

ATLAS measurements of transverse and longitudinal flow fluctuations

Phys. Rev. C 101 (2020), 024906

arXiv:2001.04201

Arabinda Behera

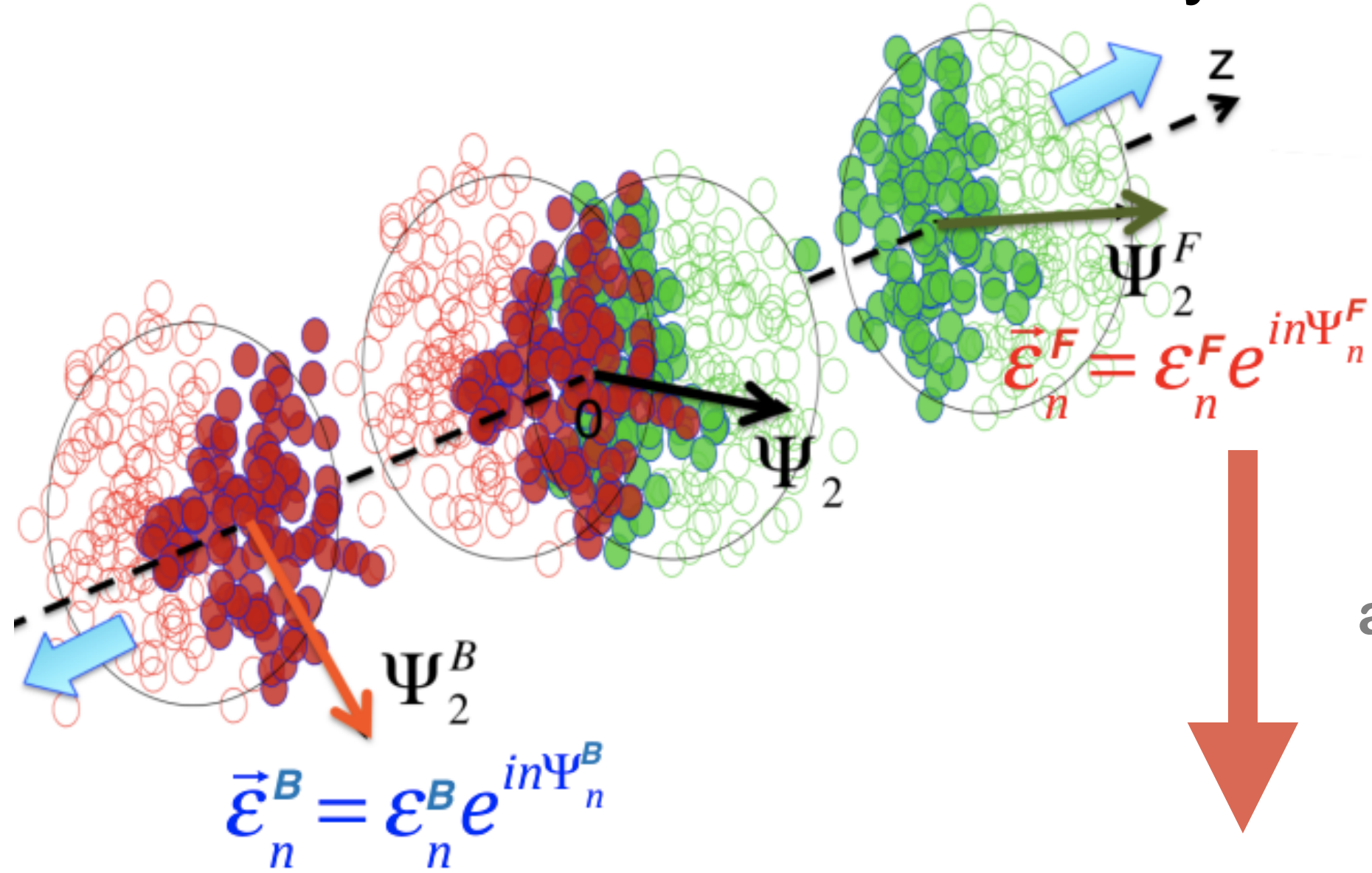
For the ATLAS Collaboration

3rd June, 2020

Hard Probes 2020



- Initial state fluctuations - FB asymmetry in eccentricities



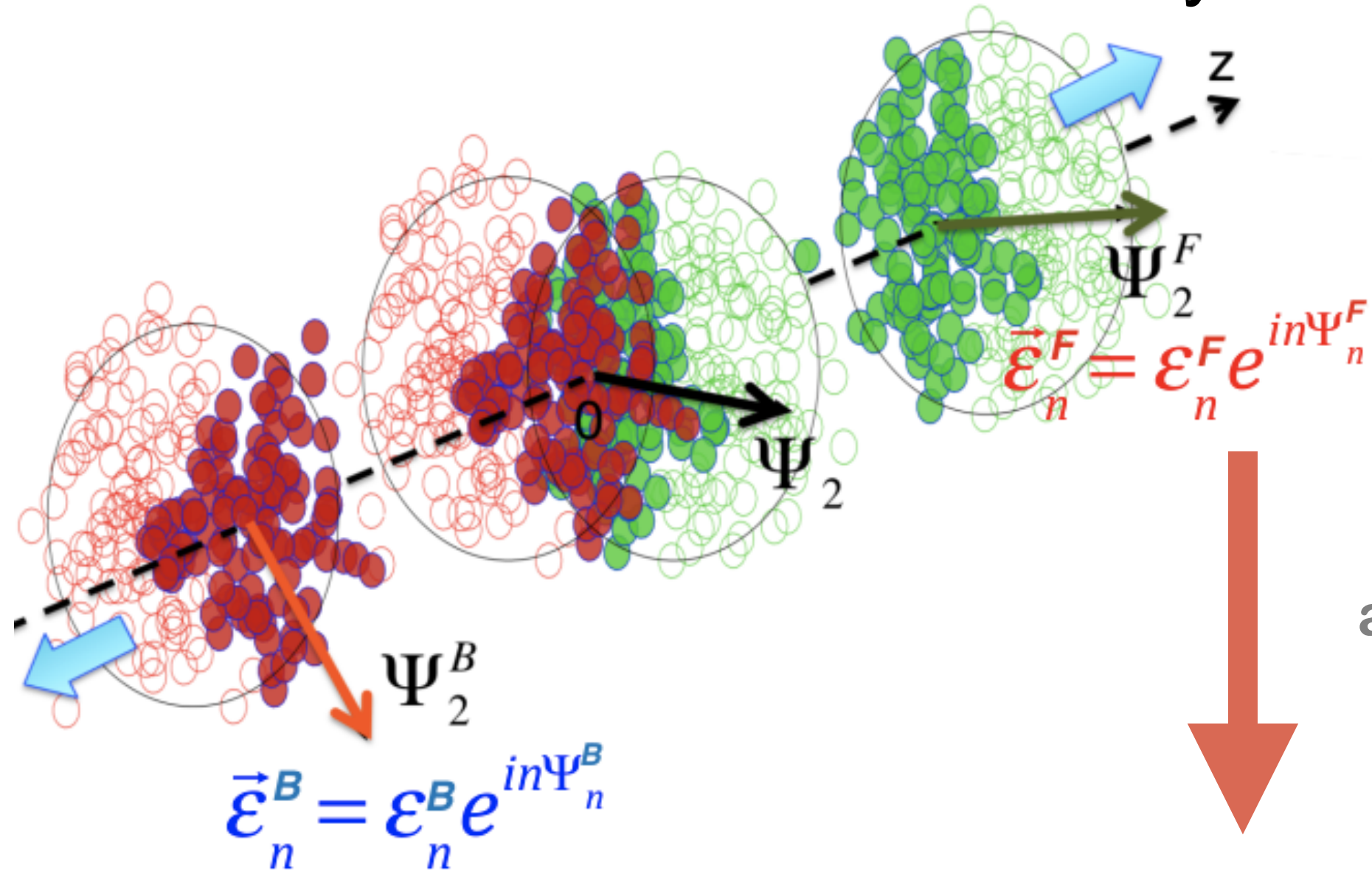
$$\vec{\epsilon}_n^F \neq \vec{\epsilon}_n^B$$

$$\vec{\epsilon}_{n+} = (\vec{\epsilon}_n^F + \vec{\epsilon}_n^B) / 2 \implies \text{Average Flow}$$

$$\vec{\epsilon}_{n-} = (\vec{\epsilon}_n^F - \vec{\epsilon}_n^B) / 2 \implies \text{Longitudinal Fluctuation}$$

arXiv:1701.02183

- Initial state fluctuations - FB asymmetry in eccentricities



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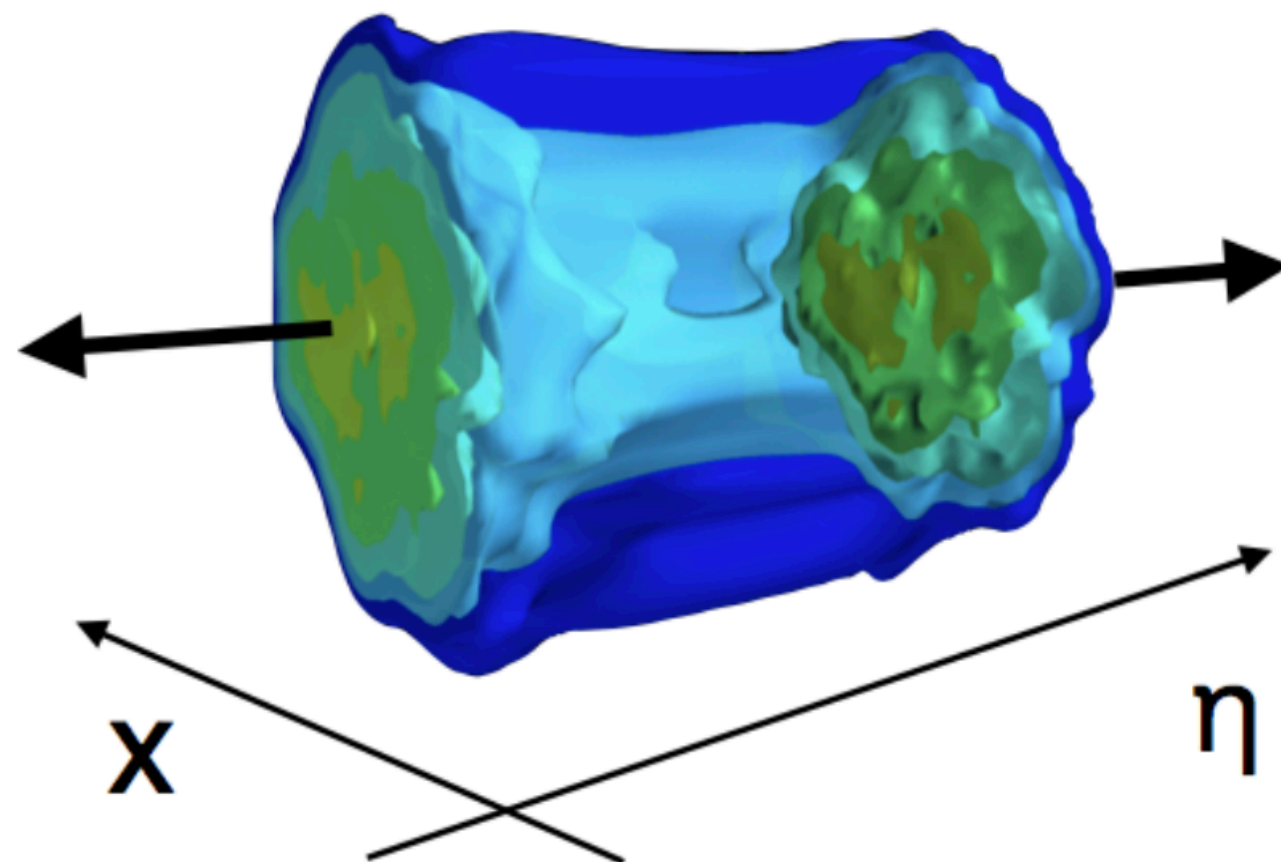
$$\vec{\epsilon}_{n+} = (\vec{\epsilon}_n^F + \vec{\epsilon}_n^B)/2 \implies \text{Average Flow}$$

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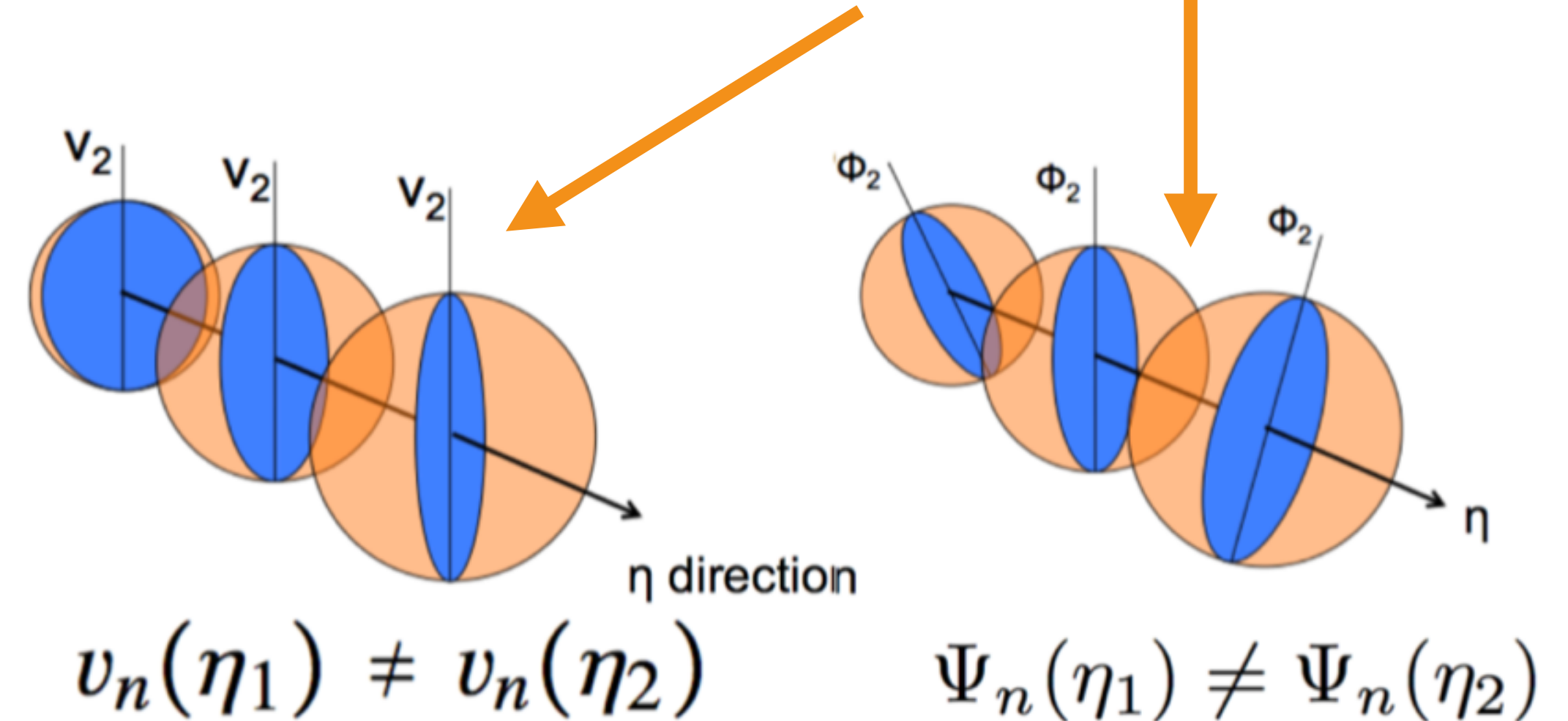
- Hydro evolution of QGP (3+1D)

$$\vec{v}_n = \kappa_n \vec{\epsilon}_n \quad \vec{v}_n(\eta) = v_n(\eta) e^{in\Psi_n(\eta)}$$



- Final State - average v_n and v_n decorrelation

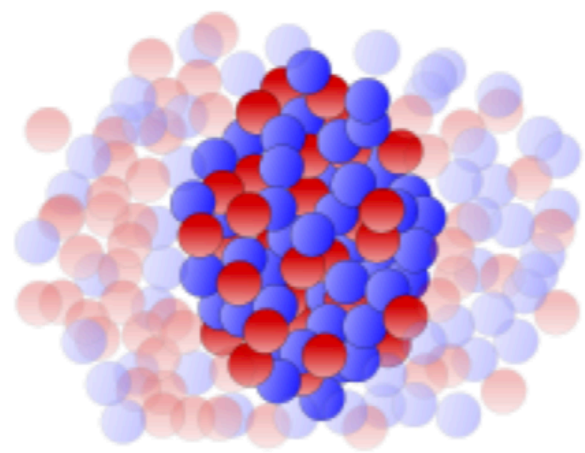
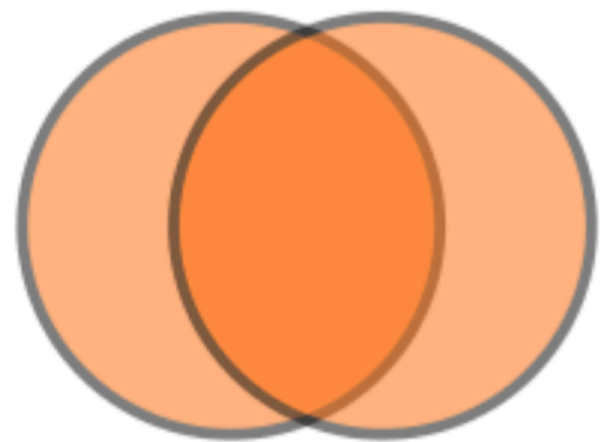
$$\vec{v}_n(\eta) \approx \vec{v}_n(0) (1 + \alpha_n \eta) e^{i\beta_n \eta}$$



Pb+Pb

$A=208, R=6.62$ fm

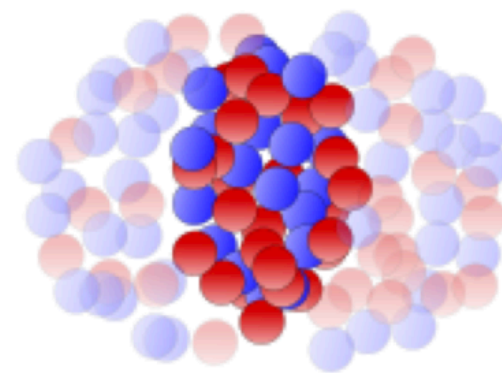
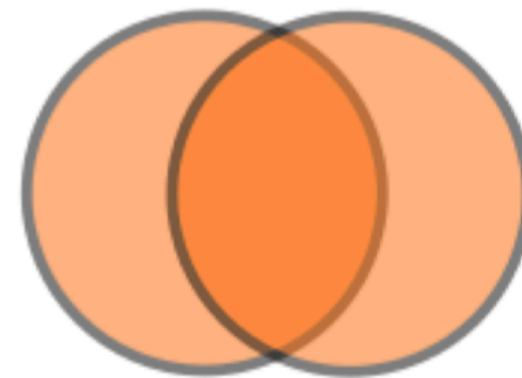
$\sqrt{s_{NN}} = 5.02$ TeV



Xe+Xe

$A=129, R=5.42$ fm

$\sqrt{s_{NN}} = 5.44$ TeV



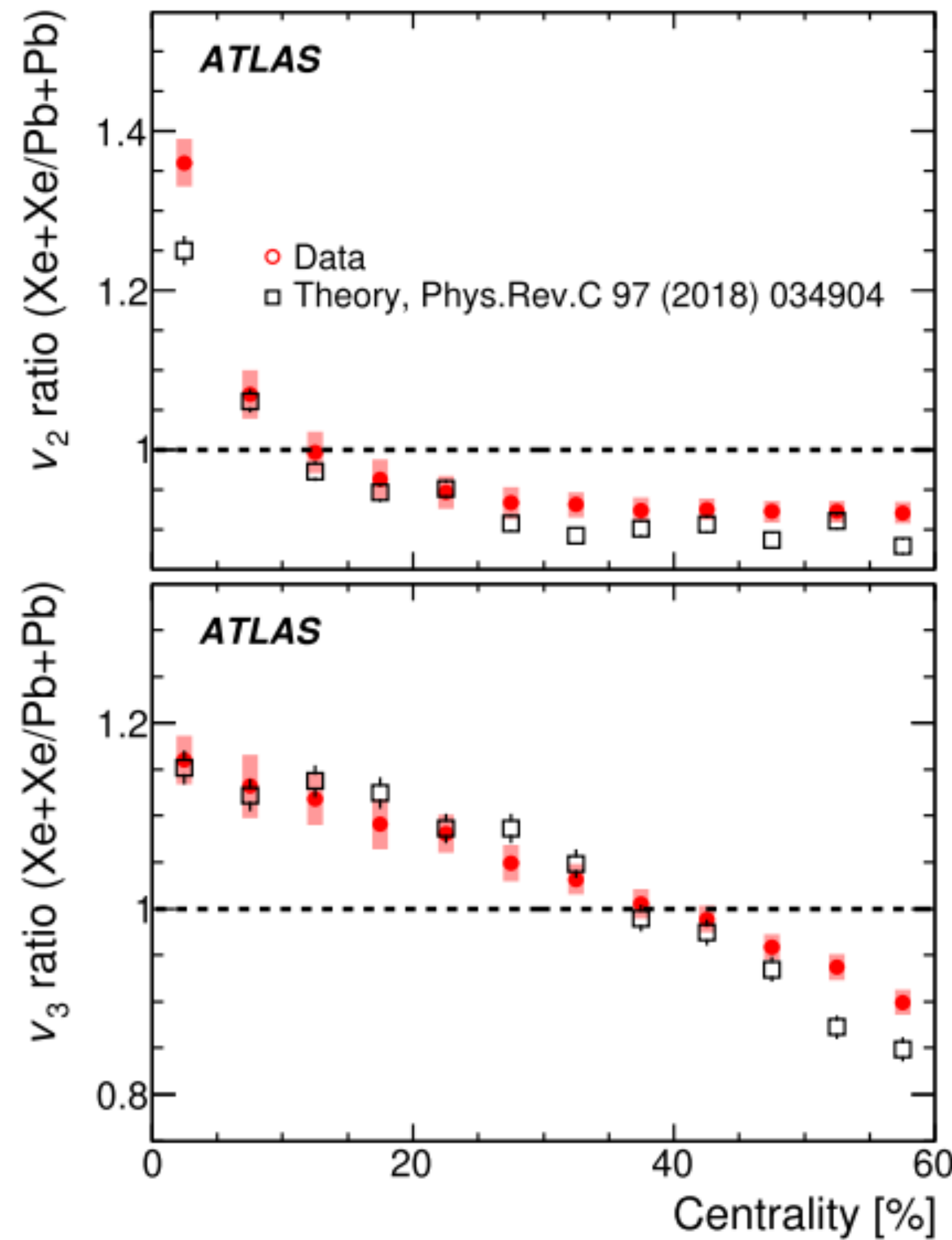
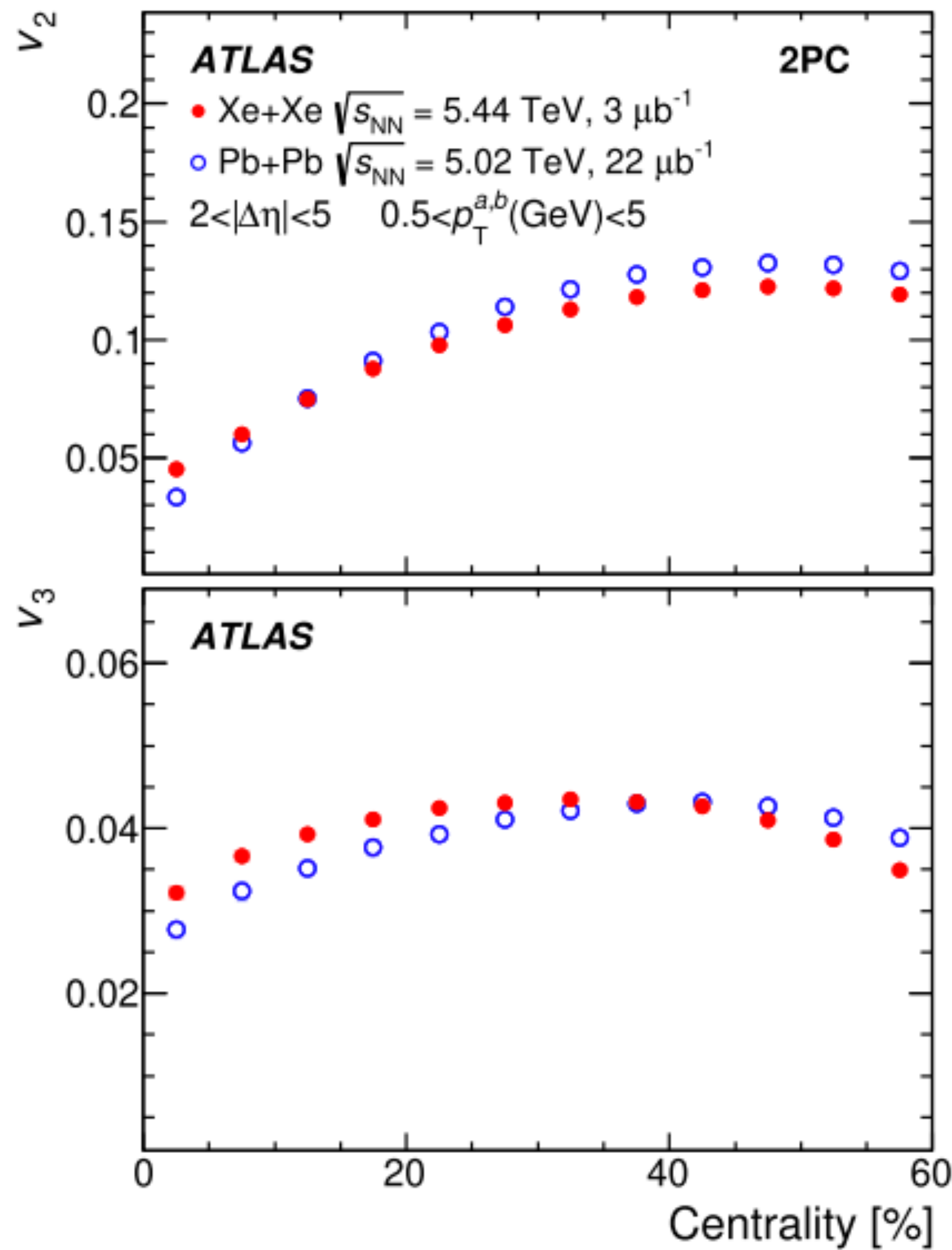
- Different factors can affect the v_n and decorrelation -
 - Initial state geometry
 - Initial momentum anisotropy
 - Shear viscosity
 - Hydrodynamic fluctuations
- How flow and decorrelation are related?
- What is the effect of nuclear deformation (Xe+Xe)?
- Is the dependence same for different harmonics?
- How to disentangle early-time and late-time effects?

◆ System-size dependence of v_n and v_n decorrelation : Xe+Xe vs Pb+Pb

◆ Ratio between the two systems - disentangle initial and final-state contributions

Centrality dependence - Xe+Xe vs Pb+Pb

- Method - 2-particle correlation technique



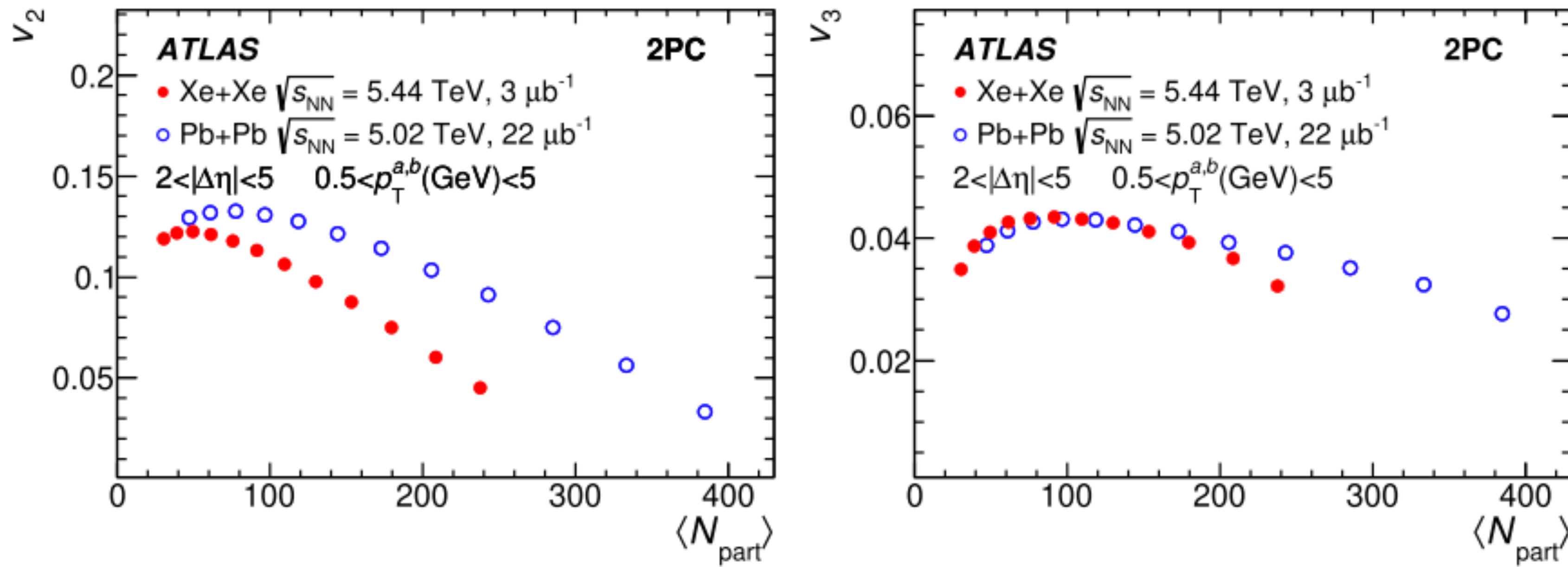
v_2

- Most Central events - Xe+Xe $v_2 > \text{Pb+Pb } v_2$
- From central to mid-central events, the ratio for v_2 decreases, becomes < 1 after 15% centrality
- Peripheral events - ratio seems to saturate

v_3

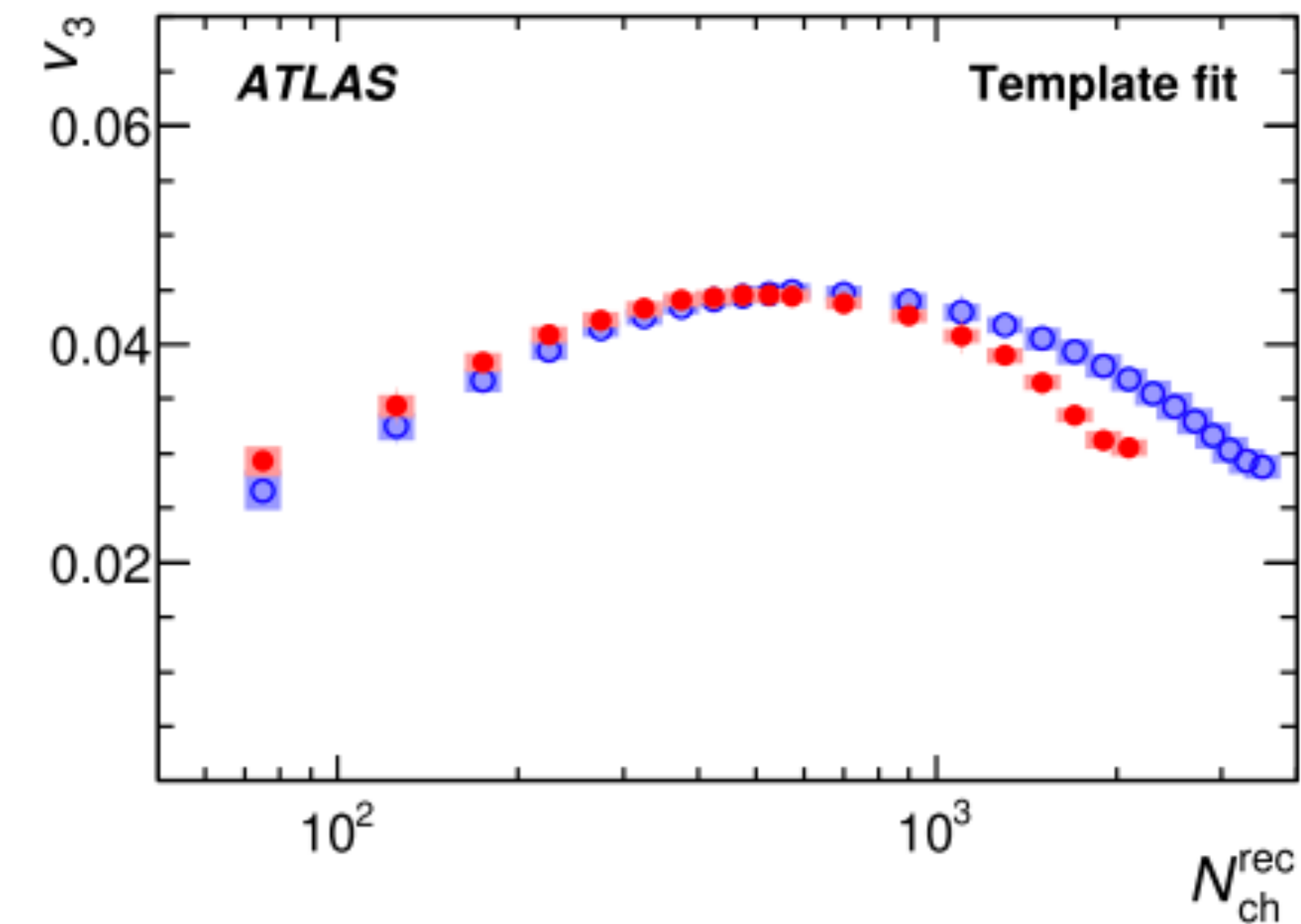
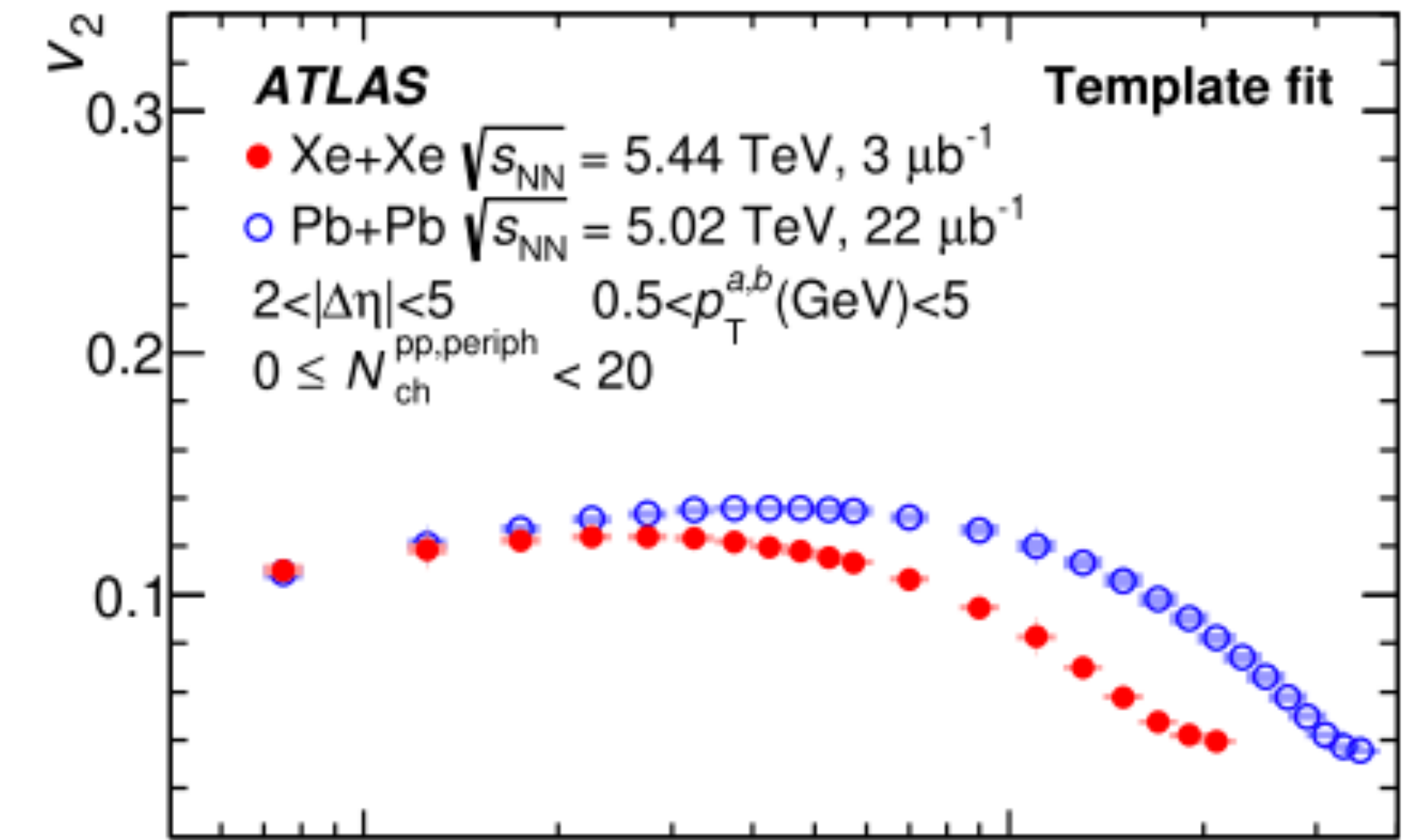
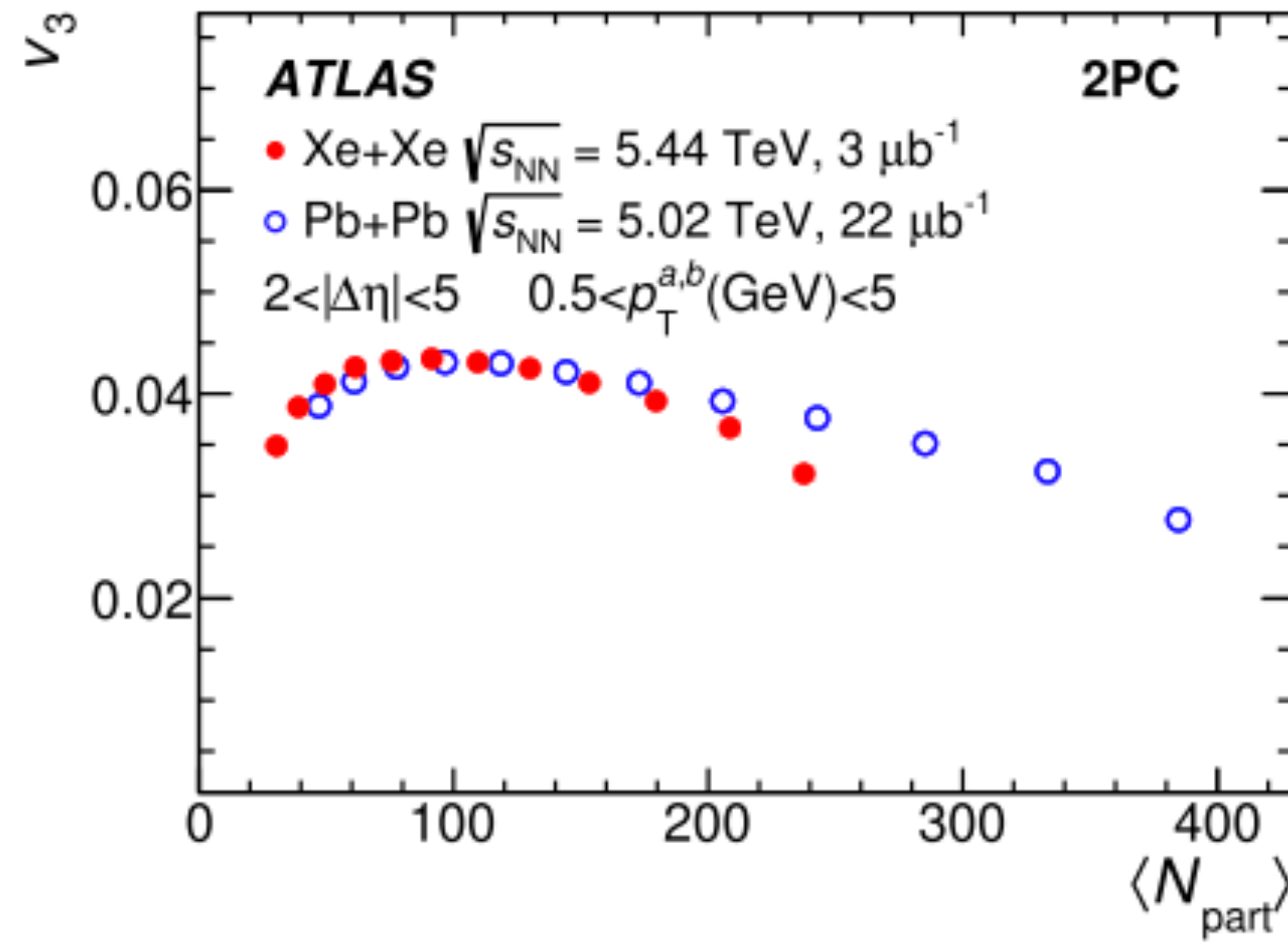
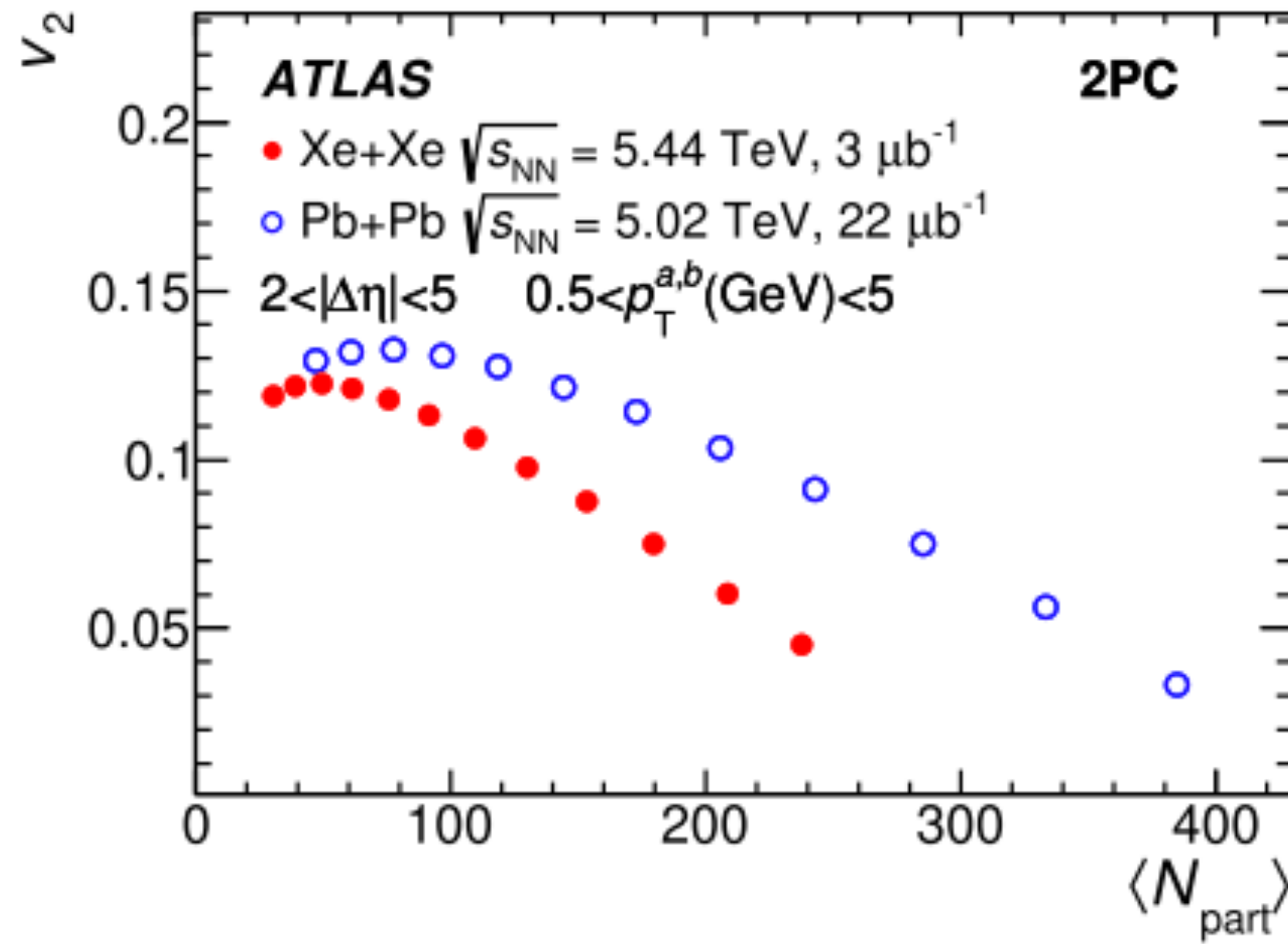
- Similar qualitative trends
- Ratio < 1 after 35% centrality

- Central - Geometry dominates, Non-central - Viscous effect dominates



- v_2 - smaller in Xe+Xe than Pb+Pb
- v_3 - similar values except in large N_{part}
- Low N_{part} - all v_n similar values between Xe+Xe and Pb+Pb

N_{part} and N_{ch} dependence - Xe+Xe vs Pb+Pb

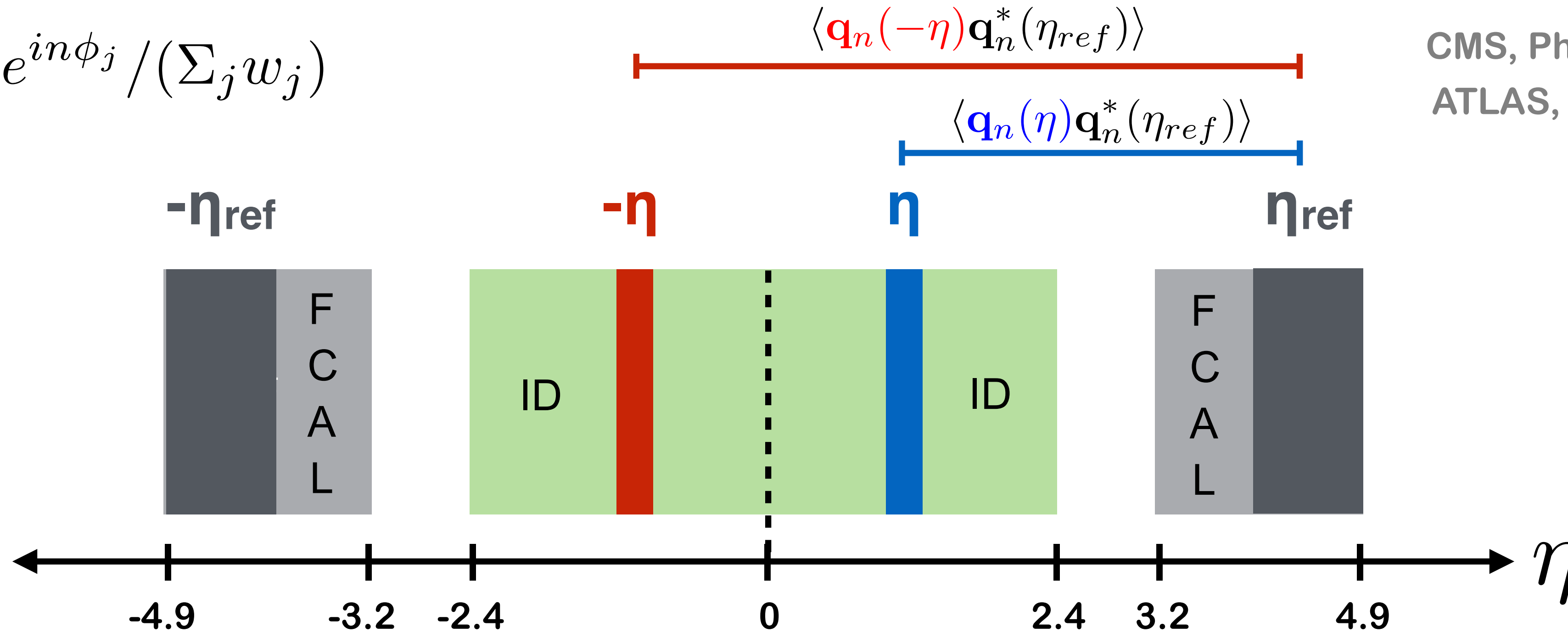


- v_2 - smaller in Xe+Xe than Pb+Pb
- v_3 - similar values except in large N_{part}
- Low N_{part} - all v_n similar values between Xe+Xe and Pb+Pb
- N_{ch} dependence - follow similar behavior as N_{part}

- Flow vector

$$\mathbf{q}_n \equiv \sum_j w_j e^{in\phi_j} / (\sum_j w_j)$$

CMS, Phys. Rev. C 92, 034911
 ATLAS, Eur. Phys. J. C (2018)
 78:142

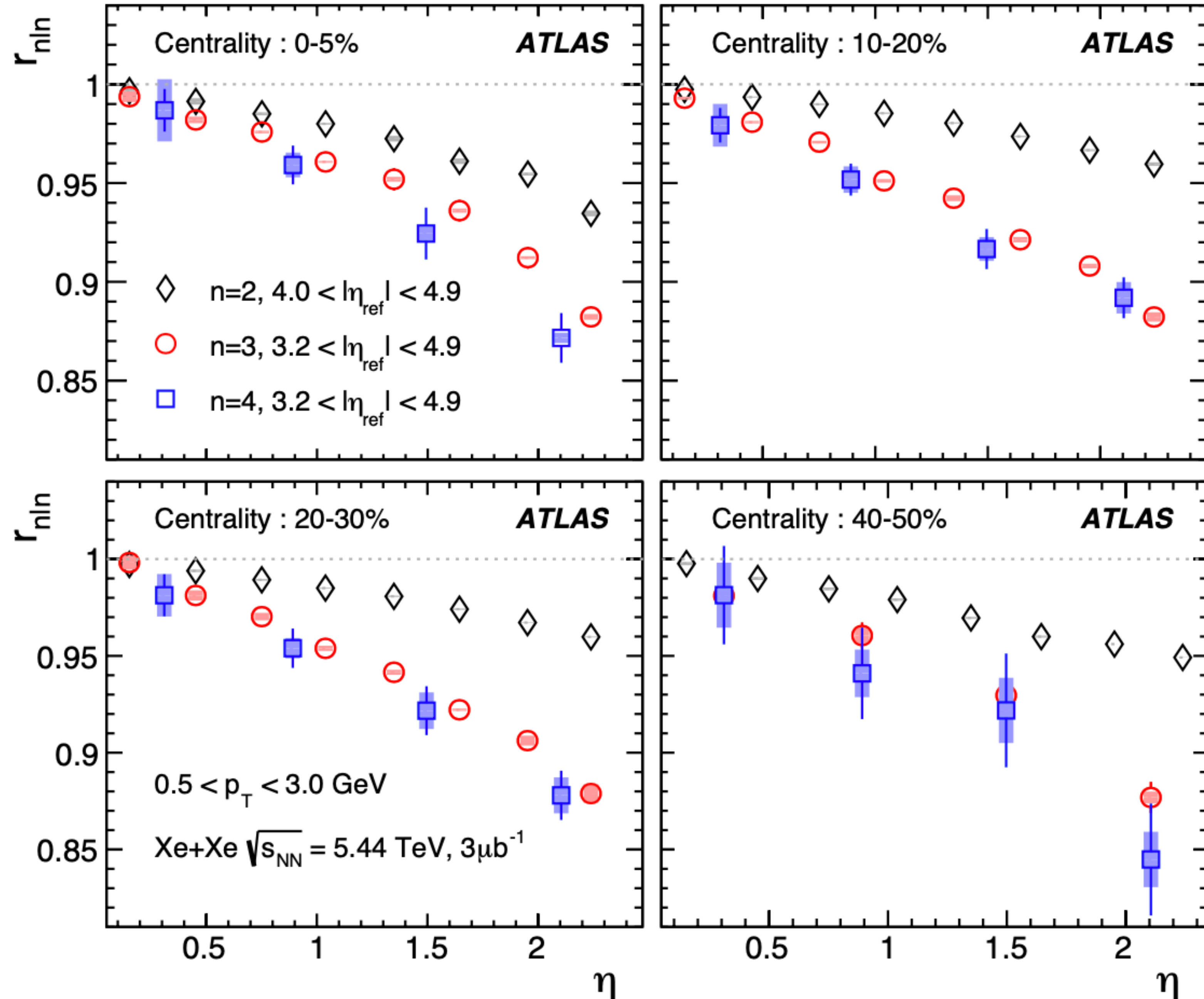


$$r_{n|n}(\eta) = \frac{\langle \mathbf{q}_n(-\eta) \mathbf{q}_n^*(\eta_{ref}) \rangle}{\langle \mathbf{q}_n(\eta) \mathbf{q}_n^*(\eta_{ref}) \rangle} = \frac{\langle v_n(-\eta) v_n(\eta_{ref}) \cos n(\Psi_n(-\eta) - \Psi_n(\eta_{ref})) \rangle}{\langle v_n(\eta) v_n(\eta_{ref}) \cos n(\Psi_n(\eta) - \Psi_n(\eta_{ref})) \rangle} \quad r_{n|n}(\eta) \leq 1$$

Measures decorrelation between $-\eta$ and η

Decorrelation in Xe+Xe

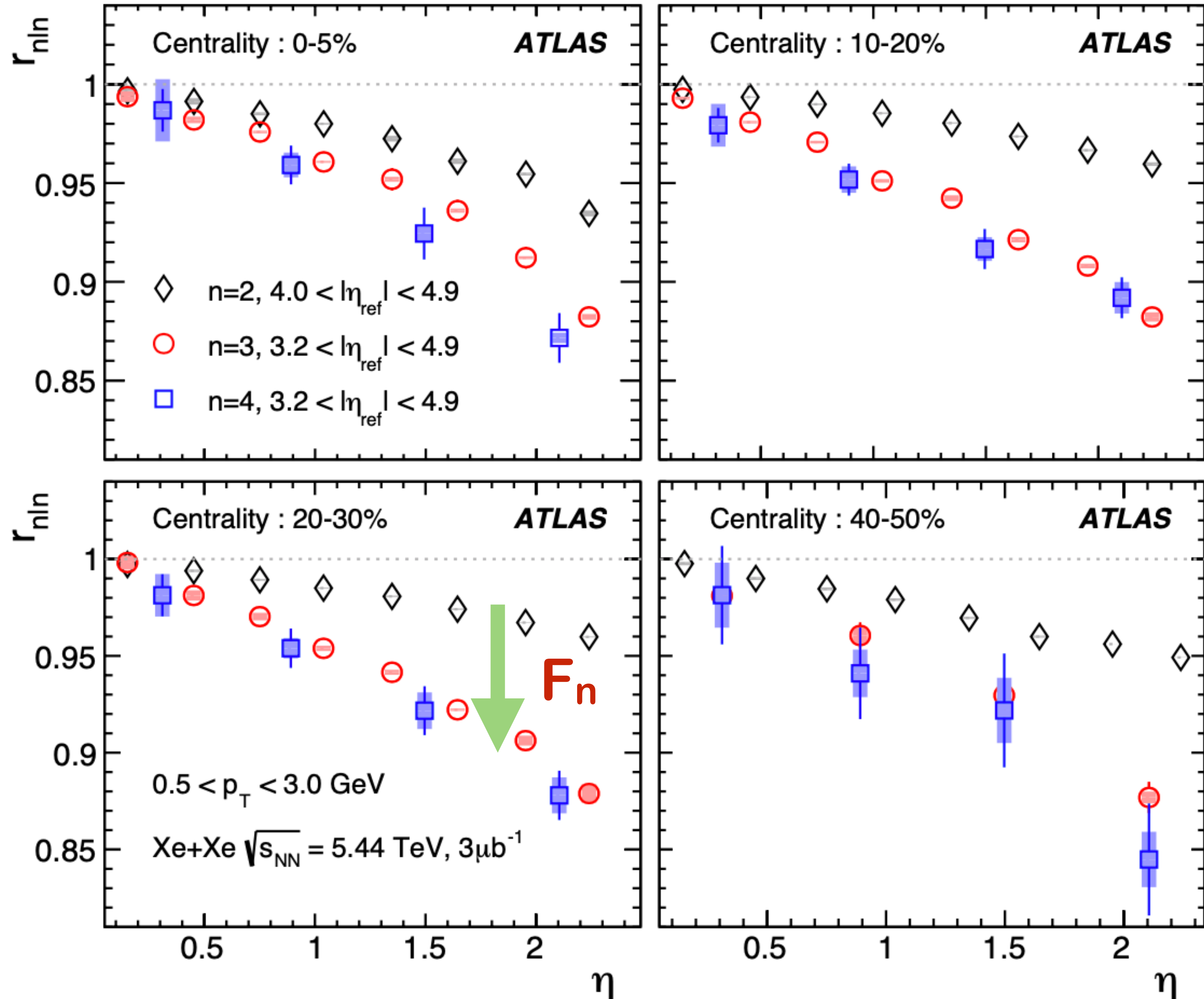
- $r_{2|2} > r_{3|3} > r_{4|4}$ - Linear decrease with η
- $r_{2|2}$ strong centrality dependence - larger in mid-central
- $r_{3|3}$ and $r_{4|4}$ - weak centrality dependence
- Similar behavior as 5 TeV Pb+Pb



- $r_{2|2} > r_{3|3} > r_{4|4}$ - Linear decrease with η
- $r_{2|2}$ strong centrality dependence - larger in mid-central
- $r_{3|3}$ and $r_{4|4}$ - weak centrality dependence
- Similar behavior as 5 TeV Pb+Pb
- Slope F_n :

$$r_{n|n}(\eta) = 1 - 2F_n\eta$$

↓
Decorrelation strength

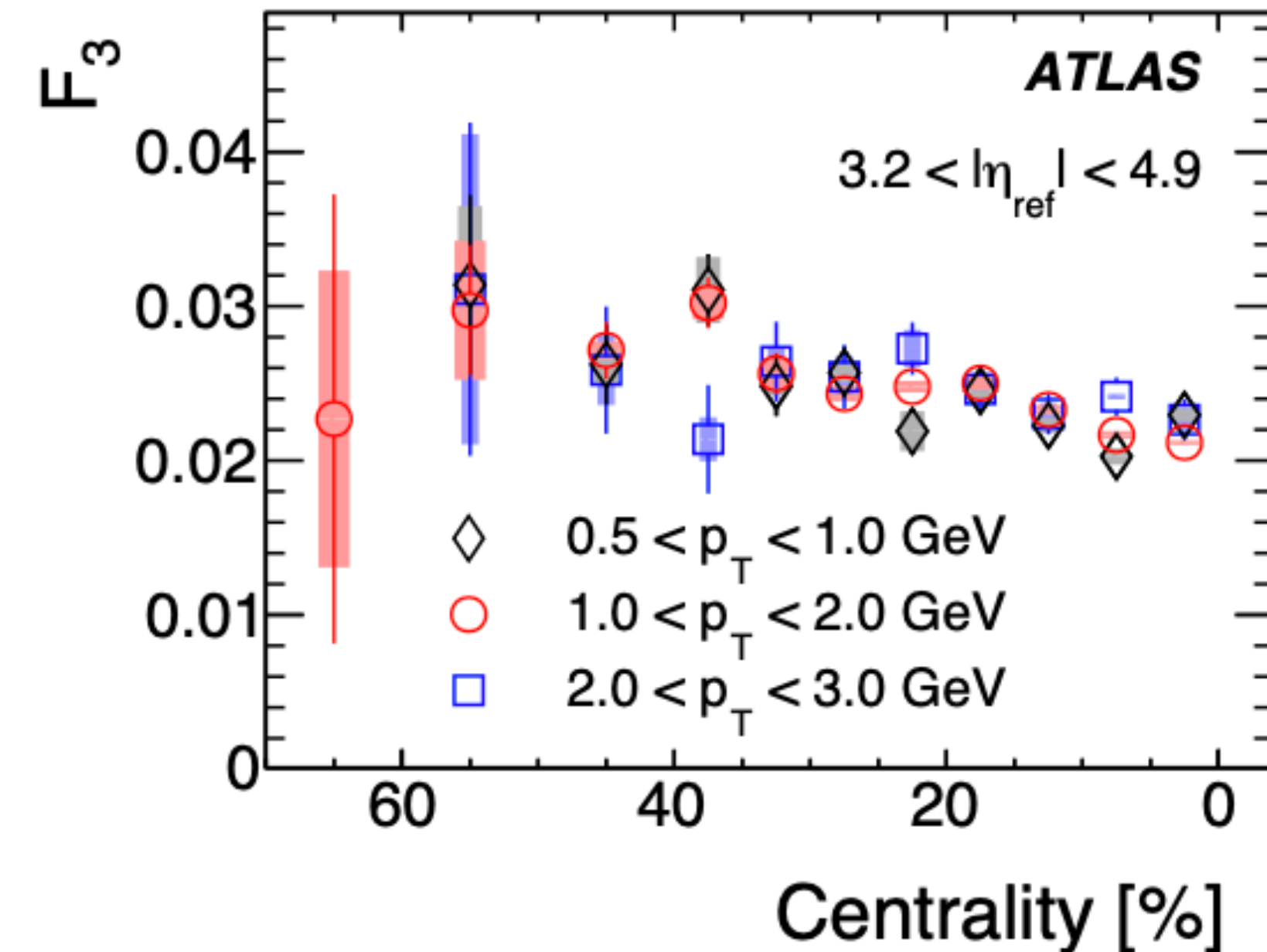
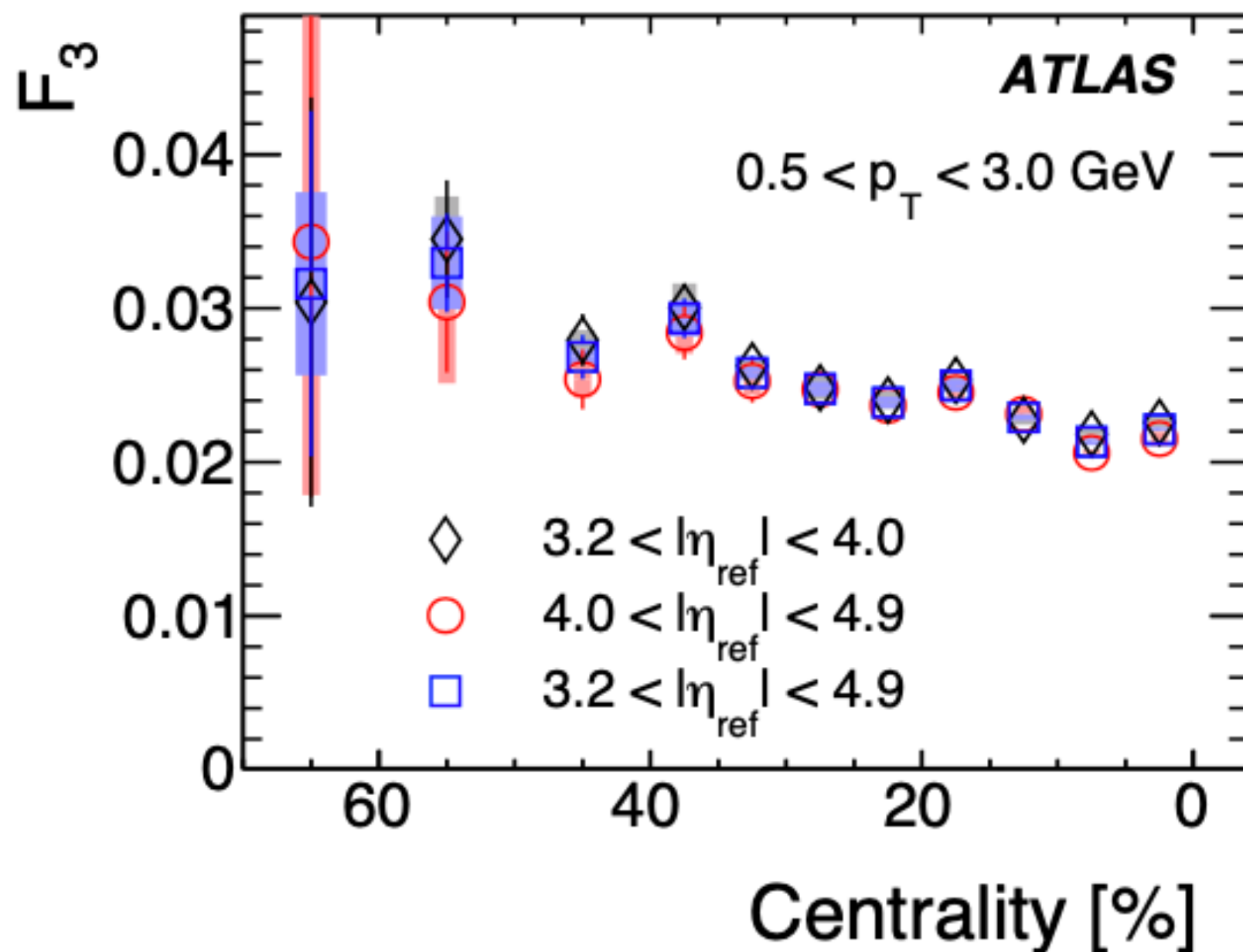
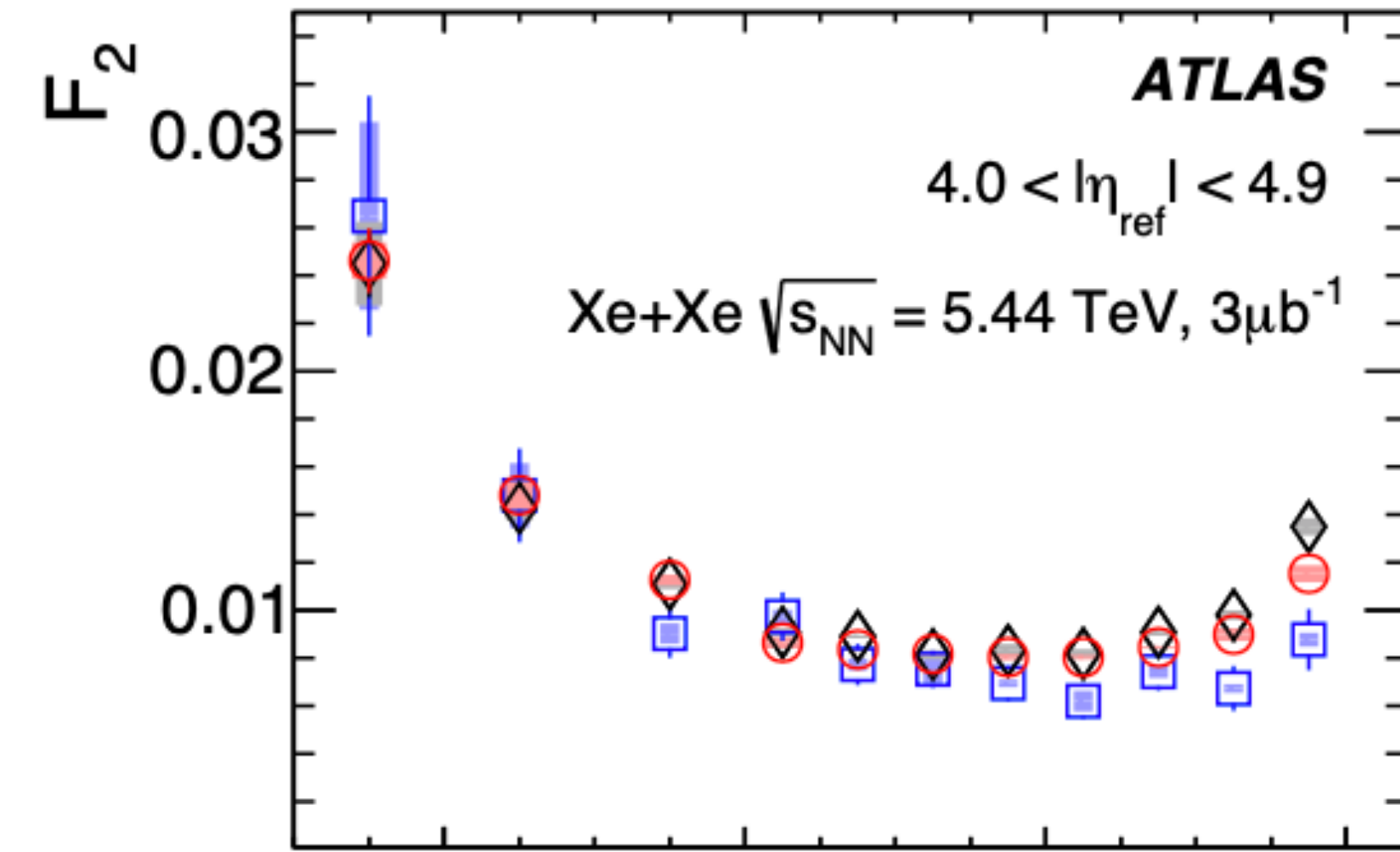
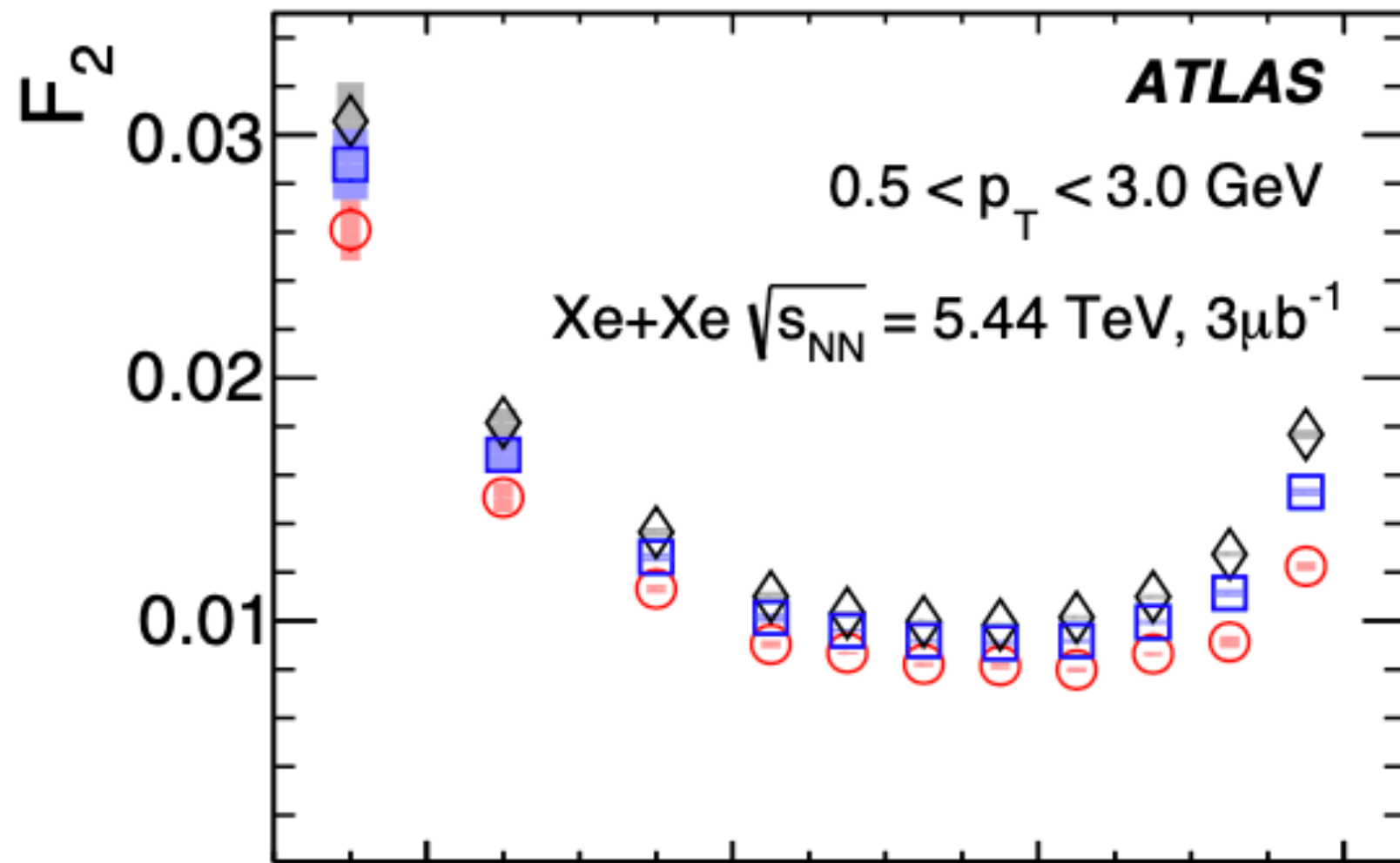


Slope $F_n - |\eta_{ref}|$ and p_T dependence

- Non-flow (dijets) at small $|\eta_{ref}|$ or large p_T

$$r_{n|n}(\eta) = \frac{\langle \mathbf{q}_n(-\eta) \mathbf{q}_n^*(\eta_{ref}) \rangle}{\langle \mathbf{q}_n(\eta) \mathbf{q}_n^*(\eta_{ref}) \rangle} \uparrow \text{NF} \Rightarrow F_n \uparrow$$

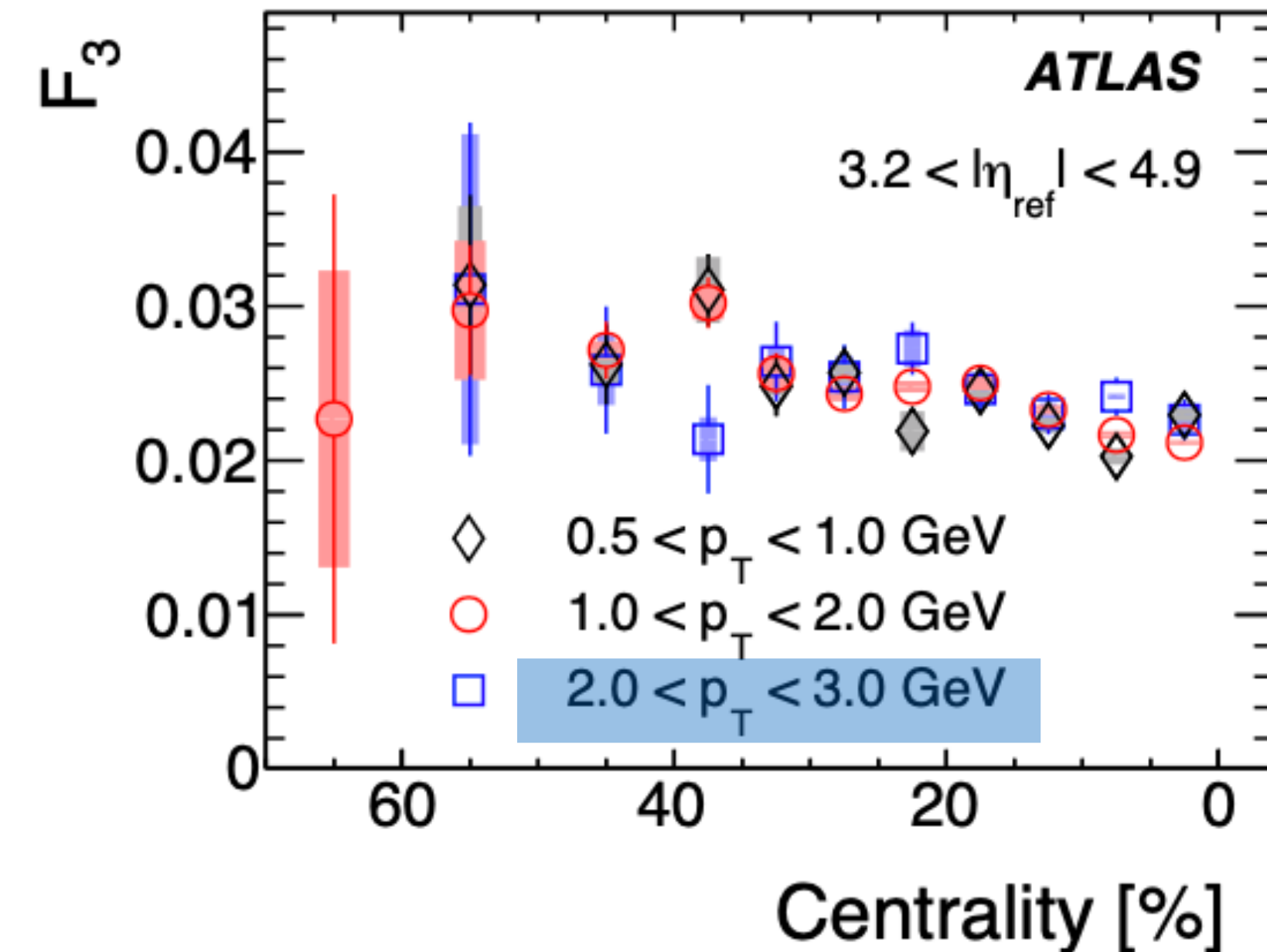
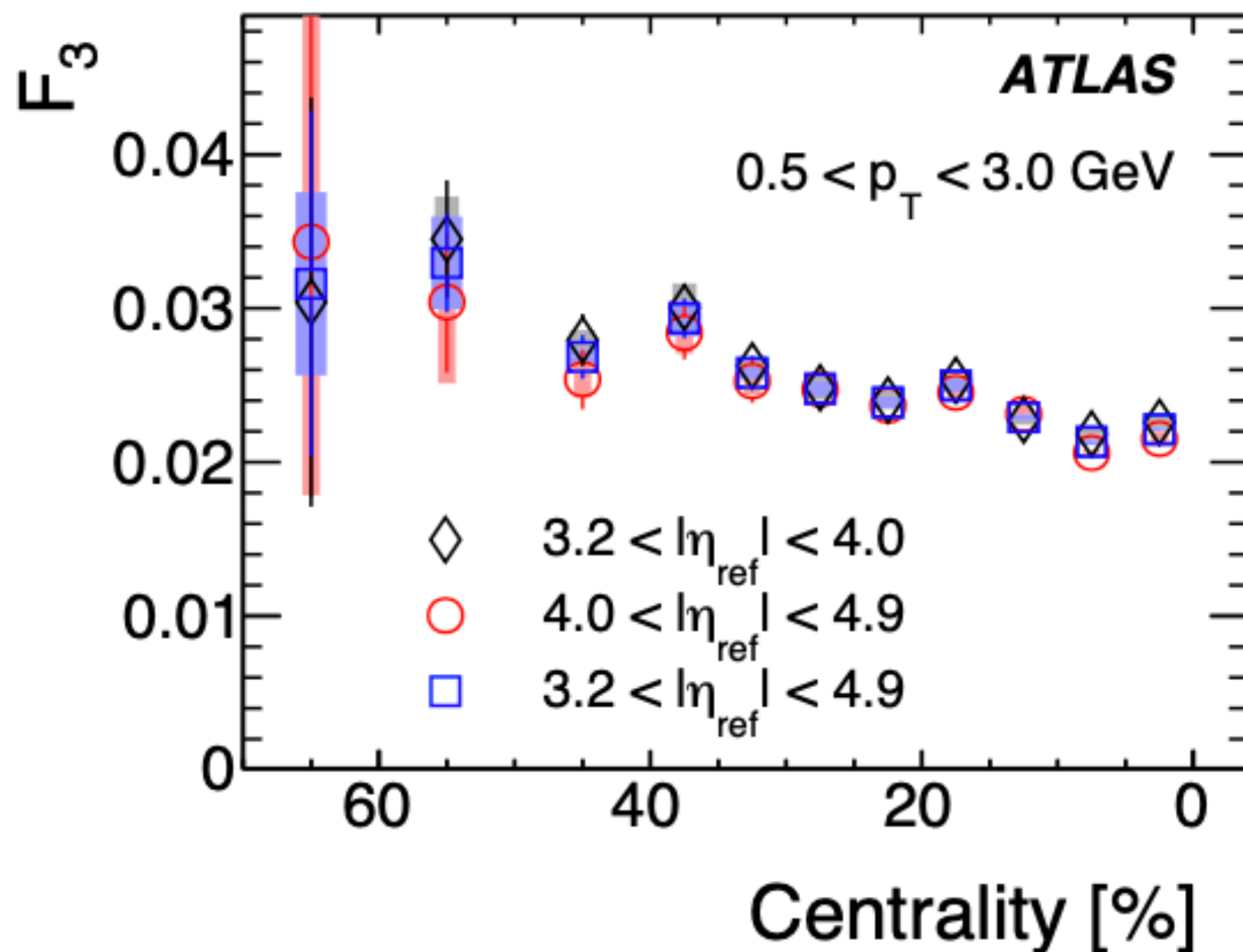
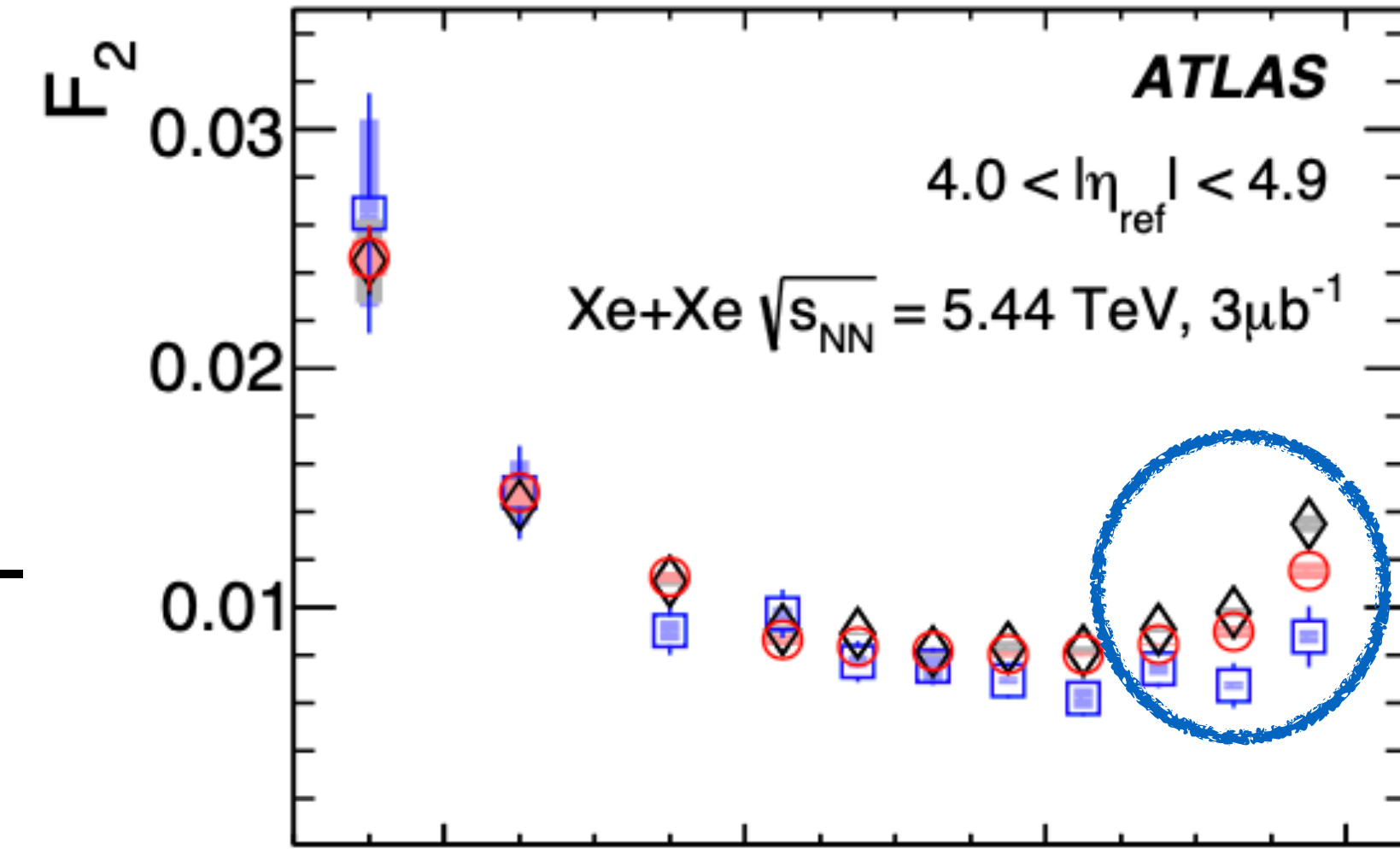
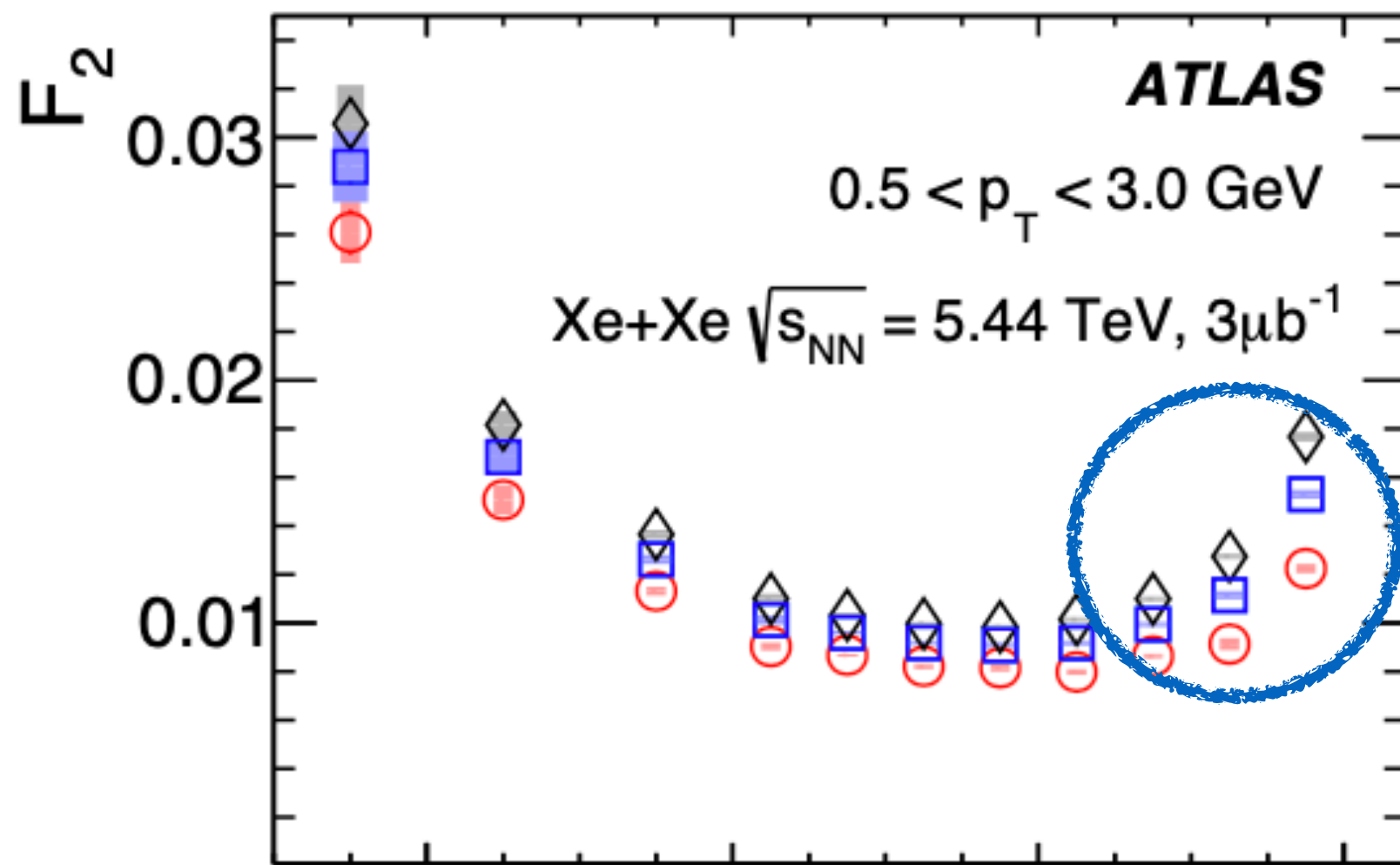
- Larger F_2 at smaller $|\eta_{ref}|$ - NF



Slope F_n - $|\eta_{ref}|$ and p_T dependence

- Non-flow (dijets) at small $|\eta_{ref}|$ or large p_T

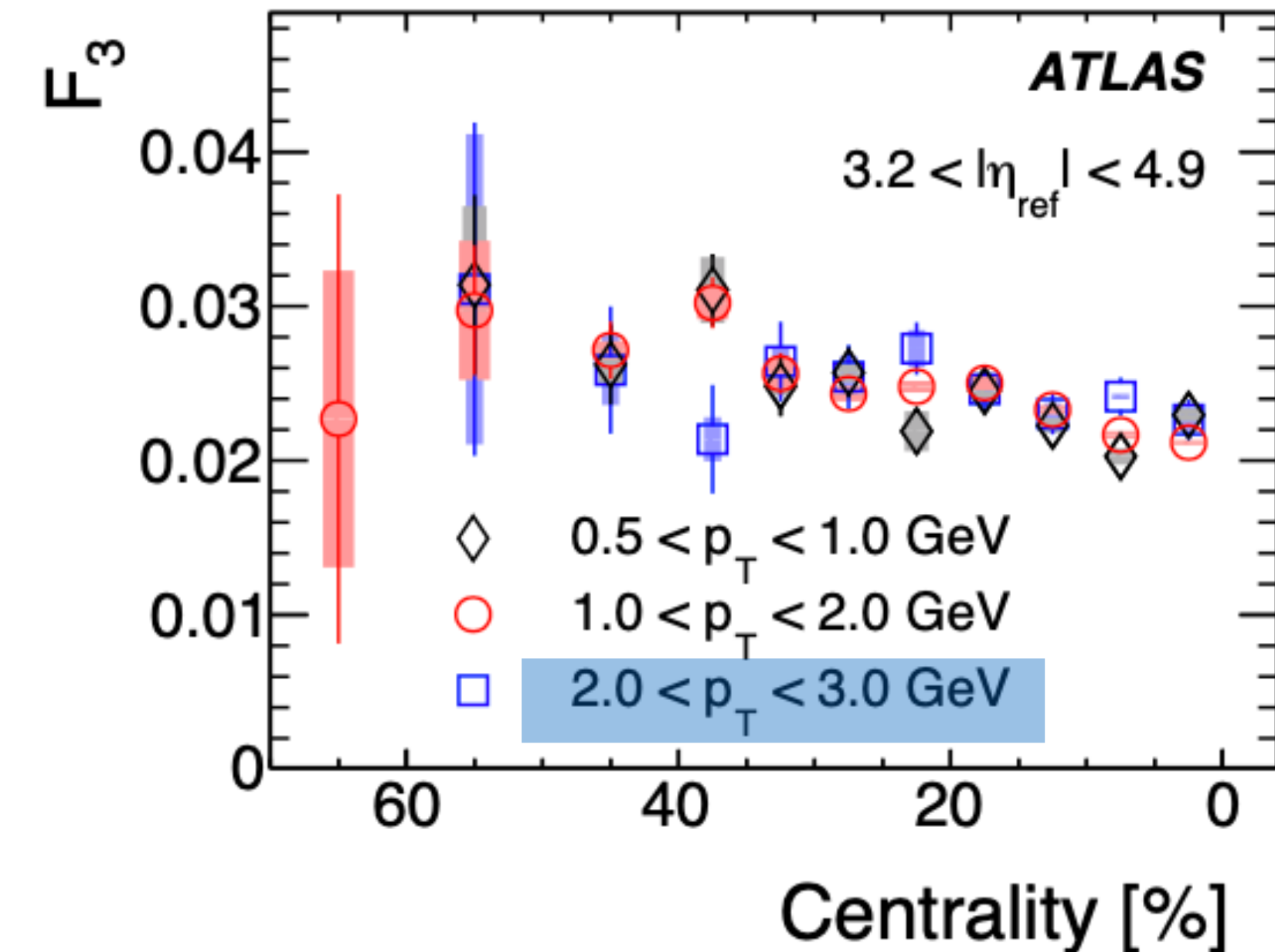
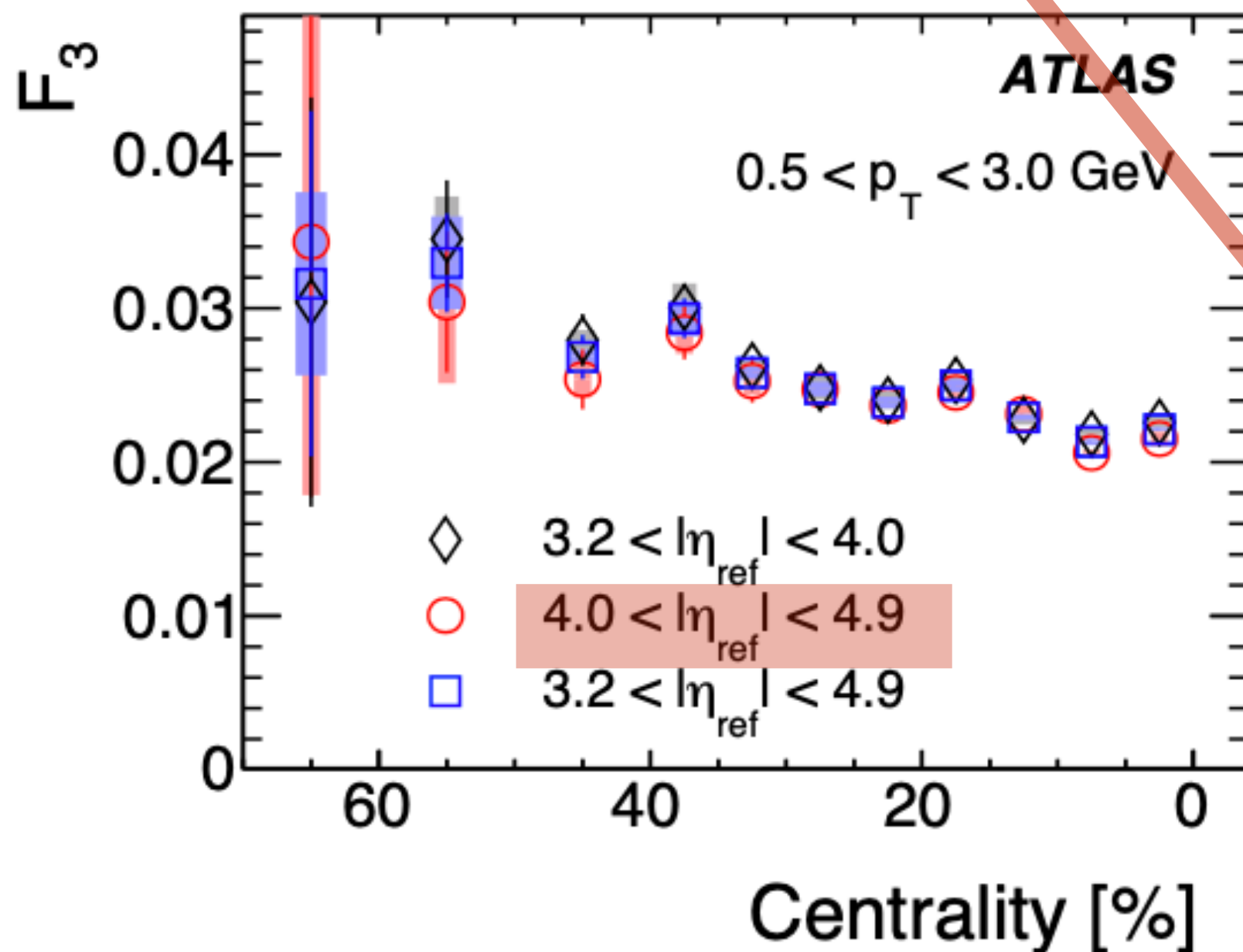
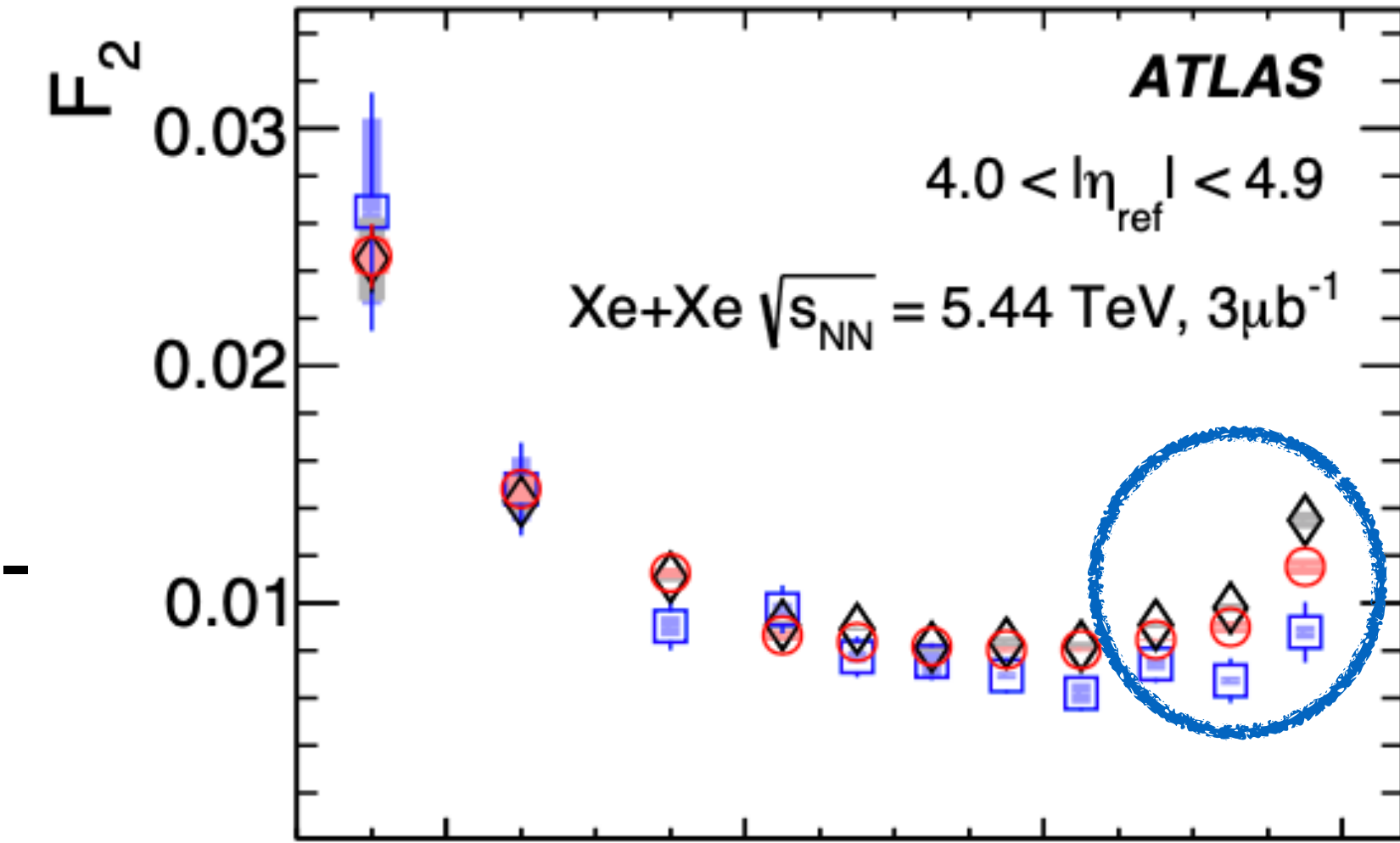
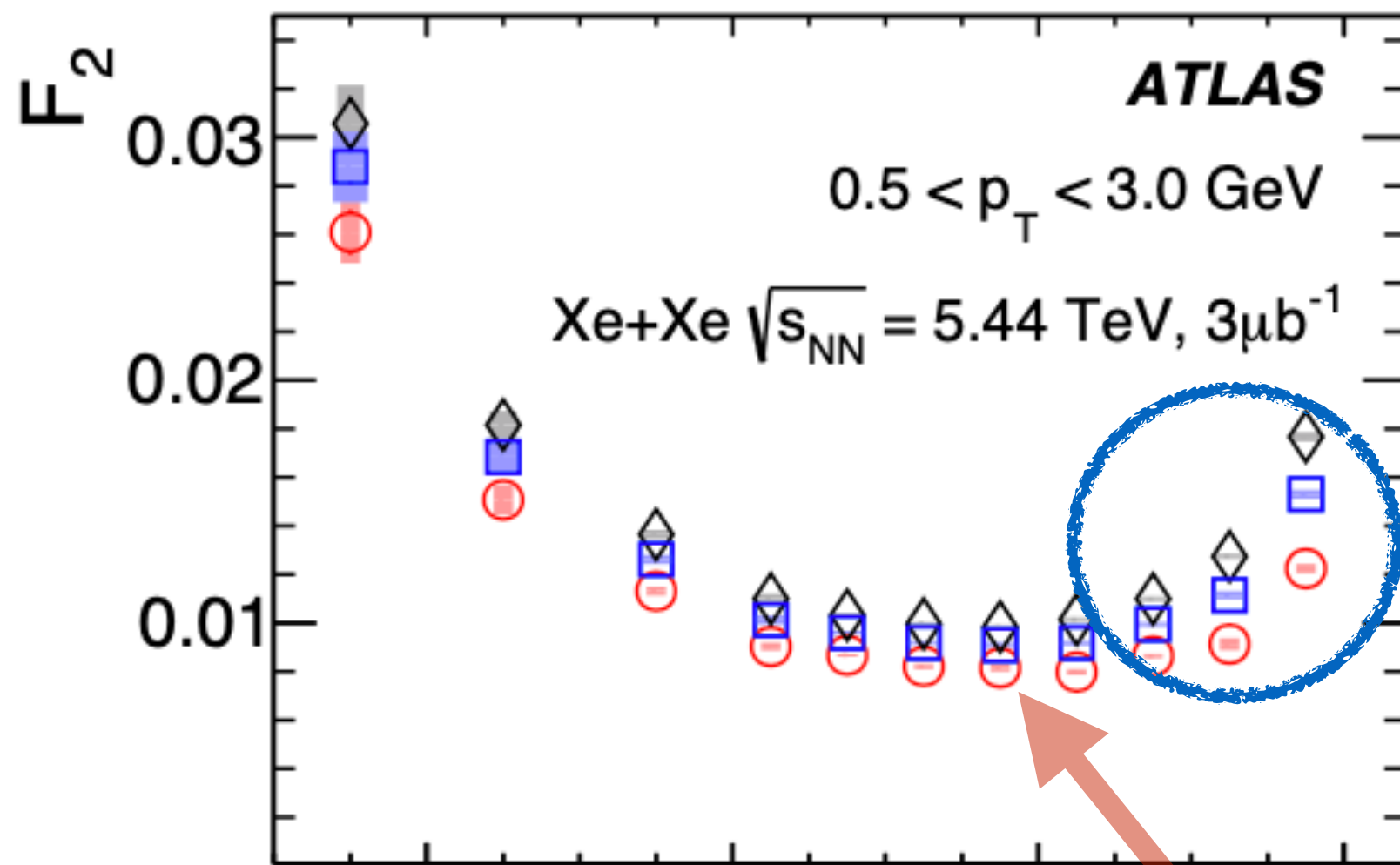
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- Larger F_2 at smaller $|\eta_{ref}|$ - NF
- Smaller F_2 at larger p_T - not NF
- Effect larger in central collisions - could be due to non-linear v_2 decorrelation
- F_3 - not affected by NF

Slope F_n - $|\eta_{ref}|$ and p_T dependence

- Non-flow (dijets) at small $|\eta_{ref}|$ or large p_T $r_{n|n}(\eta) = \frac{\langle \mathbf{q}_n(-\eta) \mathbf{q}_n^*(\eta_{ref}) \rangle}{\langle \mathbf{q}_n(\eta) \mathbf{q}_n^*(\eta_{ref}) \rangle} \uparrow \text{NF} \Rightarrow F_n \uparrow$

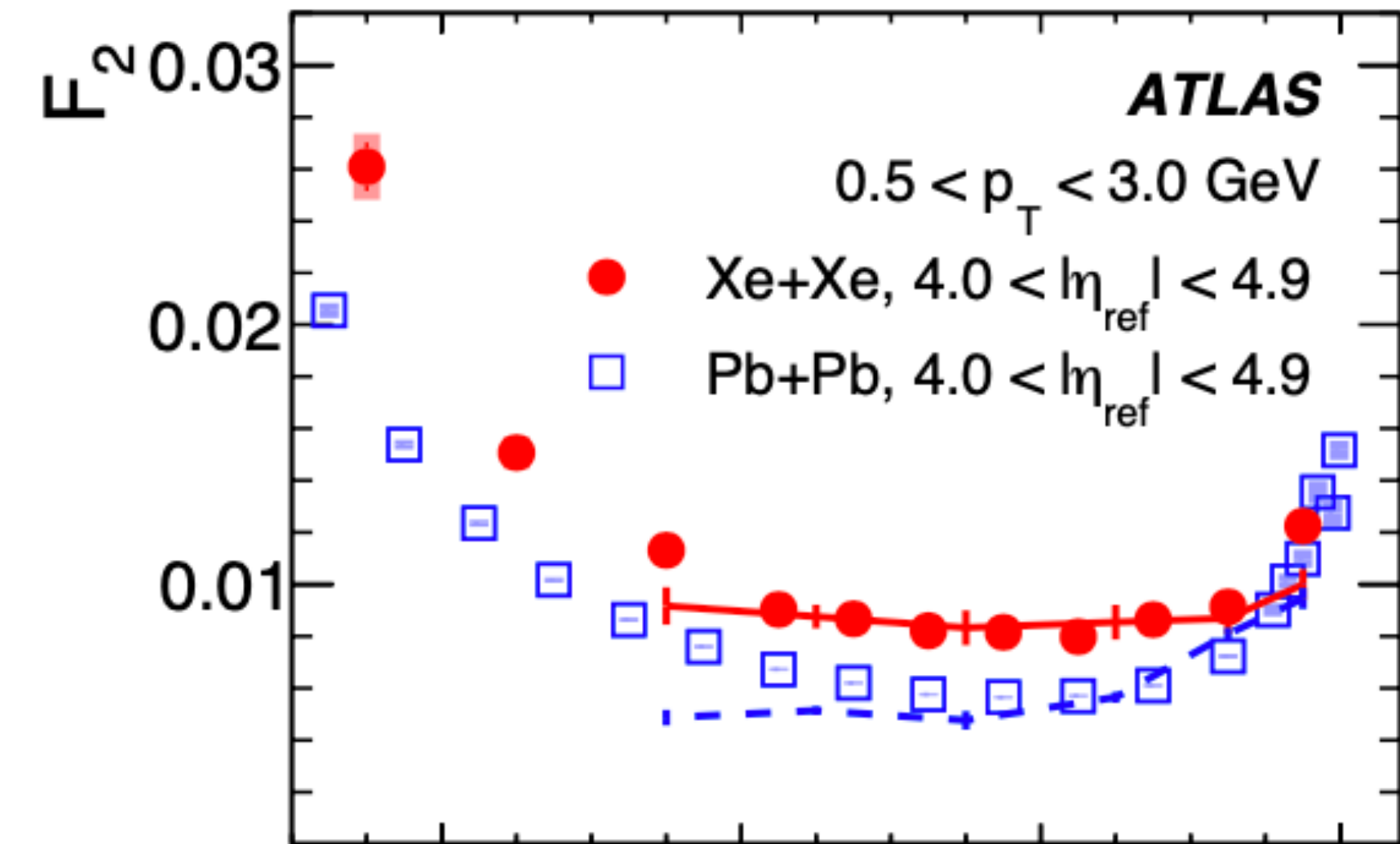


- Larger F_2 at smaller $|\eta_{ref}|$ - NF
- Smaller F_2 at larger p_T - not NF
- Effect larger in central collisions - could be due to non-linear v_2 decorrelation
- F_3 - not affected by NF
- Suppress these variations

$0.5 < p_T < 3.0 \text{ GeV}$
 $F_2 \longrightarrow 4.0 < |\eta_{ref}| < 4.9$
 $F_3 \longrightarrow 3.2 < |\eta_{ref}| < 4.9$

Comparison with Pb+Pb and Hydro

Centrality (avg. geometry)

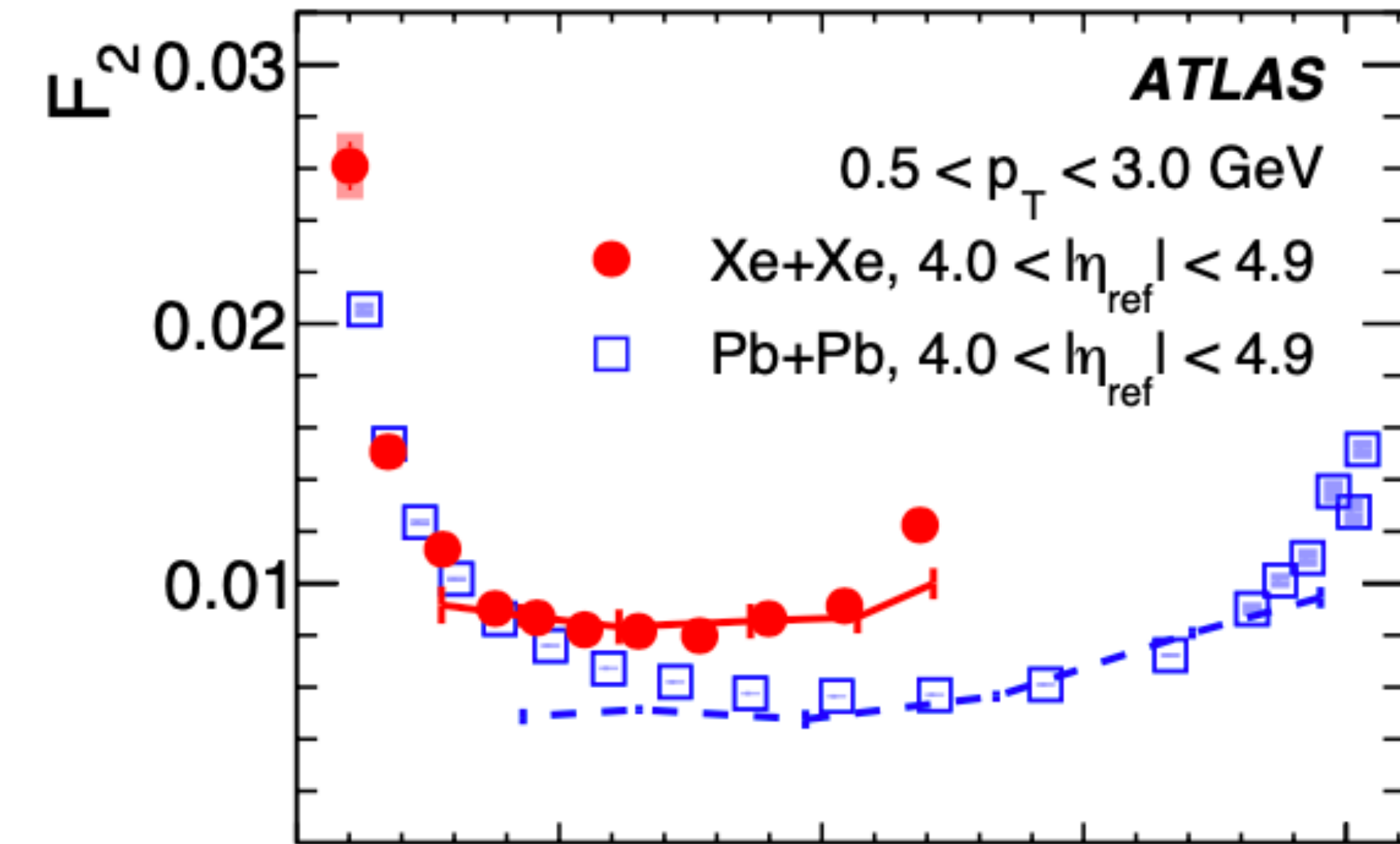


- Reverse ordering for $n=2$ and 3

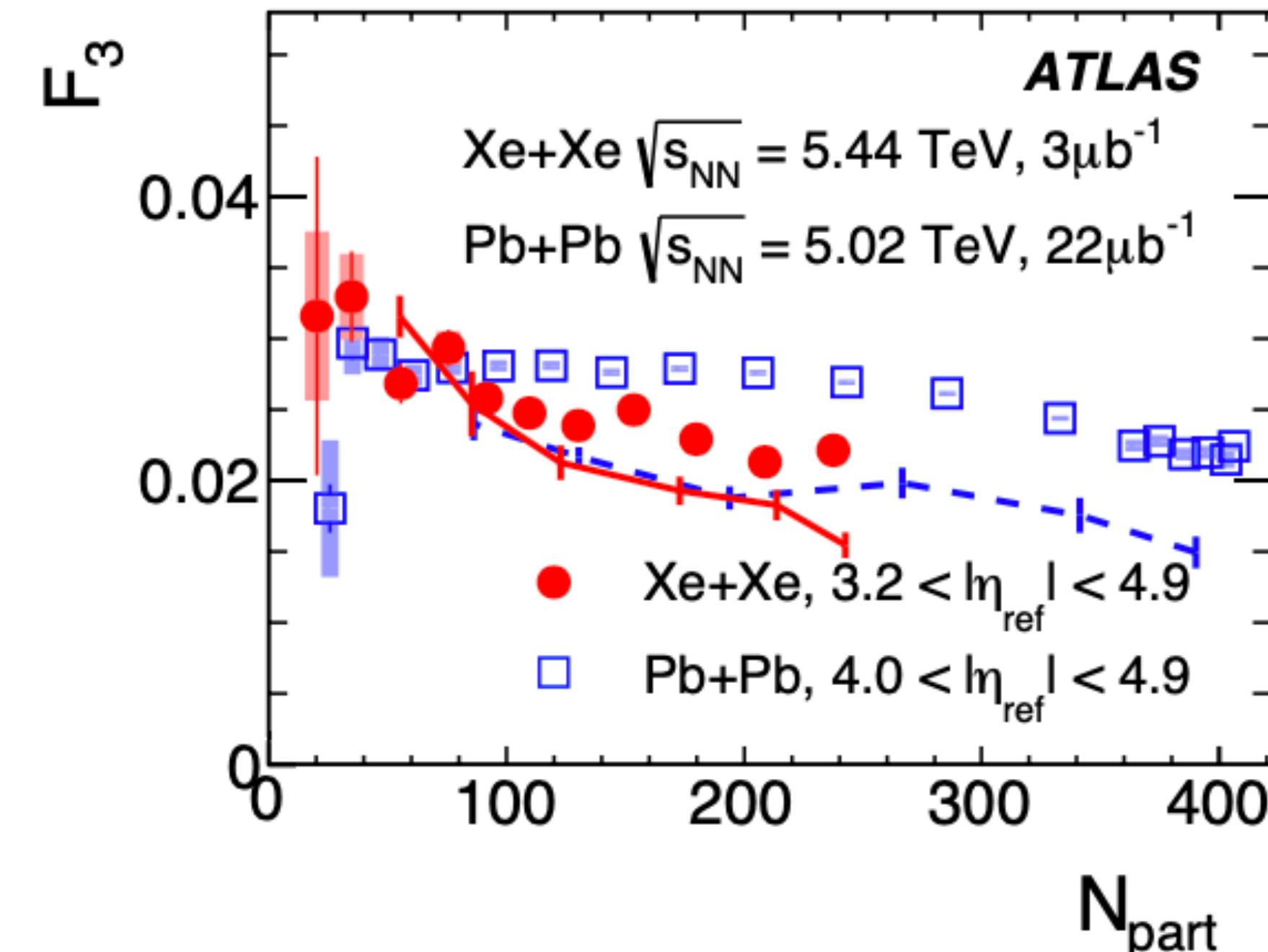
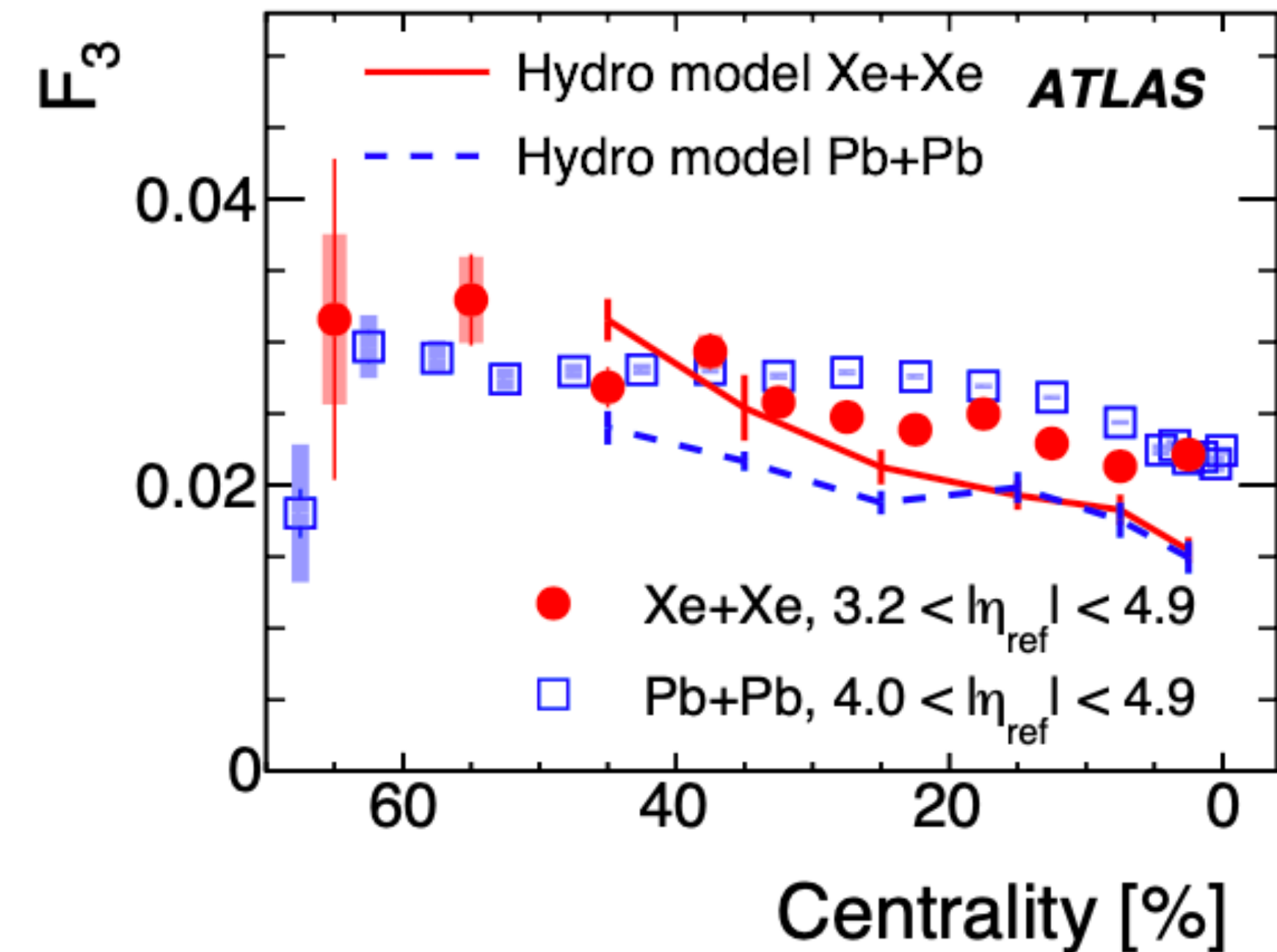
$$F_2^{XeXe} > F_2^{PbPb}$$

$$F_3^{XeXe} < F_3^{PbPb}$$

N_{part} (system-size)

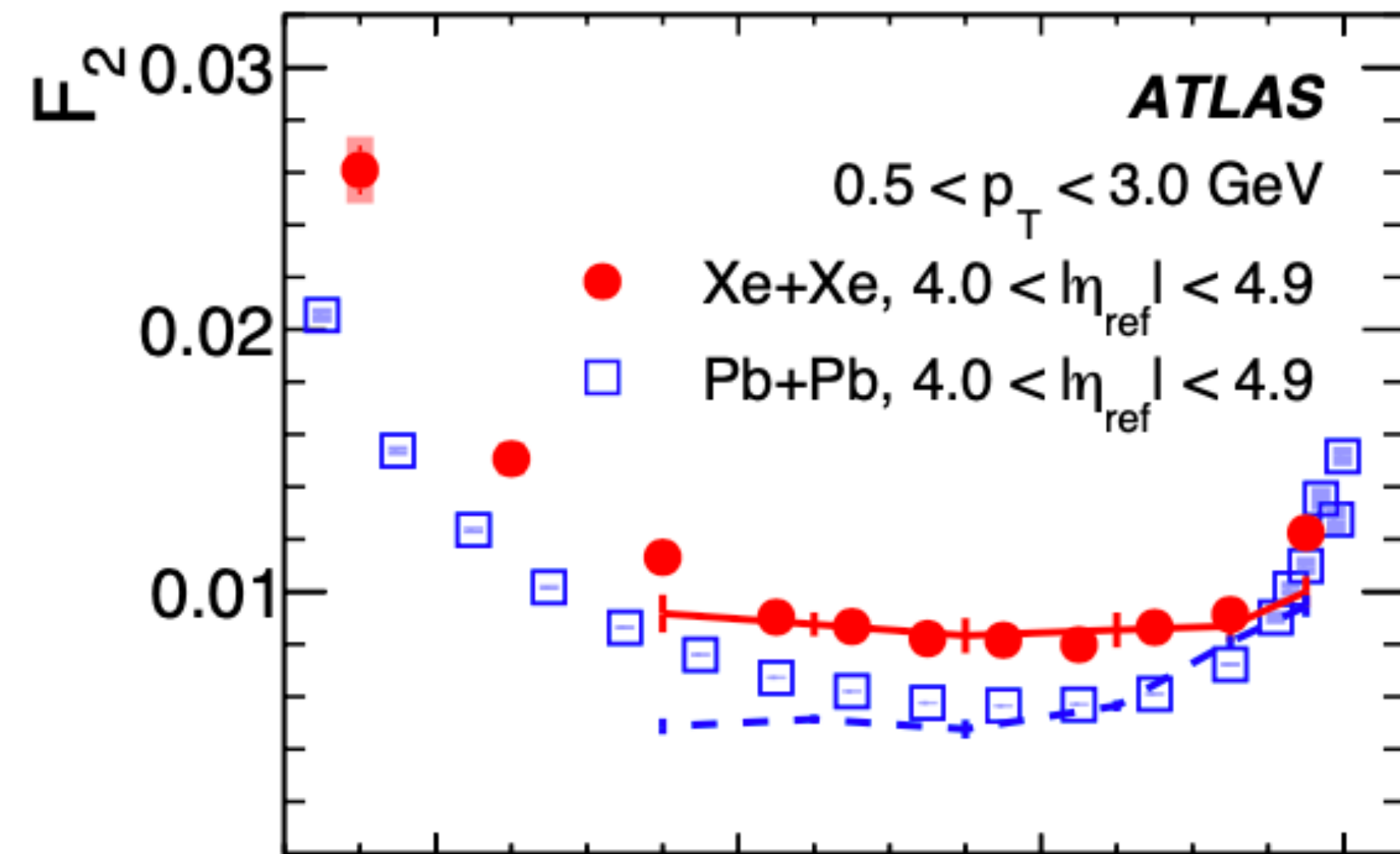


- F_2 and F_3 match - in 0-5% centrality or $N_{\text{part}} < 80$



Comparison with Pb+Pb and Hydro

Centrality (avg. geometry)

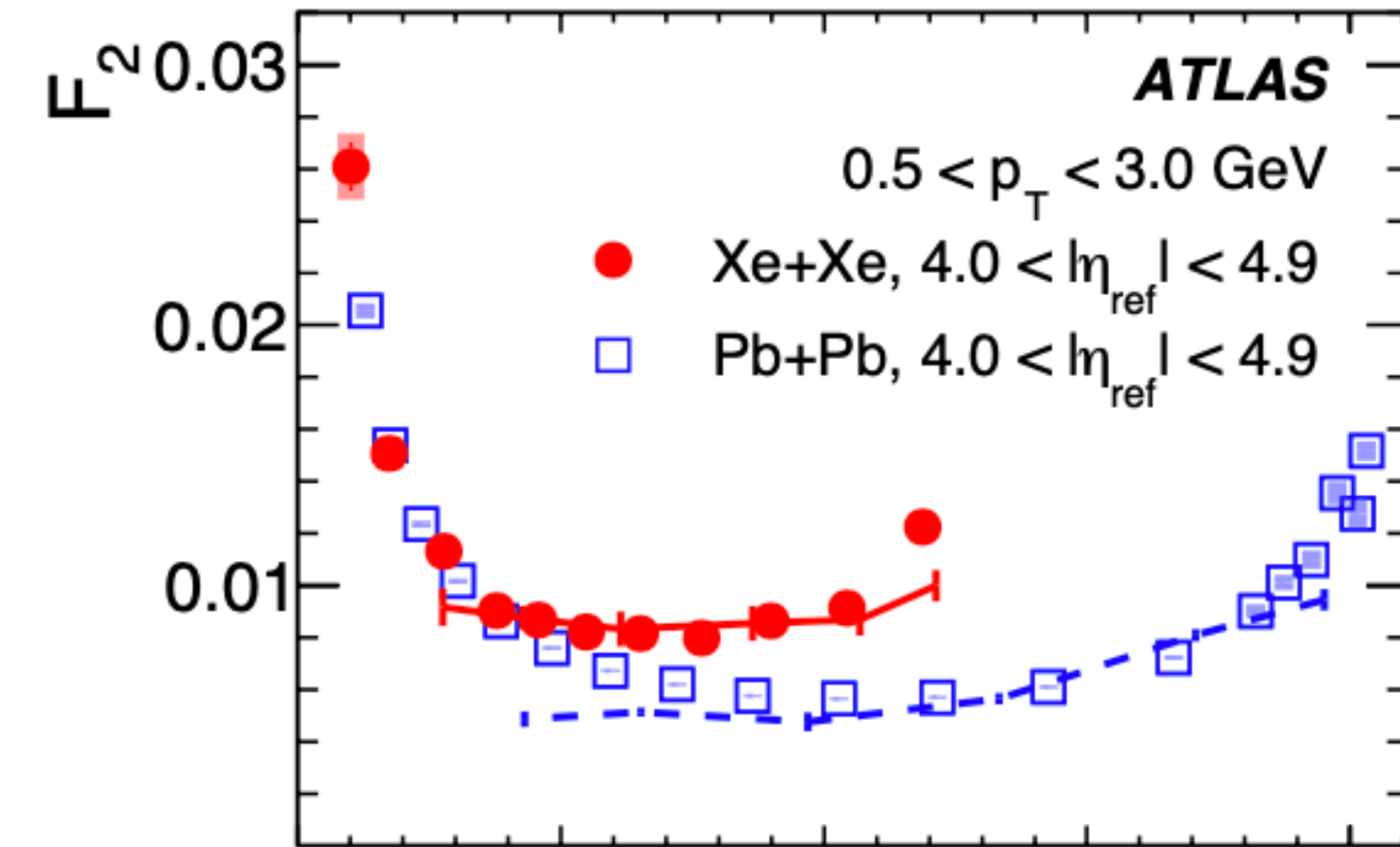


- Reverse ordering for $n=2$ and 3

$$F_2^{XeXe} > F_2^{PbPb}$$

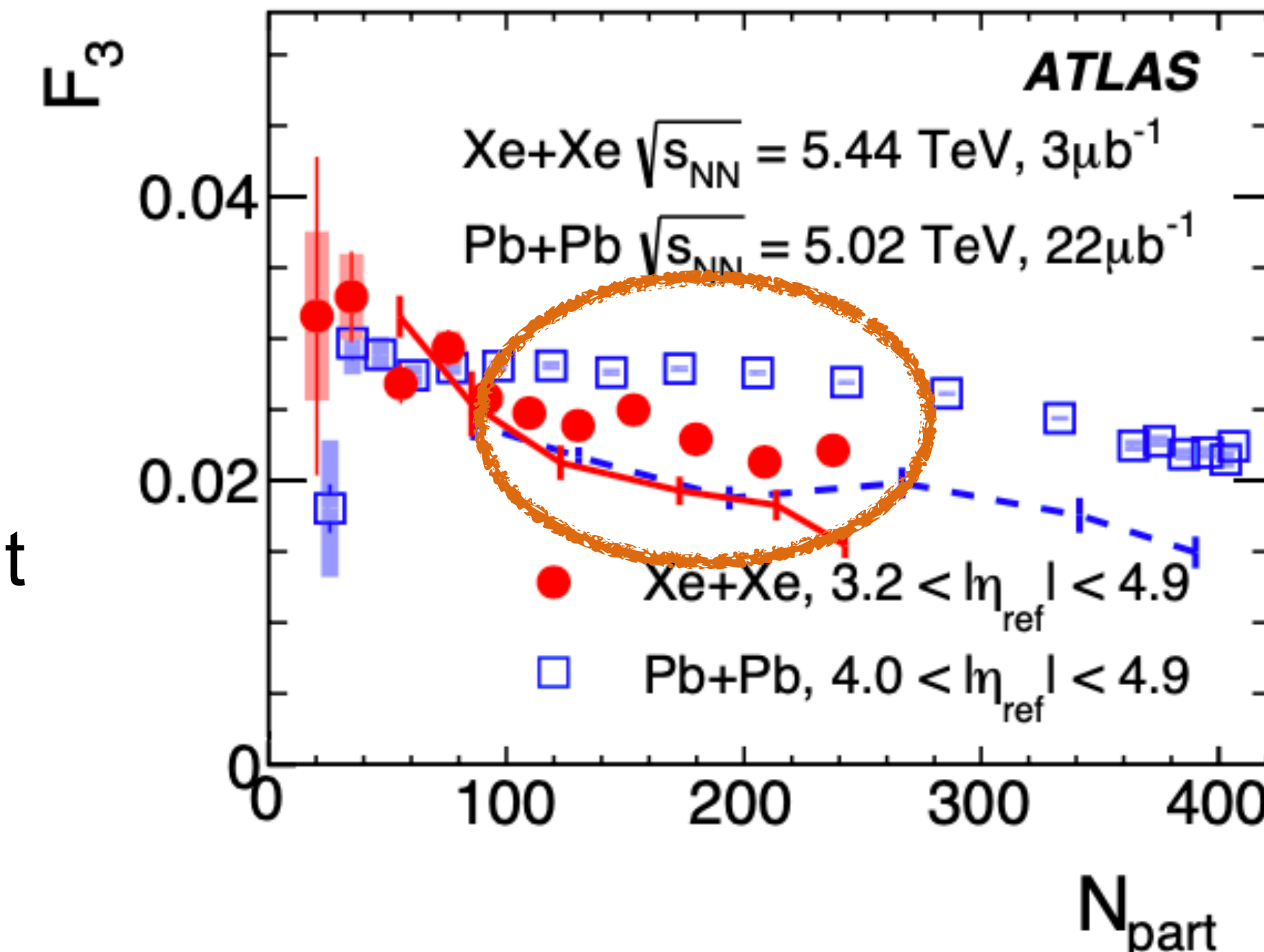
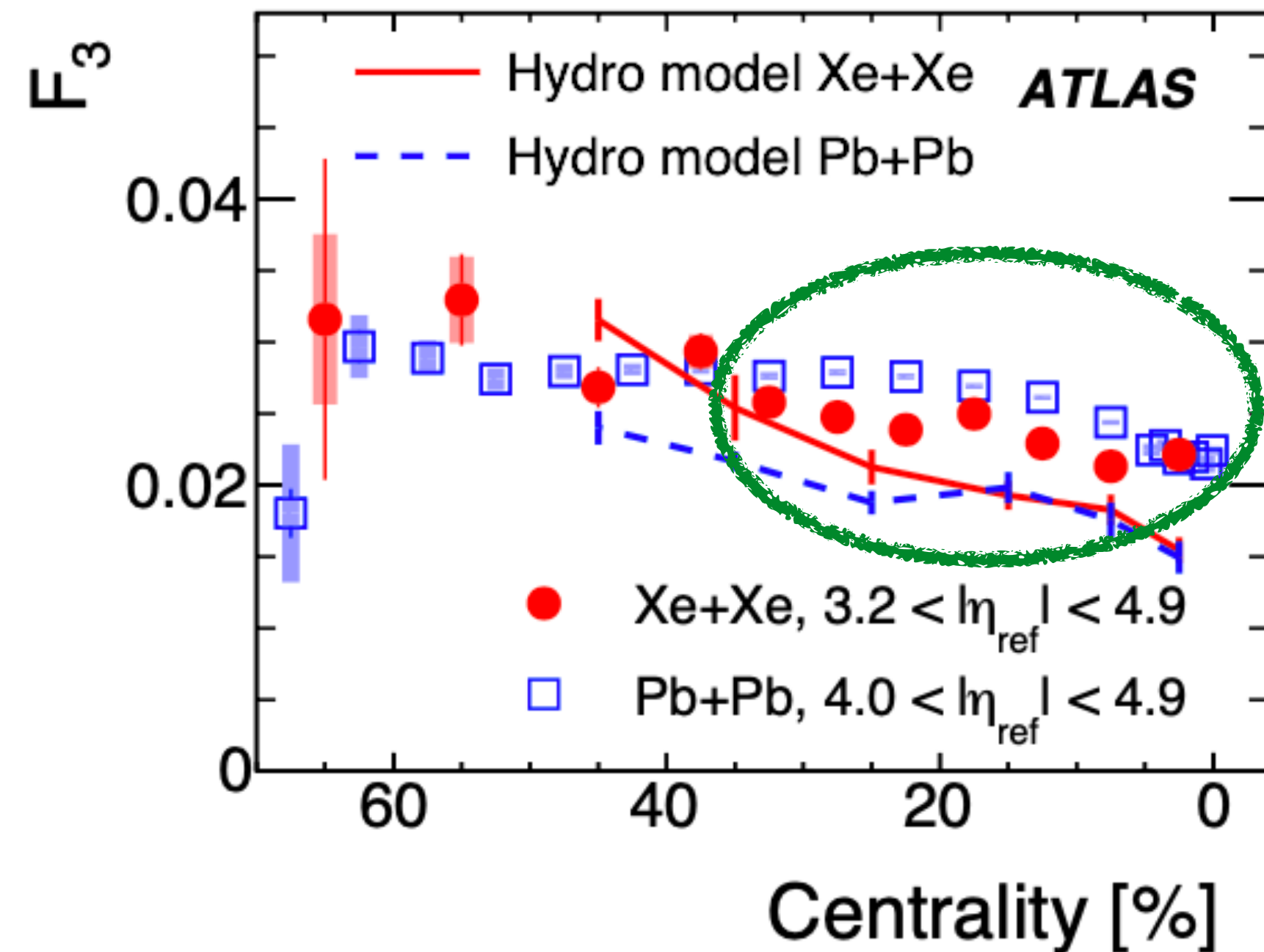
$$F_3^{XeXe} < F_3^{PbPb}$$

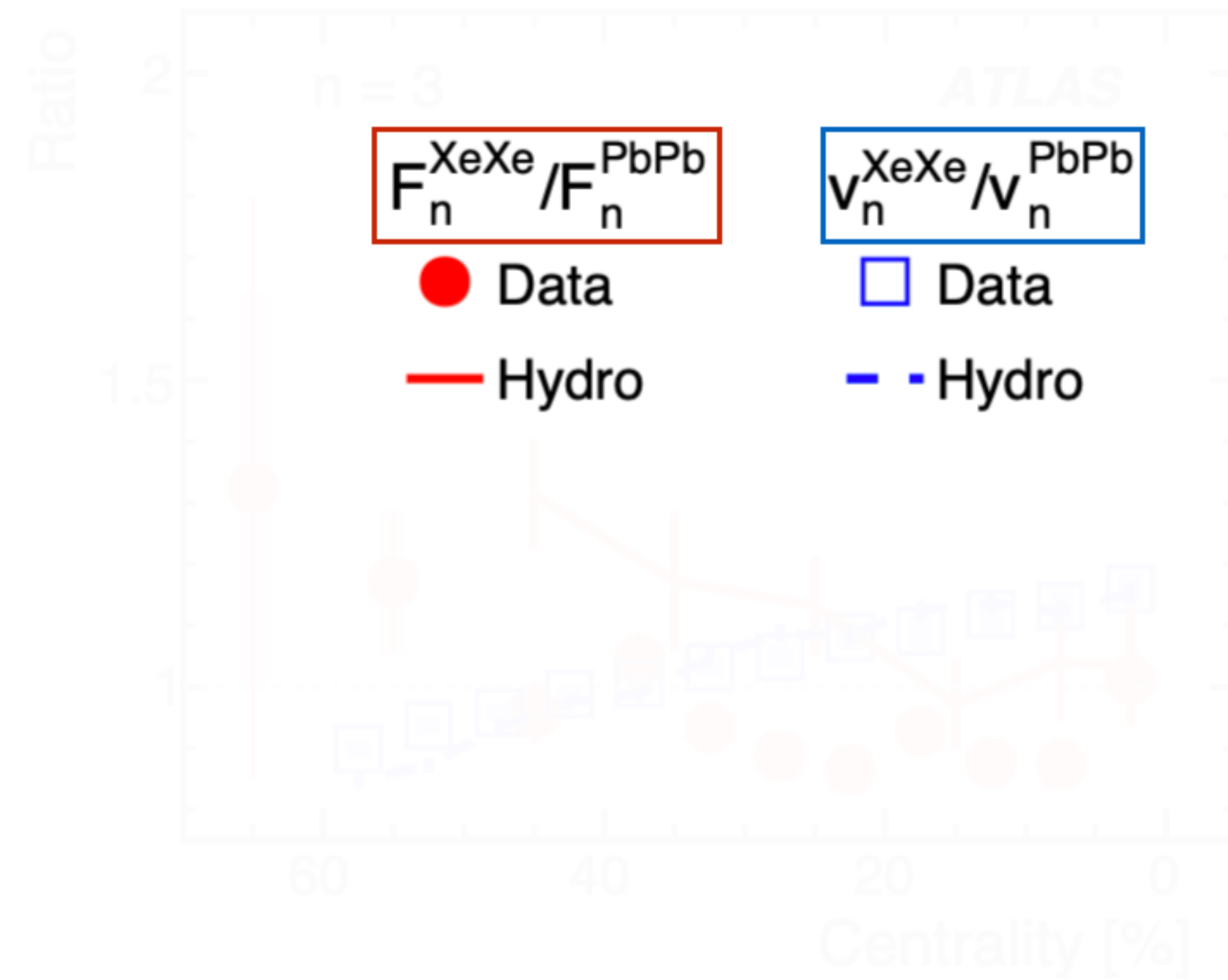
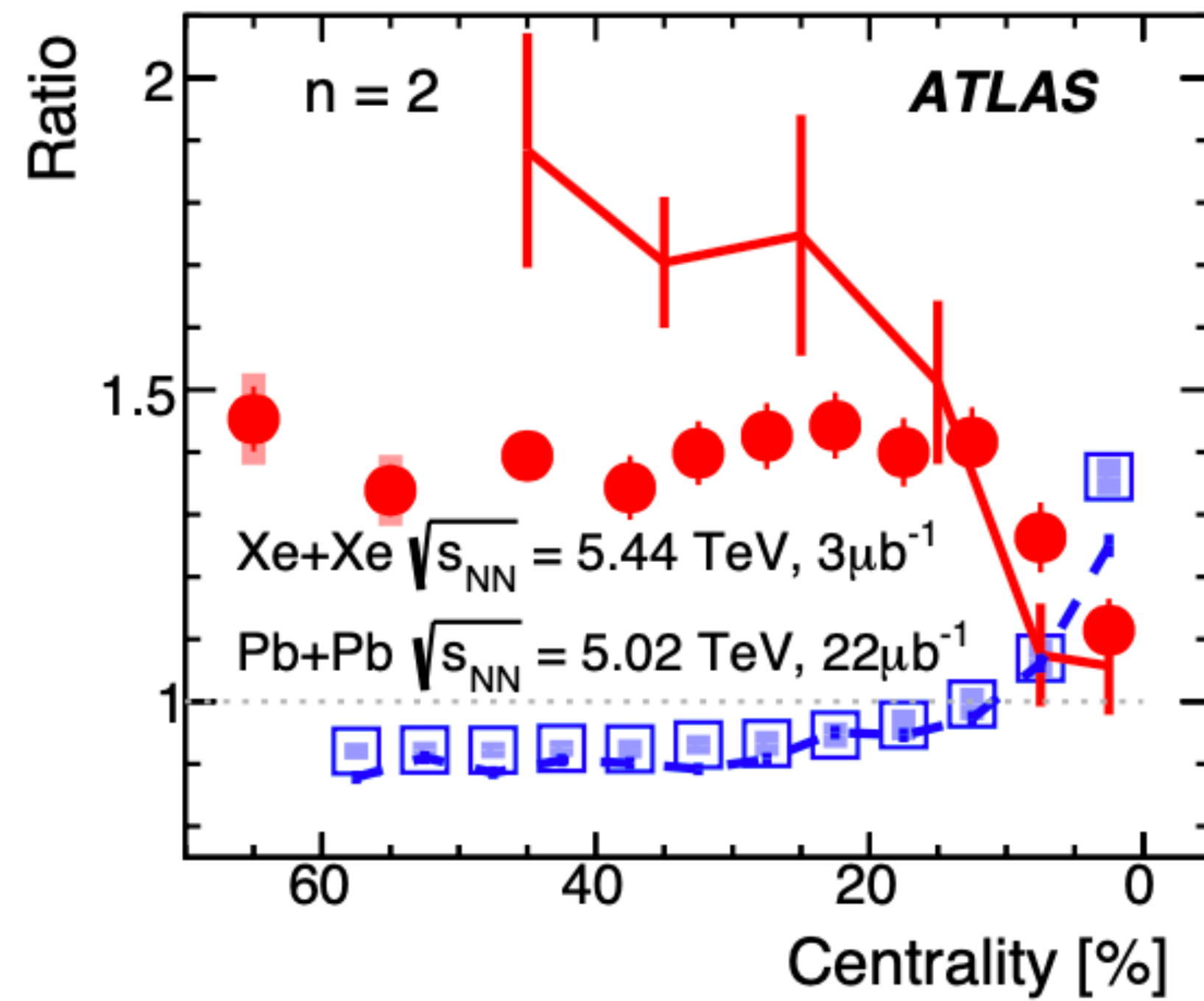
N_{part} (system-size)



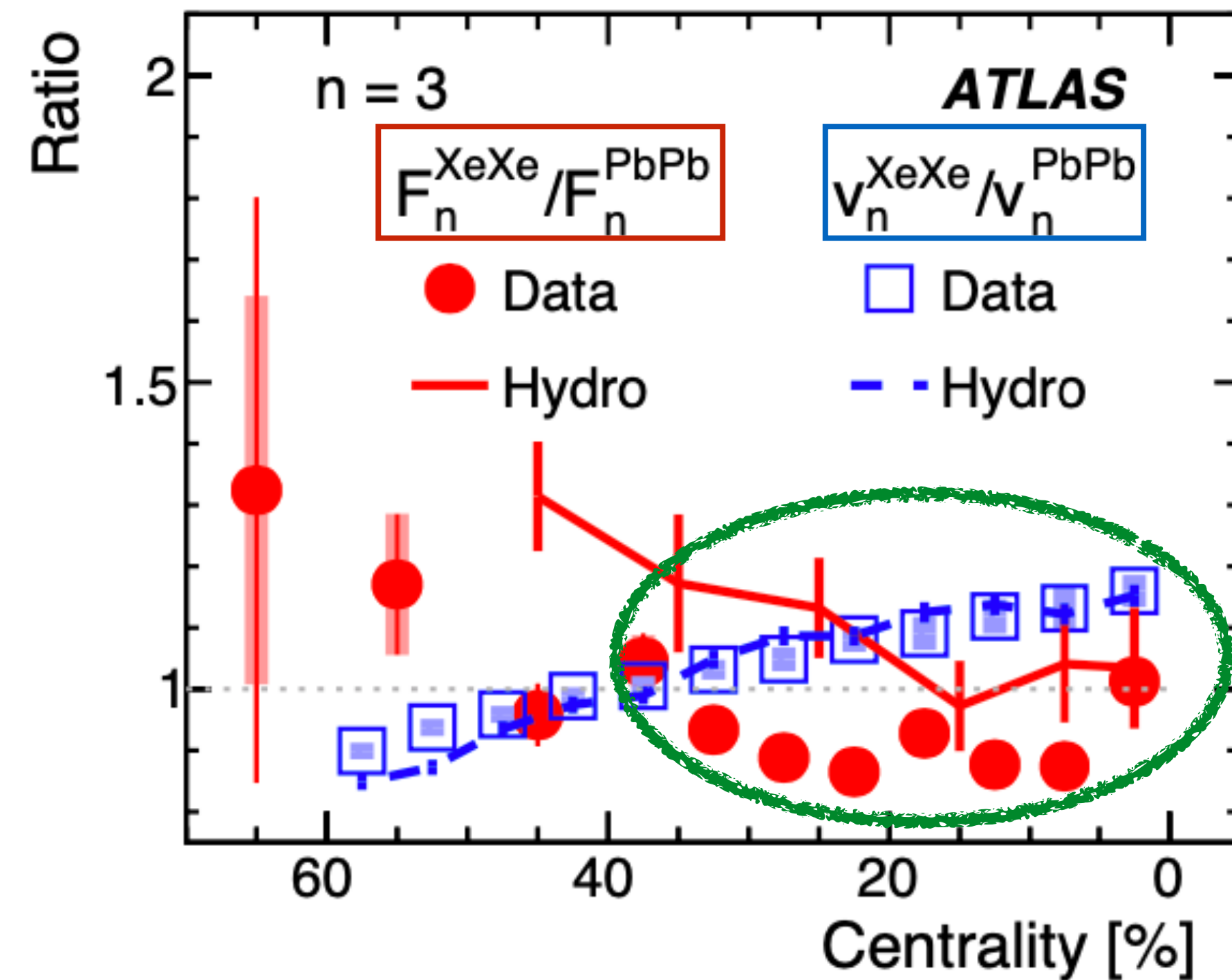
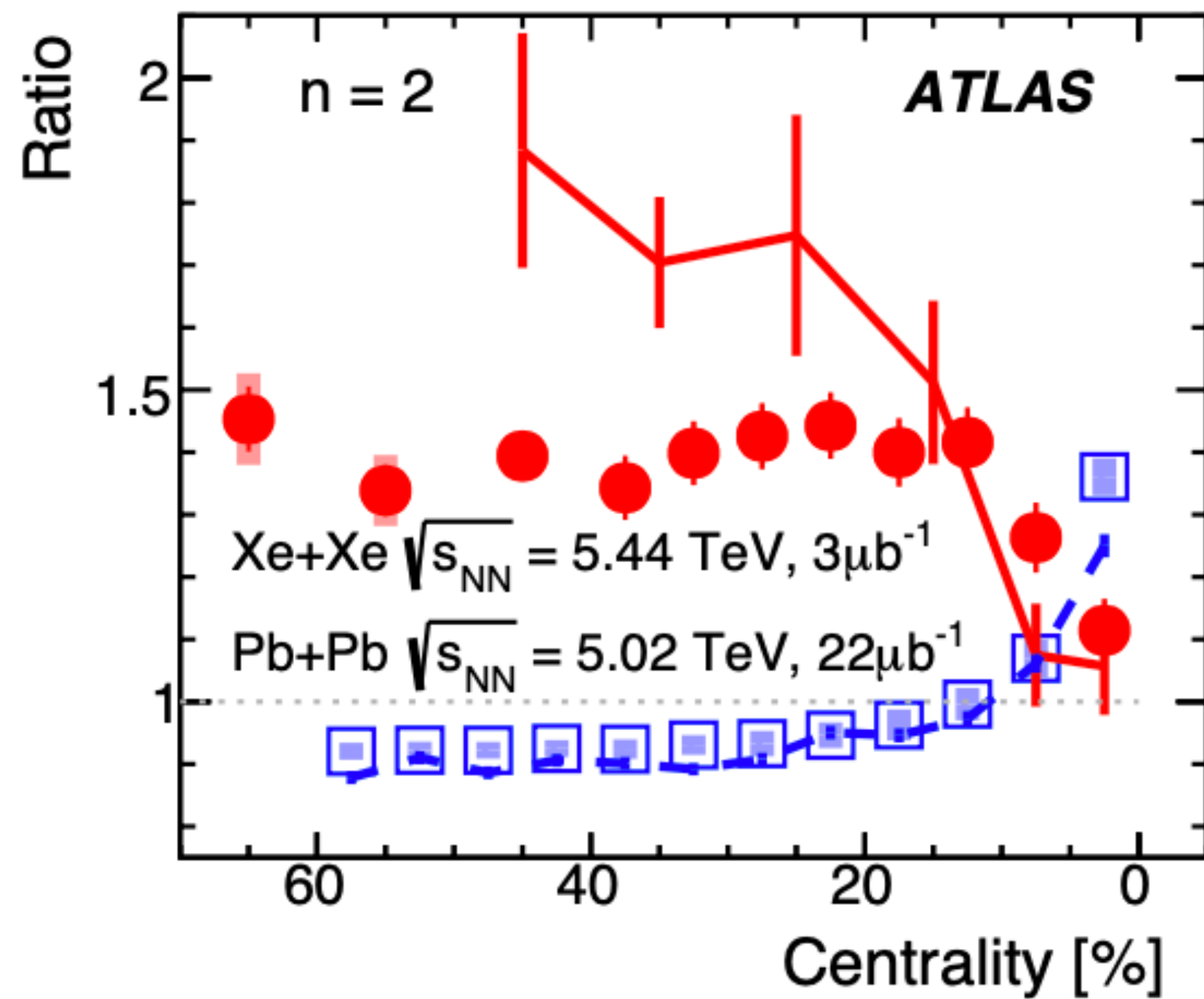
- F_2 and F_3 match - in 0-5% centrality or $N_{part} < 80$

- Hydro model works for F_2 but fails for F_3 in magnitude and splitting





- Anti-correlation between v_n and F_n - opposite centrality dependence



- Anti-correlation between v_n and F_n - opposite centrality dependence

- Reverse ordering of F_2 and F_3 ratios

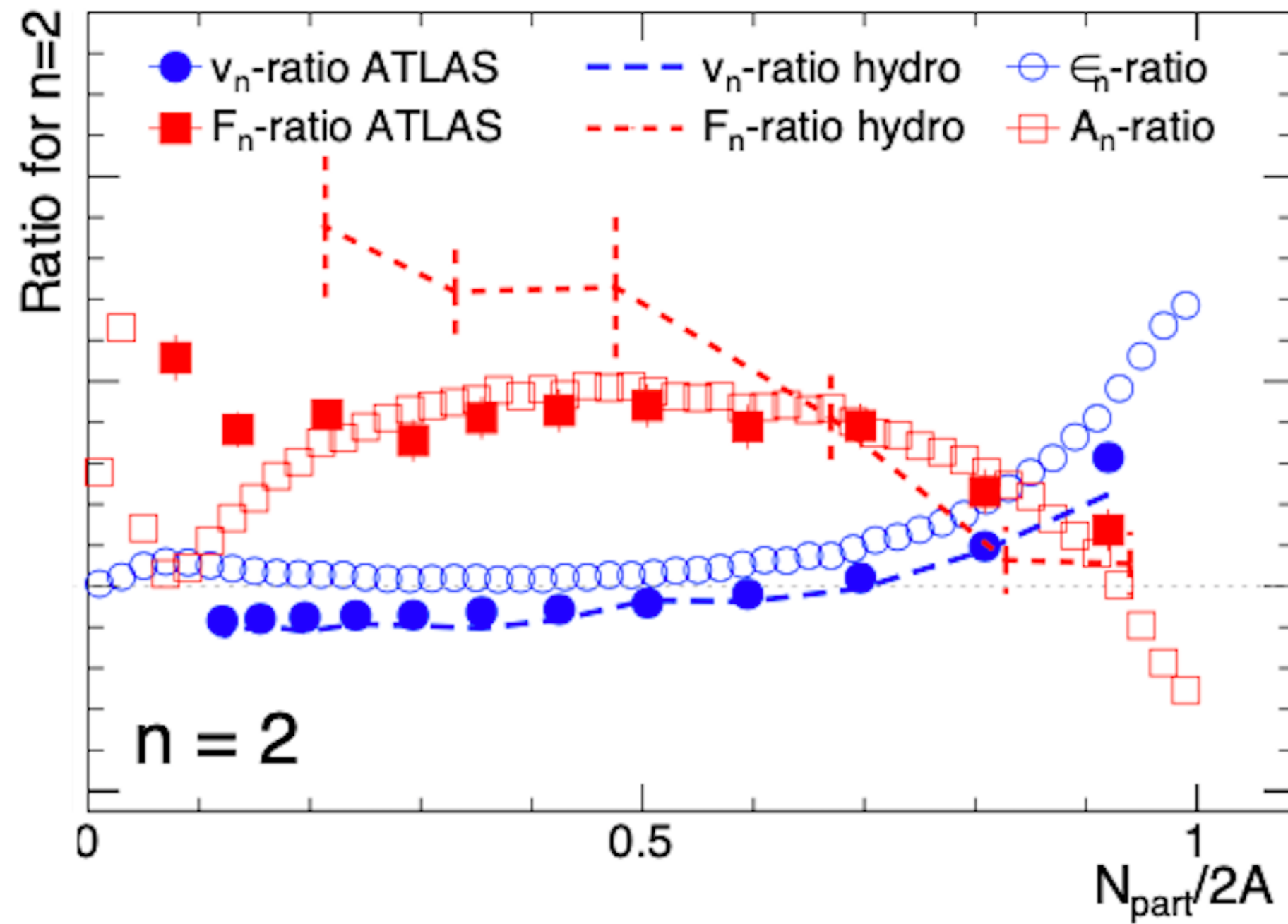
$$F_2^{\text{XeXe}} / F_2^{\text{PbPb}} > 1 \quad F_3^{\text{XeXe}} / F_3^{\text{PbPb}} < 1$$

- Hydro explains v_n ratio very well but fails to explain F_n ratio accurately

- Whats the origin of the reverse ordering ?

- Eccentricity ϵ_n and eccentricity decorrelation A_n from subnucleon-Glauber model

arXiv:2003.04340

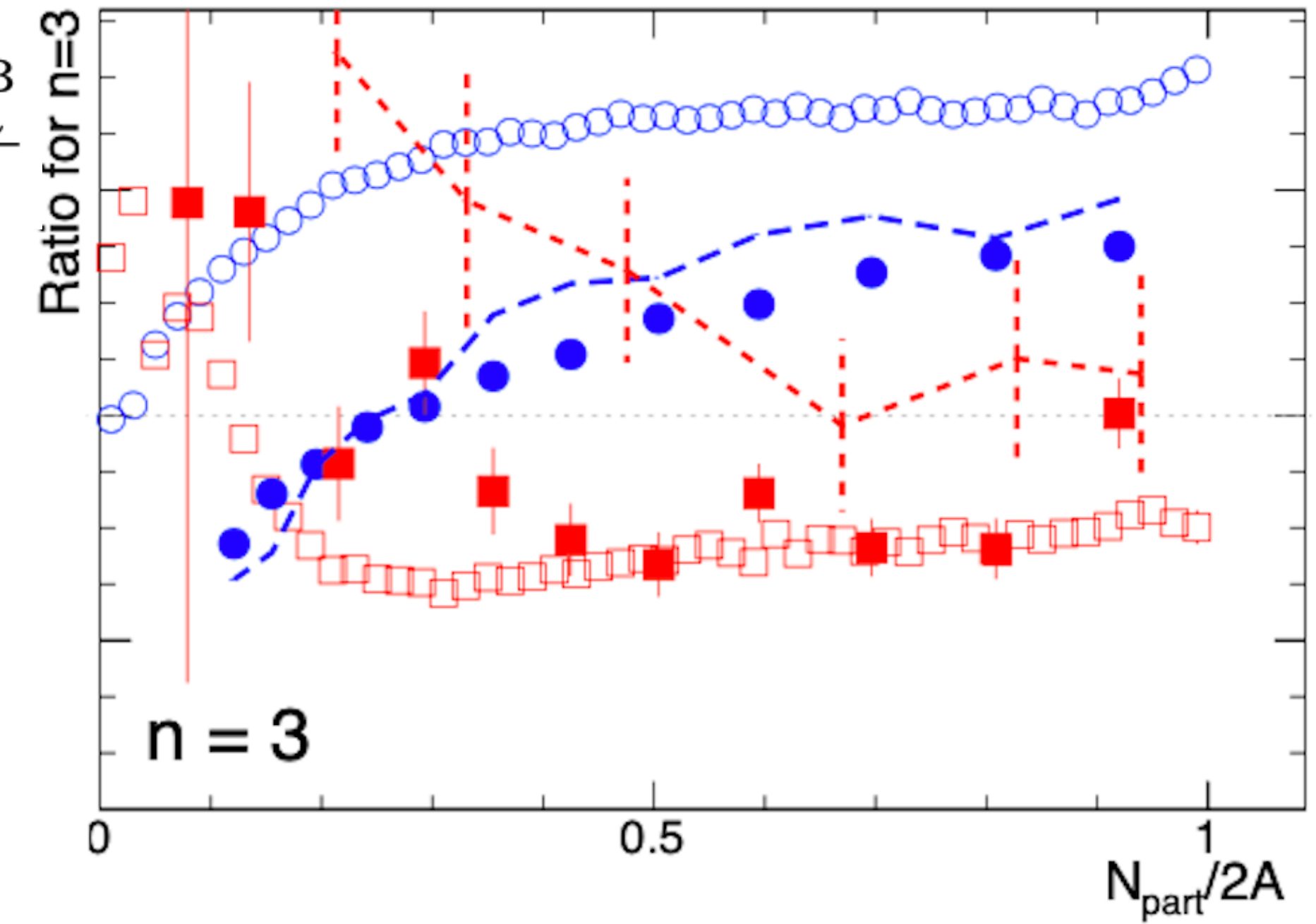


$$\epsilon_{n+} = \frac{\epsilon_n^F + \epsilon_n^B}{2} \quad \epsilon_{n-} = \frac{\epsilon_n^F - \epsilon_n^B}{2}$$

$$A_n \equiv \frac{\langle \epsilon_{n-}^2 \rangle}{\langle \epsilon_{n+}^2 \rangle + \langle \epsilon_{n-}^2 \rangle} \approx \frac{\langle \epsilon_{n-}^2 \rangle}{\langle \epsilon_n^2 \rangle}$$

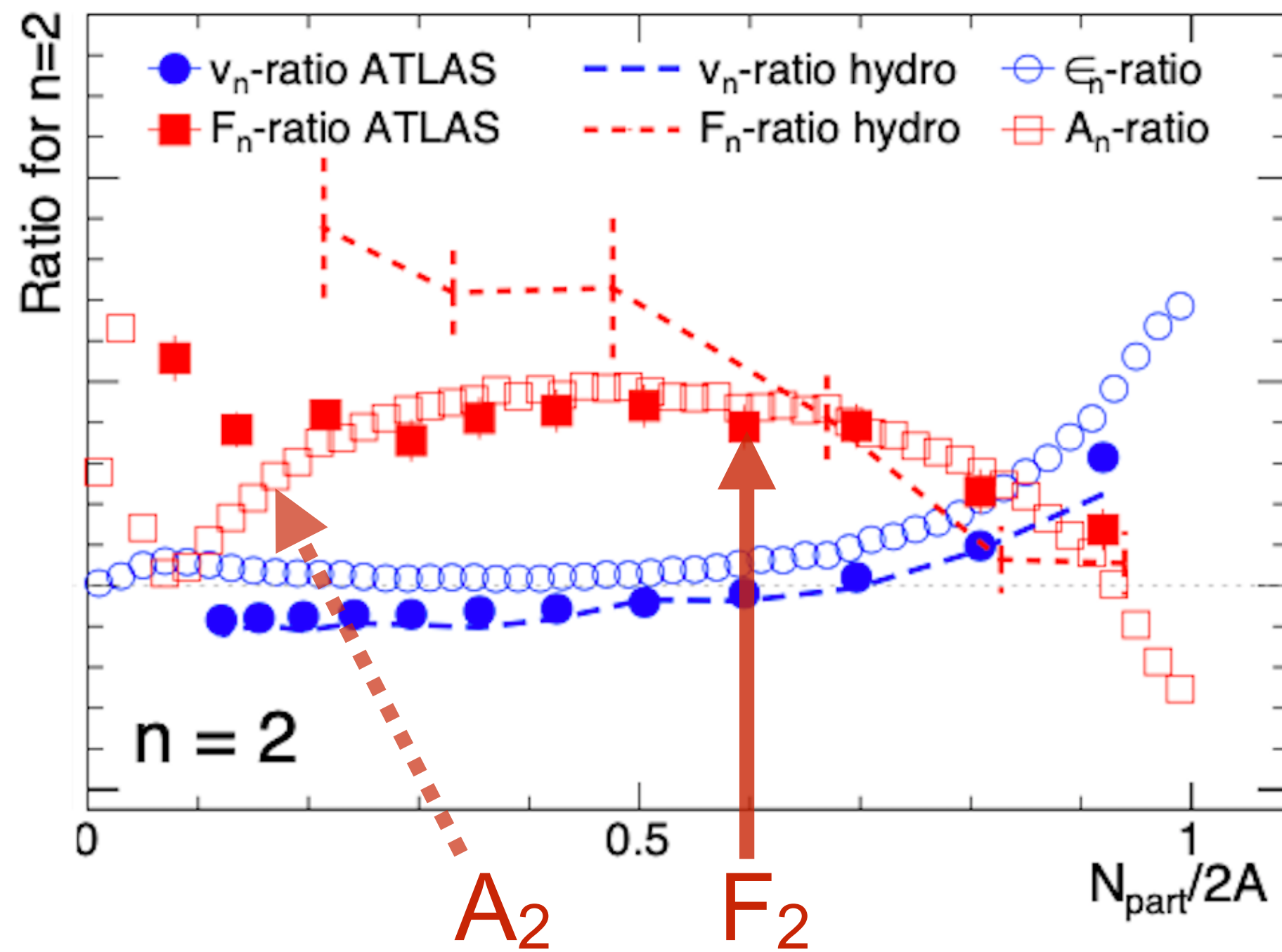
$$\epsilon_n \longrightarrow v_n$$

$$A_n \longrightarrow F_n$$

 $N_{\text{part}}/2A \longrightarrow \text{Centrality}$


- Difference between v_n and ϵ_n ratios - Late time effects

- Eccentricity ϵ_n and eccentricity decorrelation A_n from subnucleon-Glauber model arXiv:2003.04340



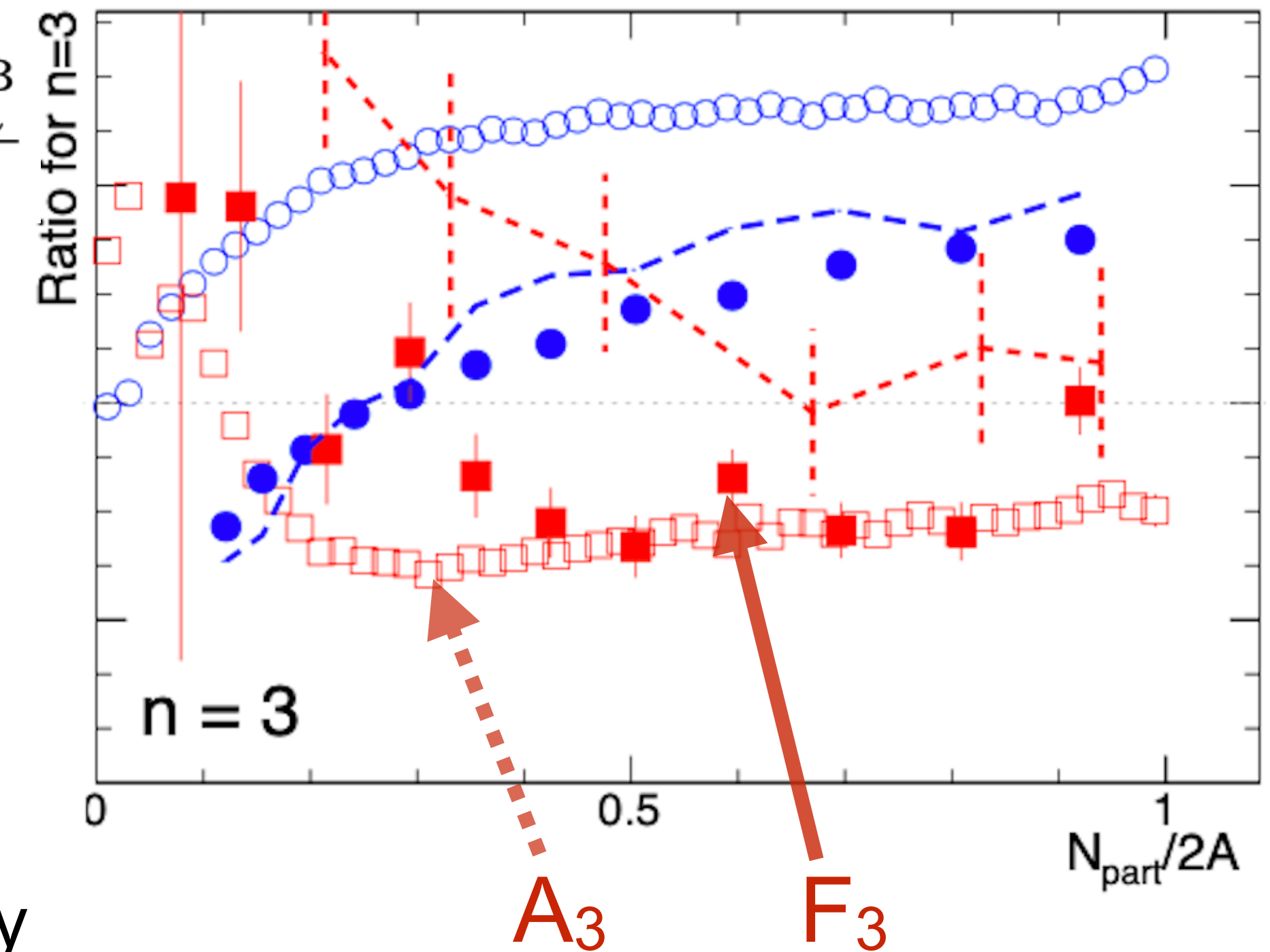
$$\epsilon_{n+} = \frac{\epsilon_n^F + \epsilon_n^B}{2} \quad \epsilon_{n-} = \frac{\epsilon_n^F - \epsilon_n^B}{2}$$

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$N_{\text{part}}/2A \longrightarrow \text{Centrality}$

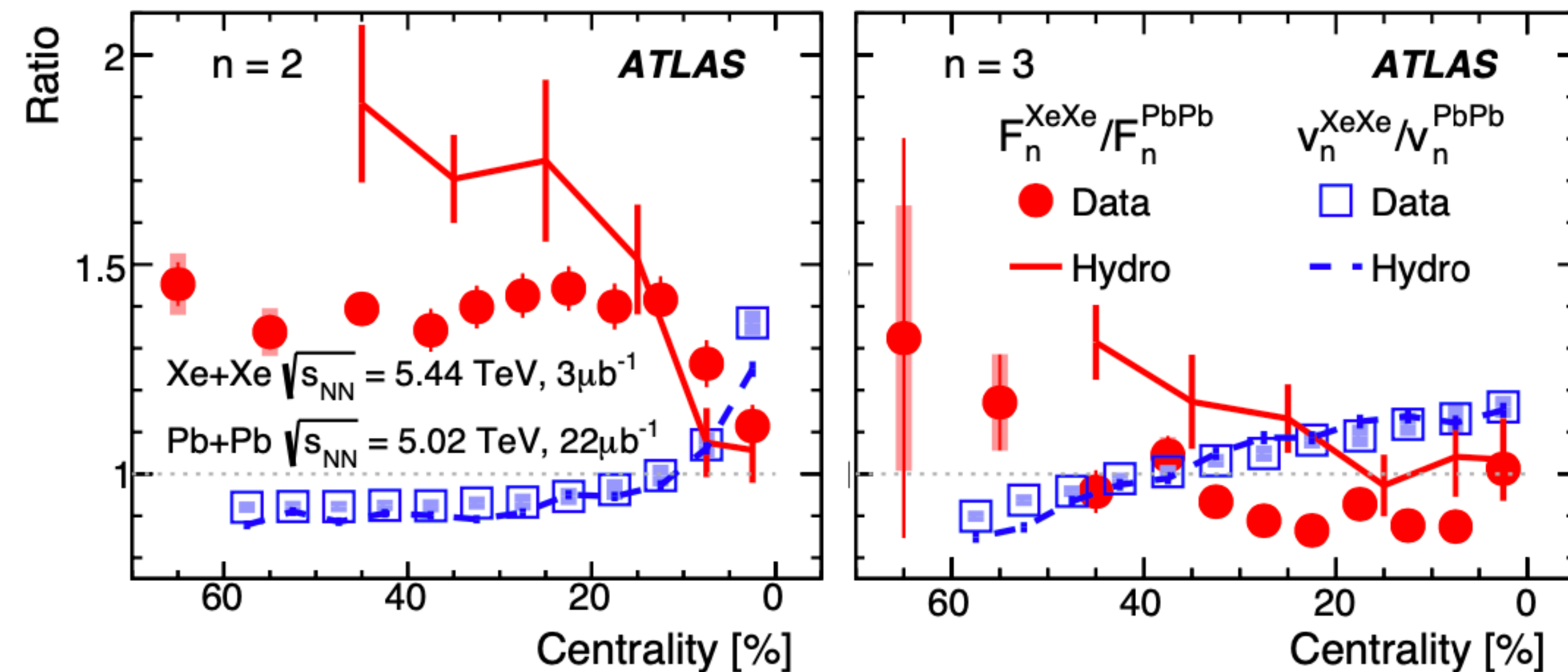


- Difference between v_n and ϵ_n ratios - Late time effects
- F_n and A_n ratios show good agreement - insensitive to late time effects
- Reverse ordering of ATLAS F_2 - F_3 ratios are explained by the Glauber A_2 - A_3 ratios
- System-size dependence of F_n - depends mainly on initial state

- Strong system-size dependence (Xe+Xe vs Pb+Pb) in ATLAS - both v_n and F_n
- Direct comparison v_n and F_n ratios - **anti-correlation**
- v_n ratios - IS Geometry dominant (central) , Viscous effects dominant (non-central)
- $n=2$ vs $n=3$ - same ordering for v_n ratios **but reverse ordering for F_n ratios**

$$F_2^{XeXe} > F_2^{PbPb}$$

$$F_3^{XeXe} < F_3^{PbPb}$$



- Current Hydro models tuned for v_n ratio cannot explain F_n ratio accurately
- Glauber Calculation - F_n ratio insensitive to late-time effects, Reverse ordering of $n=2, n=3$ due to eccentricity decorrelation
- Measurements can constrain different flow models and help develop full 3+1D initial conditions

Backup

- Glauber model calculations compared with ATLAS data and hydro calculations

arXiv:2003.04340

$$\mathcal{E}_n(\eta) = \mathcal{E}_{n+} + f_n(\eta)\mathcal{E}_{n-}, \quad \mathcal{E}_{n+} = \frac{\mathcal{E}_n^F + \mathcal{E}_n^B}{2}, \quad \mathcal{E}_{n-} = \frac{\mathcal{E}_n^F - \mathcal{E}_n^B}{2}$$

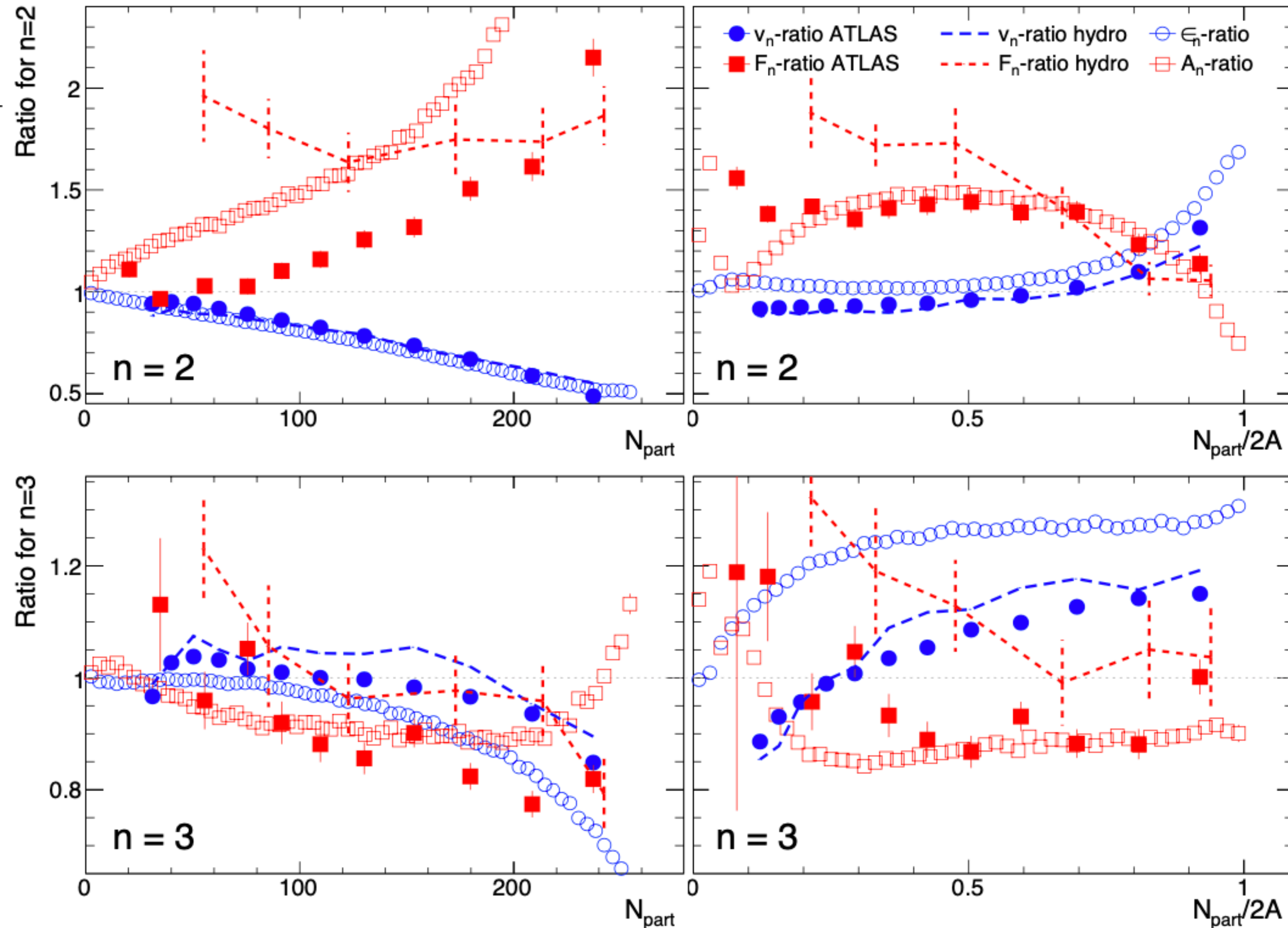
$$r_n^s(\eta) \approx 1 - 2\eta a_n A_n, \quad a_n = \left\langle \frac{\partial f_n}{\partial \eta} \Big|_{\eta=0} f_n(\eta_r) \right\rangle,$$

$$A_n \equiv \frac{\langle \varepsilon_{n-}^2 \rangle}{\langle \varepsilon_{n+}^2 \rangle + \langle \varepsilon_{n-}^2 \rangle} \approx \frac{\langle \varepsilon_{n-}^2 \rangle}{\langle \varepsilon_n^2 \rangle}$$

$$\sqrt{\langle v_n^2 \rangle} = \kappa_n \sqrt{\langle \varepsilon_n^2 \rangle} \quad F_n = \kappa'_n A_n$$

$$\frac{v_n^{A+A}}{v_n^{B+B}} = \frac{\kappa_n^{A+A}}{\kappa_n^{B+B}} \frac{\varepsilon_n^{A+A}}{\varepsilon_n^{B+B}}, \quad \frac{F_n^{A+A}}{F_n^{B+B}} = \frac{\kappa_n'^{A+A}}{\kappa_n'^{B+B}} \frac{A_n^{A+A}}{A_n^{B+B}}$$

- Opposite ordering of F_2 and F_3 ratios are well reproduced by ε_n and A_n



2-particle Correlation

$$C(\Delta\eta, \Delta\phi) = \frac{S(\Delta\eta, \Delta\phi)}{B(\Delta\eta, \Delta\phi)}$$

Integrate with $|\Delta\eta| > 2$ - remove short-range non-flow

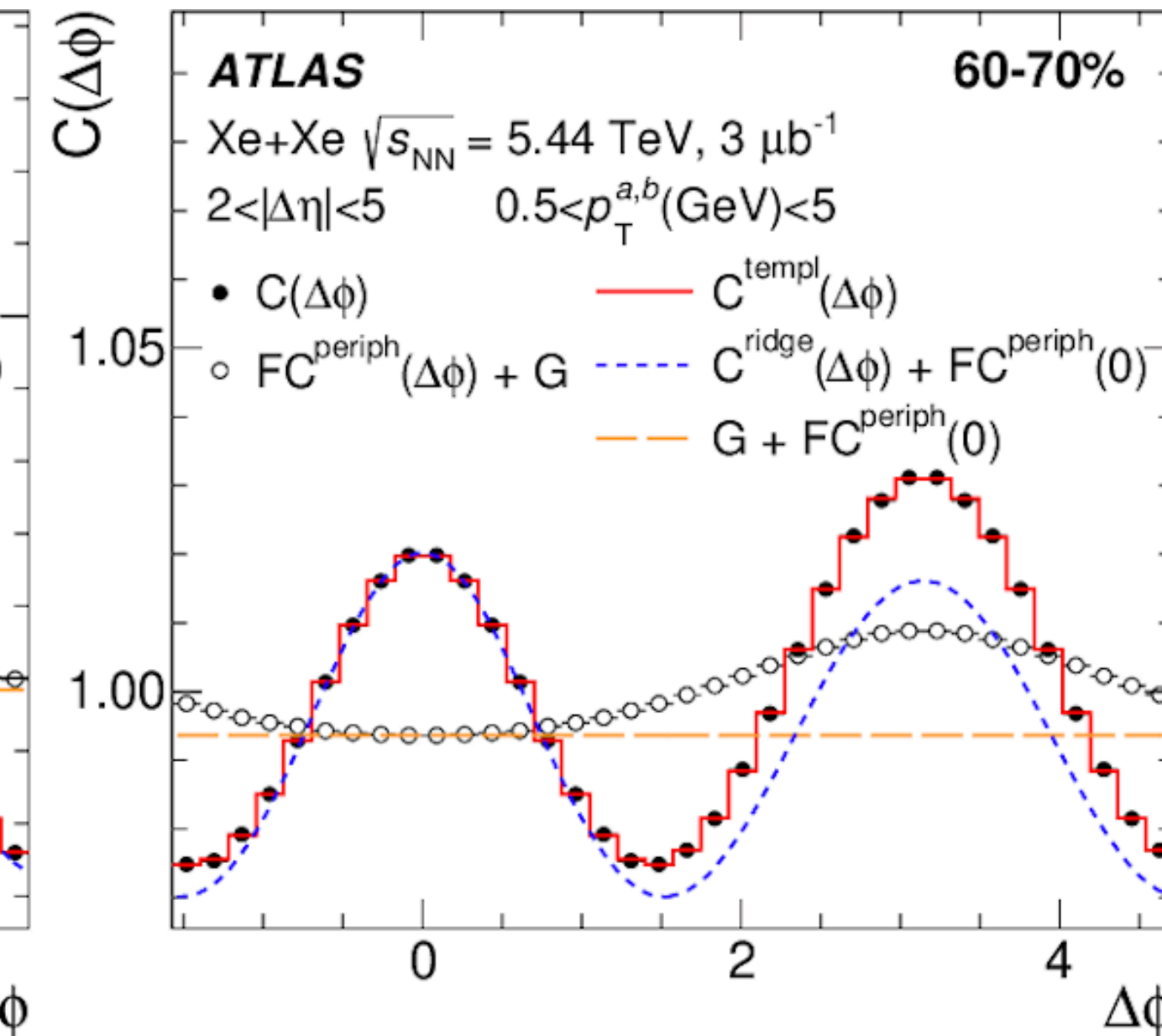
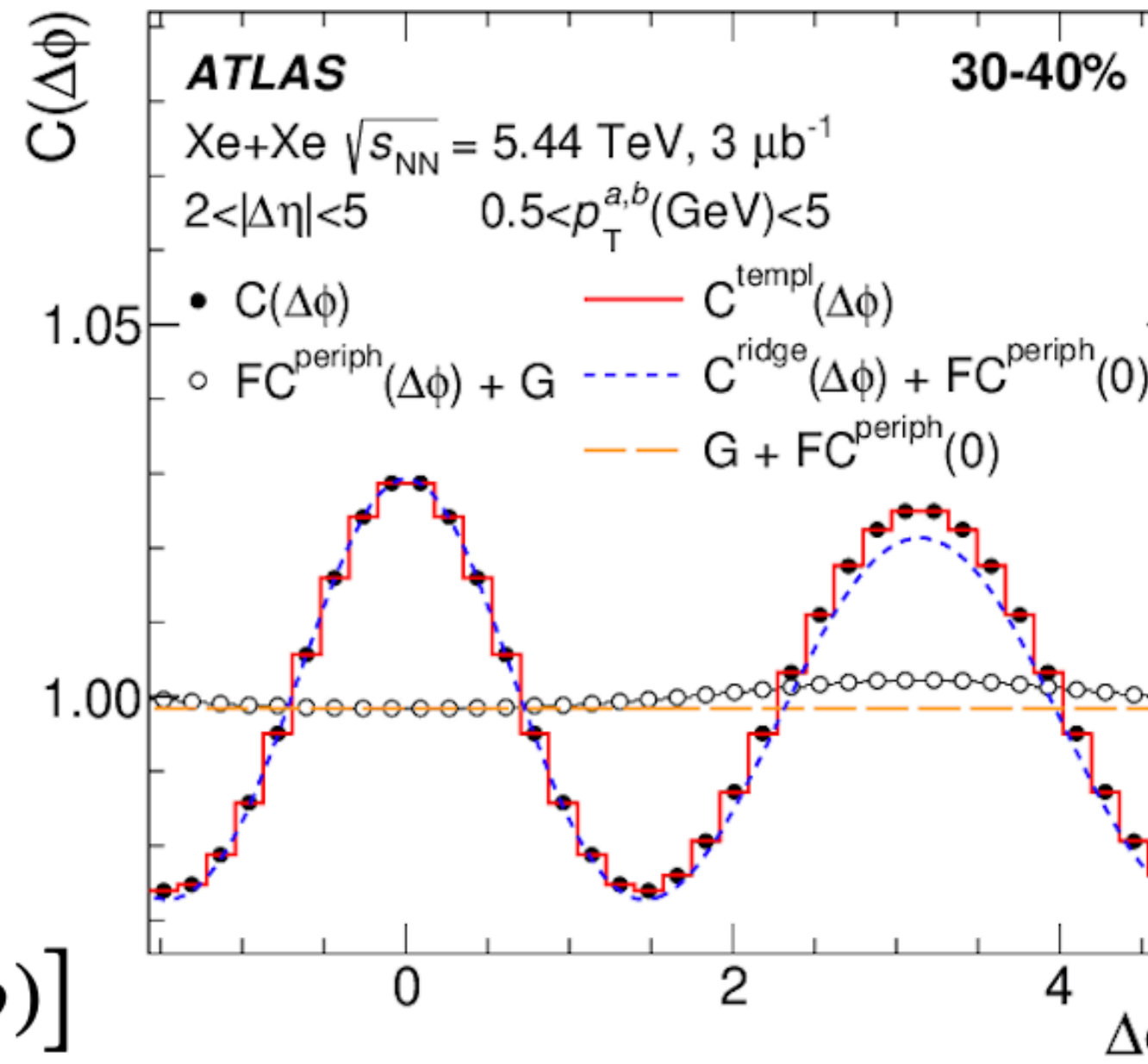
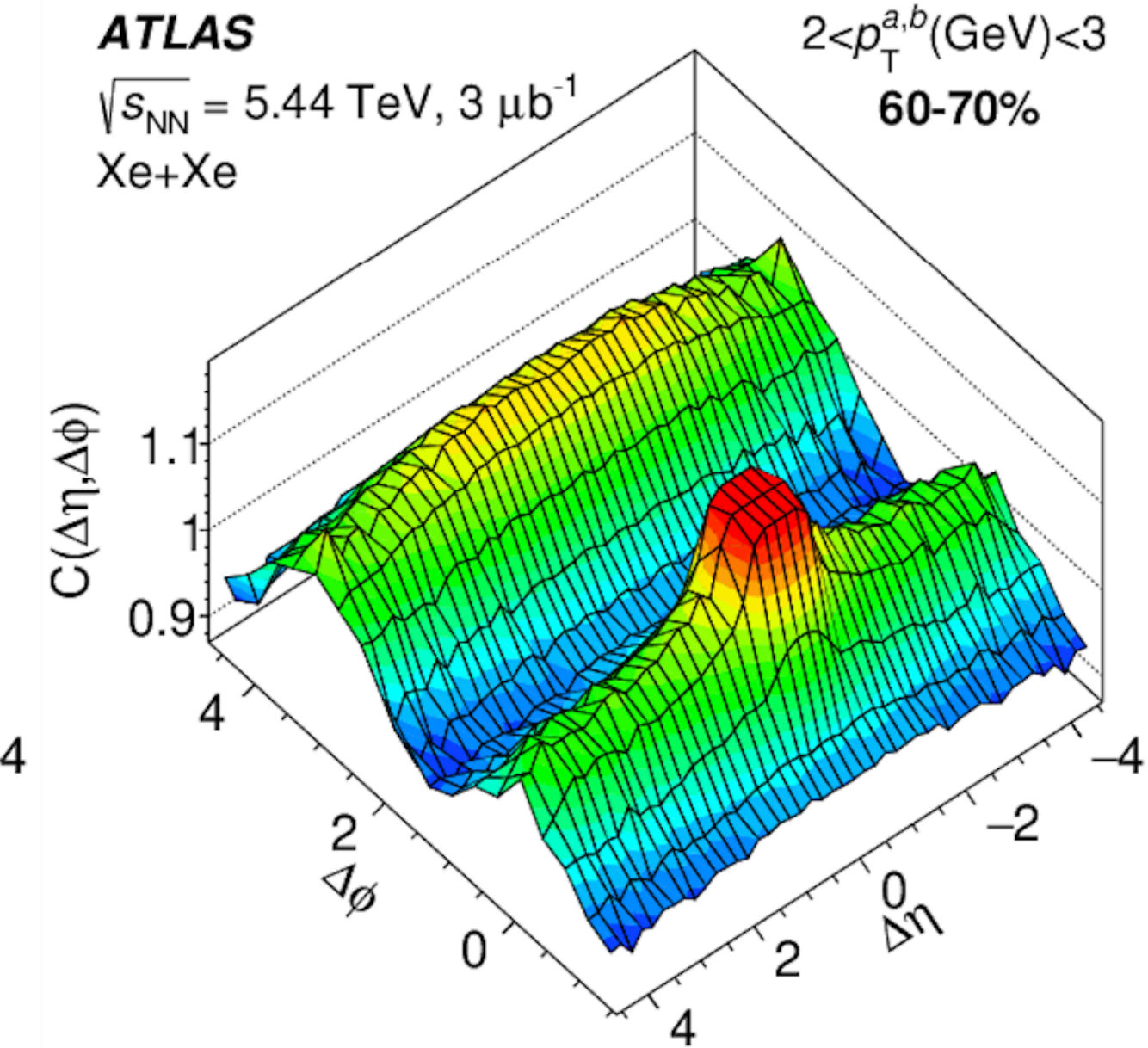
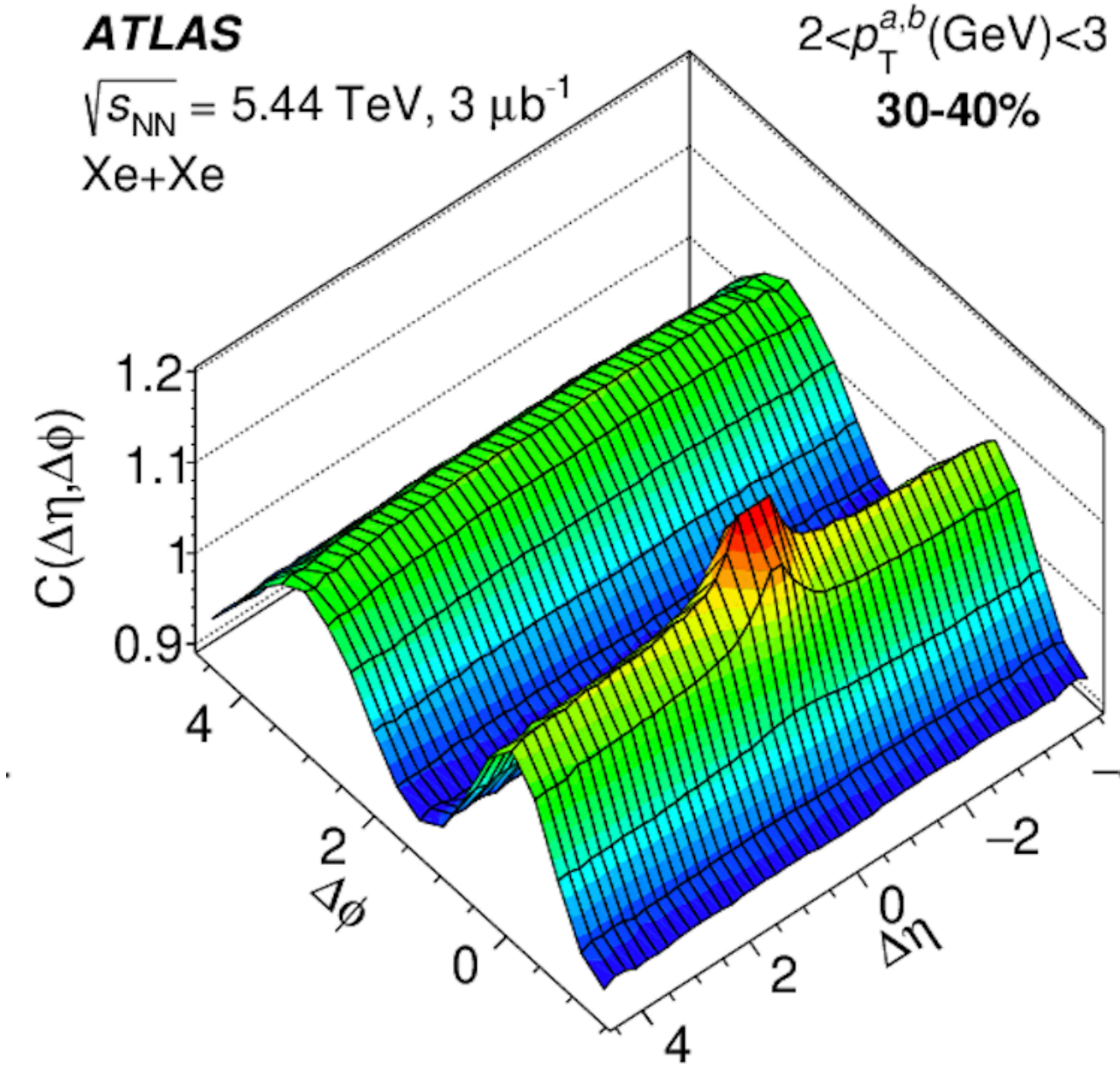
$$C(\Delta\phi) = C_0 \left[1 + 2 \sum_{n=1}^{\infty} v_{n,n}(p_T^a, p_T^b) \cos(n\Delta\phi) \right]$$

- ◆ Long-range Non-flow (dijets) - **peripheral collisions and/or at $p_T > 4\text{GeV}$**

Template fit method is also implemented

- ◆ Developed for pp collisions
- ◆ Can be applied to estimate dijet correlation bias in v_n measurements
- ◆ Significant for Xe+Xe collisions

$$\begin{aligned} C^{\text{templ}}(\Delta\phi) &\equiv FC^{\text{periph}}(\Delta\phi) + C^{\text{ridge}}(\Delta\phi) \\ &= FC^{\text{periph}}(\Delta\phi) + G \left[1 + 2 \sum_{n=2}^{\infty} v_{n,n} \cos(n\Delta\phi) \right] \end{aligned}$$



- Assume factorization

$$v_n(p_T^b) = \frac{v_{n,n}(p_T^a, p_T^b)}{v_n(p_T^a)} = \frac{v_{n,n}(p_T^a, p_T^b)}{\sqrt{v_{n,n}(p_T^a, p_T^a)}}$$

- Plots - p_T dependence of v_n obtained from 2PC and from template fitting (with and without ZYAM)
- 2PC - Large jet-bias at high- p_T in peripheral collisions, most prominent for v_2 and v_4
- Template fitting - reduces this jet bias

