EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

Proposal to the ISOLDE and Neutron Time-of-Flight Committee

Probing the largest core-breaking prediction towards 100 Sn: proton single-particle strength in $\frac{110}{5}$ Sn

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G. Zhang^{1,2}, D. Mengoni^{1,2}, S. Bottoni^{3,4}, A. Gottardo⁵, L. Gaffney⁶, M. Assié⁷,

D. Beaumel⁷, G. Benzoni⁴, Y. Blumenfeld⁷, A. Bracco^{3,4}, D. Brugnara^{1,5}, P. Butler⁶,

F. Camera^{3,4}, S. Capra^{3,4}, J. Casal^{1,2}, F.C.L. Crespi^{3,4}, G. de Angelis⁵, F. Flavigny⁸,

L. Fortunato^{1,2}, S. Freeman⁹, F. Galtarossa⁷, A. Goasduff⁵, F. Gramegna⁵, J. Ha^{1,2},

M. Labiche¹⁰, J.A. Lay Valera^{11,12}, I. Lazarus¹⁰, S.M. Lenzi^{1,2}, S. Leoni^{3,4}, T. Marchi⁵,

A. Matta ⁸, M. Mazzocco^{1,2}, R. Menegazzo², G. Montagnoli^{1,2}, D.R. Napoli⁵,

B. Olaizola¹³, R. Page⁶, G. Pasqualato^{1,2}, R. Pérez⁵, A. Pullia^{3,4}, R. Raabe¹⁴

F. Recchia^{1,2}, D.K. Sharp⁹, F. Soramel^{1,2}, J.J. Valiente Dobón⁵, I. Zanon⁵

- ¹ Dipartimento di Fisica e Astronomia, Università di Padova, Italy.
- 2 INFN, Sezione di Padova, Padova, Italy.
- 3 Dipartimento di Fisica, Università degli Studi di Milano, Italy.
- 4 INFN Sezione di Milano, Italy.
- 5 INFN, Laboratori Nazionali di Legnaro, Legnaro (Padova), Italy.
- ⁶ University of Liverpool, UK.
- 7 IJCLab, Orsay, France.
- ⁸ LPC Caen, France.
- ⁹ University of Manchester, UK.
- ¹⁰ STFC Daresbury, UK.
- ¹¹ Departamento de FAMN, Universidad de Sevilla, Spain.
- ¹² Instituto Interuniversitario Carlos I de Física Teórica y Computacional (iC1), Seville, Spain.
- ¹³ CERN, Geneva, Switzerland.
- ¹⁴ KU Leuven, Belgium.

Spokespersons: Guangxin Zhang [guangxin.zhang@pd.infn.it], Daniele Mengoni [daniele.mengoni@pd.infn.it] Contact person: Liam Gaffney [liam.gaffney@cern.ch]

Abstract:

The Sn isotopic chain is a formidable testing ground for the study of shell evolution in different nuclear models. Even though the energies of the first 2^+ states in eveneven isotopes between ¹⁰²Sn and ¹³⁰Sn are quite similar, in agreement with the seniority scheme for this valence neutron space, the significant increment of the $B(E2; 0^+ \rightarrow 2^+)$ value approaching ¹⁰⁰Sn has remained a major puzzle over decades. Recently, the result of state-of-the-art Monte Carlo shell model (MCSM) calculations on Sn isotopes, performed in a large model space including single-particle orbits below and above the $Z=50$ and $N=82$ magic numbers, has explained the anomalous trend of $B(E2; 0^+ \rightarrow 2^+)$ by considering a core breaking contribution, involving the $g_{9/2}$ proton orbits to the gds valence space. The MCSM calculations also suggested that the number of proton holes for the first 0^+ , 2^+ and 4^+ states in even-even Sn isotopes will increase when approaching $N = 60$, i.e. ¹¹⁰Sn, being maximum for the $4⁺₁$ state. In this proposal, we aim at measuring the proton singleparticle strength, by probing directly the $\Delta\ell=2$ contribution in the proton transfer to the $4₁⁺$ state in ¹¹⁰Sn. This will be achieved using the ¹⁰⁹In(³He,d)¹¹⁰Sn one-proton transfer reaction, with the light ejectiles being measured using the solenoidal spectrometer at ISOLDE.

Requested shifts: 21 shifts (7 days, 1 run over 1 year)

1 Motivations

The self-conjugated and doubly-magic ¹⁰⁰Sn nucleus has been regarded as one of the milestones in present and future nuclear structure research. Consequently, a great experimental effort has been devoted to ¹⁰⁰Sn and the nearby nuclei, employing different techniques. Up to now, only the β -decay spectroscopy of 100Sn has been experimentally studied, providing indication about a doubly-magic nature [1, 2]. Yet, single-particle energies are poorly known, due to the difficulty of studying ¹⁰¹Sn, given also the controversy regarding the spin-parity assignments of ground and first excited states [3].

The whole tin isotopic chain, from 10^2 Sn to 130 Sn (neutron number 52–80), roughly exhibits the typical behavior predicted by the generalized seniority scheme for neutrons in the $g_{\frac{7}{2}}$ and $h_{\frac{11}{2}}$ orbitals, for N < 64 and >64, respectively. See Ref. [4] for details.

The energies of the first 2^+ states in even-even isotopes between 10^2 Sn and 130 Sn are quite similar, in agreement with the prediction of the seniority scheme for this valence neutron space.

However, the experimental $B(E2; 0^+ \rightarrow 2^+)$ values suggested a *presumed* deviation from the expected parabolic behavior, as shown in Fig. 1 (a) (taken from Ref. [5]). From a theoretical point of view, various attempts have been made to explain the experimental results, in particular by including core-breaking excitations in the shell-model calculations [6, 7]. The predicted flat trend is not firmly established experimentally, because limited data are available on the extremely neutron-deficient nuclei beyond ¹⁰⁴Sn. Recently, Togashi et al. [5] claimed that Monte Carlo shell-model (MCSM) calculations have explained the $B(E2; 0^+ \rightarrow 2^+)$ trend (see the red curve in Fig. 1 and Ref. [5]) by promot-

Figure 1: (a) Calculated and measured $B(E2)$ values and (b) $E2$ matrix elements as well as their decomposition in proton and neutron contribution ratios $(\%)$. The figure is taken from Ref. [5].

ing protons and neutrons from the $g_{\frac{9}{2}}$ orbital to those at higher energy. Such state-of-art theoretical prediction opens a new perspective for the study of core-breaking effects in Sn isotopes especially for the neutron-deficient ones. In the bottom panel of Fig. 1, proton and neutron contributions to the E2 matrix elements are shown. The proton contribution becomes larger approaching the most neutron-deficient ¹⁰⁰Sn: the proton/neutron ratio of the calculated matrix element exceeds 30% for N = 50-64, with the $1g_{9/2} \rightarrow 2d_{5/2}$ excitation being dominant, and the proton contribution reaches a maximum at N=60 in $110Sn$.

Further to the $B(E2; 0^+ \rightarrow 2^+)$ measurement, which are still out of reach for the most exotic neutron-deficient Sn isotopes, a powerful alternative way to investigate the robustness of the 100 Sn core is to directly probe the proton single-particle occupation above the $Z=50$ shell closure. For the low-lying states, this would imply a partial occupation of the $g_{\frac{7}{4}}$, $d_{\frac{5}{2}}, d_{\frac{3}{2}}, s_{\frac{1}{2}}$ proton valence orbitals., if core-breaking effects are present. In particular from R_{eff}^{2} , $\frac{2}{10}$, the main core-breaking contribution is expected from proton excitation to the $d_{\frac{5}{2}}$ orbital which is predicted to reach a maximum at N=60, corresponding to ^{110}Sn , see F ig. 2. Such a contribution is also expected to become larger with increasing spin, making the 4^+ state in ¹¹⁰Sn the best candidate for testing the possible $\ell=2$ proton component of its wave function. In addition, to the best of our knowledge, this will represent the first time a neutron deficient Sn isotope is probed against the proton degree of freedom via direct reactions.

This will be possible by measuring the 109 In(3 He,d) 110 Sn single-proton transfer reaction using the ISOLDE Solenoidal Spetrometer, which will allow us to populate the 4^+ state at 2.197 MeV in ¹¹⁰Sn. Deuterons, which will be detected at backward angles, will show

Figure 2: Occupation numbers of (a) proton and (b) neutron holes in the $1g_{9/2}$ orbit. For each value of N, histograms for the 0^+ , 2^+ , and 4^+ states are shown from left to right. The maximum is expected in correspondence of the 4^+ state in 110 Sn. The figure is taken from Ref. [5].

a characteristic angular distribution depending on the $\Delta\ell$ transfer, providing information on the occupancy of proton orbitals above the Z=50 shell closure. This will shed light on possible proton excitations across the shell gap, as a consequence of core-breaking effects in ¹¹⁰Sn. In summary, we aim at:

- 1. getting insights into the possible core-breaking effects in $\frac{110}{5}$ by measuring absolute integral cross section for the proton-adding $(^{3}He,^{2}H)$ reaction, predicting to be mainly from a $\Delta \ell = 2$ transfer.
- 2. establishing the individual contribution of each proton component by measuring differential cross sections and extracting spectroscopic factors.

2 Experimental details and rate estimation

The 109 In(3 He,d)¹¹⁰Sn single-proton transfer reaction is proposed to investigate the core breaking effects, which manifests via the occupation of the valence orbitals above $Z=50$. The deuterons emitted at backward angles will be measured along the Z axis by the ISOLDE Solenoidal Spectrometer (ISS) silicon detectors.

The 109 In_{gs} isotope is produced from a UC_X target irradiated with the proton beam from the PS Booster. The expected production yield is $5x10⁹$ pps, as recently measured at the COLLAPS experiment [8]. Considering transmission efficiency and post acceleration, we expect a beam intensity on secondary target of at least $\sim 1x10^7$ pps. Possible 10÷20% contamination, presumably from ¹⁰⁹Pd, might also be present. For such $^{109}_{46}Pd, \ell=4$ would be the most probable ℓ transfer, being the Z=50 core open, but this will be largely suppressed because of the high ℓ value. As a consequence, no significant effect to the aim of the experiment is expected from Pd contaminants, considering also that the estimated statistical error on the final cross section is expected to be around 20%. Please see further below for details.

Figure 3: Realistic MC simulation of ISS energy vs Z axis position for the deuteron ejectiles. According with the calculated cross section, the 4^+ state is the most abundantly populated, followed by the 2^+ state, whilst the ground state is barely visible. Dashed lines suggest the centroid of each state: the 4^+ state could be unambiguously selected.

A beam energy of 7.38 MeV/u from HIE-ISOLDE is suited to enhance the probability of single-particle proton transfer. The transfer will populate low-lying states, in particular the 4⁺ state, and will enable the determination of the total and differential cross section, which is predicted to be dominated by the $\Delta \ell=2$ transfer.

Cross sections were calculated using a one-step DWBA approach for the $0^+, 2^+, 4^+$ states [9]. The optical parameters of the $109\text{In}+3\text{He}$ incoming and the $110\text{Sn}+d$ outgoing scattering potentials were taken from Ref. [10] and Ref. [11], respectively. The ³He wave function was obtained with a Woods Saxon potential and fitted to the $\langle ^3He|d+p\rangle$ overlap calculated with VMC techniques, using the the AV18+UX interaction (see Ref. [12]). The parameters of the potential were adjusted to match the ANC and the proton binding energy. The calculated integral cross section is 0.75 mb and 0.25 mb, for the $4⁺₁$ and $2⁺₁$ state, respectively. The total cross section calculated for the transfer to the ground state is only 0.01 mb. Unitary spectroscopic factors have been considered in all the calculations. For the secondary target, we plan to use a solid ³He target on a W substrate, developed within our collaboration, which accounts for a density of $\sim 0.3 \times 10^{18}$ atoms/cm². Substrate thickness, $0.6 \mu m$, is necessarily limited due to the beam energy straggling, and so the maximum density of the sputtered ³He. Please see Ref. [13] for any further detail.

We performed a commissioning run of a similar target prototype with an intense Zn stable beam delivered by the LNL-INFN Tandem, at a variable intensity ranging from $(0.3 \div 5) \times 10^{10}$ pps for 22 hrs. After 11 hours, during which the beam intensity was increased, the density of ³He decreased by a factor of ∼3. In conclusion, we do not expect any significant loss with an intensity of 10^7 as expected for 10^9 In beam.

At the chosen beam energy, 7.38 MeV/u, the centre of mass energy is 50% above the Coulomb barrier for W. Nevertheless, the expected reaction mechanism will be (quasi) fusion-fission process, which hinders any light-charged particle background.

Figure 4: Realistic MC simulations of the differential cross section for the 4^+ state. A deconvolution of each $\Delta\ell$ transfer component is performed, revealing the main contribution of $\Delta \ell = 2$, as a given input, in agreement with theoretical predictions.

Fig. 3 shows a realistic MC simulation where the first excited 4^+ and 2^+ states are the most abundantly populated states, whilst the ground state in ¹¹⁰Sn is barely visible.

Two other states at 2058.0(4) and 2121.04(23)-keV energy are quoted in literature. However, the first one is only reported in a FE reaction in Ref. [14] as $(0^+, 2)$ but neither spectra are shown nor the character of the state discussed; the second one is reported in Ref. [15] to be a 2^+_2 marginally produced in the ϵ capture from ¹¹⁰Sb and in Ref. [16] in a neutron transfer reaction, but the tiny cross section couldn't be quoted in the paper. Therefore, we expect these two states, if there, to be marginally produced also in our experiment.

The separation shown in the Energy vs Z plot of ISS is sufficient to gate on the 4^+ events and extract the differential cross section, which is shown in Fig. 4. Such angular distribution for the proton transferred to the 4^+ state was simulated considering the different ℓ contributions and the estimated number of counts. Spectroscopic factor are presently under calculation, consequently in our simulation, each contribution was provisionally taken proportional to the relative s,d,g composition of 4^+ wave function [17]. The relative weights used as input are 7%, 80% and 13% for $s_{1/2}, d_{5/2}+d_{3/2}, g_{7/2}$ components, respectively.

Due to the reaction kinematics, deuterons are emitted backward in the laboratory frame with a $\Delta \ell = 2$ peaked at ∼120 degrees in the lab frame. This implies that, including also the presence of the target and Si detection system bore, an angular coverage of $3 \div 30$ degrees in the CM frame remains available for the detection, see Fig. 4. Notwithstanding this limit, the result of the simulation gives for the deconvolution of the $\Delta \ell=2$ component a value of 80%, as a given input, with a relative error of ∼20%. This error sets also the sensitivity of the measurement to the core breaking effect, giving a lower detectable limit of ∼100 µb. For the extraction of the absolute cross section we plan to use a downstream Faraday cup and a monitor detector for the ³He elastically scattered.

Summary of requested shifts: in total we request 21 shifts of beam on target, to be able to estimate a core breaking contribution with an error of 20% corresponding to a lower limit of \sim 100 µb.

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Appendix

DESCRIPTION OF THE PROPOSED EXPERIMENT

The experimental setup comprises: (name the fixed-ISOLDE installations, as well as flexible elements of the experiment)

HAZARDS GENERATED BY THE EXPERIMENT (if using fixed installation:) Hazards named in the document relevant for the fixed [COLLAPS, CRIS, ISOLTRAP, MINIBALL + only CD, MINIBALL + T-REX, NICOLE, SSP-GLM chamber, SSP-GHM chamber, or WITCH] installation.

Additional hazards:

Hazard identification:

Average electrical power requirements (excluding fixed ISOLDE-installation mentioned above): [make a rough estimate of the total power consumption of the additional equipment used in the experiment]