

Recent ATLAS measurements of correlations in Pb+Pb and Xe+Xe collisions



*Krzysztof Woźniak, IFJ PAN, Kraków, Poland
on behalf of the ATLAS Collaboration*

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ON HIGH ENERGY PHYSICS

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PRAGUE, CZECH REPUBLIC

Collisions of nuclei - creation of Quark-Gluon Plasma

Large energy released in collisions of nuclei leads to creation of strongly interacting dense matter - **Quark-Gluon Plasma**

Collective phenomena - particle flow

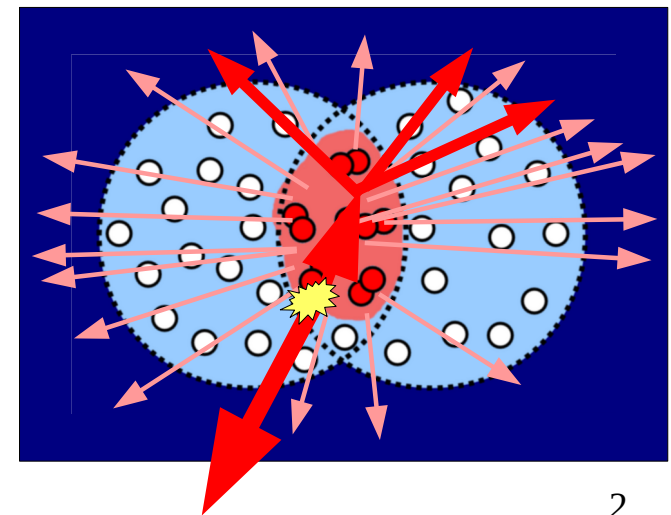
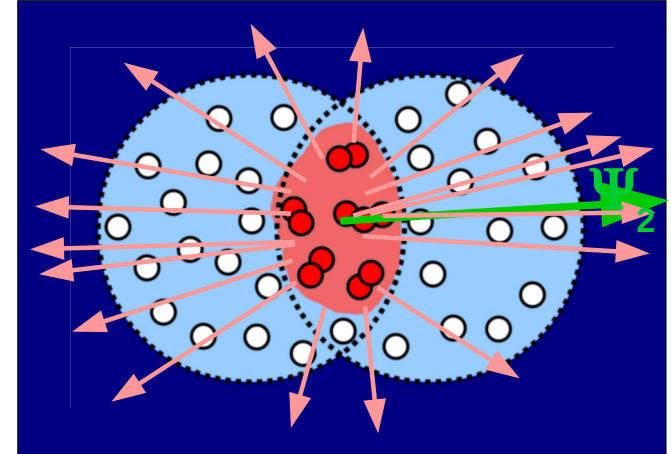
- ▶ Area of the overlap of nuclei has an elongated shape
- ▶ In this region Quark-Gluon Plasma is created
- ▶ QGP expands and its components, quark and gluons, behave collectively
- ▶ More particles are produced in the event plane Ψ

this talk and the talk by Adam Trzupek

Strong interactions of partons in QGP

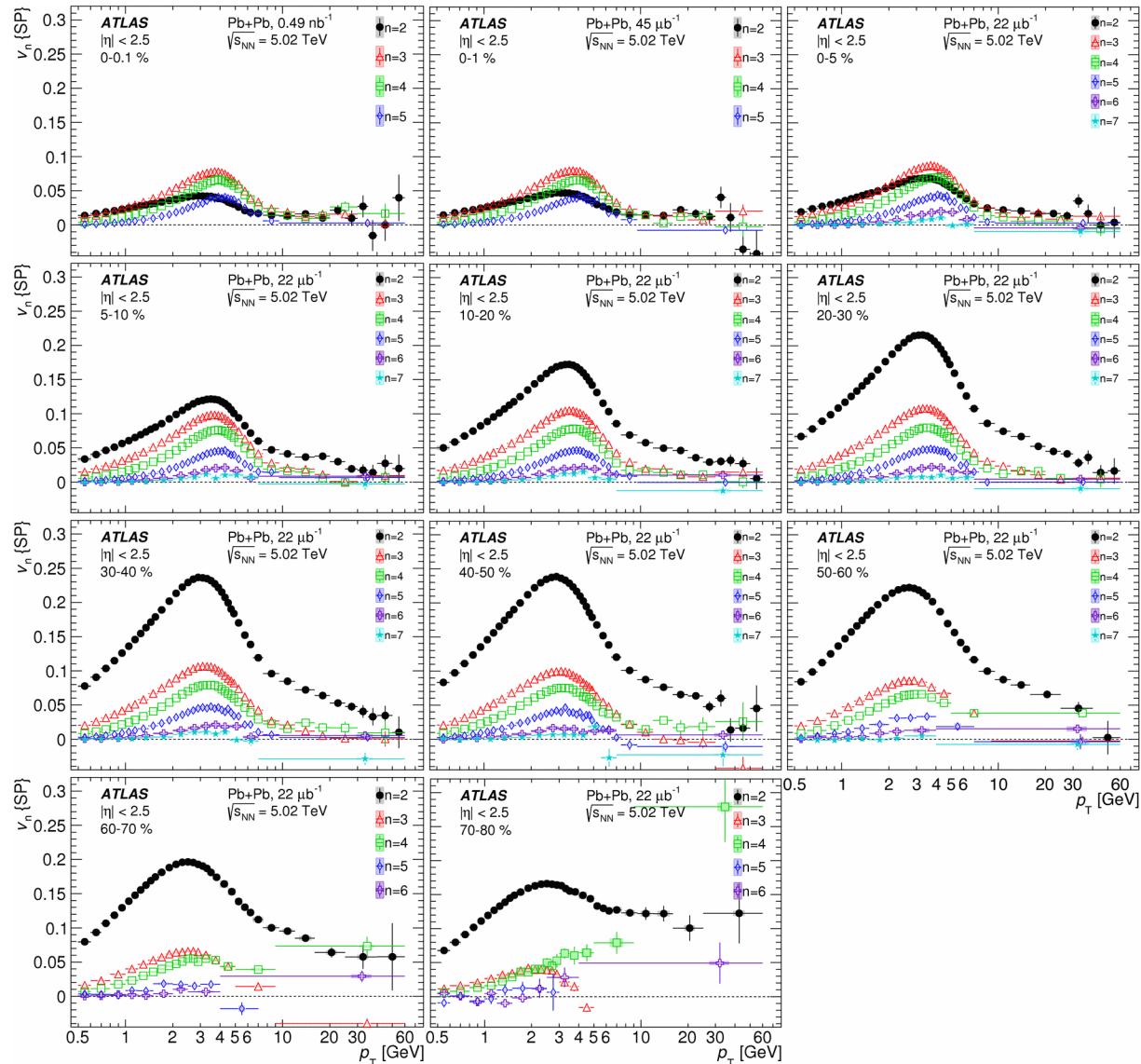
- ▶ Partons created in a hard scattering travel through QGP
- ▶ Interactions with the dense matter causes that even the most energetic partons loose energy and may be stopped inside QGP area

see the talk by Helena Santos



Detailed measurements of harmonics v_2-v_7 :

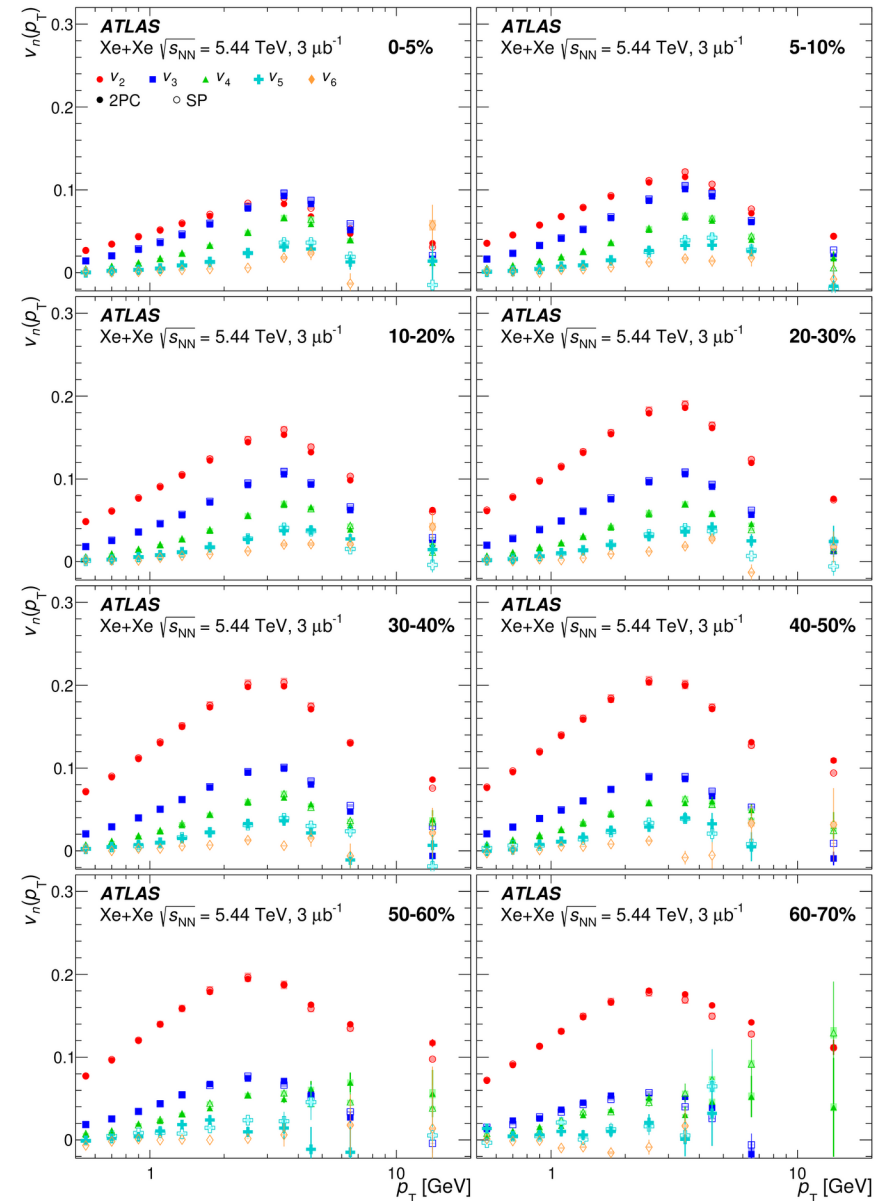
- ◆ p_T dependence
 - an increase up to ~ 3 GeV
 - a decrease above ~ 3 GeV
- ◆ centrality dependence - harmonics are the largest in mid-central (20-50%) events
- ◆ harmonics up to v_7 are non-zero
- ◆ various methods of calculations of flow harmonics



ATLAS, Eur. Phys. J. C 78 (2018) 997.

Flow measurement in Xe+Xe collisions

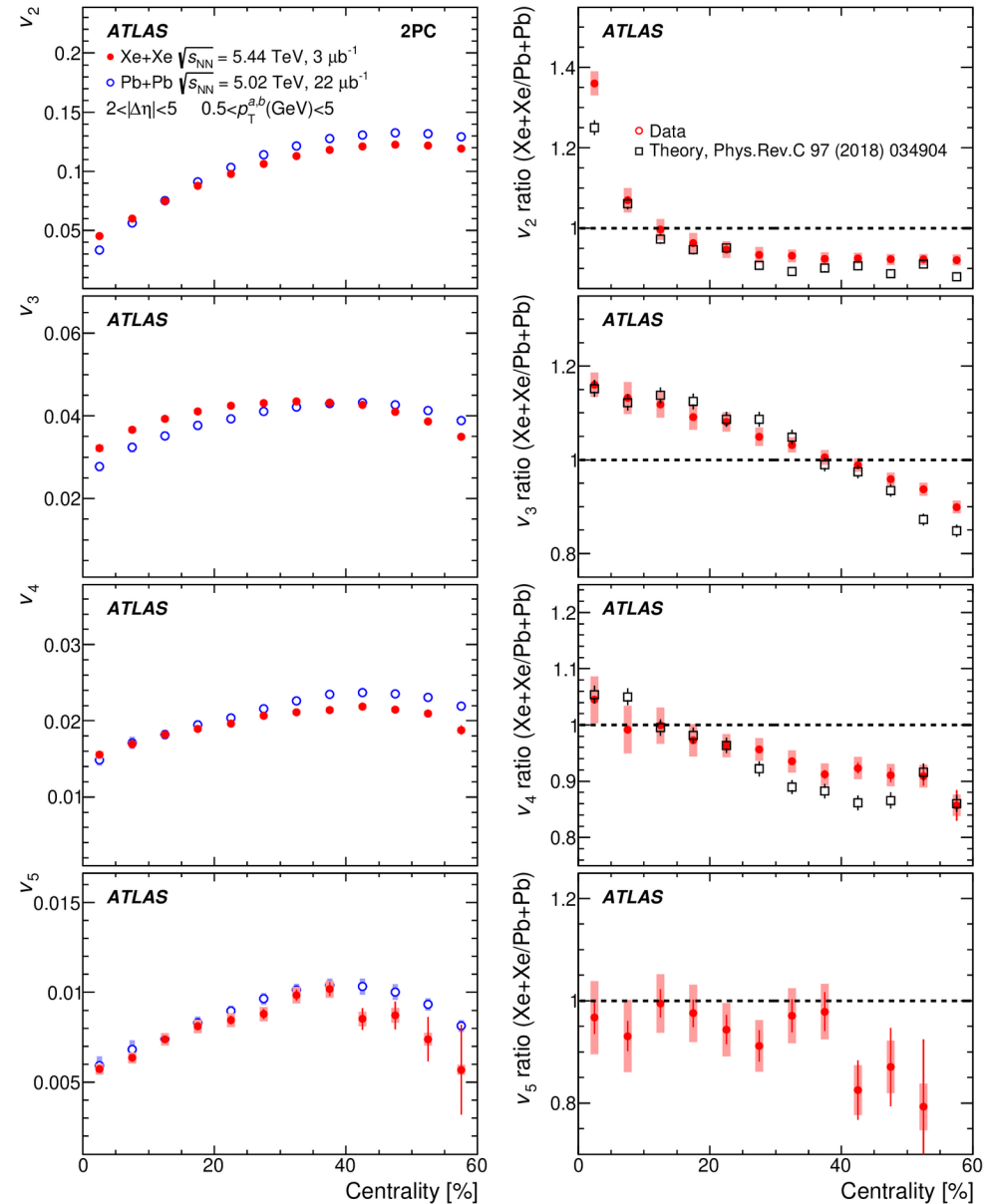
- ◆ harmonics v_2 to v_6
- ◆ at the first look p_T dependence similar to that for other collision systems



ATLAS, Phys. Rev. C 101 (2020) 024906

Comparison of flow measurement in Xe+Xe and Pb+Pb collisions

- ▶ elliptic flow larger in Xe+Xe in most central collisions and smaller in peripheral collisions
- ▶ similar, but less strong trend with centrality for v_3 and v_4
- ▶ slightly lower v_5 in Xe+Xe than in Pb+Pb collisions in all centralities

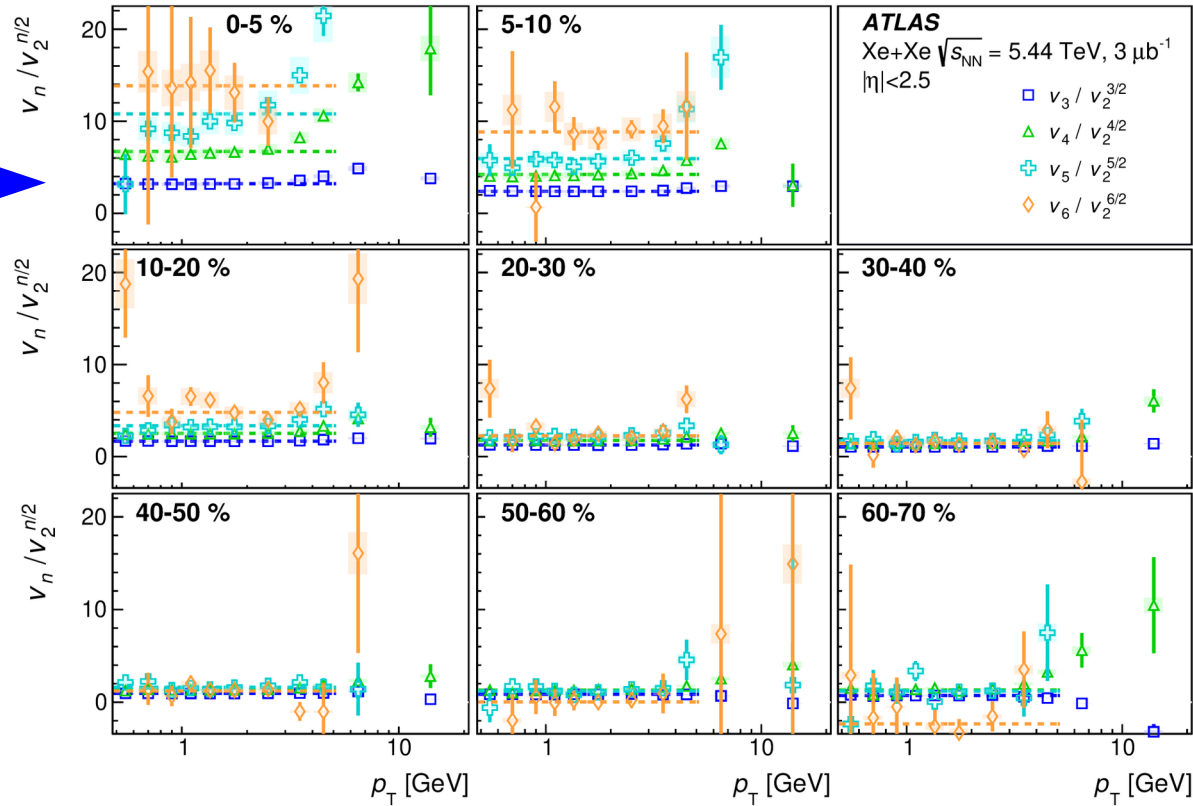


ATLAS, Phys. Rev. C 101 (2020) 024906

Scaling of flow harmonics

scaling by v_2 : $V_n / (V_2)^{n/2}$

constant as a function of p_T



ATLAS, Phys. Rev. C 101 (2020) 024906

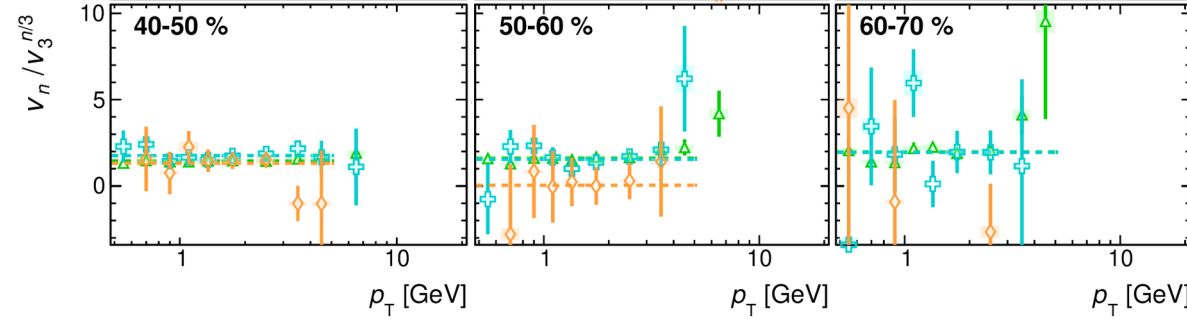
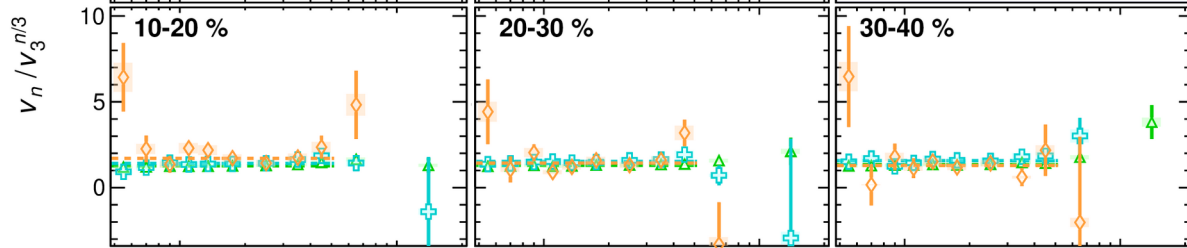
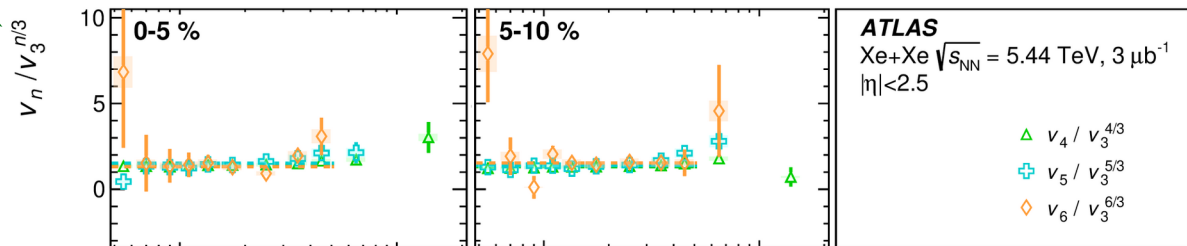
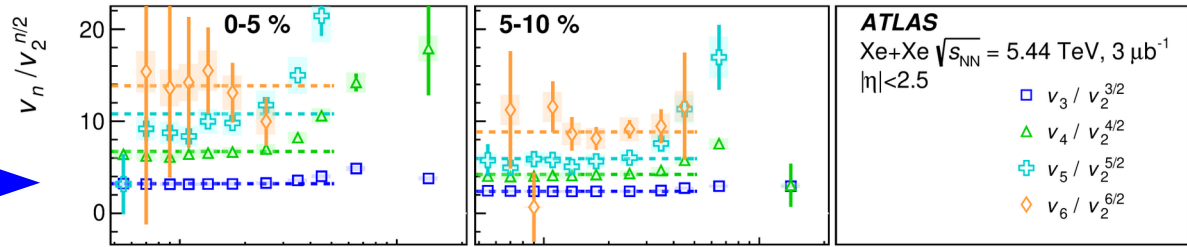
Scaling of flow harmonics

◆ scaling by v_2 $v_n / (v_2)^{n/2}$

◆ scaling by v_3 $v_n / (v_3)^{n/3}$

◆ constant as a function of p_T

◆ similar values of $v_n / (v_3)^{n/3}$ for different v_n



ATLAS
Xe+Xe $\sqrt{s_{NN}} = 5.44$ TeV, $3 \mu\text{b}^{-1}$
 $|\eta| < 2.5$

- $v_3 / v_2^{3/2}$
- △ $v_4 / v_2^{4/2}$
- + $v_5 / v_2^{5/2}$
- ◇ $v_6 / v_2^{6/2}$

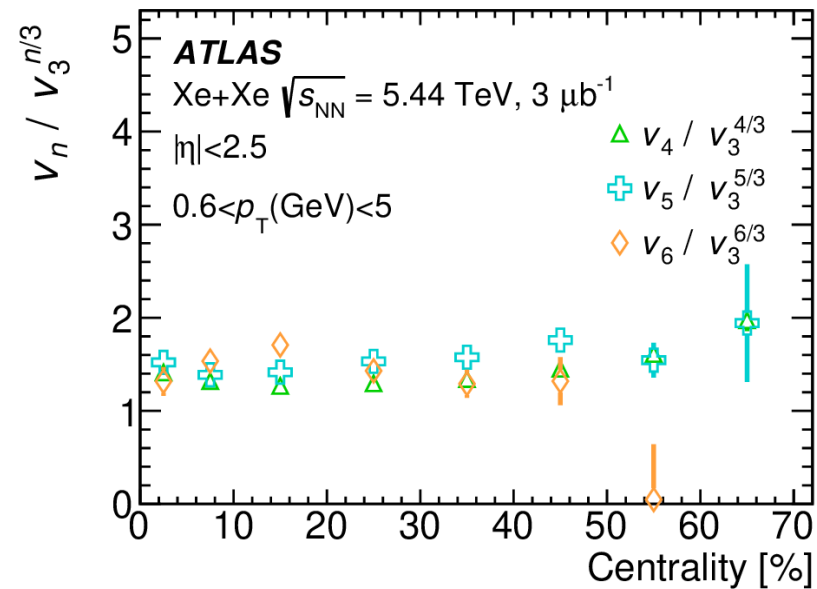
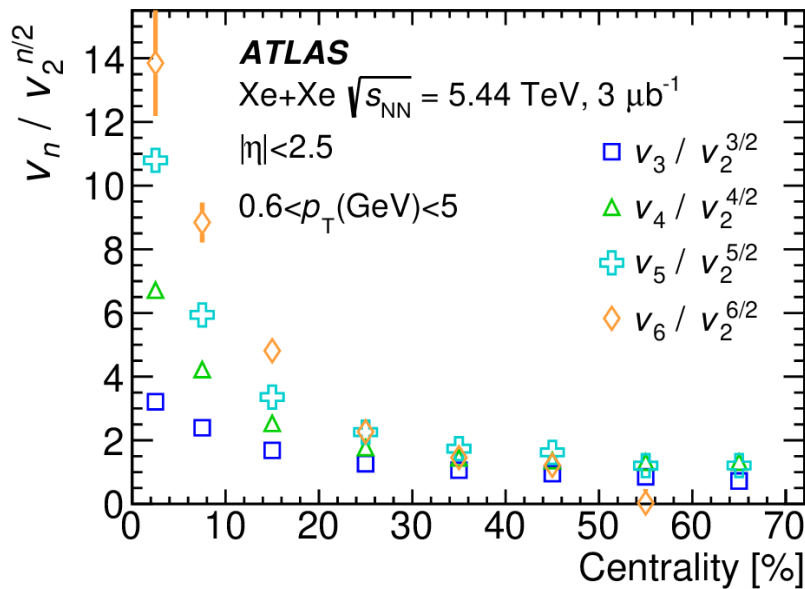
ATLAS
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- △ $v_4 / v_3^{4/3}$
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ATLAS, Phys. Rev. C 101 (2020) 024906

Scaling of flow harmonics

- ◆ scaling by v_2 $v_n / (v_2)^{n/2}$
- ◆ scaling by v_3 $v_n / (v_3)^{n/3}$
- ◆ constant as a function of p_T
- ◆ scaling by v_3 gives almost the same value for all higher harmonics

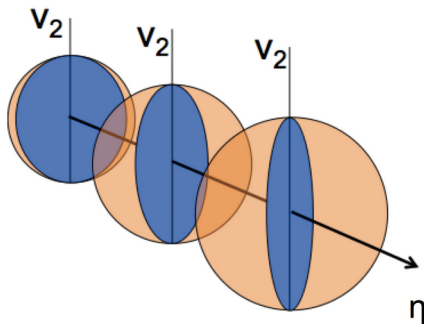


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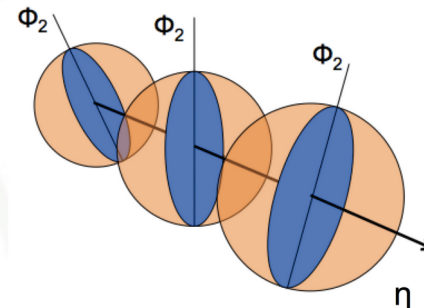
Longitudinal flow decorrelations

Measurement of flow harmonics fluctuations as a function of pseudorapidity

Change of flow magnitude



Change of event plane (twist)



Correlators using flow vectors calculated in different η intervals

$$r_{n|n;k} = \frac{\langle \mathbf{q}_n(-\eta)^k \mathbf{q}_n^{*k}(\eta_{\text{ref}}) \rangle}{\langle \mathbf{q}_n(\eta)^k \mathbf{q}_n^{*k}(\eta_{\text{ref}}) \rangle}; \quad \mathbf{q}_n = \sum_i w_i e^{in\phi_i} / \sum_i w_i$$

$$r_{n|n;k} = \frac{\langle [v_n(-\eta)v_n(\eta_{\text{ref}})]^k \cos(kn(\Phi_n(-\eta) - \Phi_n(\eta_{\text{ref}}))) \rangle}{\langle [v_n(\eta)v_n(\eta_{\text{ref}})]^k \cos(kn(\Phi_n(\eta) - \Phi_n(\eta_{\text{ref}}))) \rangle}$$

$$R_{n,n|n,n} = \frac{\langle \mathbf{q}_n(-\eta_{\text{ref}})\mathbf{q}_n(-\eta)\mathbf{q}_n^*(+\eta)\mathbf{q}_n^*(\eta_{\text{ref}}) \rangle}{\langle \mathbf{q}_n(-\eta_{\text{ref}})\mathbf{q}_n(+\eta)\mathbf{q}_n^*(-\eta)\mathbf{q}_n^*(\eta_{\text{ref}}) \rangle}$$

Note:

in the absence of any systematic modifications of flow harmonics

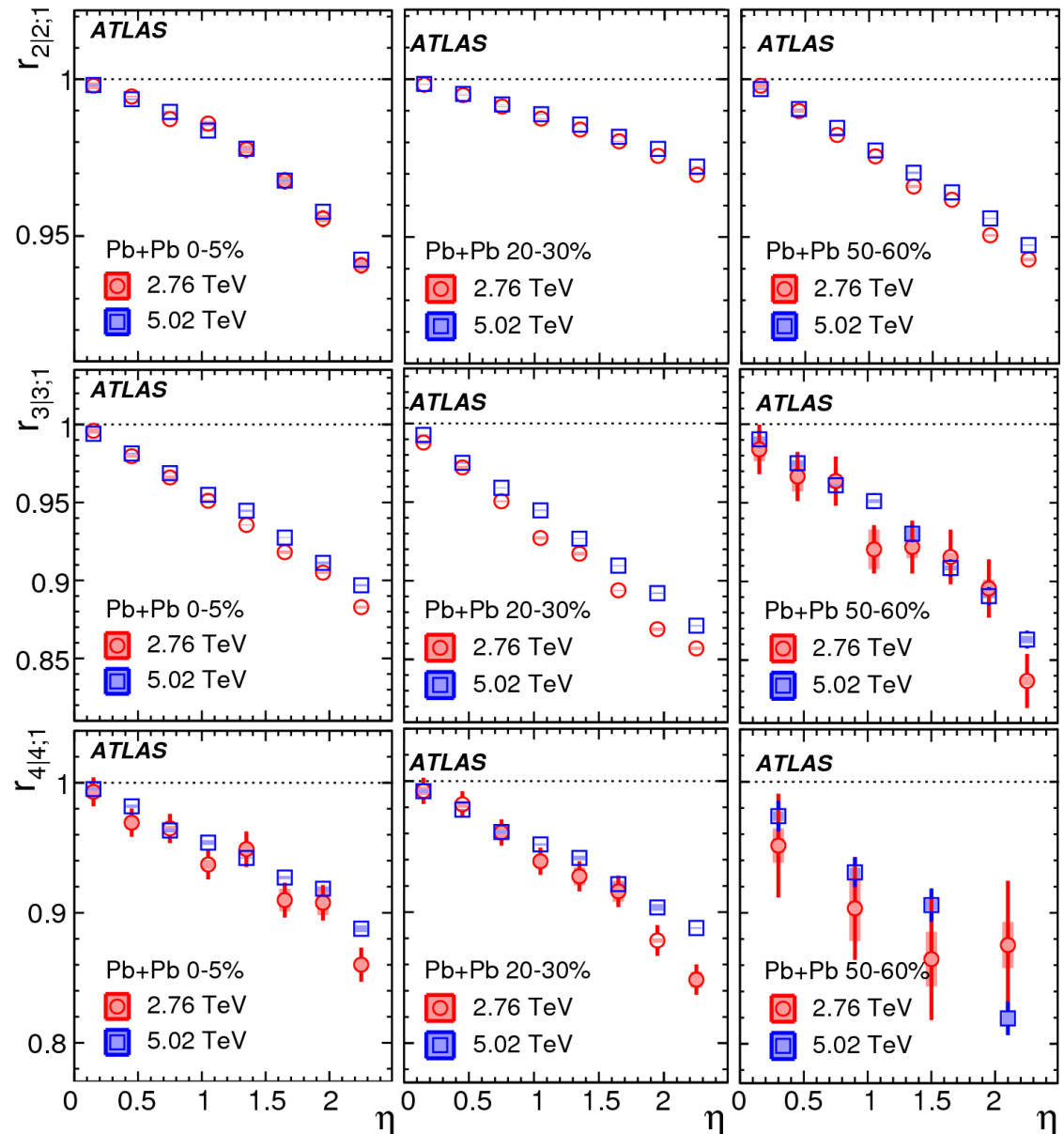
$r = 1$ and $R = 1$

Measured in Pb+Pb and Xe+Xe collisions

ATLAS, Eur. Phys. J. C 76 (2018) 142 and arxiv:2001.04201

Correlator r

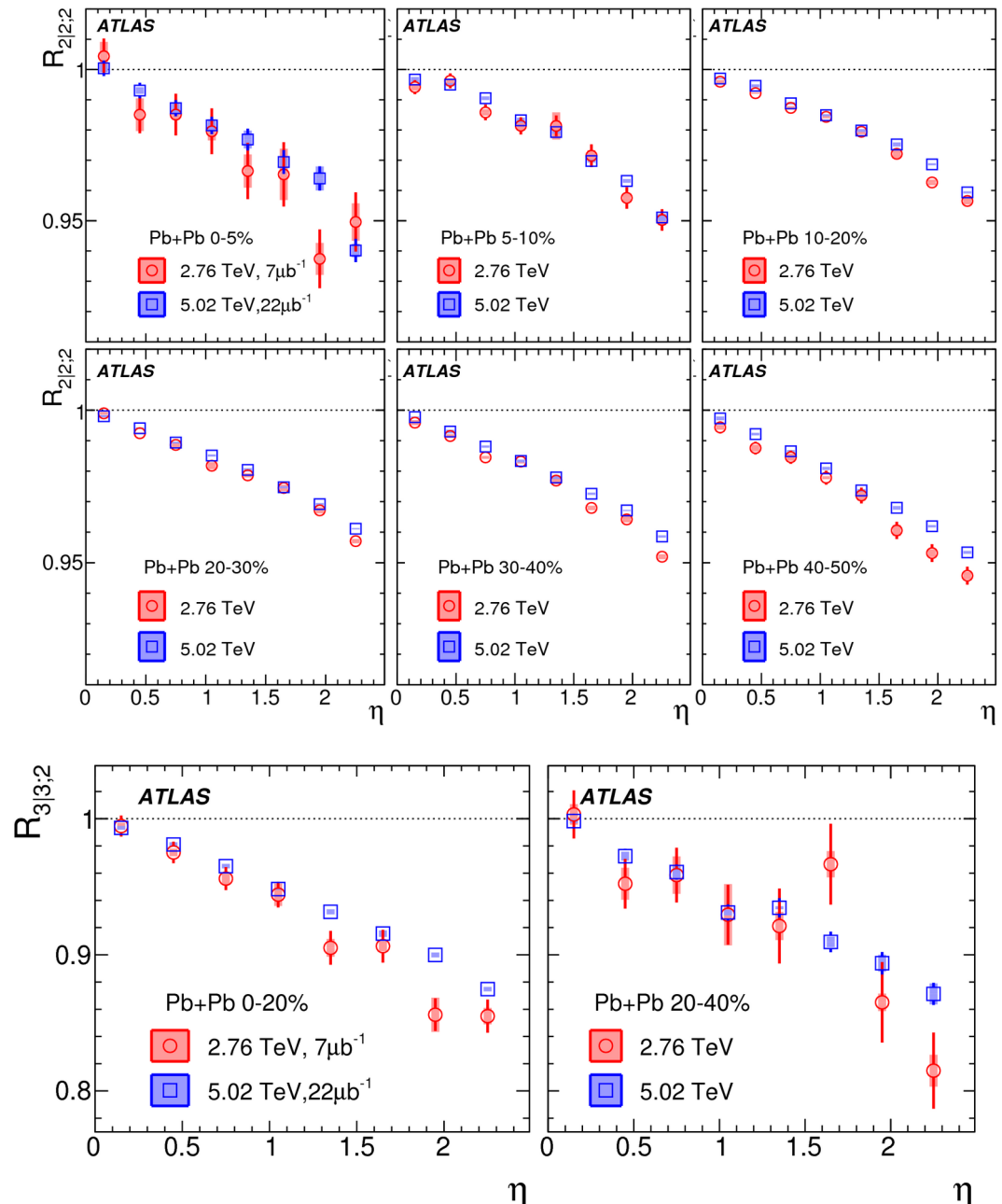
- ◆ flow decorrelations are approximately linear in η
- ◆ elliptic decorrelation $r_{2|2;1}$ depends on centrality - it deviates from 1 more in central and peripheral collisions than in semi-central collisions
- ◆ higher order correlators, $r_{3|3;1}$ and $r_{4|4;1}$, departure from 1 is even larger but similar for all centralities
- ◆ no difference for two Pb+Pb collision energies



ATLAS, Eur. Phys. J. C 76 (2018) 142

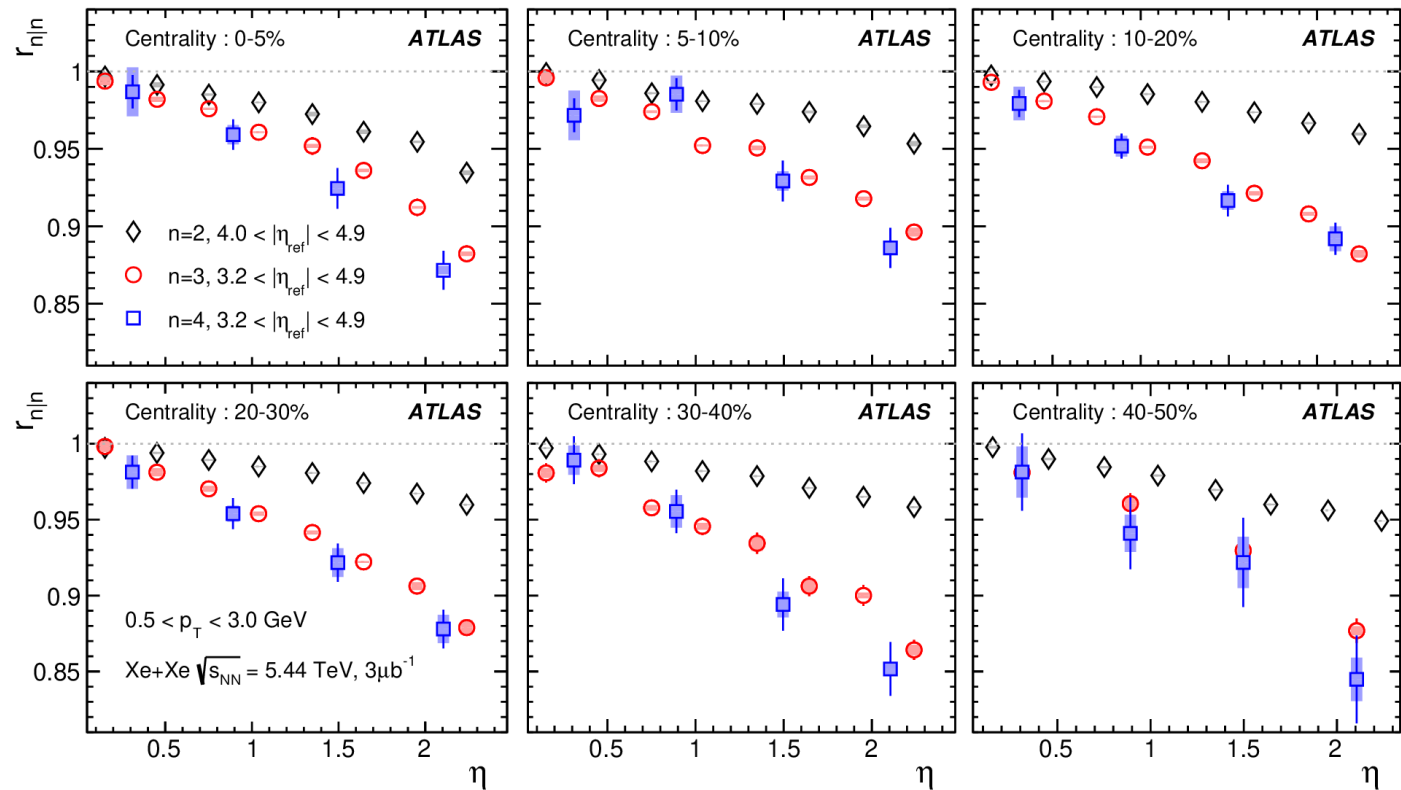
Correlator R

- ◆ flow decorrelations are approximately linear in η
- ◆ elliptic decorrelation $R_{2|2;1}$ depends on centrality - it deviates from 1 more in central and peripheral collisions than in semi-central collisions
- ◆ higher order correlator, $R_{3|3;2}$, deviates from 1 even more but in a similar way for different centralities
- ◆ no difference for two collision energies



Similar trends in Xe+Xe collisions:

- ◆ flow decorrelations are approximately linear in η
- ◆ $r_{2|2;1}$ depends on centrality - it is larger in central and peripheral collisions than in semi-central collisions
- ◆ decrease of higher order correlators, $r_{3|3;1}$ and $r_{4|4;1}$, are larger then for $r_{2|2;1}$



ATLAS, arxiv:2001.04201

- flow decorrelations are approximately linear in η

$$r_{n|n;k} = 1 - 2 F_{n;k}^r \eta$$

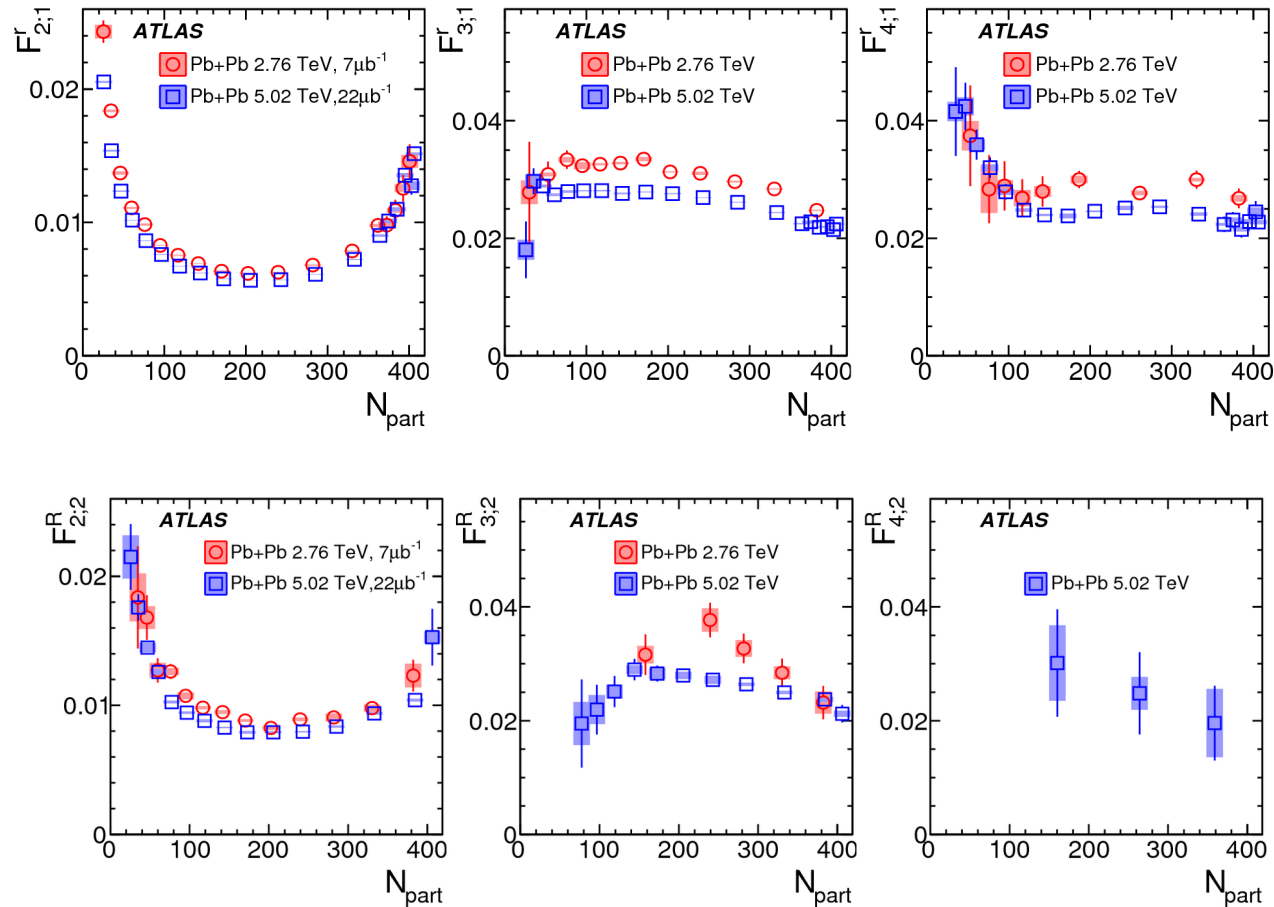
$$R_{n|n;k} = 1 - 2 F_{n;k}^R \eta$$

- decorrelation “strength” parameters:

$$F_n^r = \frac{\sum_i (1 - r_{n|n}(\eta_i)) \eta_i}{2 \sum_i \eta_i^2}$$

$$F_n^R = \frac{\sum_i (1 - R_{n|n}(\eta_i)) \eta_i}{2 \sum_i \eta_i^2}$$

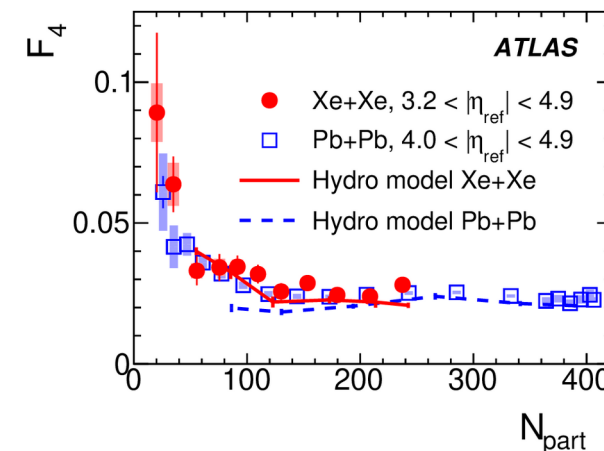
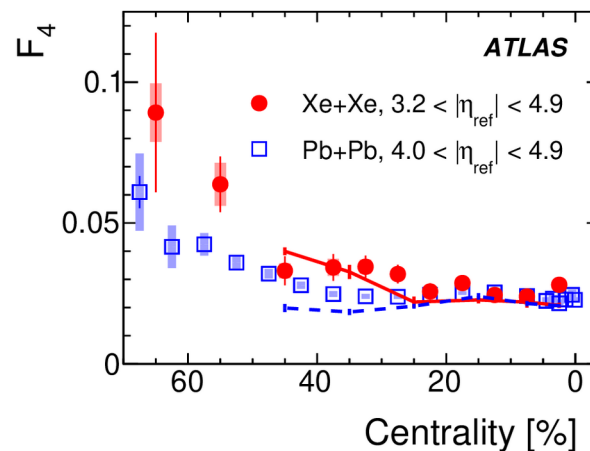
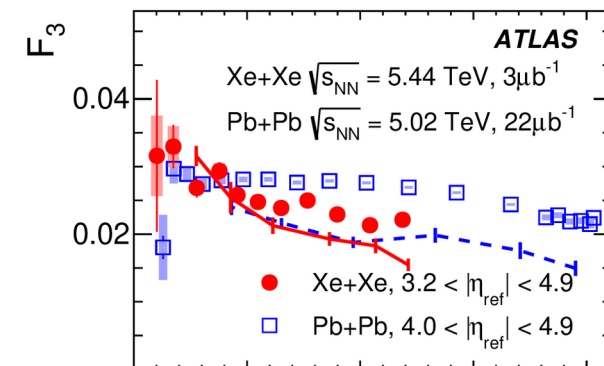
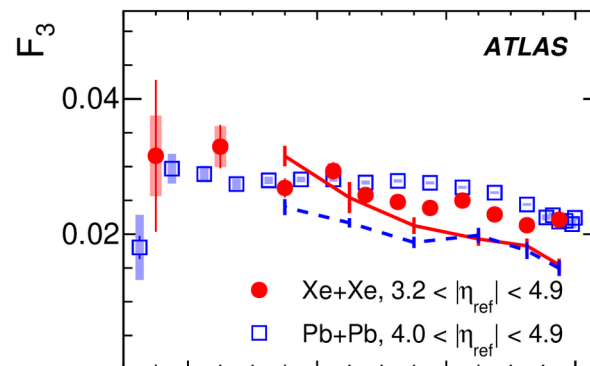
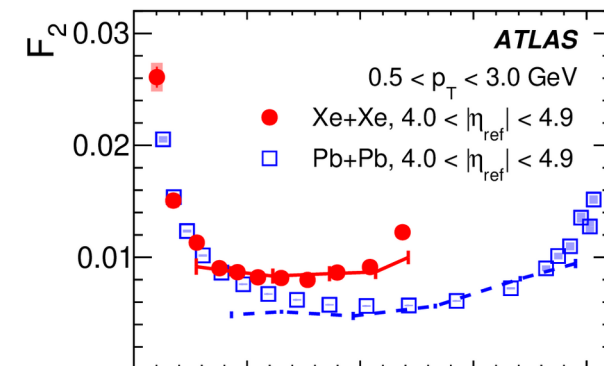
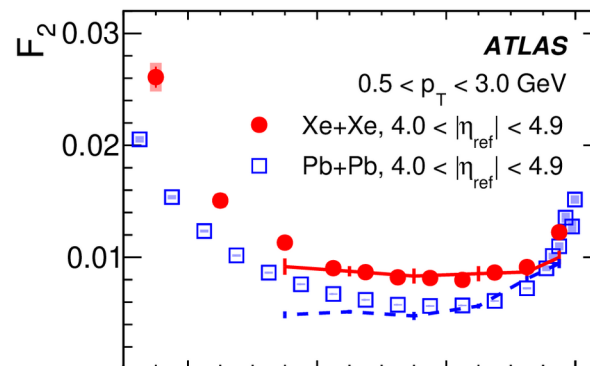
- F_2^r and F_2^R strongly depend on centrality (here N_{part})



- decorrelation “strength” parameter:

$$F_n = \frac{\sum_i (1 - r_{n|n}(\eta_i)) \eta_i}{2 \sum_i \eta_i^2}$$

- F_2 larger in Xe+Xe collisions than in Pb+Pb collisions
- F_3 smaller in Xe+Xe collisions than in Pb+Pb collisions



◆ decorrelation “strength”

parameters:

$$F_n^r = \frac{\sum_i (1 - r_{n|n}(\eta_i)) \eta_i}{2 \sum_i \eta_i^2}$$

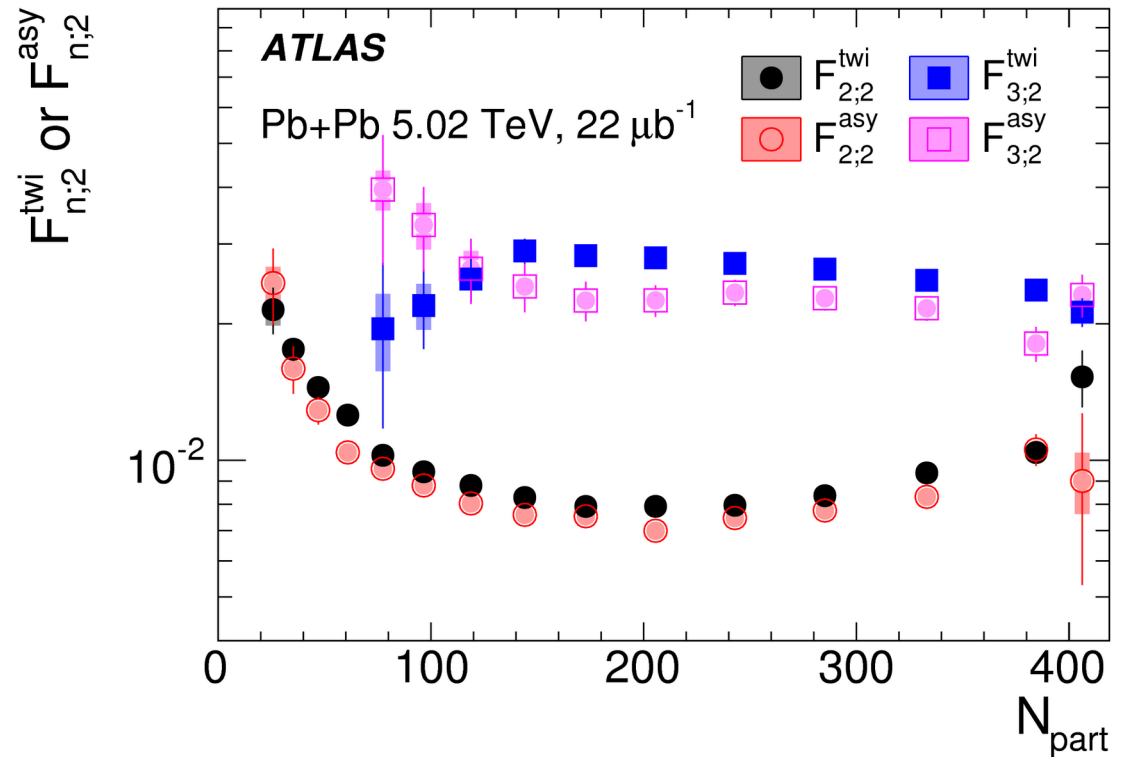
$$F_n^R = \frac{\sum_i (1 - R_{n|n}(\eta_i)) \eta_i}{2 \sum_i \eta_i^2}$$

◆ connection with twist and asymmetry fluctuations:

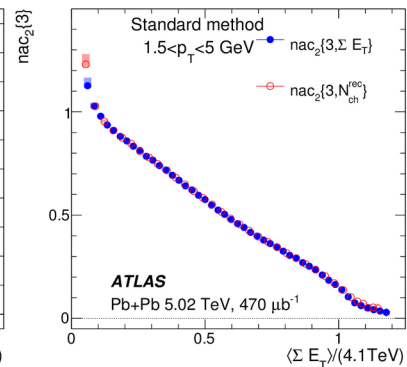
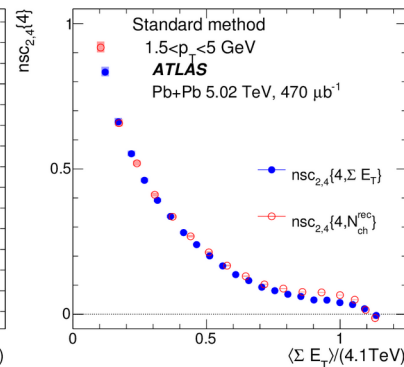
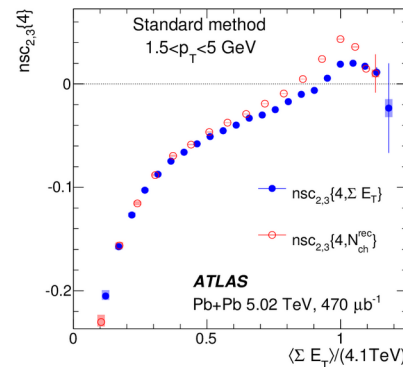
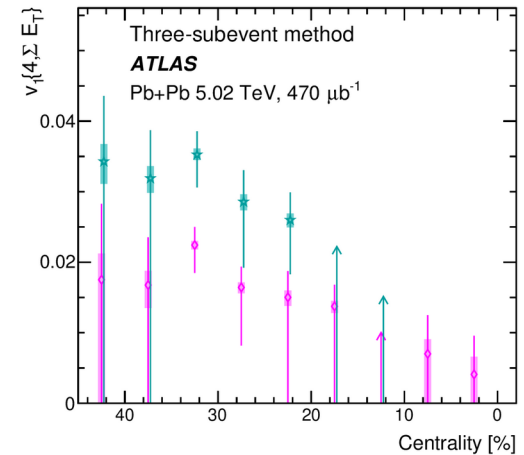
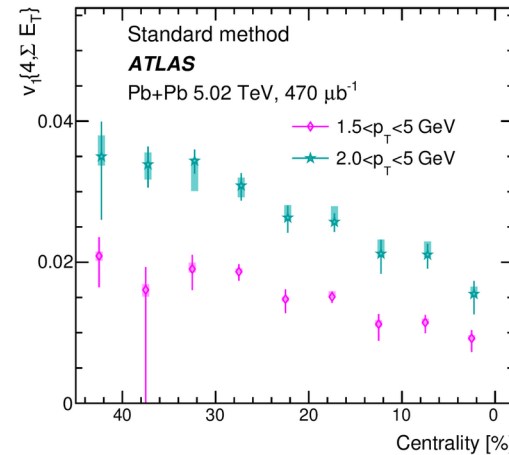
$$F_{n;2}^{twi} = F_{n;2}^R$$

$$F_{n;2}^{asy} = F_{n;2}^r - F_{n;2}^R$$

◆ similar size of these two types of decorrelations



- study of event-by-event fluctuations of harmonics, $p(v_n)$ and $p(v_n, v_m)$, using multiparticle cumulants calculated for different centralities and p_T ranges
- negative $c_1\{4\}$ (thus positive $v_1\{4\}$) found
- negative $c_4\{4\}$ implicates a non-linear contribution to v_4 proportional to v_2^2
- negative asymmetric cumulant $nsc_{23}\{4\}$ reflects anticorrelation between v_2 and v_3 , while positive $nsc_{24}\{4\}$ is consistent with non-linear contributions to v_4
- centrality fluctuations lead to additional v_n fluctuations



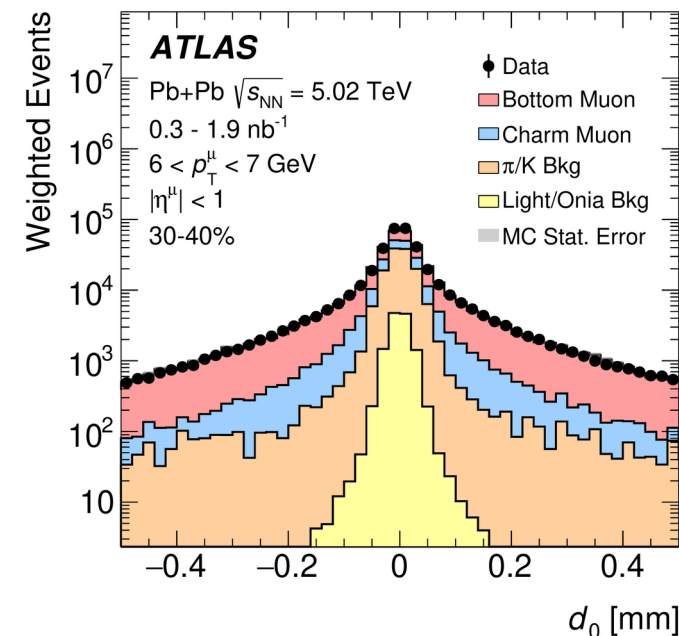
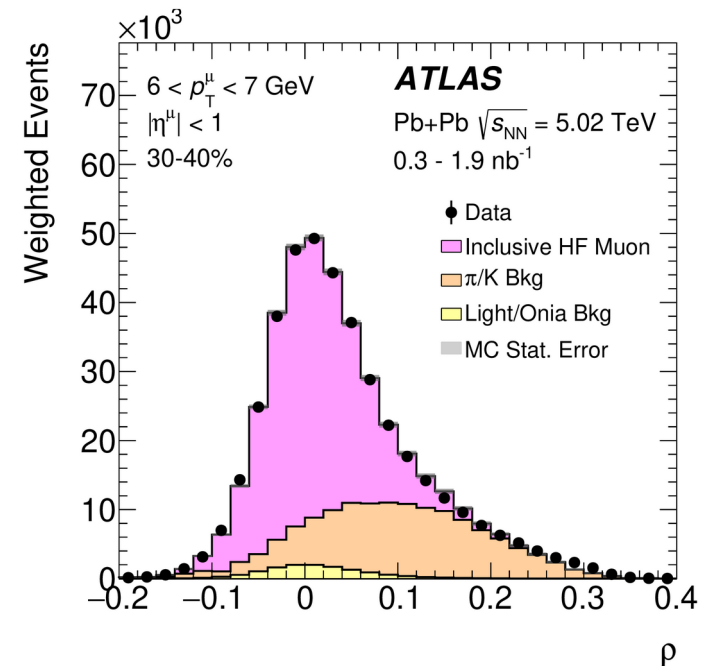
ATLAS, JHEP 01 (2020) 51.

Reconstruction of muons from charm and bottom hadron decays and separation from those from π/K background using:

- ◆ imbalance between momentum measured in the inner detector and in the muon spectrometer

$$\rho = (p^{\text{ID}} - p^{\text{MS}}) / p^{\text{ID}}$$
- ◆ transverse impact parameter d_0

Different distributions of these variables enable separation of muons from different decays



v_2 :

- elliptic flow of bottom muons is smaller than that of charm muons
- v_2 increases in more peripheral collisions

ATLAS

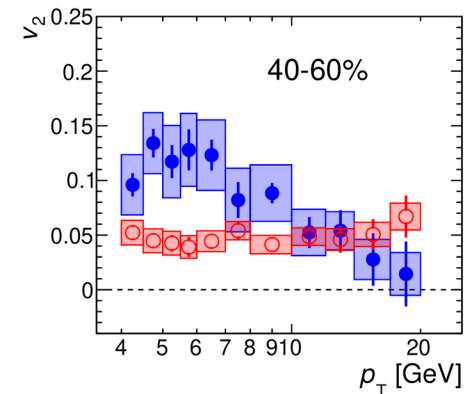
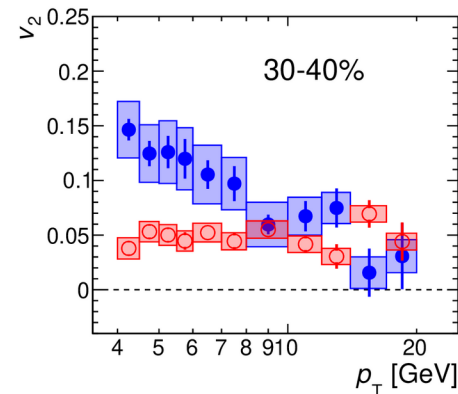
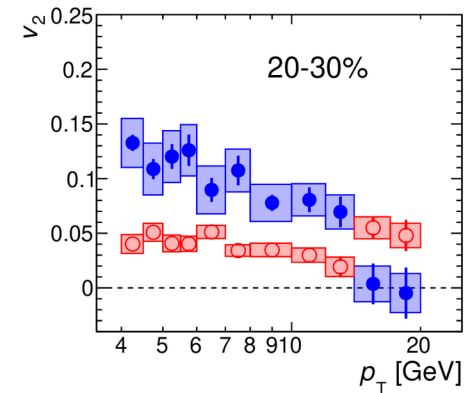
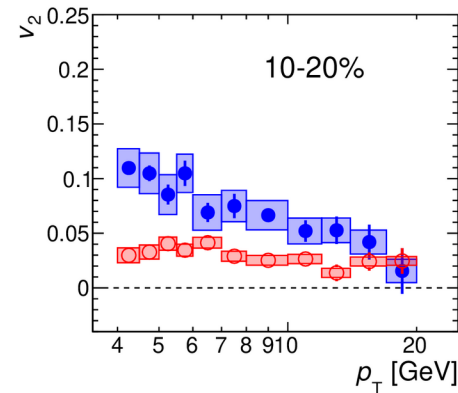
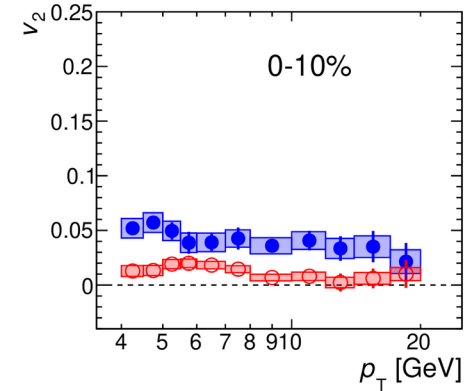
Pb+Pb $\sqrt{s_{NN}} = 5.02$ TeV

0.3 - 1.9 nb⁻¹

$|\eta^\mu| < 2$

● charm muon

○ bottom muon



ATLAS, arxiv:2003.03565

V_3 :

- also triangular flow of bottom muons is smaller than that of charm muons

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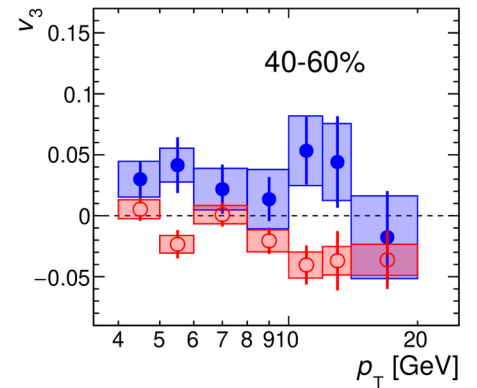
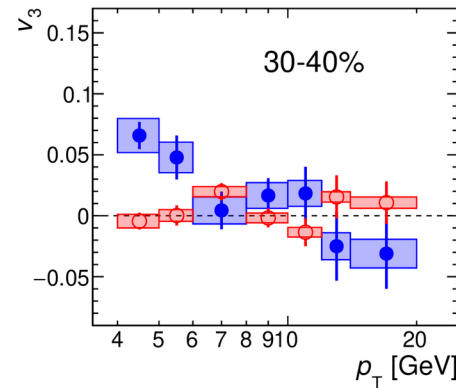
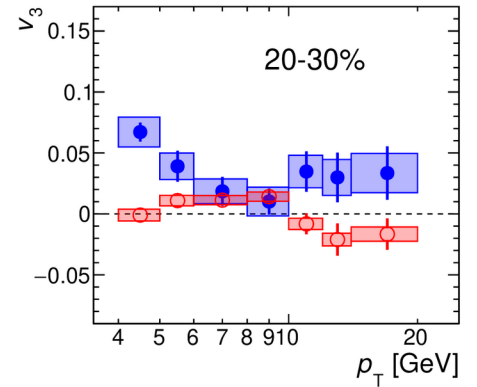
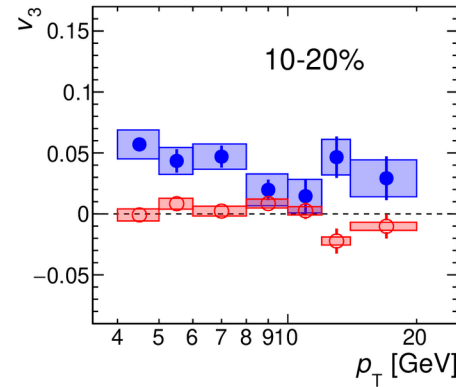
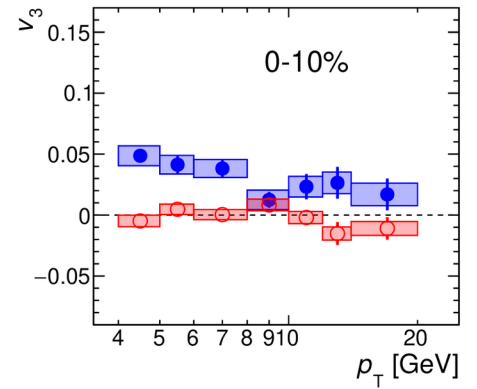
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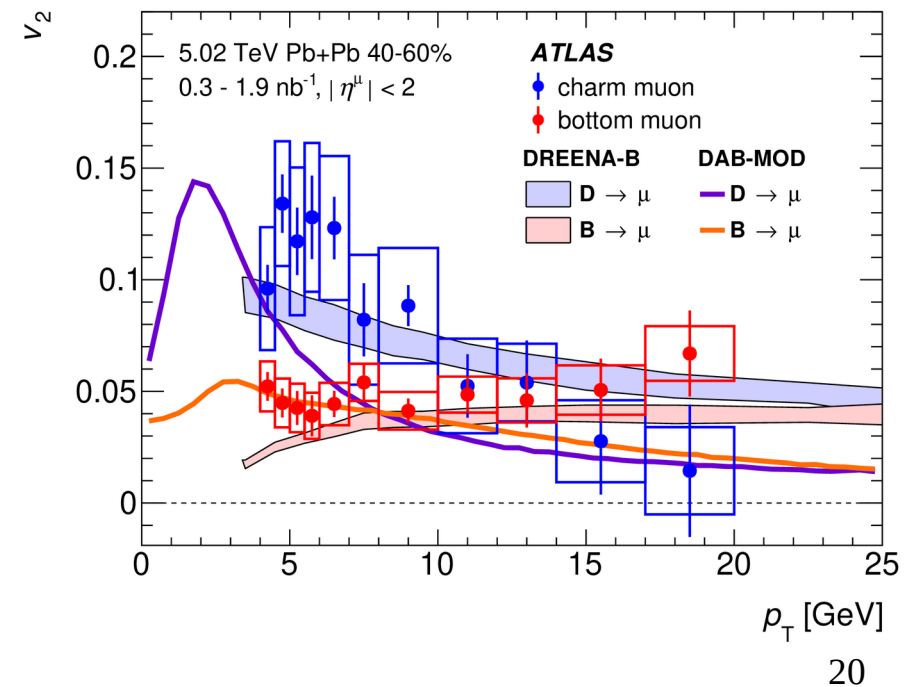
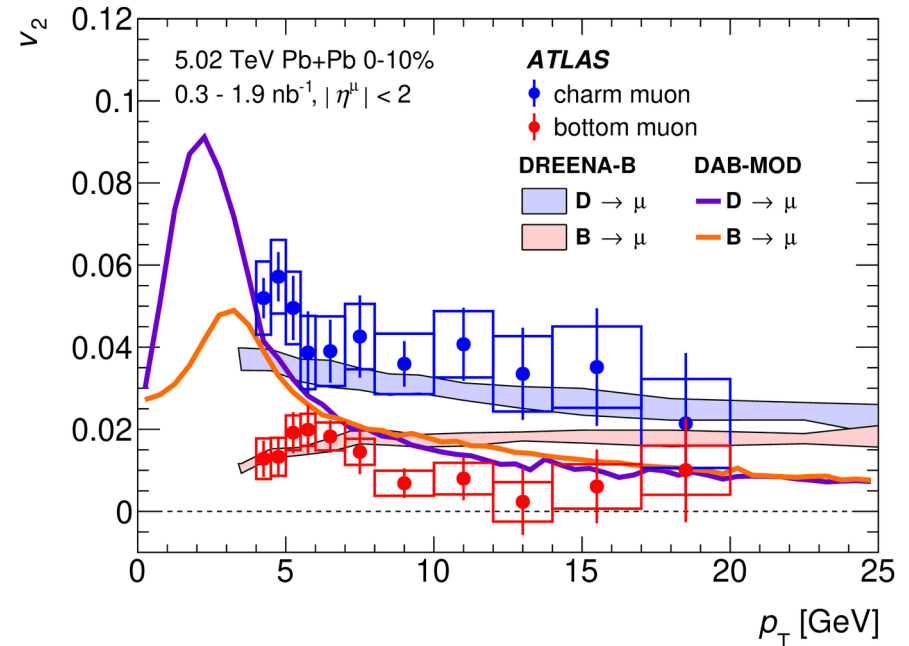
● charm muon

○ bottom muon



ATLAS, arxiv:2003.03565

Transverse momentum dependence of v_2 has similar trend as predicted by the models, differences may shed the light on details of QGP expansion and the energy-loss mechanism.



ATLAS, arxiv:2003.03565

Summary

Correlation studies in Pb+Pb and Xe+Xe collisions:

- detailed analysis of azimuthal correlations in Pb+Pb and Xe+Xe collisions,
- scaling of $v_n / (v_k)^{n/k}$ ratio,
- fluctuations of harmonics as a function of pseudorapidity (longitudinal decorrelations),
- analysis of event-by-event fluctuations with multiparticle cumulants,
- flow of charm and bottom hadrons.

Deeper insight into properties and expansion of QGP and the effects of fluctuations of initial conditions in the collisions

Backup

Pearson correlation coefficient R , substituted by ρ to be not distorted at small multiplicity

$$R = \frac{\text{cov}(v_n\{2\}^2, [p_T])}{\sqrt{\text{Var}(v_n\{2\}^2)}\sqrt{\text{Var}([p_T])}}$$

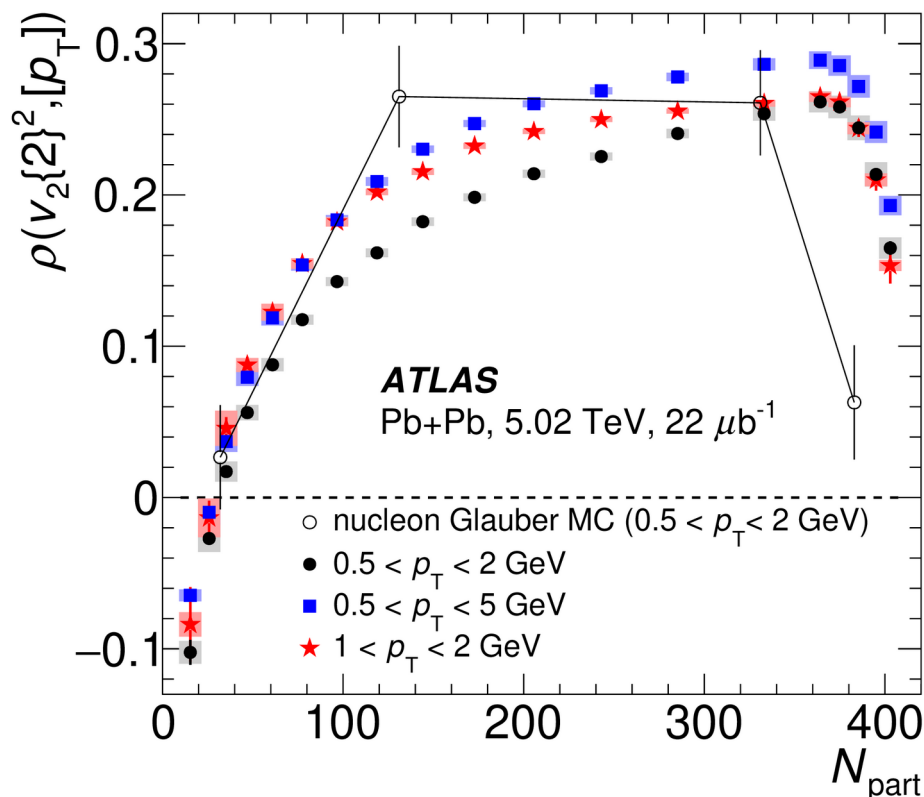


$$\rho = \frac{\text{cov}(v_n\{2\}^2, [p_T])}{\sqrt{\text{Var}(v_n\{2\}^2)_{\text{dyn}}}\sqrt{c_k}}$$

in which dynamical variances are used

Harmonics are calculated in different pseudorapidity intervals ($0.75 < |\eta| < 2.5$) than the mean p_T of the event ($|\eta| < 0.5$)

ATLAS, Eur. Phys. J. C 79 (2019) 985



- Negative correlation in peripheral collisions
- Positive correlation for more central collisions
- Qualitatively in agreement with hydrodynamics 1+3D
(P. Božek, Phys. Rev. C93 (2016) 044908)

Suppression of contributions from non-flow effects

Flow estimation methods

- ◆ event plane
- ◆ two-particle correlations
- ◆ scalar product

$$v_n = \langle \cos(n(\phi - \Phi_n)) \rangle$$

$$v_{n,n} = \langle \cos(n(\phi_a - \phi_b)) \rangle$$

$$v_n = \text{Re}(\langle q_n^N Q_n^{P*} \rangle / \sqrt{\langle Q_n^N Q_n^{P*} \rangle}); \quad q_n, Q_n = (1/\Sigma_j w_j) \Sigma_j w_j e^{in\phi_j}$$

Multi-particle correlations

- ◆ standard cumulants
- ◆ symmetric cumulants
- ◆ asymmetric

$$\langle \langle \{2k\}_n \rangle \rangle = \langle \langle e^{in(\phi_1 + \dots + \phi_k - \phi_{k+1} - \dots - \phi_{2k})} \rangle \rangle = \langle v_n^{2k} \rangle$$

$$\langle \langle \{4\}_{n,m} \rangle \rangle = \langle \langle e^{in(\phi_1 - \phi_2) + \Im(\phi_3 - \phi_4)} \rangle \rangle = \langle v_n^2 v_m^2 \rangle$$

$$c_n\{4\} = \langle \langle \{4\}_n \rangle \rangle - 2 \langle \langle \{2\}_n \rangle \rangle^2$$

$$SC_{n,m}\{4\} = \langle \langle \{4\}_{n,m} \rangle \rangle - \langle \langle \{2\}_n \rangle \rangle \langle \langle \{2\}_m \rangle \rangle$$

$$ac_2\{3\} = \langle \langle \{3\}_n \rangle \rangle = \langle \langle e^{i(n\phi_1 + n\phi_2 - 2n\phi_3)} \rangle \rangle^2$$

Subevent methods - particles selected from different regions in pseudorapidity

- ◆ two-subevents
- ◆ three-subevents
- ◆ four-subevents

$$SC_{n,m}^{2a|2c}\{4\} = \langle \langle \{4\}_{n,m} \rangle \rangle_{2a|2c} - \langle \langle \{2\}_n \rangle \rangle_{a|b} \langle \langle \{2\}_m \rangle \rangle_{a|b}$$

$$SC_{n,m}^{a,b|2c}\{4\} = \langle \langle \{4\}_{n,m} \rangle \rangle_{a,b|2c} - \langle \langle \{2\}_n \rangle \rangle_{a|c} \langle \langle \{2\}_m \rangle \rangle_{b|c}$$

$$SC_{n,m}^{a,b|c,d}\{4\} = \langle \langle \{4\}_{n,m} \rangle \rangle_{a,b|c,d} - \langle \langle \{2\}_n \rangle \rangle_{a|c} \langle \langle \{2\}_m \rangle \rangle_{b|d}$$

Many different methods tested in order to obtain "real" flow