results of the tracking simulations including beam–beam effect and synchrotron radiation for different ERL designs

ERL size	1/3 C_{LHC}	1/4 C_{LHC}	1/5 C_{LHC}
$\gamma \varepsilon_x^{inj}$ [$\mu\text{m mrad}$]	25.4	22.7	15.1
$\Delta p/p$ at IP	0.021%	0.029%	0.041%
Transmission	99.9%	98.9%	98.4%
Energy recovery	97.9%	96.7%	95.4%

deceleration turns are following in the same lattice structure, established via a RF phase shift in the highest energy return arc 6.

The objectives are: obtain the required transverse emittance at the IP; collide with the proton beam; minimise the emittance growth; taking into account eventual optics mismatch and distortion due to the non-linear beam–beam effect; decelerate the electron beam during the energy recovery process and guarantee minimum particle losses, while the energy spread will reach levels of a few percent in the last deceleration step.

The synchrotron radiation for each ERL circumference that has been studied varies significantly and has a strong influence on the energy recovery efficiency, see the results Table 6.

The results of the tracking simulations and the obtained emittance growth during the three turn beam acceleration agree nicely with the analytical calculations. After the interaction region the particles increasingly gain energy spread that creates a deviation from the design optics. The optics mismatch results in an extra emittance growth and ultimately leads to beam losses during the deceleration phase. The results of the emittance growth for the largest LHeC circumference studied, 1/3 of LHC, can be found in Fig. 22.

The recent tracking studies demonstrated that the electron beam quality can be preserved until the IP in order to meet a normalised transverse emittance of 30 mm mrad at the interaction point. Then, it is followed by a strong non-linear beam–beam interaction and finally decelerated over 3 turns to be dumped at 500 MeV. The tracking results of all the

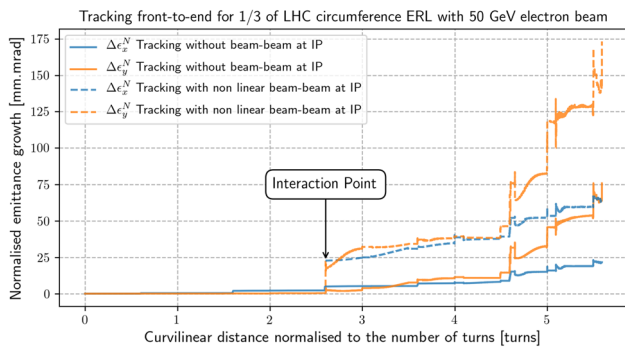


Fig. 22 Emittance growth along the curvilinear coordinate for the largest ERL design, corresponding to 1/3 of the LHC circumference

ERL circumferences studied, also including the synchrotron radiation and the beam–beam disruption, give an excellent transmission of close to 100%, see Table 6. The energy recovery efficiency is mainly constrained by the synchrotron light losses in the arcs that need to be compensated by extra RF cavities. It can be noted that the ERL designs that consider smaller machine circumferences require a smaller horizontal injection emittance that will potentially not allow enough margin for further studies including magnet field errors and misalignment.

5 Joint HI physics and a detector

The initial studies on a common IR for electron-hadron and hadron-hadron scattering, as presented above, suggest to indeed think of its possible joint use. Subsequently, a heavy-ion physics programme is sketched which would be enabled by combining *eh* and *hh* opportunities. The section also presents a sketch of how the LHeC detector would need to be modified when used also for *hh* collisions. Would this promising direction be followed, a joint design concept would be a next natural step, regarding obviously the beam pipe, particle ID, hadronic calorimetry, end cap and other design features.

5.1 Heavy ion physics

The physics opportunities provided by the availability in the same detector of DIS off nuclei, proton-nucleus and nucleus-nucleus collisions, *eA*, *pA* and *AA* respectively, are immense (see e.g. the discussions in [87] and refs. therein):

- On the one hand, as extensively discussed in [1,2], *eA* collisions at high energies at the LHeC will reveal the partonic structure of nuclei and the QCD dynamics in hitherto unexplored kinematic regions of high energies

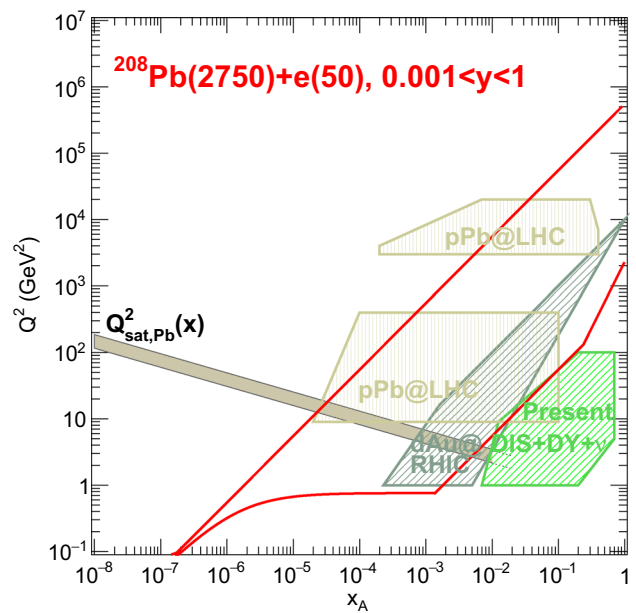


Fig. 23 Kinematic plane studied in *ePb* collisions at the LHeC [2] (solid red lines) together with the regions explored in present analyses [91]: DIS and DY fixed target data (hatched area in green), hadron production in *dAu* collisions at RHC (hatched area in grey) and Run 1 dijet and EW boson studies in *pPb* collisions at the LHC (hatched upper region in brown). Also shown in the hatched upper region in brown are the expected coverage from dijets in Run 2 [92] and from EW bosons in future Runs [88], and in the hatched lower region in brown the expectations from Run 2 D-meson analyses [93] and from DY and photon studies in future LHC Runs [88,94]. LHeC will not only access the saturation region at small *x*, “below” the Q^2_{sat} line and still small α_s , but covering 5 orders of magnitude in Q^2 can uniquely unravel the parton dynamics in nuclei

and parton densities. This is the region of relevance for *pA* and *AA* collisions at the LHC and beyond.

- On the other hand, the proposed heavy-ion (HI) detector to be installed in IP2 during LS4 to operate during subsequent LHC Runs [4] aims to provide outstanding tracking capabilities in the soft region down to tens of MeV and, due to fast timing, large possibilities for PID beyond *dE/dx* [59], and to be able to work and record minimum bias collisions at the largest *AA* achievable luminosity.
- The combination of outstanding track reconstruction, extended to ~ 1 degree in the backward and forward directions, and particle ID in the soft sector, with electromagnetic and hadronic calorimetry and muon detection makes this a general purpose detector for *pp*, *pA* and *AA* collisions, with larger capabilities for QCD than ATLAS and CMS and larger acceptance than LHCb.

Such detector configuration will be suited to explore the new possibilities for physics with ions after LS4 discussed in [88] as, for example, the larger luminosities provided by ions lighter than Pb (O, Ar, Kr) to analyse the presently

least understood stage of hadronic collisions, the initial one [89, 90], by using hard probes. In the following we elaborate on such possibilities.

5.1.1 Nuclear structure

The kinematic $x - Q^2$ range⁸ to be explored in eA collisions at the LHeC and during future pA Runs at the LHC is shown in Fig. 23 and compared with that of the set of data used in present analyses of collinear nuclear parton densities. LHC data will cover most of the kinematic region also covered by the LHeC (note also that the region between the lower and upper hatched regions in brown can be analysed by DY studies at LHCb), but the extraction of nuclear parton densities in pA and AA collisions relies on the validity of collinear factorisation down to rather low values of x and transverse momenta where other dynamics beyond leading twist perturbative factorisation could be at work. These new dynamics are strongly suggested by the finding at the LHC that many observables in pp and pA behave in a similar manner to that in AA , where they are interpreted as signatures of the existence of the Quark-Gluon Plasma (QGP) – the *small system problem*, see [95] and refs. therein, and [96–98] for recent studies in DIS at HERA. Besides, even assuming collinear factorisation to hold, the sensitivity to different flavours varies strongly when moving in the kinematic plane.

DIS offers fully constrained kinematics through the reconstruction of the electron angle and energy, a cleaner theoretical environment where factorisations can be proven [99] and perturbative calculations and resummations can be pushed to very high orders, and the possibility of full flavour decomposition through the combination of NC and CC and heavy flavour tagging. As shown in [1, 2], these opportunities will be fully exploited at the LHeC, where the nuclear PDFs can be determined with unprecedented precision without requiring prior knowledge of proton PDFs. This leads to a complete change of nuclear PDF physics, beyond the huge kinematic extension, as flavour by flavour and for the gluon one will determine the famous R ratios as a function of x and Q^2 which will enable the necessary distinction of nuclear effects, such as anti-shadowing [100] for example, from the appearance of non-linear parton dynamics.

Factorisation schemes exist beyond collinear factorisation, such as high-energy factorisation, TMD, . . . [99], and, eventually, the breaking of linear evolution when parton densities become large enough with decreasing x or increasing mass number of the colliding objects. Our current understanding of non-linear QCD dynamics views them as den-

sity effects, making both ep and eA equally important inputs to check such explanation. The combination of inclusive, diffractive and exclusive (vector mesons and photons) studies at the LHeC [1, 2] will establish the correct factorisation and dynamics in the different kinematic regions. Then, with the relevant non-perturbative information (PDFs, GPDs, TMDs, . . .) available, the validity of the corresponding factorisation will be checked in pA [88], thus elucidating the mechanism of particle production in high-energy nuclear collisions.

Finally, the possibility of accelerating ions lighter than Pb will clarify the dependence of parton densities on the mass number. Therefore, it will eliminate the need of interpolations, based on assumed factorisation of the mass number dependence, between different nuclear species in global fits. This will greatly reduce the theoretical uncertainties inherent to the interpolation procedure.

5.1.2 Soft physics

Soft heavy-ion physics deserves attention for a number of reasons. The low transverse-momentum spectrum of di-leptons may be sensitive to the restoration of chiral symmetry at high temperatures. Low energy photon spectra may be influenced by dynamics at the initial state. Physics studied beyond LS4 would focus on aspects that may not, or not fully, be answered before [101]: concerning the existence of charm-nuclei, hadronisation and thermalisation of beauty via its production and flow, production of multi-charm baryons, temperature evolution of the fireball via transverse momentum differential di-lepton measurements or $\rho - a_1$ mixing.

While the standard description of these collective features [95] is done in the framework of relativistic hydrodynamics, and the comparison with data used to extract QGP properties, it is known that hydrodynamics works well in out-of-equilibrium situations. In fact, it is currently believed that hydrodynamics is the long wavelength limit of quantum field theories. How this macroscopic description emerges from the microscopic QCD dynamics off the highly out-of-equilibrium initial conditions is the hottest topic in the field. To clarify this, it is crucial to establish the proper factorisation at work in pp , pA and AA collisions at high energies and the dynamics in the initial stages prior to the application of hydrodynamics. It is here where the contribution from ep and eA collisions in similar kinematic regions – at the LHeC – becomes key, as DIS is the ideal system to elucidate these aspects. It will also contribute to reducing the uncertainties in the extraction of QGP properties from the comparison of data with hydrodynamic calculations, those regarding the initial conditions and the initial stage dynamics [102]. Additional information on the initial stages of the collision and the thermalisation problem can be accessed using hard probes that

⁸ Note that in DIS Q^2 denotes the squared virtuality of the exchanged photon, while in hadronic collisions it designates the kinematic variable (transverse momentum, energy, mass, . . .) squared which enters as factorisation scale in the PDFs.

demand high transverse momentum particle detection and calorimetry.

5.1.3 Hard probes: vector mesons and jets

Among the hard probes, heavy-quarkonium production processes have always been a subject of special interest in high-energy physics. They involve both perturbative and non-perturbative aspects of QCD, corresponding to the production of the heavy-quark pair and its non-perturbative evolution. In addition to the hadronic experiments, the LHeC is particularly suited to study electro- and photo-production of quarkonia up to high energies [1]. These processes, which involve a highly virtual photon for electro-production or a real one for photo-production, provide unique opportunities for the study of the quarkonium production mechanism and the perturbative QCD calculation reliability. Moreover, the high gluon densities involved in these processes offer the opportunity to have an insight into the gluon generalised parton distribution in nuclei, the role of color correlations, and the color-dipole nature of quarkonia in unexplored regions far beyond those at HERA and the EIC.

At low transverse momentum, the proposed HI detector [4] at IP2 will offer the possibility to study separately the prompt and non-prompt quarkonium production with the identification of the contribution from excited states by detecting low energy photons. Such separation will allow a better characterisation of the QGP [103], based up to now on the anomalous nuclear dependence of quarkonium hadroproduction. Besides, the capabilities of the detector can have a great impact on the field of hadron spectroscopy, opening the possibility to measure the photoproduction of X, Y, Z states. Such studies demand an understanding of the production mechanism of quarkonia which presents large uncertainties until now, and of the effects of conventional, cold nuclear matter on quarkonia yields, both of the nuclear modification of parton densities but also of possible absorption or final state effects. Note that quarkonia are suppressed also in pA collisions, which constitutes one of the pieces of the small system puzzle. eA collisions at the LHeC, with the possibility of varying the nuclear species, are a crucial part of the here sketched hh/eh programme to clarify all these questions.

Of special interest is the physics of high transverse momentum particles and of jets, named *jet quenching*, usually employed in HI collisions as tools to analyse the QGP properties [104] but of great interest in QCD and SM and for searches of BSM as well. The consideration of a complete detector, adding electromagnetic and hadronic calorimetry and muon detection to a superb inner tracking device, opens numerous possibilities for studies of jet substructure, hadrochemistry and EM radiation within jets, heavy flavoured tagged jets, etc.

It is to be noted that jet quenching is the only observation in HI that has not been found in small systems. eA collisions at the LHeC offer the opportunity to study the influence of nuclear matter on jets [105], with abundant yields at high transverse momentum [1], thus contributing to the understanding of the small system puzzle and of the physics of jets in HI collisions. A related subject is the use of high transverse momentum particles and jets to understand the initial stage of hadronic collisions [89, 90, 106], an aspect that will benefit greatly from the possibility of varying the nuclear size of the colliding hadrons which provide larger centre-of-mass energies and luminosities [88].

5.1.4 Ultrapерipheral collisions

Ultrapерipheral collisions (UPC), in which one or both of the colliding hadrons act as sources of large fluxes of quasi-real photons, are a hot topic at the LHC [107]. They offer the possibility of studying photo-production, being in that sense complementary to DIS in which the photon virtuality can be controlled and varied. They have been exploited until now through studies of exclusive vector meson production and dijets with the aim studying nuclear PDFs, and of dimuons and of two-particle correlations in the search of collective effects in systems smaller than pp .⁹ UPC have also been used to study light-by-light scattering [109]. All these possibilities can be further exploited in a new detector which besides tracking, calorimetry and muon detection, will provide photon, electron, proton and nucleus detection in the very backward and forward regions. eA offers similar opportunities but in a much better controlled setup, with the additional possibility of further constraining the photon distribution inside the electron, see e.g. [110] and refs. therein. Further, the determination of the nuclear PDFs in inclusive processes in eA would verify the numerous assumptions underlying their extraction in UPC.

5.2 Detector concept for eh and hh collisions

As described above in Sect. 4.6, the new accelerator optics is able to provide collisions for eh and hh configurations in the same interaction point. As a consequence and if confirmed by further study, IP2 could indeed house one, common multipurpose detector serving for all of these, mostly related physics programs, of ep , pp , eA , pA and AA interactions, with high precision and large acceptance, and the unique advantage for cross-calibration of performance and physics.

Exploiting the LHeC detector described above, a first design of a detector suitable for both eh and hh collisions at IP2 is presented in Fig. 24. The detector is symmetrised in

⁹ In this respect, the ATLAS Collaboration claims the observation of azimuthal asymmetries [108] in γ Pb collisions.

the two beam directions, which is preferred for hh operation. The rapidity coverage of the electron side is extended from -4.6 to -5.2 by adding two additional wheel tracking layers and a deeper and symmetric calorimetry. The solenoidal magnetic field for normal eh running is chosen to be 3 T, while lower strengths may also be considered.

The barrel, from inside out, comprises an inner Silicon tracker, surrounded by the electromagnetic (LAr) calorimeter and a combined solenoid and dipole magnet section, at roughly 2 m radius, followed by a hadronic calorimeter and muon system, both outside the 3 T solenoidal field. It extends to 4.67 m radius, which fits into the L3 magnet support structure, see [1] for a study on the installation of an LHeC detector in the IP2 hall.

The LHeC central tracker provides momentum resolution of about 1% for $\eta < 3$ and p_T below a few GeV or about 10% for $\eta < 4.5$ and $p_T < 5$ GeV. Detailed acceptance parameters are listed in Table 1, where now the backward detector parameters are to be replaced by those in forward direction. The calorimetry also provides measurements for full rapidity coverage with fine segmentation, allowing to measure lepton, jets and energy flow to high precision as is illustrated in Sect. 3.2.2. The radial excursion of the LHeC barrel muon detector, as mentioned, is about 1.3 m.

Work on the LHeC detector concept continues, and more detailed simulations would be a next step to study its performance. The detail of the shape of the beam pipe for accommodating the synchrotron radiation fan and hence the shape of the circular-elliptical pixel layer are being worked out. A further optimisation may also involve the size of strip layers for occupancy in hh collisions and thickness of the calorimeters in backward direction. From this concept to a technical design, work is ahead for a couple of years, linked also to the FCC feasibility study and the FCC-eh detector. Novel ideas, as are summarised in the ECFA Detector R&D Roadmap, such as on high resolution Silicon time measurements and integrated readout and wafer configurations are directions which will influence the LHeC detector concept.

6 Summary

A detailed design has been presented and updated for the introduction of a programme of high energy electron-hadron scattering in a future phase of running of the CERN Large Hadron Collider. The design is based on collisions at interaction point IP2, utilising one of the LHC hadron beams, and assumes concurrent running with hadron-hadron collider experiments.

The electron beam as designed previously uses two superconducting linear accelerators of length around 900 m each, arranged in a racetrack configuration with three separate return arcs, allowing acceleration to an energy of 50 GeV

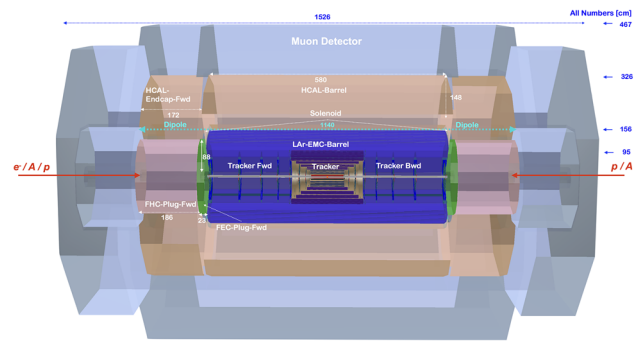


Fig. 24 Side view of a first design of the LHeC detector for both eh and hh collisions, where the detector coverage of the backward (electron) direction is extended to match that for the forward (hadron) direction. The overall size in beam direction increased from 13.75 to 15.26 m while the diameter is unchanged as compared to the new eh detector version presented above

before bringing the beams into collision with the LHC hadrons. A key feature of the electron accelerator design is energy recovery, for which plans for a lower energy, prototype ERL facility (PERLE) are well advanced.

The resulting LHeC experiment offers new sensitivity to a broad and original programme of physics at the energy frontier also complementing the existing LHC hadron-hadron experiments and their upgrades. Main physics topics are presented with a view on their requirements on the detector design. Highlights of the high luminosity LHeC ep programme picked out in this document include Higgs physics, in a joint ep and pp coupling analysis, studies of single top-quark production with correspondingly high precision on the Wtb vertex and competitive sensitivity to physics beyond the standard model across a range of processes that benefit from initial state leptons.

In terms of hadron structure, the LHeC allows the extraction of the complete set of parton densities with unprecedented precision, extending onto a new kinematic region, at low Bjorken x where new dynamics are expected and to high x where new physics may reside and factorisation be tested independently of power and nuclear corrections. In eA mode, the LHeC offers unique sensitivity to nuclear parton densities and exploits their enhanced sensitivity to low x effects over those of the proton, as well as complementing the relativistic heavy ion collision programme at the LHC and RHIC by providing cold-matter baselines for the understanding of quark-gluon plasma effects and contributing to a range of topics with hard probes as well as soft physics and ultra-peripheral collisions.

The ambitious physics programme is matched by a hermetic, compact, high performance LHeC detector design based around a strong (3 T) central solenoid for the precision measurement of high transverse momentum charged particles. Inner detectors based on depleted MAPS silicon sensors will provide tracking and vertexing at the highest

possible precision with a modest material budget. The pivotal importance of scattered electron detection and measurement in a DIS experiment is matched through electromagnetic calorimeter designs based on cold liquid argon with lead absorbers, building on technologies used successfully in previous experiments. The need for a high quality hadron response, from high transverse momentum jets to the inclusive measurement of the hadronic final state for kinematic reconstruction, is met using a steel/scintillating tile solution. The importance of forward and very forward (and backward) instrumentation is recognised by implementing detector components throughout the range $|\eta| < 5$ and by incorporating beamline instrumentation in the outgoing hadron and electron directions into the interaction region design from the outset.

Optimisations of the Interaction Region lead to improved performance parameters such as reduced synchrotron backgrounds enabling the high luminosity goals to be pursued. New studies have been presented on beam–beam effects and front-to-end tracking calculations been made for different circumferences of the LHeC racetrack.

A first configuration has been developed of an IR which may alternately serve eh and hh collisions while the other experiments at the HL-LHC stay in hh collision mode concurrently. The feasibility of realising a common eh/hh IR implies that one may realistically consider an apparatus which would permit both electron-hadron and hadron-hadron collisions to be registered.

The addition of ep and eA experiment capabilities to the LHC accelerator infrastructure, in combination with ongoing pp and AA programmes, deepens the sensitivity to new physics in the existing programme and adds TeV scale DIS physics in the quest for finding new particles, dynamics or symmetries beyond the Standard Model. Whether as part of a multi-purpose apparatus or operating in standalone mode, the LHeC offers fundamental, new perspectives on an energy frontier physics, detector and accelerator landscape in the 2030s which may indeed differ from today’s perceptions.

Acknowledgements This paper relies on a decade of work and collaboration with the hundreds of authors of the 2012 [1] and 2020 [2] LHeC design papers. The work has been accompanied by the CERN Directorate and guided by an International Advisory Committee chaired by em. DG of CERN, Herwig Schopper. Very fruitful discussions are acknowledged which some of us had with John Jowett, Luciano Musa and further members of the ALICE Collaboration. Special thanks are due to the organisers of the Off-Shell Conference, Kristin Lohwasser, Matthias Schott and colleagues, who encouraged and supported this paper to be written. JGM gratefully acknowledges the hospitality of the CERN theory group. We further acknowledge financial support by: Xunta de Galicia (Centro singular de investigación de Galicia accreditation 2019-2022); the “María de Maeztu” Units of Excellence program MDM2016-0692 and the Spanish Research State Agency under project FPA2017-83814-P; Fundação para a Ciência e a Tecnologia (Portugal) under project CERN/FISPAR/0024/2019; the South African Department of Science and Innovation and the National Research Foun-

ation; the UK Science and Technology Facilities Council; U.S. DOE under contracts DE-AC05-06OR23177 and DE-SC0012704; European Union ERDF; the European Research Council under project ERC-2018-ADG-835105 YoctoLHC; MSCA RISE 823947 “Heavy ion collisions: collectivity and precision in saturation physics” (HIEIC); and European Union’s Horizon 2020 research and innovation programme under grant agreement No. 824093.

Data Availability Statement This manuscript has no associated data or the data will not be deposited. [Authors’ comment: The manuscript contain discussions on the physics and accelerator and detector considerations about a new experimental proposal. The relevant data are contained in the tables and figures. Concerning the latter, if numerical values are required, they can be obtained from the authors by request.]

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Funded by SCOAP³.

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