

Theory requirements for SM Higgs and EW precision physics at the FCC-ee

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Abstract. High precision experimental measurements of the properties of the Higgs boson at ~ 125 GeV as well as electroweak precision observables such as the W -boson mass or the effective weak leptonic mixing angle are expected at future e^+e^- colliders such as the FCC-ee. This high anticipated precision has to be matched with theory predictions for the measured quantities at the same level of accuracy. We briefly summarize the status of these predictions within the Standard Model (SM) and of the tools that are used for their determination. We outline how the theory predictions will have to be improved in order to reach the required accuracy, and also comment on the simulation frameworks for the Higgs and EW precision program.

1 Introduction

With the discovery of the Higgs boson, all possible elements of the Standard Model (SM) have been experimentally confirmed and tested in great depth. On the other hand, observational evidence for dark matter and the matter-antimatter asymmetry require physics beyond the SM. One promising way to probe such new physics is through precision measurements of the properties of the Higgs boson. A complementary way is to measure electroweak precision (pseudo-)observables (EWPO) with highest precision. These are the avenues pursued by several proposals for a future e^+e^- collider. In particular, the FCC-ee concept is designed to run at $\sqrt{s} = 250$ GeV as a Higgs factory, and at $\sqrt{s} \sim M_Z, 2M_W$ for high-precision EWPO measurements. In this way it can improve (indirect) probes for new physics by several orders of magnitude compared to existing bounds [1–3].

The anticipated experimental accuracy of an observable has to be matched with a theory prediction of at least the same level of accuracy to make maximum use of the experimental data. For the accurate study of the properties of the Higgs boson, precise predictions for the various partial decay widths, the branching ratios (BRs) and the Higgs-boson production cross sections along with their theoretical uncertainties are indispensable. For the EWPO on the one hand the SM prediction has to reach the level of the experimental accuracy. Similarly, the extraction of these quantities from experimental data must be equally well under control. All these types of uncertainties (various theoretical ones and experimental) must be taken into account when deriving constraints on new physics from the data.

Several sources of theory uncertainties have to be distinguished. The *intrinsic* uncertainties are due to missing higher orders in the perturbative expansion of the SM (or BSM) prediction for an observable. The *parametric* uncertainties are due to the imperfect experimental knowledge of the SM input parameters (as well as theory uncertainties induced in their extraction from data). The extraction of a quantity from a cross section or an asymmetry requires the theory prediction of this cross section or asymmetry to at least the same order of precision¹.

In this paper the current status and future implications of theory uncertainties on (SM) Higgs-boson physics and EWPO's will be summarized. While we focus on the FCC-ee, they are valid for all future high precision e^+e^- colliders running above the HZ threshold (such as ILC, CLIC, or CEPC). We will use anticipated FCC-ee precisions to illustrate the impact of theory uncertainties. We build on [4], where also a discussion of the determination of the SM input parameters can be found. In addition, we review Monte Carlo (MC) event generator frameworks which are indispensable for simulations in the perturbative and non-perturbative regime, building upon the experience gained

¹ Theoretical predictions for the past electron collider experiments (LEP/SLC) were at least a factor 3 more precise than experimental data.

Quantity	FCC-ee	Current intrinsic unc.	Projected intrinsic unc.
M_W [MeV]	0.5	4 ($\alpha^3, \alpha^2\alpha_s$)	1
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	0.6	4.5 ($\alpha^3, \alpha^2\alpha_s$)	1.5

Table 1. Estimated precision for the direct determination of M_W and $\sin^2 \theta_{\text{eff}}^\ell$ at FCC-ee [2, 3] (column two), current intrinsic uncertainty (column three) and projected future uncertainty when leading 3-loop corrections become available (column 4).

from the LEP era [5–7] and the extensive simulations for CEPC, ILC and CLIC. Work on dedicated FCC-ee MC samples has started.

2 Pseudo-observables versus realistic observables

The quantities that can be directly measured in experiments are cross sections, line shape observables, forward–backward asymmetries etc., called “realistic observables” [8, 9]. The obtained results depend on the specific set of experimental cuts, however detector inefficiencies are removed and kinematic cuts are brought to simplified form using MC event generators. In order to determine idealized quantities like masses, partial widths or couplings, or the effective electroweak mixing angle, from the primarily measured realistic observables, a “QED unfolding” procedure is applied. This procedure involves manipulations like subtracting photon-exchange or interference terms, subtracting box-diagram contributions, unfolding higher-order QED or QCD corrections, etc. These secondary quantities are therefore called “pseudo-observables”. The Higgs-boson observables and EWPOs on which we focus here are pseudo-observables. It should be stressed that the role of EWPOs is to represent and encapsulate experimental data in the *model-independent* way². Fitting EWPOs to the SM and its extensions comes as an independent next step. In order to reach the required theoretical accuracy special care has to be taken in the application of these unfolding procedures – in particular the procedure of the extraction of EWPOs from data causes certain bias, partial loss of information on physics in the realistic data. The unfolding procedure for the LEP experiments is described in Refs. [8, 9]. A discussion of the additional problems in this respect arising in beyond the SM (BSM) models can be found in [10]. In Refs. [8, 9] it was proven that the bias induced in the extraction of EWPOs (QED deconvolution) was smaller than the LEP experimental errors. This will not be true in the FCC-ee era where experimental errors will be smaller by up to two orders. The very survival of the technique of the model independent EWPOs as the bridge between the experiment and the theory at FCC-ee will require significant improvement of the precision of the trivial but large QED effects. Ref. [11] covers this issue in a great detail, listing challenges and outlining solutions. In addition, the definition of EWPOs and the algorithm of their extraction will have to be modified. This is elaborated in Sect. 3 in Ref. [12]. Two major changes are anticipated. One is that we may be forced to remove from EWPOs not only QED but also first order pure electroweak corrections, in order to achieve sharper resolution in the search of the BSM effects in the new EWPOs. Another important change will be the increased role of Monte Carlo generators in extracting EWPOs from data – this is because semianalytic programs like ZFITTER and TOPAZ0 of LEP era will not be able to handle precisely enough detector geometry (cut-offs).

3 Electroweak precision observables

The most important electroweak precision observables (EWPOs) are related to properties of the Z and W bosons. The Z -boson properties are determined from measurements of $e^+e^- \rightarrow f\bar{f}$ on the Z -pole. To isolate the physics of the Z -boson, the typical set of pseudo-observables is defined in terms of the de-convoluted cross-section $\sigma_f(s)$, where the effect of initial- and final-state photon radiation and from s -channel photon and double-boson (box) exchange has been removed. The customary set of pseudo-observables are total and partial Z -boson cross sections with their subsequent decay to quarks or leptons, forward-backward and left-right asymmetries, from which the effective weak leptonic mixing angle, $\sin^2 \theta_{\text{eff}}^f$, is extracted, as well as the W -boson mass, M_W , see Ref. [13]. The expected (pure) experimental precision at FCC-ee of M_W and $\sin^2 \theta_{\text{eff}}^\ell$ is summarized in Tab. 1 based on Refs. [2, 3].

The quantities listed in Tab. 1 can be predicted within the SM in terms of the input parameters G_F , $\alpha(M_Z)$, $\alpha_s(M_Z)$, M_Z , M_H and m_t . The radiative corrections in these predictions are currently known including complete two-loop corrections. In addition, approximate three- and four-loop corrections of $\mathcal{O}(\alpha_t^3)$, $\mathcal{O}(\alpha_t^2\alpha_s)$, $\mathcal{O}(\alpha_t\alpha_s^2)$ and $\mathcal{O}(\alpha_t\alpha_s^3)$ are available (with $\alpha_t = y_t^2/(4\pi)$ and y_t being the top Yukawa coupling), see [14] for a review. The theory uncertainties from missing higher-order corrections are given in the third column of Tab. 1 [15–18]. Also indicated are the main sources for the respective uncertainties. In order to match the FCC-ee precision the theory predictions will have to be improved substantially. It was estimated in [4, 12, 19] that the intrinsic uncertainties will be reduced by

² For instance $\sin^2 \theta_{\text{eff}}^f$ is not a SM parameter – it parametrizes *data* on the ratio of the vector to axial couplings in $e^+e^- \rightarrow f\bar{f}$.

complete $\mathcal{O}(\alpha_s^2)$ corrections, fermionic $\mathcal{O}(\alpha^2\alpha_s)$ corrections, double-fermionic $\mathcal{O}(\alpha^3)$ corrections, and leading four-loop corrections enhanced by the top Yukawa coupling, resulting in the projected intrinsic uncertainties shown in column four of Tab. 1.

A crucial point here is that the determination of the pseudo-observables in Tab. 1 from experimental data also requires theory input for the removal of initial-state and final-state photon radiation and s -channel photon exchange and box contributions. The theory uncertainty from missing higher QED orders is estimated to amount to a few times 0.01% [5, 8, 9, 20] for the Z -peak cross-section and total width measurements. In order not to be dominated by this uncertainty, it will need to be reduced by about a factor 10 for the FCC-ee. This will require the calculation of non-leading log contributions to $\mathcal{O}(\alpha^3)$ corrections, $\mathcal{O}(\alpha^3L^2)$ and $\mathcal{O}(\alpha^4L^4)$ contributions (with $L \equiv \log s/m_e^2$) as well as an improved treatment of fermion pair production from off-shell photons.

The W -boson mass will be determined from a threshold scan near the W -pair threshold, $\sqrt{s} \approx 161$ GeV. It is foreseen that the experimental uncertainty at FCC-ee for this measurement is about 0.5 MeV [1–3]. At the point of highest sensitivity, an uncertainty of the cross-section measurement of 0.1% translates to an uncertainty of ~ 1.5 MeV on M_W [21]. Therefore a theoretical prediction for the process $e^+e^- \rightarrow 4f$ with an accuracy of $\Delta\sigma \sim 0.01\%$ is desirable, including effects of off-shell W bosons, which become important near threshold.

The currently best calculations are based on complete one-loop results for $e^+e^- \rightarrow 4f$ [22] and partial higher-order effects for the total cross section from an effective field theory framework [23, 24]. The resulting theory uncertainty on M_W is estimated to be about 3 MeV [24]. This result must be improved by complete 2-loop corrections to $e^+e^- \rightarrow W^+W^-$ and to $W \rightarrow f\bar{f}'$ (based on the effective field theory framework). In addition, a more accurate description of initial-state radiation will be important, which includes universal contributions from soft and collinear photon radiation (see [5, 8, 20] for a review), as well as hard photon radiation. For the latter, a proper matching and merging procedure needs to be employed to avoid double counting [23, 25], see [4] for more details. There it was estimated that a theory induced systematic uncertainty of $\Delta M_W \lesssim 0.60$ MeV can be feasible.

4 The SM Higgs boson

4.1 Higgs-boson production

The very narrow width of the Higgs boson allows for a factorization of all cross-sections with resonant Higgs bosons into production and decay parts to very high precision if the Higgs boson can be fully reconstructed. In this case, finite-width effects and off-shell contributions are of relative size $\Gamma_H/M_H \sim 0.00003$ and thus not relevant. If the Higgs boson is not fully reconstructable (e.g. in $H \rightarrow W^{(*)}W^{(*)} \rightarrow 2\ell 2\nu$) Higgs off-shell contributions have to be taken into account (which is straightforward at NLO).

At the FCC-ee with $\sqrt{s} = 240$ GeV (or other e^+e^- machines near this center-of-mass energy), the Higgs-boson production cross-section is strongly dominated by $e^+e^- \rightarrow ZH$, and $e^+e^- \rightarrow \nu\bar{\nu}H$ contributes less than 20% [1, 26]. For these two processes full one-loop corrections in the SM are available [27, 28]. For the dominating ZH production mode they are found at the level of ~ 5 –10%. Leading two-loop corrections to this cross section were evaluated at $\mathcal{O}(\alpha\alpha_s)$ [29, 30], which were found to amount to ~ 1 –2%. This number has to be compared to the anticipated experimental accuracy of 0.4% [1–3]. It becomes clear that with a full two-loop calculation of $e^+e^- \rightarrow ZH$ the intrinsic uncertainty will be sufficiently small. Computational techniques for $2 \rightarrow 2$ processes at the two-loop level exist, and it is reasonable to assume that, if required, this calculation within the SM can be incorporated for the FCC-ee Higgs precision studies (see also [31] for recent progress in this direction.)

For WBF production, the calculation of the full two-loop corrections will be significantly more difficult, since this is a $2 \rightarrow 3$ process. However, one may assume that a partial result based on diagrams with closed light-fermion loops and top-quark loops (in a large- m_t approximation) can be achieved, which should reduce the intrinsic theory uncertainty to below the 1% level. Given the fact that the WBF process is less crucial than the HZ channel for the Higgs physics program FCC-ee with $\sqrt{s} = 240$ GeV, this will probably be adequate for most practical purposes. It would be desirable to have a complete $2 \rightarrow 4$ calculation at two loops, including both processes with their interference, but very likely factorized processes with on-shell projections will be sufficient.

4.2 Higgs-boson decay

The current intrinsic and parametric uncertainties for the various Higgs-boson decay widths are given in Tab. 2 (see [4]). The status of the intrinsic uncertainties was evaluated following [32–35].

Also the parametric uncertainties can play a non-negligible role for the evaluation of the partial widths. The most important parameters are the bottom quark mass and the strong coupling constant. In Ref. [33] the current uncertainties of α_s and m_b have been assumed to be $\delta\alpha_s = 0.0015$ and $\delta m_b = 0.03$ GeV. Additionally, $\delta m_c =$

Partial width	intr. QCD	intr. electroweak	total	para. m_q	para. α_s	para. M_H
$H \rightarrow b\bar{b}$	$\sim 0.2\%$	$< 0.3\%$	$< 0.4\%$	1.4%	0.4%	–
$H \rightarrow c\bar{c}$	$\sim 0.2\%$	$< 0.3\%$	$< 0.4\%$	4.0%	0.4%	–
$H \rightarrow \tau^+\tau^-$	–	$< 0.3\%$	$< 0.3\%$	–	–	–
$H \rightarrow \mu^+\mu^-$	–	$< 0.3\%$	$< 0.3\%$	–	–	–
$H \rightarrow gg$	$\sim 3\%$	$\sim 1\%$	$\sim 3.2\%$	$< 0.2\%$	3.7%	–
$H \rightarrow \gamma\gamma$	$< 0.1\%$	$< 1\%$	$< 1\%$	$< 0.2\%$	–	–
$H \rightarrow Z\gamma$	$\lesssim 0.1\%$	$\sim 5\%$	$\sim 5\%$	–	–	2.1%
$H \rightarrow WW \rightarrow 4f$	$< 0.5\%$	$< 0.3\%$	$\sim 0.5\%$	–	–	2.6%
$H \rightarrow ZZ \rightarrow 4f$	$< 0.5\%$	$< 0.3\%$	$\sim 0.5\%$	–	–	3.0%

Table 2. Current intrinsic and parametric uncertainties in the various Higgs-boson decay width calculations, see text and Refs. [32–35]. “–” indicates a negligible source of uncertainty.

decay	projected intr.	para. m_q	para. α_s	para. M_H	FCC-ee prec. on g_{HXX}^2
$H \rightarrow b\bar{b}$	$\sim 0.2\%$	0.6%	$< 0.1\%$	–	$\sim 0.8\%$
$H \rightarrow c\bar{c}$	$\sim 0.2\%$	$\sim 1\%$	$< 0.1\%$	–	$\sim 1.4\%$
$H \rightarrow \tau^+\tau^-$	$< 0.1\%$	–	–	–	$\sim 1.1\%$
$H \rightarrow \mu^+\mu^-$	$< 0.1\%$	–	–	–	$\sim 12\%$
$H \rightarrow gg$	$\sim 1\%$	–	0.5% (0.3%)	–	$\sim 1.6\%$
$H \rightarrow \gamma\gamma$	$< 1\%$	–	–	–	$\sim 3.0\%$
$H \rightarrow Z\gamma$	$\sim 1\%$	–	–	$\sim 0.1\%$	–
$H \rightarrow WW$	$\lesssim 0.3\%$	–	–	$\sim 0.1\%$	$\sim 0.4\%$
$H \rightarrow ZZ$	$\lesssim 0.3\%^\dagger$	–	–	$\sim 0.1\%$	$\sim 0.3\%$
Γ_{tot}	$\sim 0.3\%$	$\sim 0.4\%$	$< 0.1\%$	$< 0.1\%$	$\sim 1\%$

[†] From $e^+e^- \rightarrow HZ$ production

Table 3. Projected intrinsic and parametric uncertainties for the partial and total Higgs-boson decay width predictions (see text). The last column shows the target of FCC-ee precisions on the respective coupling squared.

0.025 GeV, $\delta m_t = 0.85$ GeV and $\delta M_H = 0.24$ GeV [36] have been taken into account. The effect on the various partial widths has been evaluated as in Ref. [32] and is shown in the three right columns of Tab. 2.

When comparing the combined intrinsic and parametric uncertainties with the target precision of FCC-ee [1–3], see the last column in Tab. 3, it is clear that improvements are necessary. Concerning the intrinsic theory uncertainty, the available predictions for the $f\bar{f}$ and $\gamma\gamma$ channels are already sufficiently precise to match the expected FCC-ee experimental uncertainty. With available calculational techniques, the evaluation of complete two-loop corrections to $H \rightarrow f\bar{f}$ can be achieved. This would reduce the uncertainty of the electroweak contributions to less than 0.1%. Similarly, the complete NLO corrections to $H \rightarrow Z\gamma$ can be carried out with existing methods, resulting in an estimated precision of about 1%. More theoretical work is needed for $H \rightarrow WW, ZZ, gg$, which are currently limited by QCD uncertainties. For $H \rightarrow gg$ the calculation of massless four-loop QCD diagrams, will be required, which may be within reach [37], reducing the intrinsic uncertainties to the level of about 1%. For $H \rightarrow WW, ZZ$, the required QCD corrections are essentially identical to those for $e^+e^- \rightarrow WW$, and it is straightforward to improve them to a practically negligible level. Further significant progress would require the calculation of two-loop electroweak corrections, which for a $1 \rightarrow 4$ process is beyond reach for the foreseeable future. Here it should be noted, however, that the HZZ coupling will be mostly constrained by the measurement of the $e^+e^- \rightarrow HZ$ production process at FCC-ee with $\sqrt{s} = 240$ GeV, rather than the decay $H \rightarrow ZZ^*$, see the discussion in section 4.1, leading to a remaining intrinsic uncertainty of less than 0.3%.

Also shown in Tab. 3 are the projected parametric uncertainties, assuming FCC-ee precisions. For inputs, we use $\delta\alpha_s = 0.0002$ and $\delta m_t = 50$ MeV³, $\delta M_H \sim 10$ MeV, and $\delta m_b \sim 13$ MeV and $\delta m_c \sim 7$ MeV, see [4] for details.⁴ It becomes clear that the intrinsic and parametric uncertainties will be able to match the high anticipated experimental accuracy at FCC-ee.

³ This assumes a precision from an e^+e^- top threshold scan. At the highest FCC-ee energy, 365 GeV, the top Yukawa coupling can also be inferred from the top threshold dependence on electroweak loop corrections.

⁴ Note that the numbers in Tab. 3 do not take into account correlations between the uncertainties in the Higgs production and decay processes or between different decay processes, in particular entering via Γ_{tot} . Their impact can only be evaluated when the full experimental correlation matrix is known, see the discussion in [4].

5 Monte-Carlo and exclusive predictions

To determine systematic uncertainties, to properly extract parameters from the Higgs measurements and to find deviations from SM predictions, it is indispensable to have Monte Carlo (MC) simulations. The challenge compared to LEP/SLC times is the strongly increased experimental precision and much higher luminosities which demands much better simulations. The precise description of QED effects, both regarding the correct normalization of cross sections as well as exclusive multi-photon radiation will be one of the highest priorities. In the recent years, there has been steady progress in dealing with the calculation of higher logarithmic and finite orders in the QED radiation, enhanced by the very small electron mass, see Refs. [38–42]. These very advanced calculations are only partially applicable for the exclusive description of photon radiation which is necessary if cuts are to be applied on those photons. At the level of LO matrix elements, there is a matching procedure between explicit matrix element photons and the electron PDFs in collinear factorization in the context of WHIZARD (cf. below) which has been recently shown to work unexpectedly well [43, 44].

From the LEP era there are many dedicated MC programs which feature soft photon resummation in the exclusive form. Most of them YFS2/KORALZ [45], BHLUMI [46], BHWIDE [47], YFSWW [48] are based directly on the Yennie-Frautschi-Suura (YSF) resummation scheme [49]. A more powerful coherent exclusive exponentiation (CEEX) at the amplitude level was introduced in [20] and is so far implemented only in KKMC [50] program for two-fermion processes $e^+e^- \rightarrow f\bar{f}$, $f = \mu, \tau, \nu, q$. A generic semi-automatic implementation of CEEX for a wider class of processes would be a very desirable development, and might pave the path towards a consistent inclusion of the non-soft higher order contributions over the complete photon phase space, merging together genuine electroweak corrections with QED correction calculated to much higher orders, inclusion of spin correlations, initial-final state interferences and more. The ultimate goal is to match non-soft QED corrections up to forth order and genuine SM electroweak corrections up to second order.

Multi-purpose event generators allow to simulate the full spectrum of SM (and BSM) processes at FCC-ee. They provide (for the moment) a lower level of precision for the two-fermion processes, but allow to simulate them in the same framework as four-, six- and event eight-fermion processes. Due to the much higher precision and better detectors than at SLC/LEP these processes have to be included already at the 250 GeV stage, and not only if triple electroweak resonances or top pairs are kinematically accessible. In this context a lot of experience has been gained from the large full SM event samples generated for ILC, CLIC and CEPC using WHIZARD [51, 52]. This supports leading-logarithmic QED collinear factorization for QED photons, automated NLO QCD corrections for arbitrary processes, while NLO electroweak corrections are under way. For experimental studies, crucial observables are the correct numbers of neutral and charged hadrons. These can only be correctly simulated if the color-flow assignments are correctly made, which includes that the invariant masses of the shower systems from the most prominent processes, $e^+e^- \rightarrow W^+W^-, ZZ$ are preserved; in a full electroweak matrix element they are, however, quantum mechanically entangled. The MC hence needs to determine the correct probability of the resonant subprocesses and hand them over to parton shower and hadronization [53, 54]. It is important to note that the solution of the long-standing problem of the soft photon resummation to narrow width semi-stable charged resonances like W^\pm was recently outlined in Ref. [55]. It is also very important to take photon-induced backgrounds from Weizsäcker-Williams QED splittings into account, particularly $\gamma\gamma \rightarrow$ hadrons. WHIZARD has a special treatment for the top threshold run [56] that is used by the FCC-ee collaboration. It also supports all different kinds of beam spectra, from simple parameterized files, over Gaussian beam spreads to a sophisticated beamstrahlung setup for linear and circular lepton machines [57].

To search for deviations in the Higgs sector from the SM, WHIZARD supports a large setup of BSM models like supersymmetry, extended Higgs models like 2HDM, Higgs singlet extensions, SMEFT, composite and Little Higgs models, which can be generically extended to its interface to automated Lagrangian model tools, cf. e.g. [58, 59]

6 Outlook and deliverables

The most important steps to be taken in the future to match the high anticipated accuracy of FCC-ee in the context of Higgs physics and EWPO with SM theory predictions are the following.

- Improved unfolding techniques to go from observables to pseudo-observables.
- Calculations for the EWPO: complete $\mathcal{O}(\alpha\alpha_s^2)$ corrections, fermionic $\mathcal{O}(\alpha^2\alpha_s)$ corrections, double-fermionic $\mathcal{O}(\alpha^3)$ corrections, and leading four-loop corrections enhanced by the top Yukawa coupling.
- To extract the W -boson mass: complete 2-loop corrections to $e^+e^- \rightarrow W^+W^-$ and to $W \rightarrow f\bar{f}'$ (based on the effective field theory framework).
- Development on matching scheme between EW corrections and radiated/resummed photons.
- Full two-loop corrections for $e^+e^- \rightarrow ZH$ and $e^+e^- \rightarrow VVH$
- Full two-loop corrections for $H \rightarrow f\bar{f}(+\gamma)$.
- (Mostly) automatized MC generation including NLO QCD and EW corrections.

- Improvements by roughly one order of magnitude in the determination of the most important SM parameters (m_t , m_b , α_s , $\Delta\alpha_{\text{had}}$, M_H , ...)
- Complete simulations of relevant SM processes at different energy stages in full detector simulation

Finally, the following should be kept in mind. The SM constitutes the model in which highest theoretical precision for the predictions of EWPO and Higgs-boson observables can be obtained. This concerns their prediction as well as their extraction. However, as soon as BSM physics will be discovered, an evaluation of the predictions of the EWPO and the Higgs-boson observables in any preferred BSM model will be necessary. The corresponding theory uncertainties, both intrinsic and parametric, can then be larger (see, e.g., [10,60] for the Minimal Supersymmetric SM). A dedicated theory effort (beyond the SM) would be needed in this case.

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