



# Uncertainties and challenges in jet reconstruction in ATLAS

Aparajita Dattagupta on behalf of the ATLAS collaboration

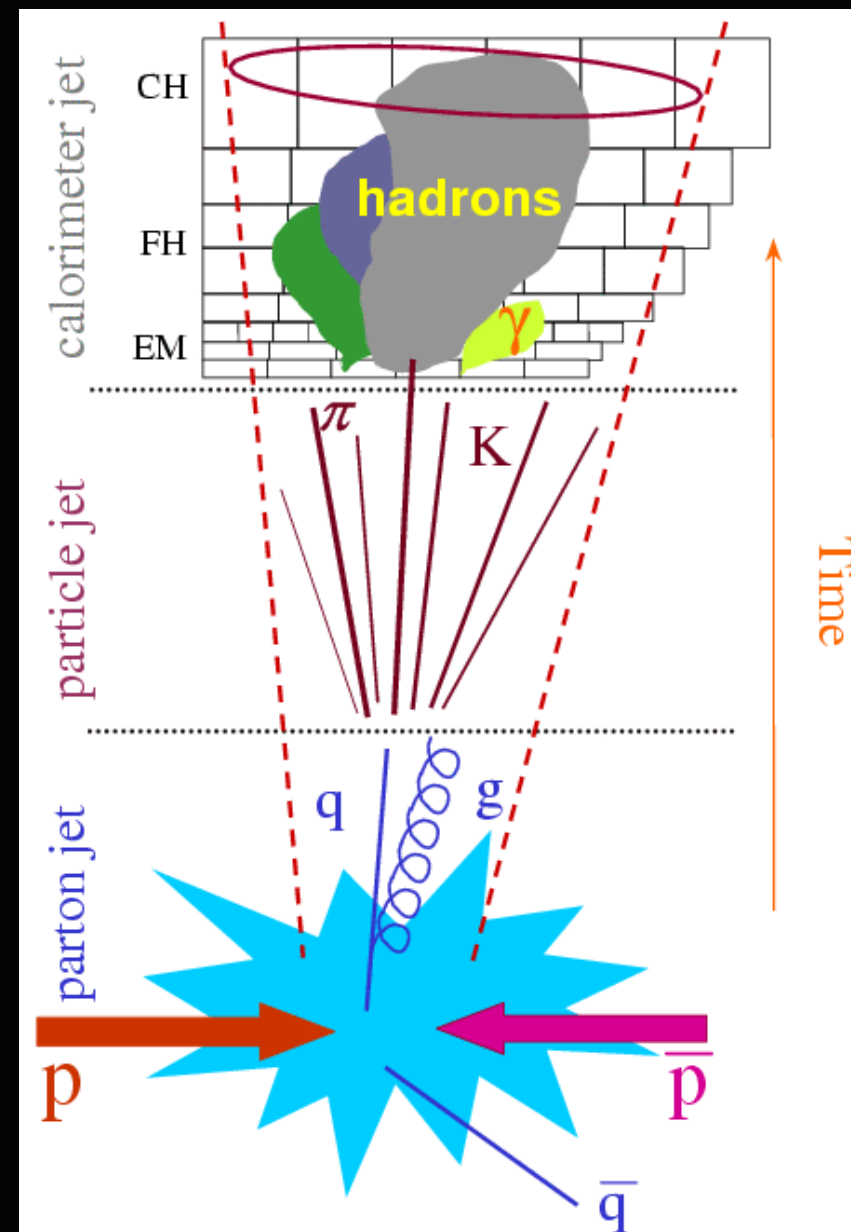
University of Oregon

QCD@LHC

27th - 31st August 2018, Dresden, Germany

# Introduction

- Jets are a useful tool to represent the hadronic energy present in a given proton-proton collision event at the LHC.
- Inelastic proton-proton collisions result in production of quarks and gluons which undergo parton showering and hadronization, as they cannot exist in isolation due to color confinement.
  - Observed as collimated streams of particles depositing energy in the calorimeters, reconstructed as jets.
- Jets are important to almost all analyses at ATLAS
  - Reconstructed and used in order to enhance signal selection in an analysis either by vetoing them or by requiring their presence
- Given their importance, it is essential to calibrate the jet energies to the correct scale and to properly take into account the uncertainties on these calibration procedures.
- In this talk, the different jet reconstruction and calibration procedures and the performance of these methods will be discussed.





# Which jet algorithm to use?

---

- Recommendations based on experience from CDF and D0 experiments:
  - Most commonly used algorithm for reconstructing jets on ATLAS is the anti-kt algorithm.
- Theoretical considerations:
  - Infrared and collinear safety.
  - Invariance under boost.
- Experimental considerations:
  - Independent of detector technology.
  - Control over pile-up effects with increasing luminosity.
  - Should capture well, the decay of the initiating particle.
  - Should be easy to calibrate.
  - Not very computing intensive.

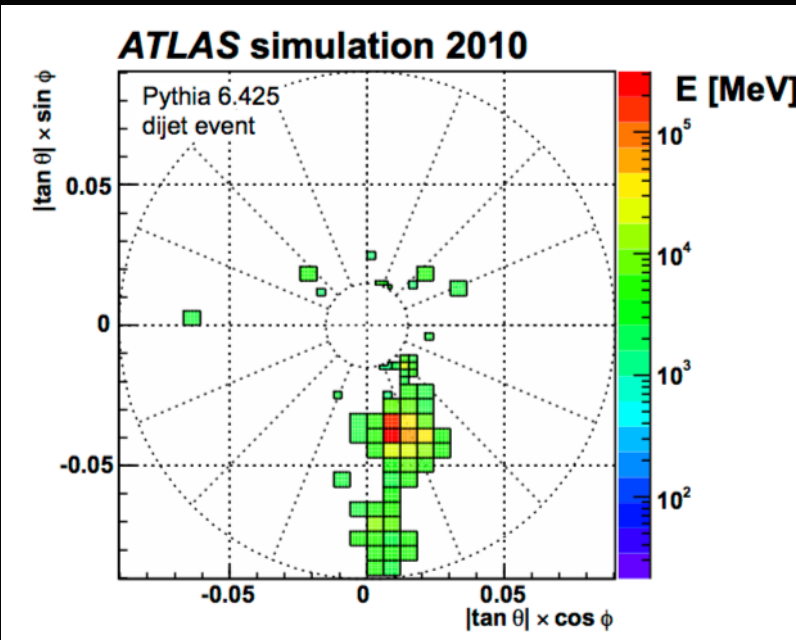


- Jets on ATLAS can be reconstructed from various inputs:
  - Standard calorimeter jets built from topoclusters.
  - Track jets built using the tracking information.
  - pFlow and TrackCaloCluster jets are examples of jets that use both tracking and calorimeter information.
  - Particle-level jets: Reference for the simulation-based jet calibration.
- Most commonly used jets are reconstructed using the anti-kt algorithm, radius parameters can be:
  - $R=0.4$ : Used in almost all analyses on ATLAS using quark/gluon-initiated jets.
  - $R=1.0$ : Used mainly to capture decay products of hadronically decaying massive particles like a top quark or W/Z/H boson (boosted topologies).
  - Variable-radius (variable-R) jets are also increasingly being used at high  $p_T$  regimes where it is shown that the hadronic decay products are contained in a smaller area than  $R=1.0$ .

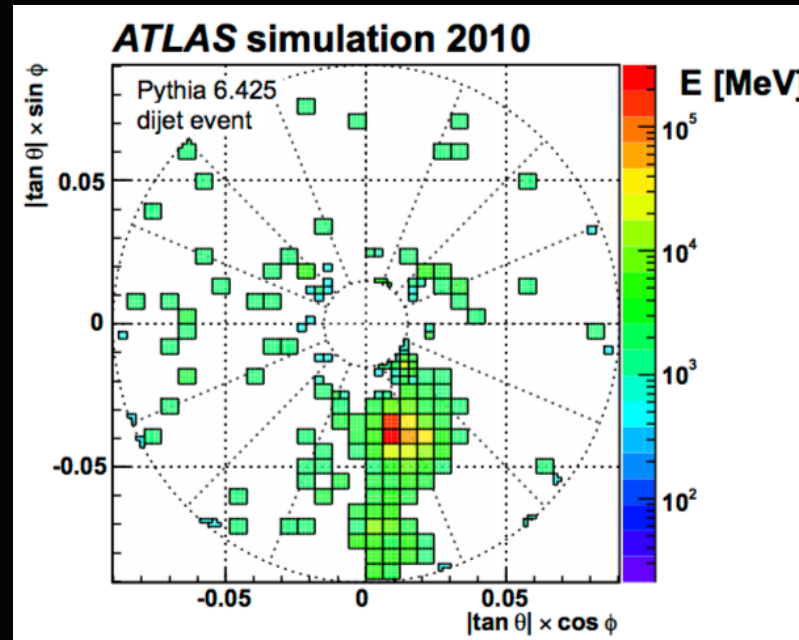


- Jets can be reconstructed from 3D topological clusters of calorimeter cells (topoclusters)
  - Calibrated or uncalibrated topoclusters (LC and EM scale).
- Cells are clustered together based on their signal significance or ratio of the cell signal to the average (expected) noise.

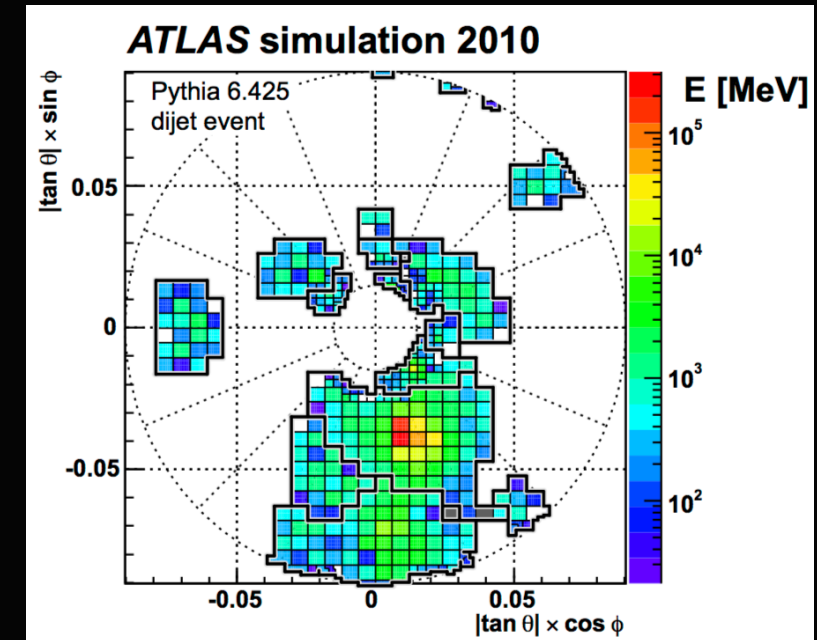
$$S_{\text{cell}}^{\text{EM}} = \frac{E_{\text{cell}}^{\text{EM}}}{\sigma_{\text{noise,cell}}^{\text{EM}}}$$



Cells with signal Significance > 4  
seed topoclusters



Cells with signal significance > 2  
control topocluster growth

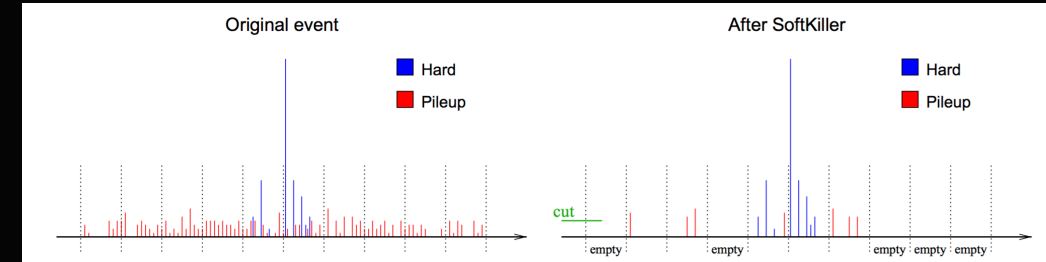


All clustered cells: topoclusters  
surrounding layer

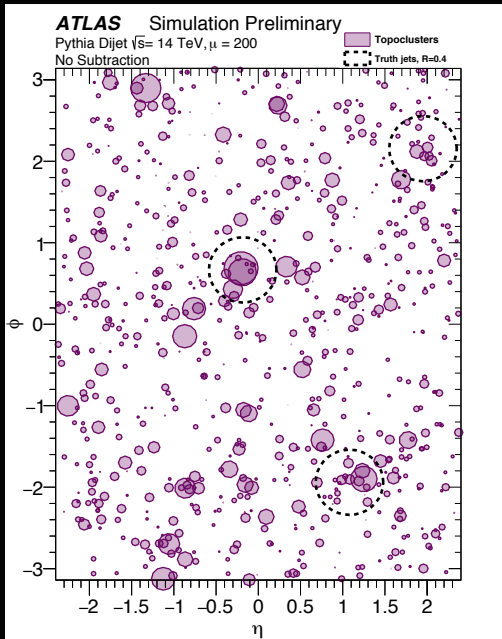
Stages of topocluster formation: Based on signal significance of calorimeter cells



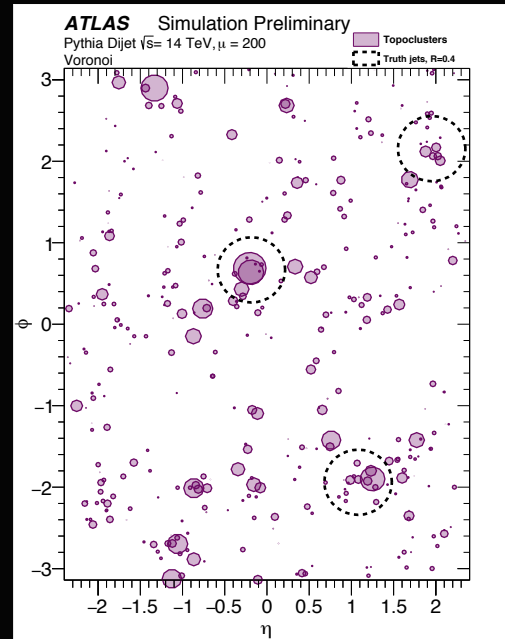
- Constituent subtraction (CS): “ghost particles” ( $E = \epsilon$ ) added uniformly in an event and clustered alongside cells: Number of clustered ghosts is proportional to the area.
  - Topoclusters are corrected based on  $N_{ghost}$  and event pile-up density ( $Rho$ ).
- Voronoi area: eta-phi area closest to each cluster:
  - $Rho$  subtracted from each cluster according to voronoi area.
  - All clusters with low significance above noise are removed.
- Softkiller (SK):
  - Clusters below an event-specific  $p_T$  cut after CS or VS are rejected to further remove pileup.
- CS+SK found to be the best performing one.



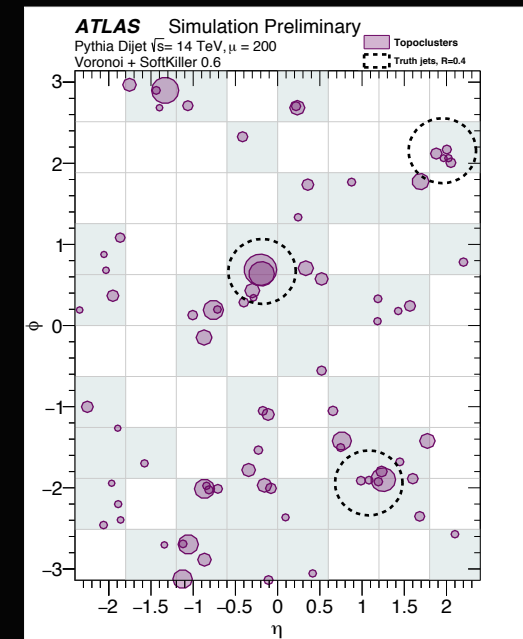
No subtraction



After Voronoi subtraction

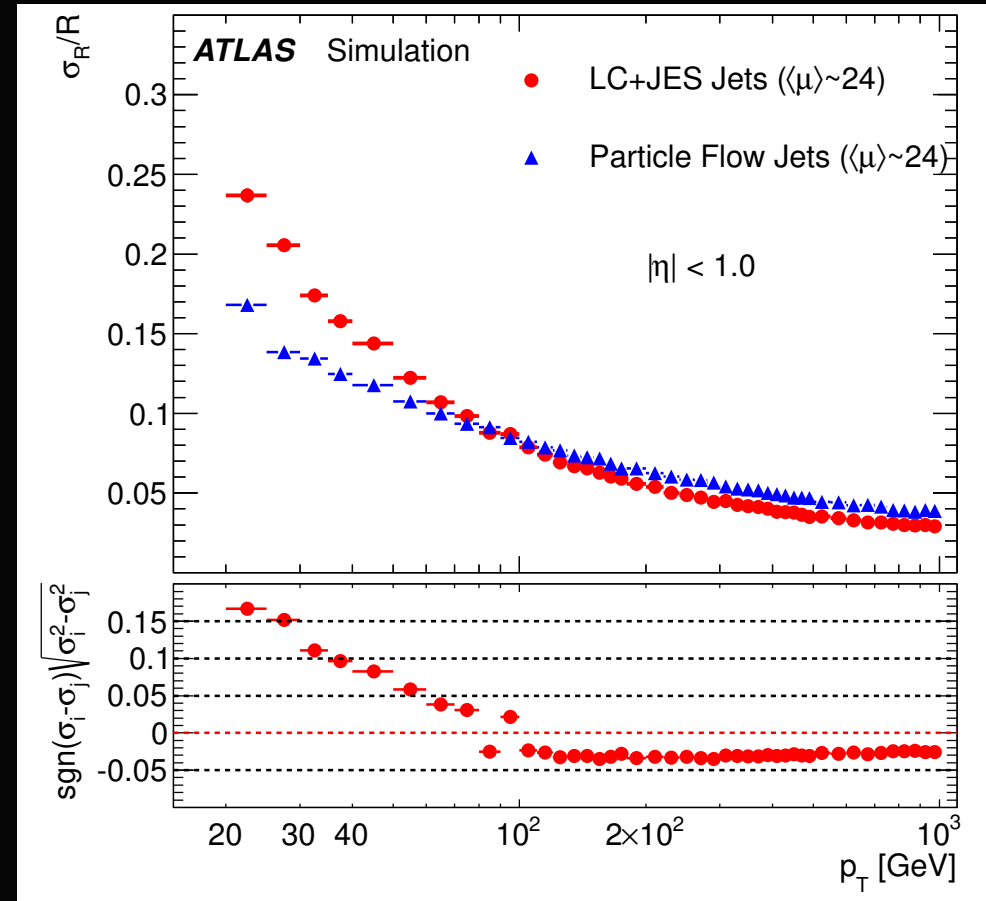
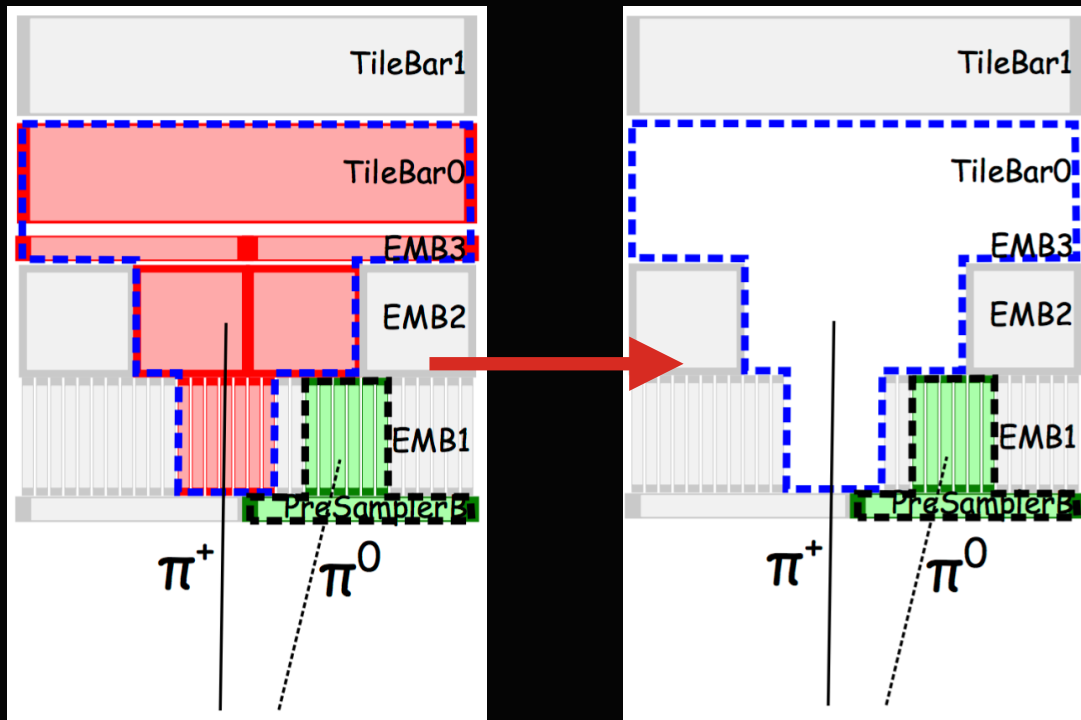


After Voronoi + SoftKiller subtraction





- Tracks matched to topo-clusters, removing charged energy while keeping neutral component. Subtraction of energy is done cell-by-cell
- Reconstructed using the anti-kt algorithm with radius parameter 0.4: inputs are topo-clusters surviving the energy subtraction step and the selected track
- Improved resolution at low  $p_T$  compared to LC+JES jets due to better tracker resolution at low  $p_T$





# Improvements in jet mass reconstruction for R=1.0 jets

Mass is a commonly used variable in the identification of hadronically decaying massive particles such as top and W/Z/H.

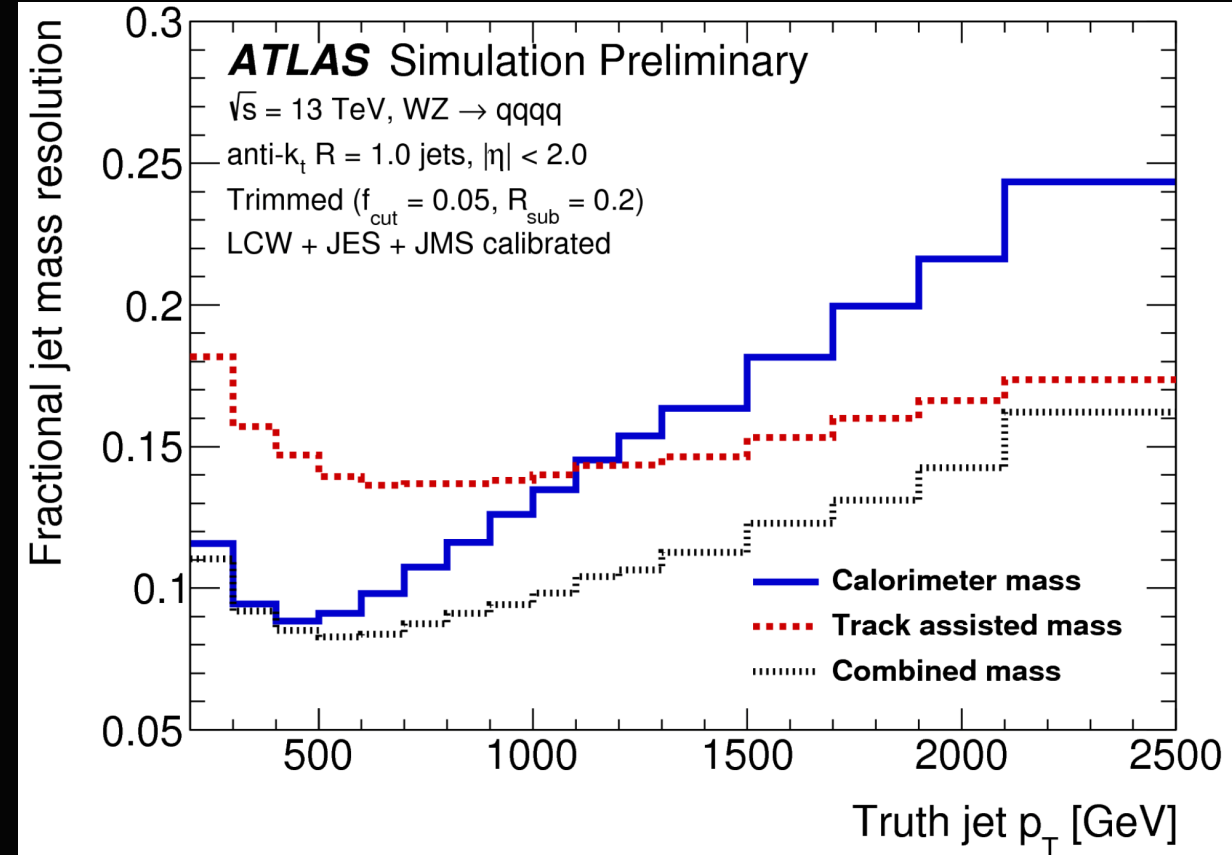
Calorimeter mass: Resolution degrades at high pt, shower size ( $1/p_T$ ) becomes comparable to calo granularity

$$m^{\text{calo}} = \sqrt{\left(\sum_{i \in J} E_i\right)^2 - \left(\sum_{i \in J} \vec{p}_i\right)^2}$$

Track-assisted mass: Tracking information can be used to maintain performance beyond the calo granularity limit.  $Pt(\text{calo})/Pt(\text{track})$  corrects for neutral energy not accounted for in track jet reconstruction.

$$m^{\text{TA}} = \frac{p_T^{\text{calo}}}{p_T^{\text{track}}} \times m^{\text{track}}$$

Combined mass uses an inverse resolution weighted combination of  $m_{\text{Calo}}$  and  $m_{\text{TA}}$



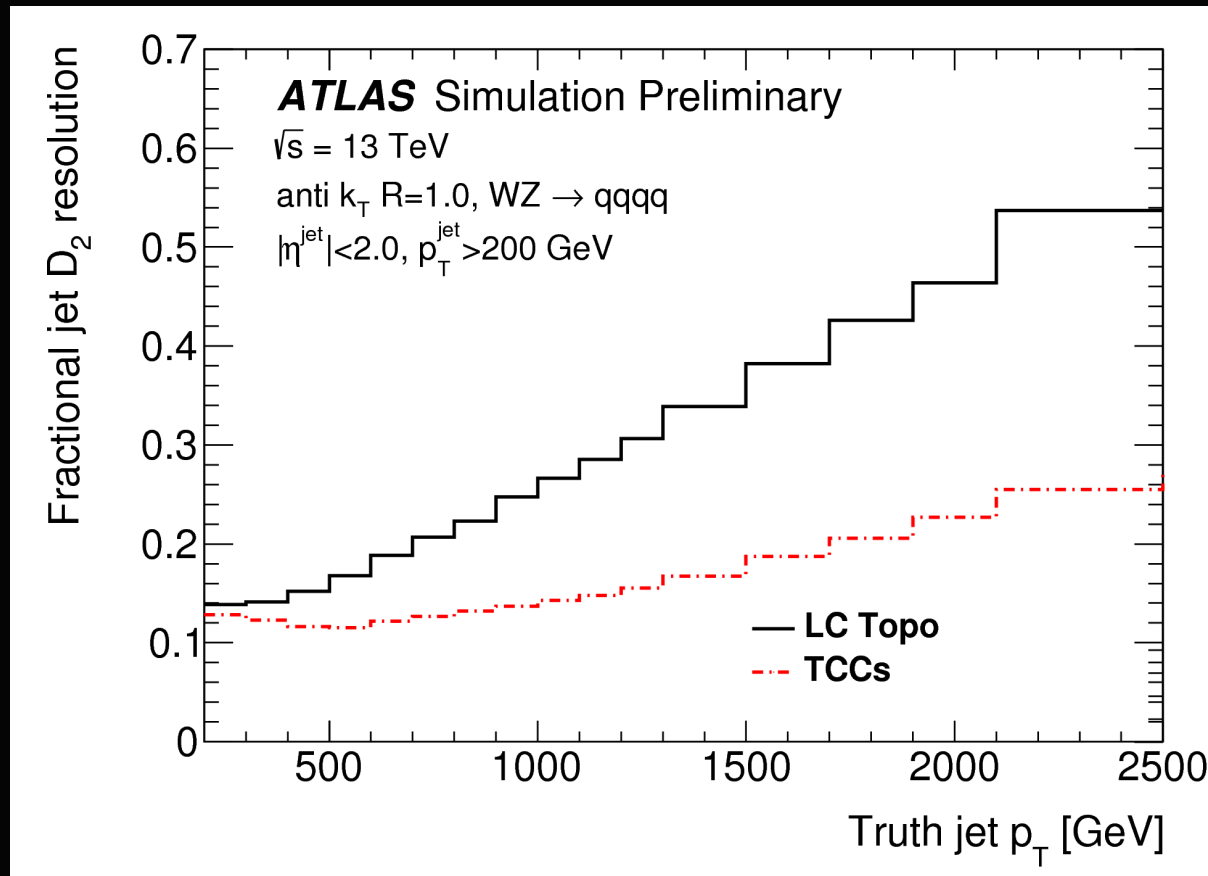
$$w_{\text{det}} = (\sigma[m_{\text{det}}])^{-2}$$

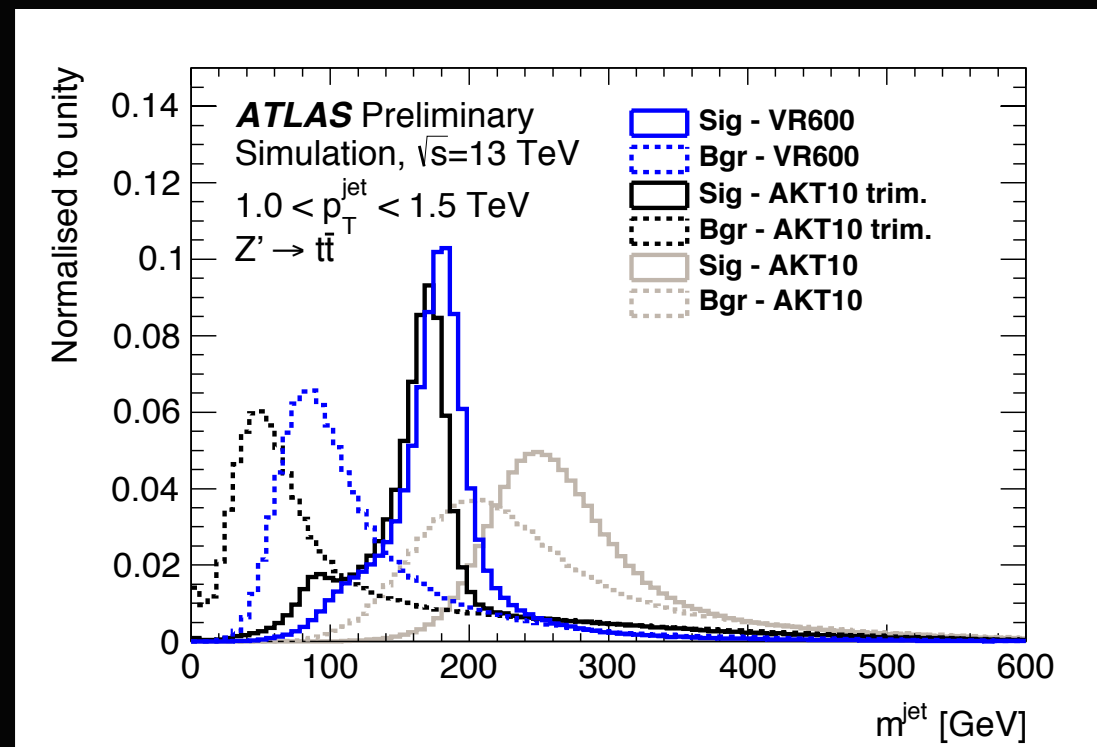
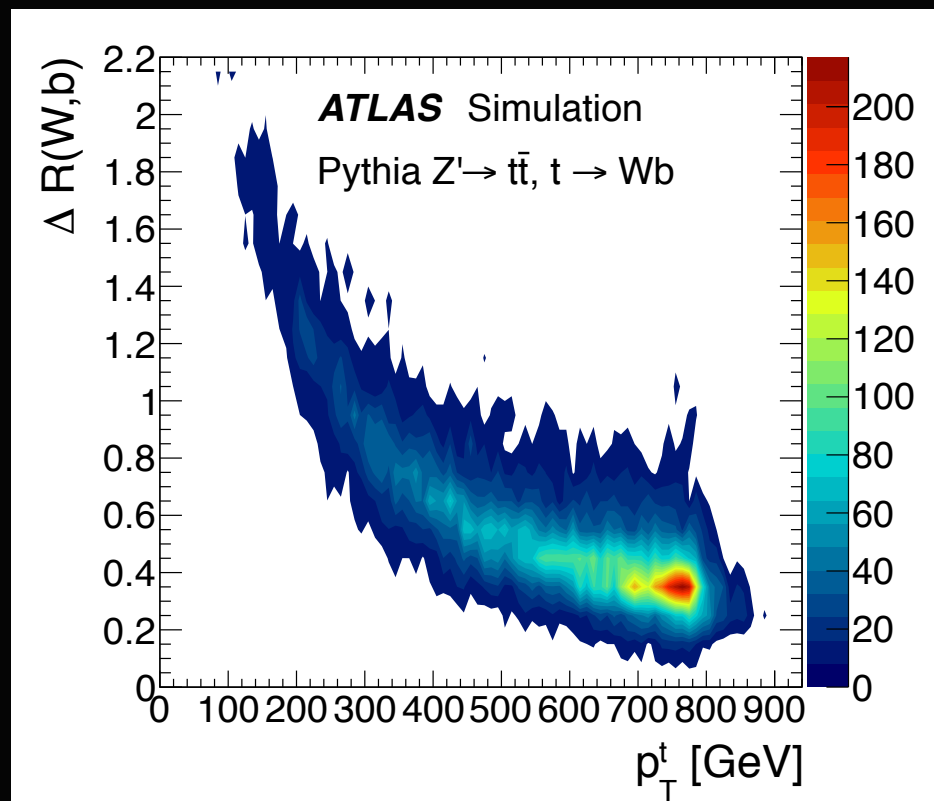
$$m_{\text{comb}} = \left( \sum_{\text{TA,calo}} w_{\text{det}} m_{\text{det}} \right) / \left( \sum_{\text{TA,calo}} w_{\text{det}} \right)$$





- Define a jet using spatial coordinates of the tracker and the energy scale of the calorimeter: Use superior angular resolution of tracker and energy resolution of calorimeter
- Neutral TCCs (unmatched topo-clusters) jets have a similar pile-up dependence to standard topocluster jets, additional pile up removal techniques applied.
- Improved resolution of large-R jet mass for very high pT and substructure observables sensitive to the substructure of a W/Z-boson jet.





- For high  $p_T$  jets (top jets in left plot)  $R < 1$  is sufficient.
- For VR jets, radius parameter scales with  $1/p_T$
- $\rho$  determines how fast the effective jet size decreases with jet  $p_T$ .

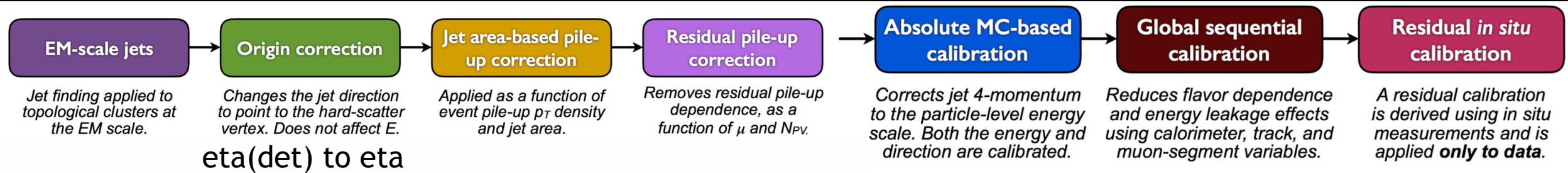
$$R_0 \longrightarrow R_{eff}(p_{T,i}) = \frac{\rho}{p_{T,i}}$$

- Min and max  $R$  values prevents the jets from becoming too large at low  $p_T$  and from shrinking below the detector resolution at high  $p_T$



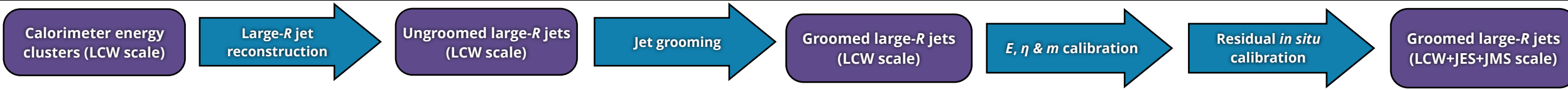
# Calibration procedure for R=0.4 (small-R) and R=1.0 jets (large-R)

[Phys. Rev. D 96 \(2017\) 072002](#)



## R=0.4 jets

[CERN-EP-2018-191](#)



Large-R jets reconstructed using the anti-kt algorithm with  $R = 1.0$ .

removes pile-up and the underlying event

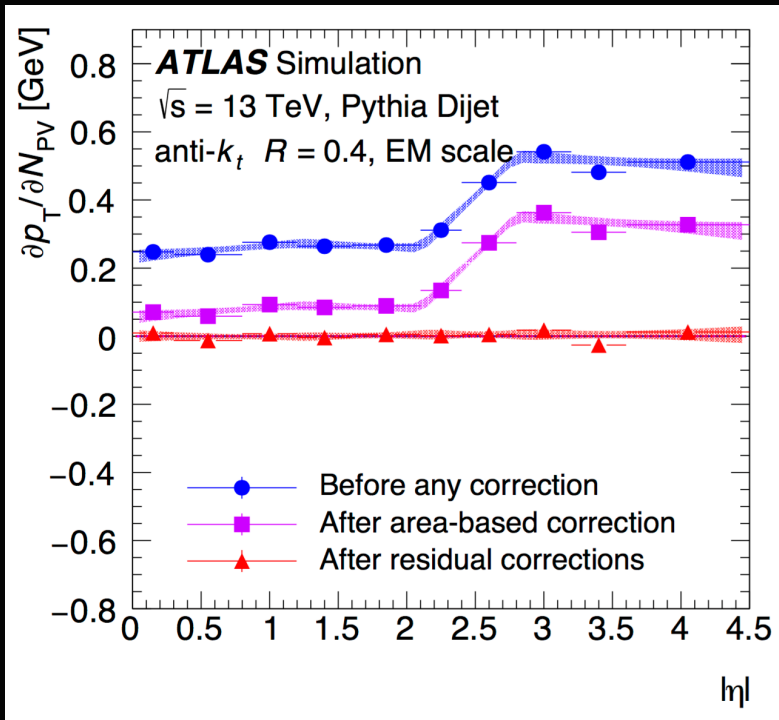
reconstructed jet  $E, \eta$  and  $m$  corrected using MC to the particle jet scale

Residual correction applied to data

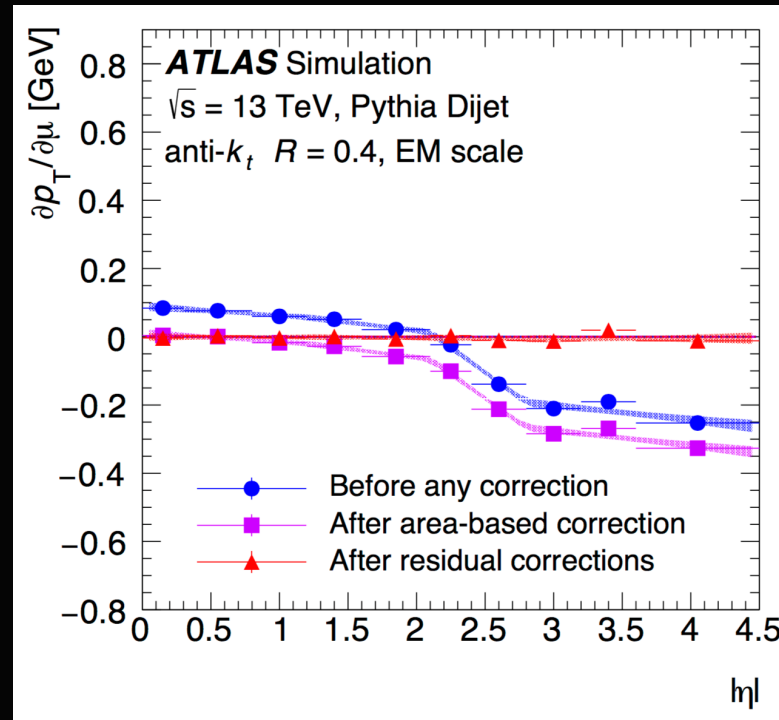
## R=1.0 jets



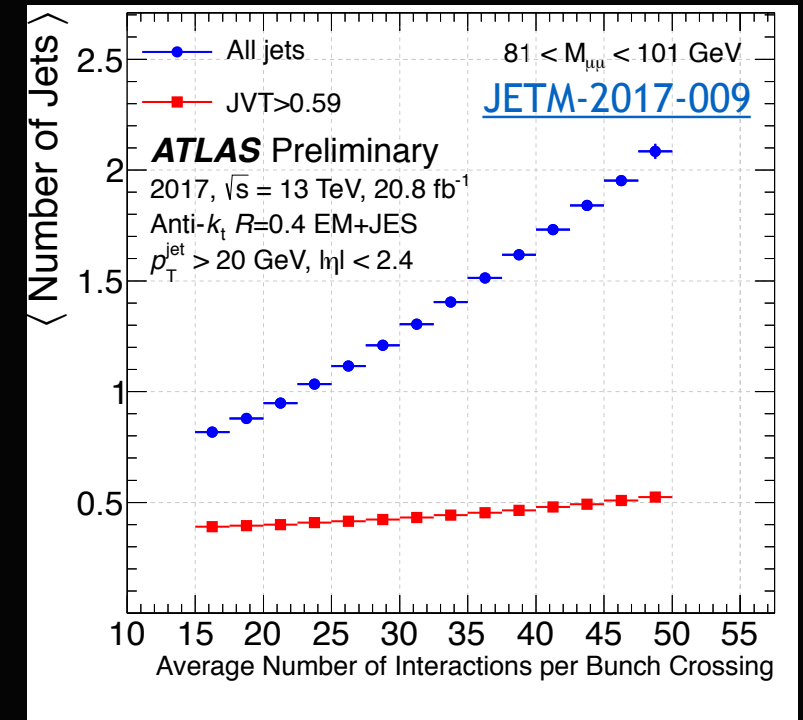
- Particles from pile-up collisions can add additional jets that are not from the hard-scatter
- Can overlap with hard scatter jets, altering their energy
- Effect of pile-up reduced by applying per-jet corrections based on pile-up density in the event and JVT (connects jets to pile-up vertices using tracking information) cuts for small- $R$  jets.



In-time pile-up correction



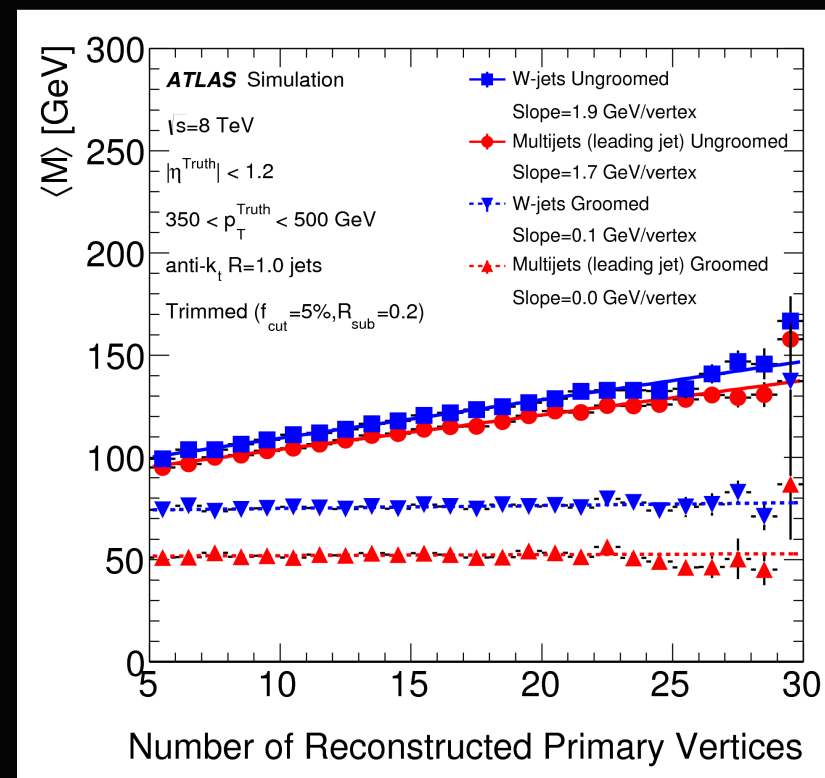
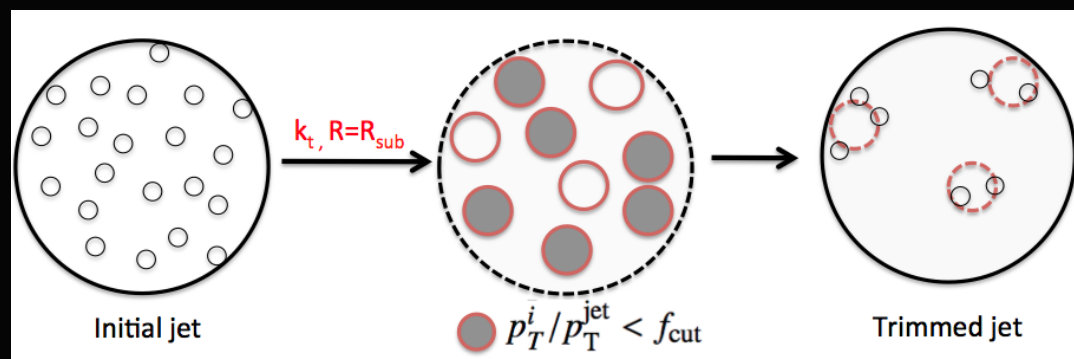
Out-of-time pile-up correction





# Pile-up correction for R=1.0 jets: “grooming”

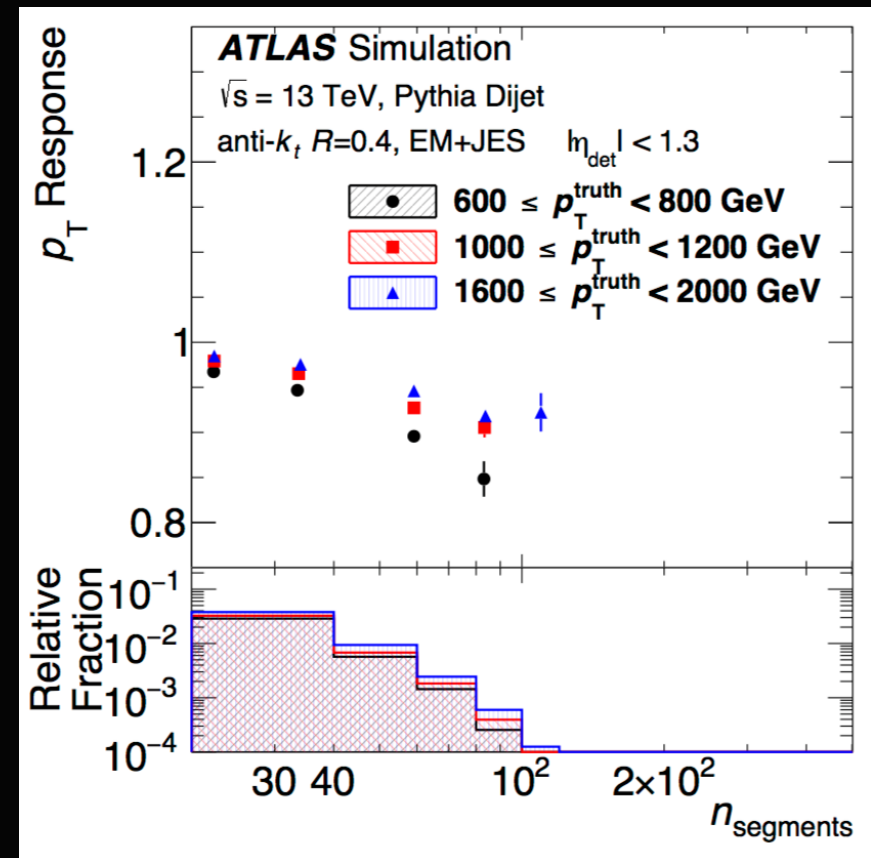
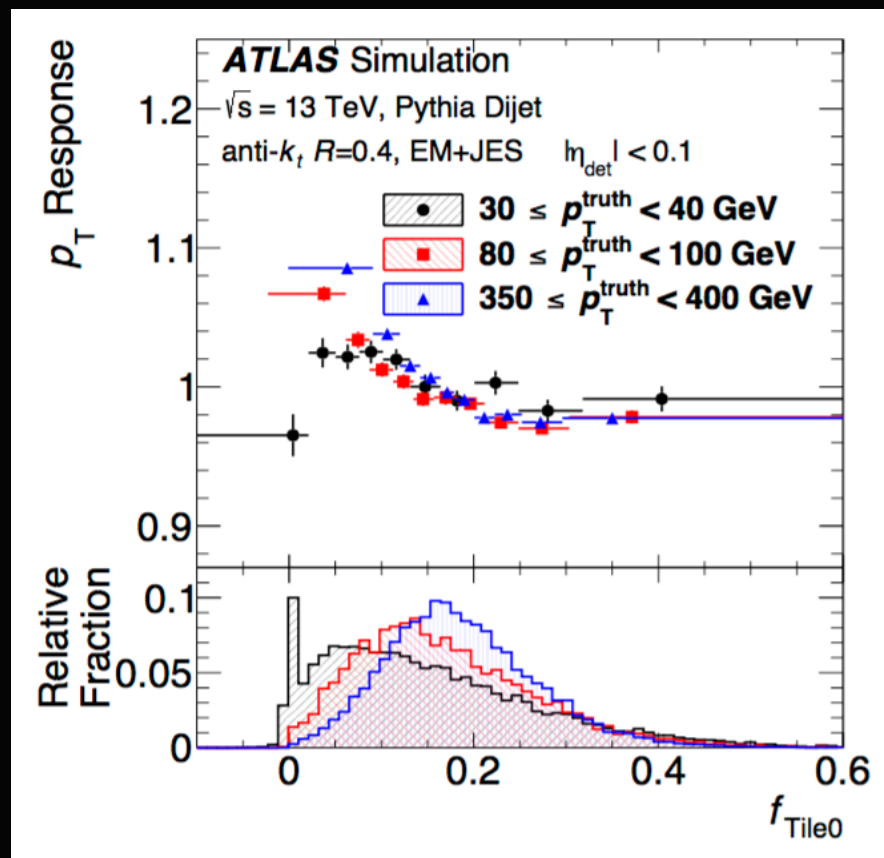
- Large-R jets due to a larger area are more susceptible to pile-up effects.
- Grooming techniques are applied to correct for these effects which can alter substructure features.
- The trimming procedure with parameters of  $f_{\text{cut}} = 5\%$  and  $R_{\text{sub}} = 0.2$  is used on ATLAS.
  - Constituents of the original anti-kt jet are reclustered using the kt algorithm with a distance parameter of  $R_{\text{sub}}$ .
  - Resulting kt sub-jet is removed if the  $p_T$  is less than  $f_{\text{cut}}$  of the large-R jet  $p_T$ .
  - Jet is rebuilt from the remaining constituents.
- Soft drop grooming has also been studied and found to have good performance.



Grooming techniques can improve the reconstruction of the jet mass helping discriminate for e.g. a W-boson initiated jet from a light quark/gluon jet, the distribution for which will peak at lower masses.

# Global sequential calibration (GSC) for R=0.4 jets

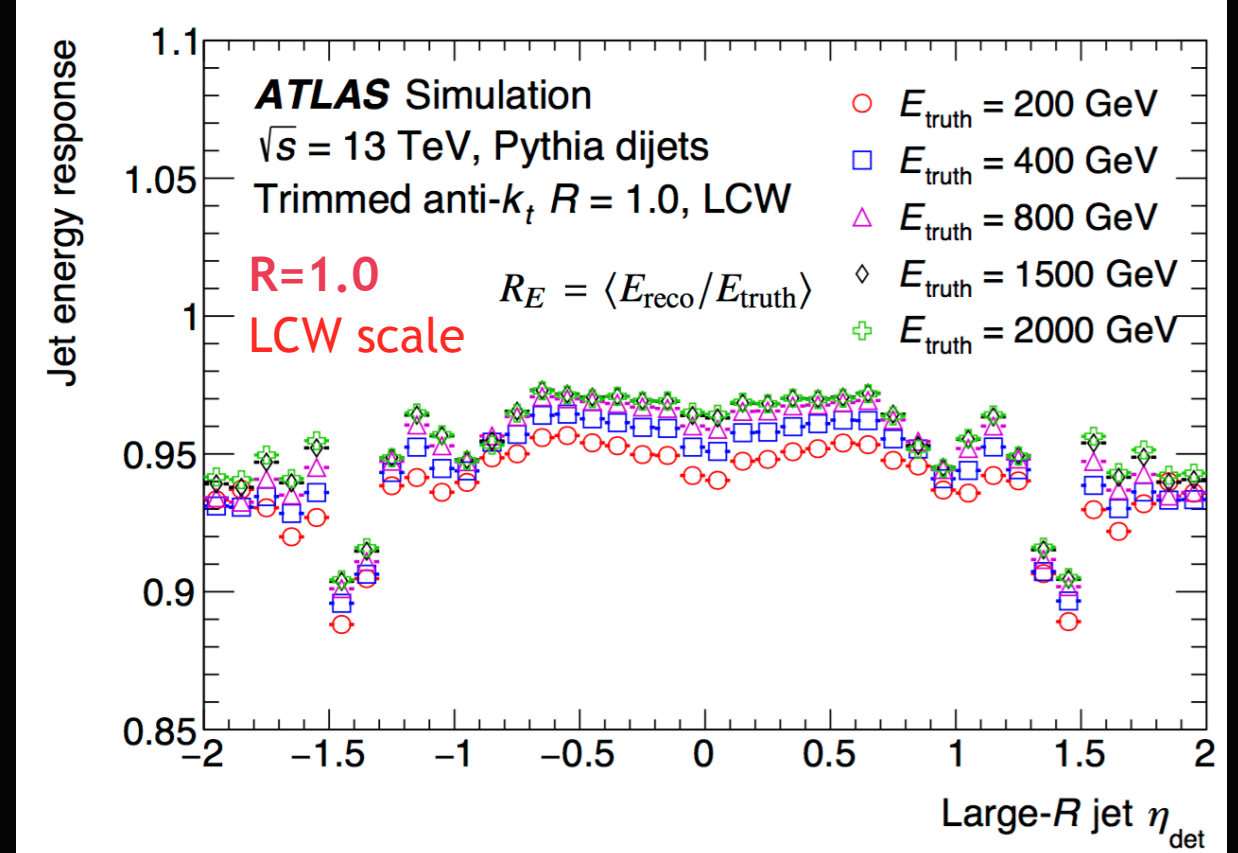
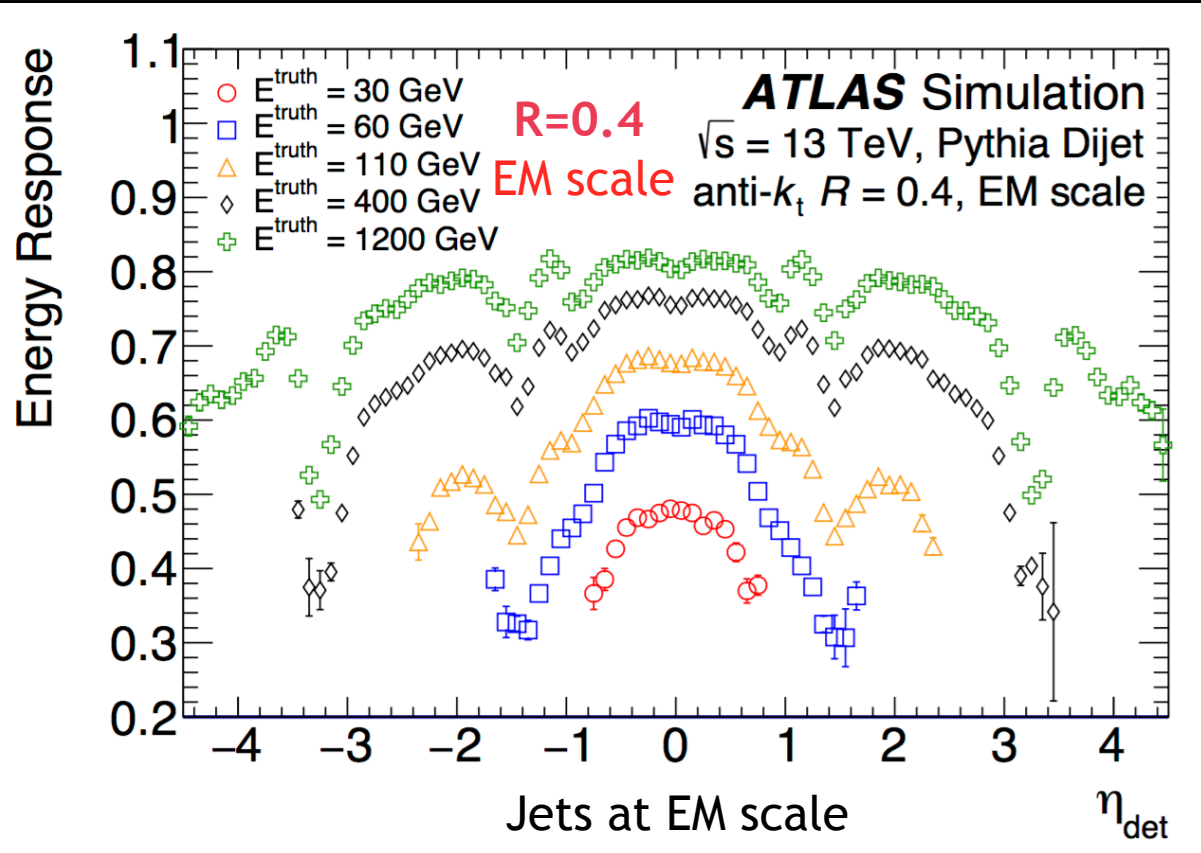
- Correct jet response according to jet shower depth, track variables and muon punch-through to characterize fluctuations in the jet particle composition, distribution of energy within the jet.
- This correction applied in MC improves jet energy resolution and jet flavor dependence.
- Average JES is unchanged





- Average response determined from a Gaussian fit to the core of the response distribution  $R_E = \langle E_{\text{reco}}/E_{\text{truth}} \rangle$
- JES correction factor  $c_{\text{JES}}$  is determined as a function of the jet energy and pseudorapidity  $\eta_{\text{det}}$ .
- Large-R jet energy, mass, eta, and pT after applying the correction factor (Phi is not changed)

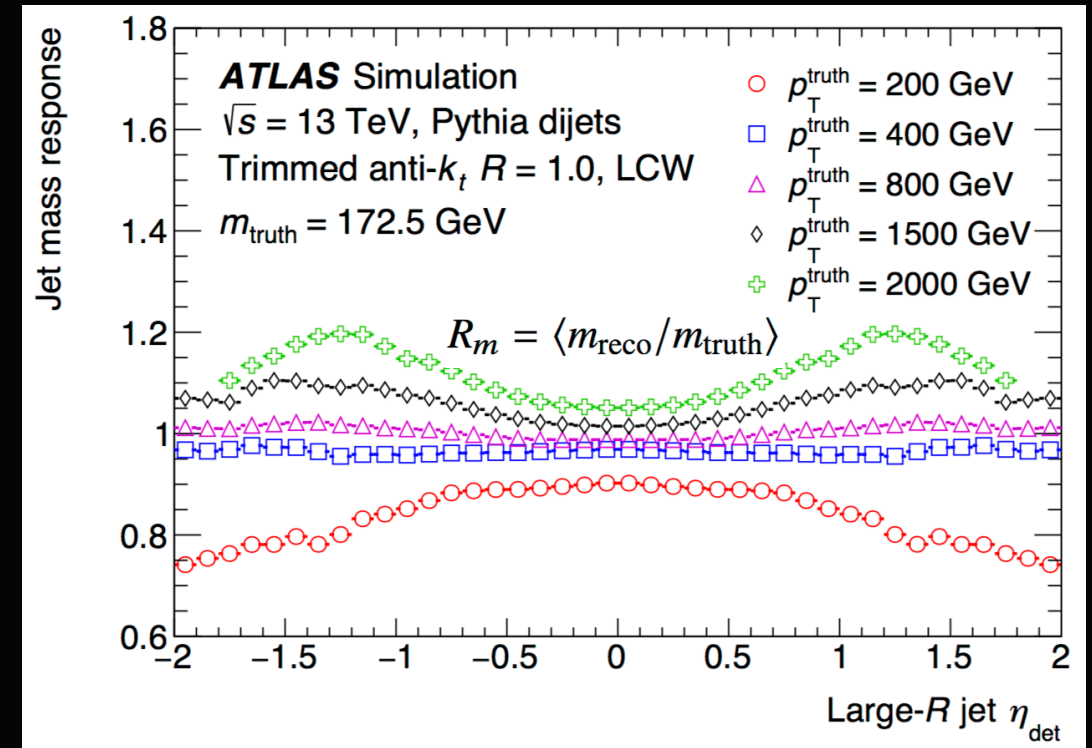
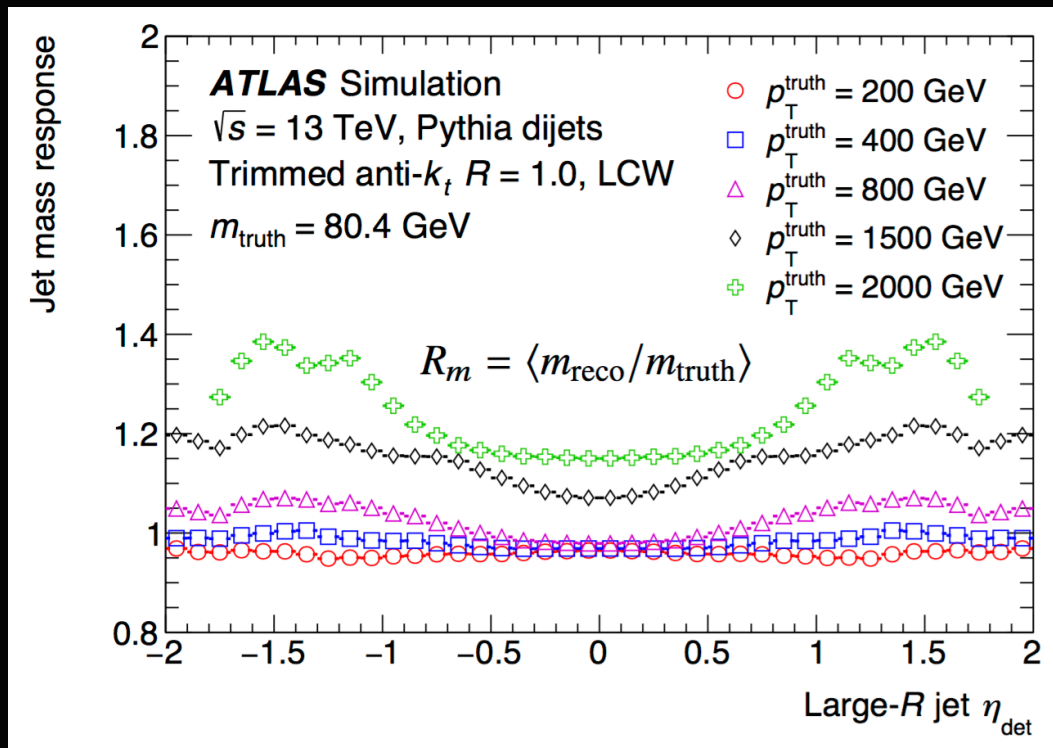
$$E_{\text{reco}} = c_{\text{JES}} E_0, \quad m_{\text{reco}} = c_{\text{JES}} m_0, \quad \eta_{\text{reco}} = \eta_0 + \Delta\eta, \quad p_{\text{T}}^{\text{reco}} = c_{\text{JES}} |\vec{p}_0|/\cosh(\eta_0 + \Delta\eta)$$





- Jet mass more sensitive to soft, wide angle radiation and topocluster merging and splitting than pT. Jet mass response measured in dijet events for different truth pT:  $R_m = \langle m_{\text{reco}}/m_{\text{truth}} \rangle$
- Correction (cJMS) applied after Jet energy response correction, large-R jet energy kept fixed and pT allowed to vary. cJMS varies from ~1-1.5.

$$E_{\text{reco}} = c_{\text{JES}} E_0, \quad m_{\text{reco}} = c_{\text{JES}} c_{\text{JMS}} m_0, \quad \eta_{\text{reco}} = \eta_0 + \Delta\eta, \quad p_{\text{T}}^{\text{reco}} = c_{\text{JES}} \sqrt{E_0^2 - c_{\text{JMS}}^2 m_0^2} / \cosh(\eta_0 + \Delta\eta).$$



Mass response for two representative values of the truth mass: W boson mass (left) and top quark mass (right)



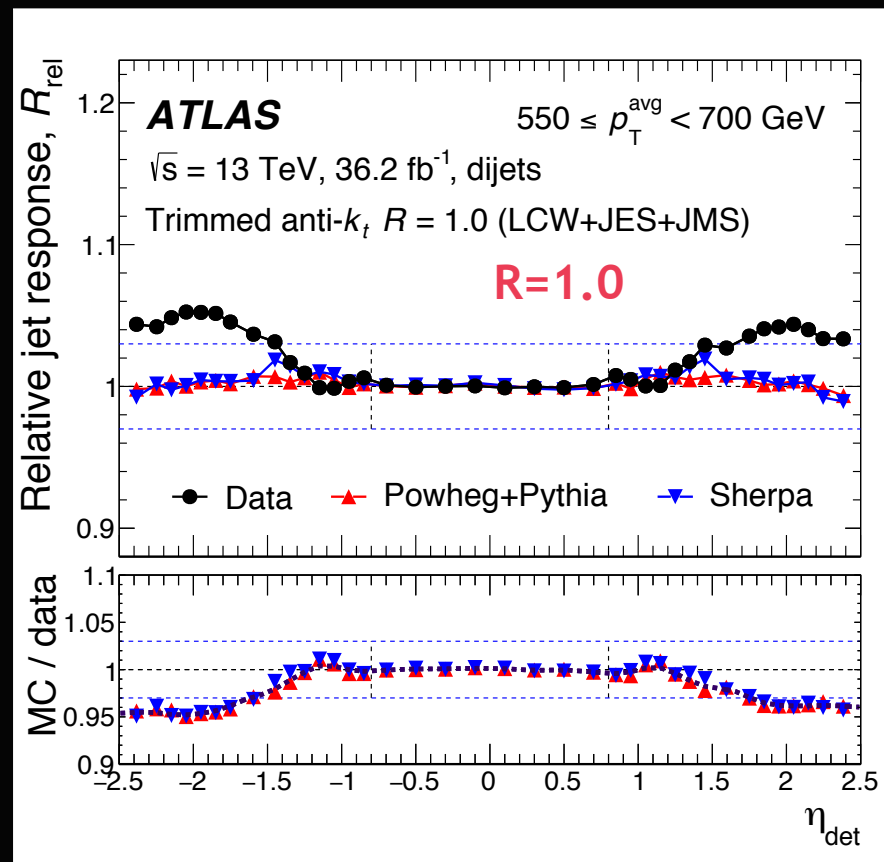
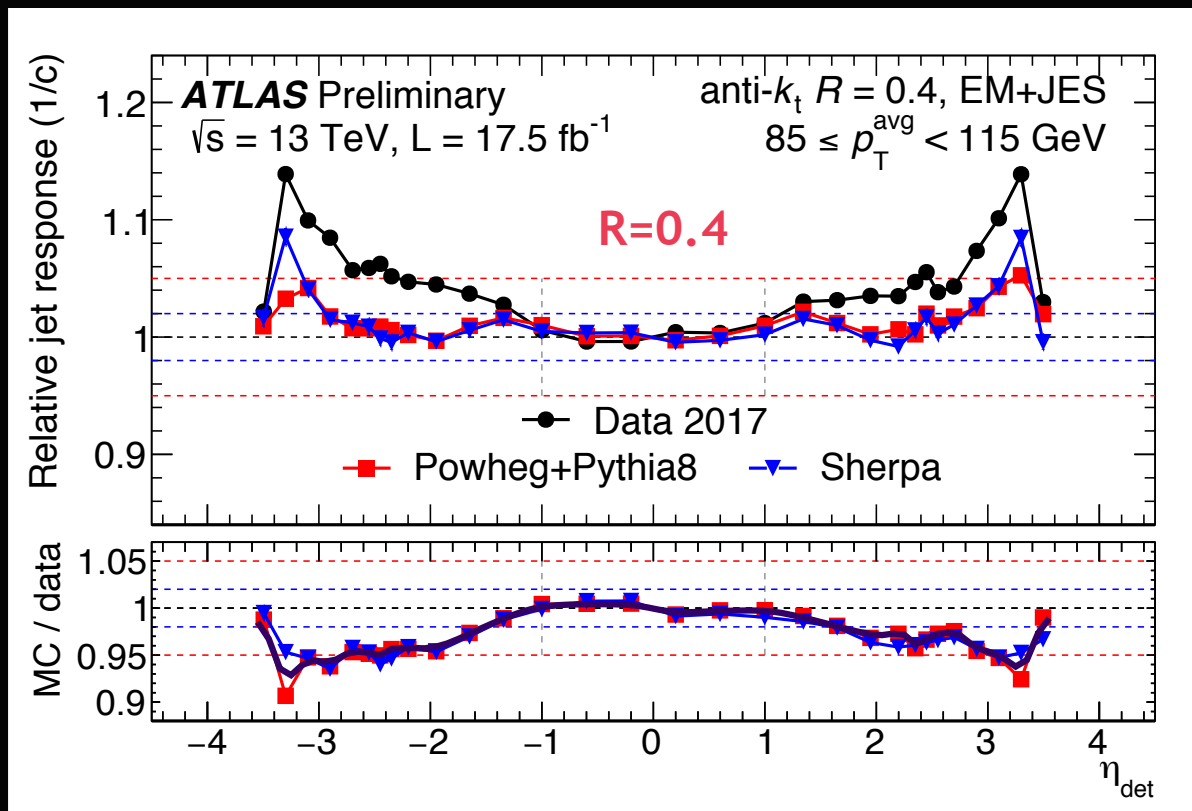


$$\mathcal{A} = \frac{p_T^{\text{probe}} - p_T^{\text{ref}}}{p_T^{\text{avg}}}$$

- Correct for the average pT asymmetry between central (eta<0.8) reference jets and forward (eta>0.8) probe jets with DeltaPhi > 2.5 as a function of eta(probe) based on the agreement between data and MC.

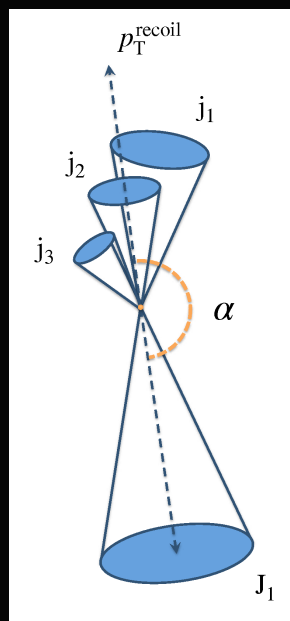
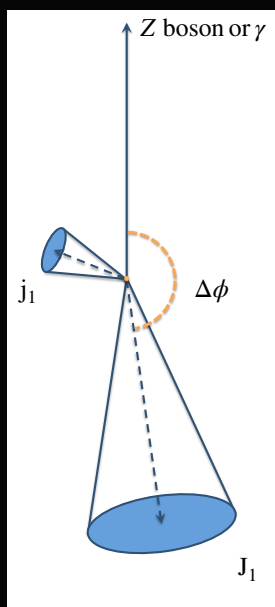
$$\left\langle \frac{p_T^{\text{probe}}}{p_T^{\text{ref}}} \right\rangle \approx \frac{2 + \langle \mathcal{A} \rangle}{2 - \langle \mathcal{A} \rangle}$$

- Relative jet response with respect to the reference region is studied given that the asymmetry is Gaussian.  $\langle \mathcal{A} \rangle$ : mean value of the Gaussian asymmetry distribution for a bin of pTavg and ηdet.



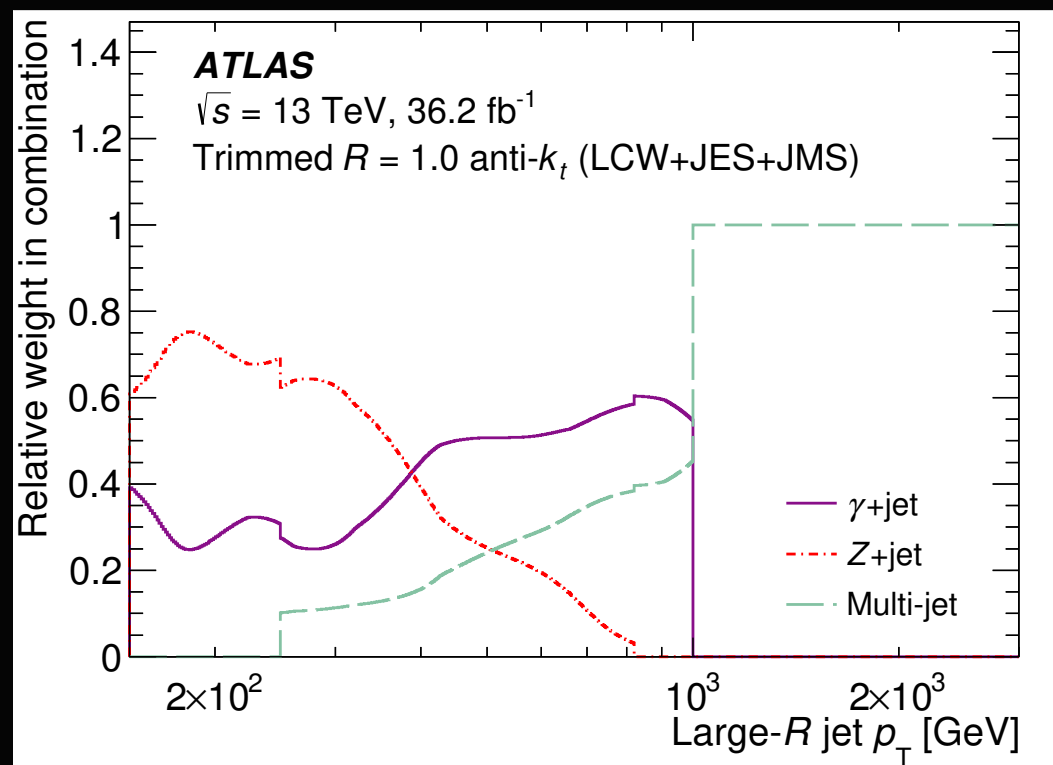


- Measure differences in average  $p_T$  balance between the jet and reference object in data and MC after the MC-based calibrations.
- Require jet to recoil against a well measured reference object
  - Gamma+jet, Z+jet and Multijet have different  $p_T$  reaches
  - Multijet:  $p_T$  balance between jet and recoiling system of calibrated small-R jets

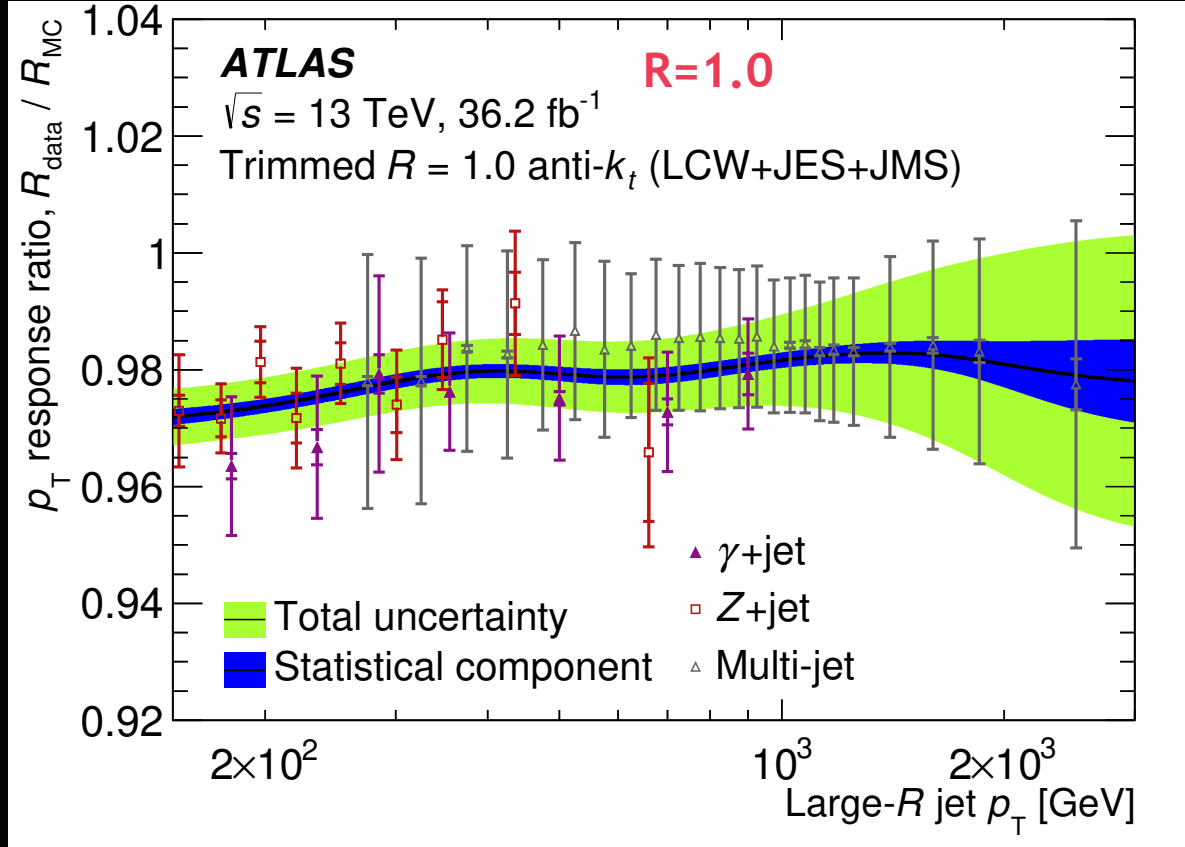
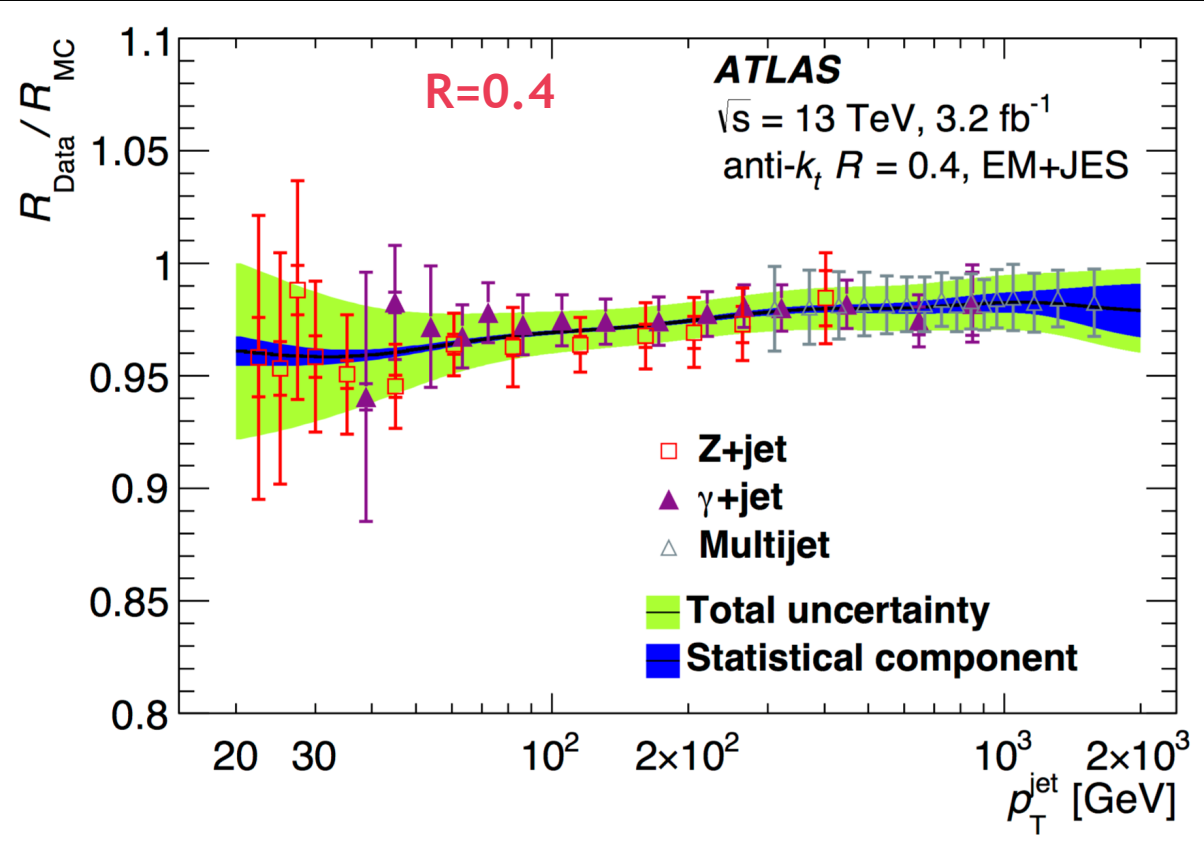


Relevance of a calibration reference:

- Energy scale uncertainty
- Trigger rate and purity
- Production cross-section



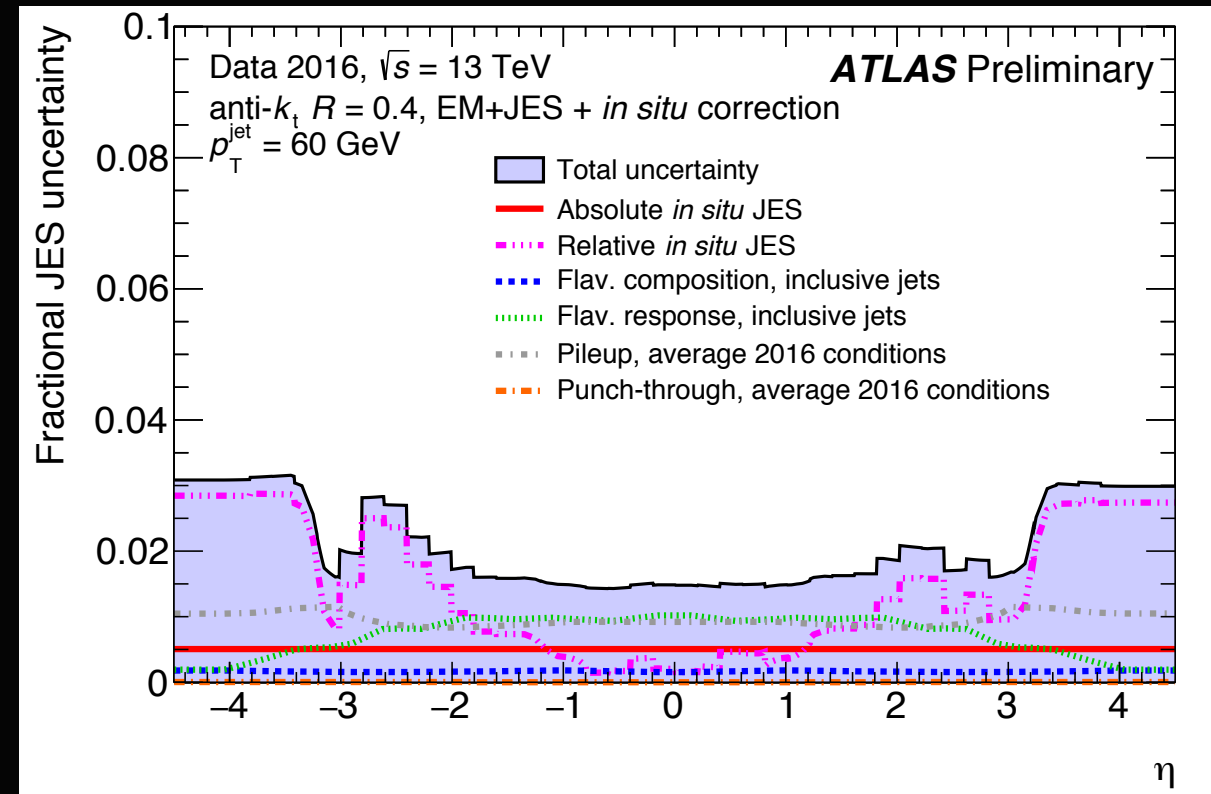
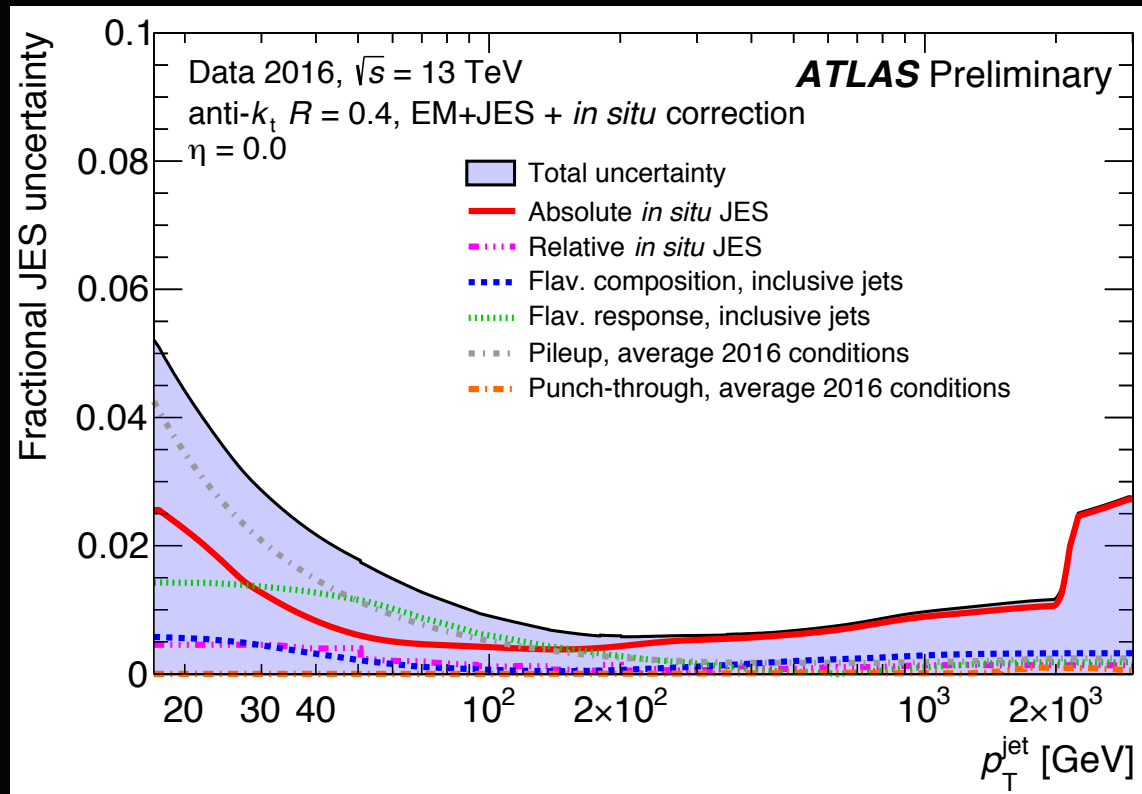
Weights for each measurement are derived by chi-square minimization using statistical and systematic uncertainties of finely-binned scale measurements





# R=0.4 jet energy scale uncertainties

- The full JES uncertainties contain *in situ* uncertainties and additional uncertainties for the modeling of pile-up, flavor composition and response differences between generators, and single particle response at the highest  $p_T$ .
- At low  $p_T$ : the pile-up uncertainties > flavor response of gluon jets > photon energy scale > single particle uncertainties.
- At high  $|\eta|$  modeling issues of the balance between forward and central jets dominate.

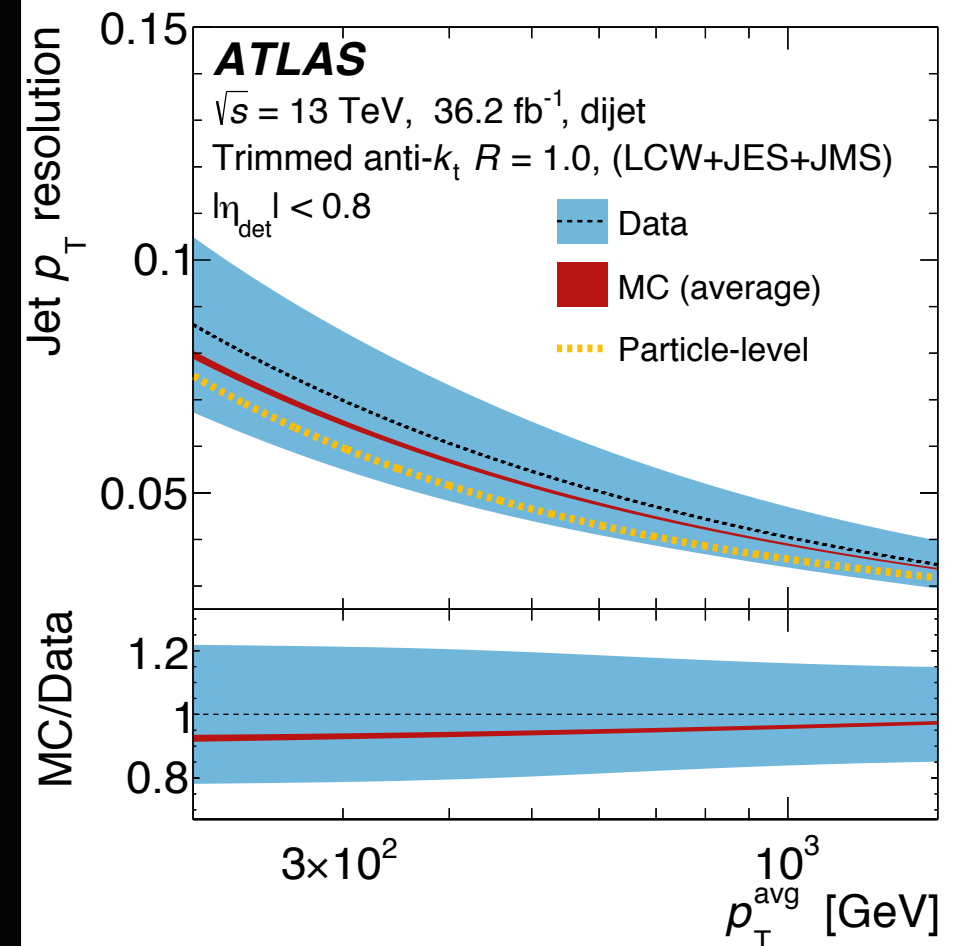
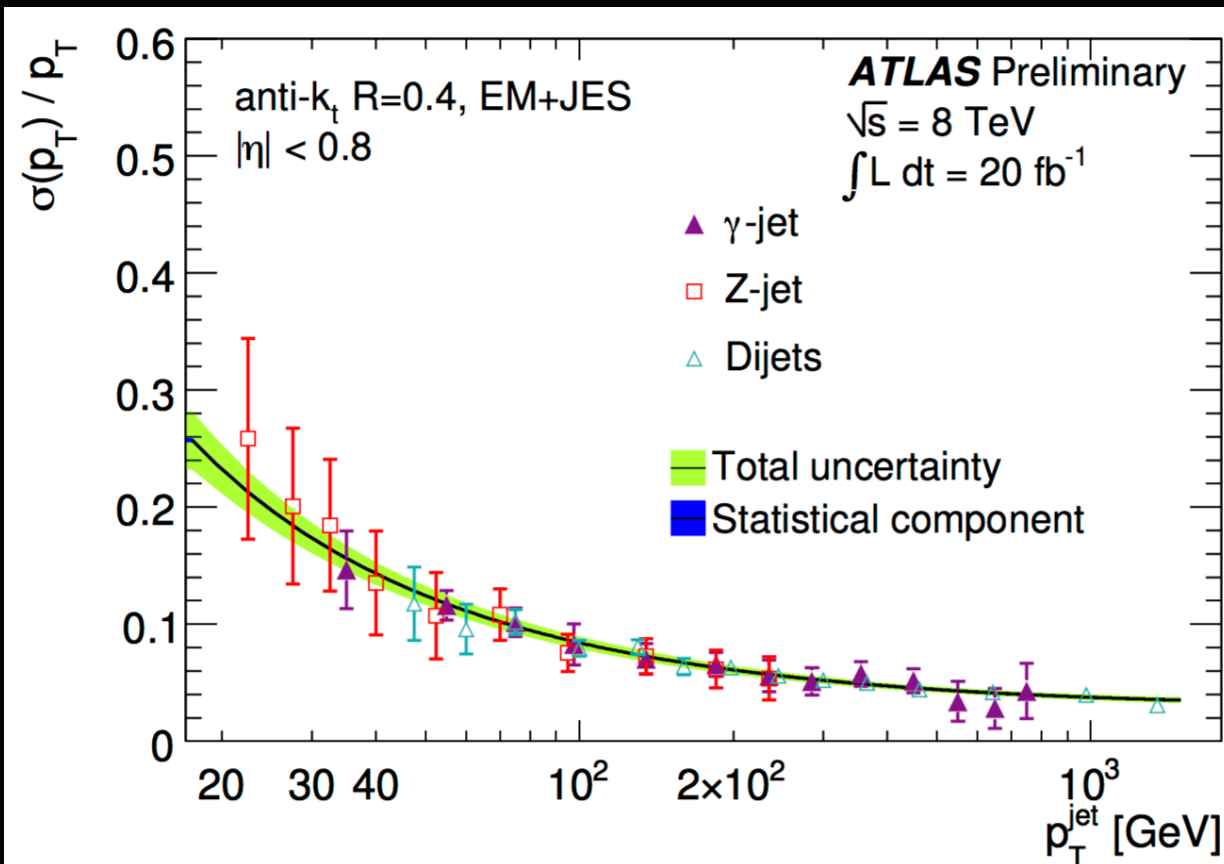




# pT resolution

For  $R=0.4$  jets, insitu pT resolution is studied by combining measurements in gamma+jet, Z+jet and Dijet events. For  $R=1.0$  jets, dijet events are used.

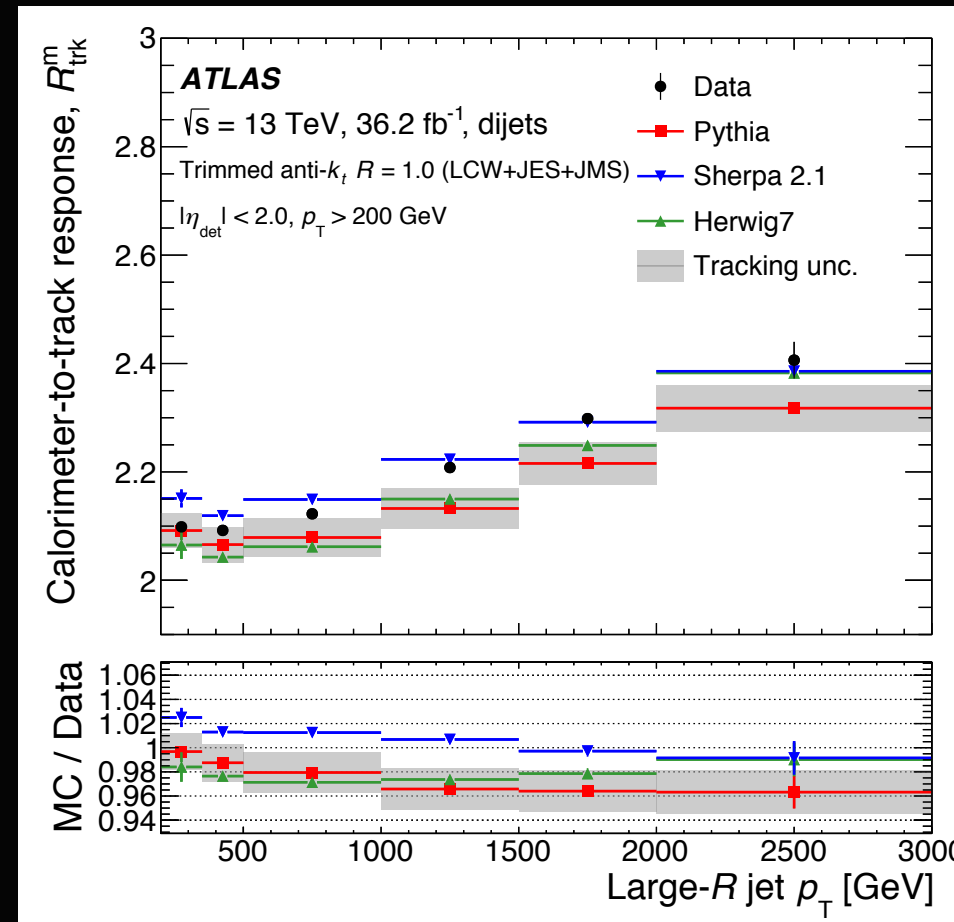
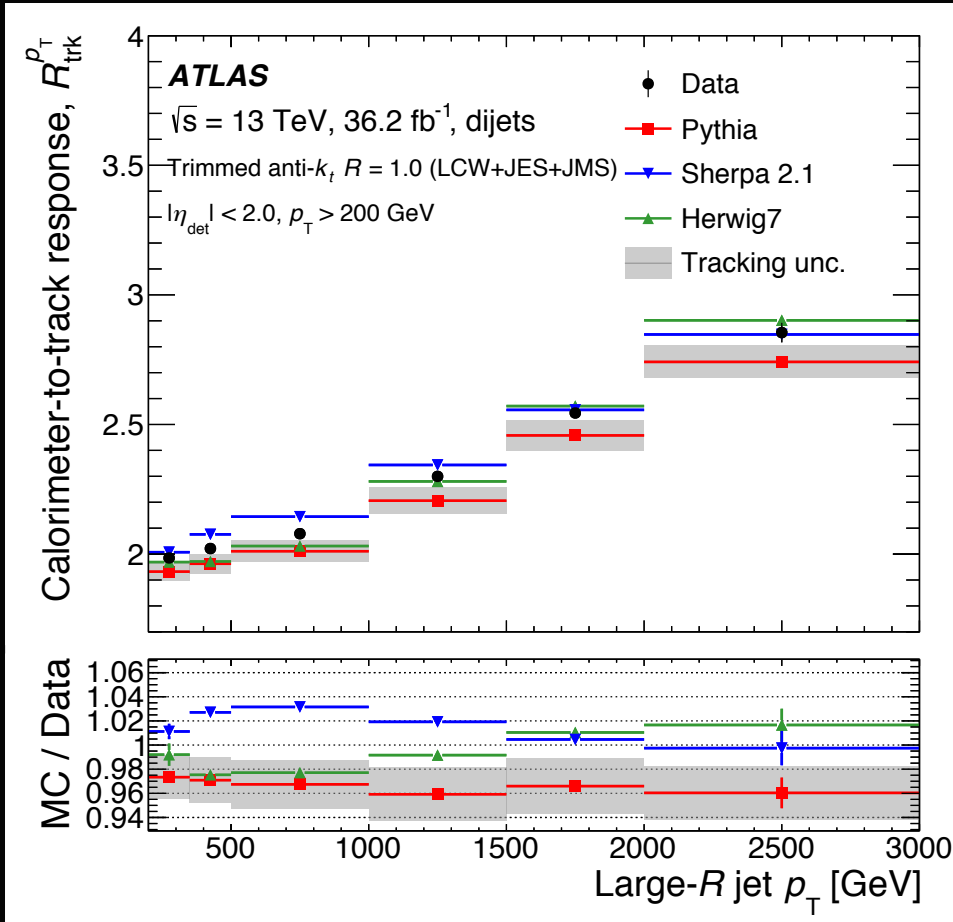
$$JER \equiv \left( \frac{\sigma[p_T]}{p_T} \right) = \sigma_{A,det} / \sqrt{2}$$





- Average calorimeter-to-track jet response is proportional to the average calorimeter-to-truth jet response.
- Comparing Rtrk in simulation and data is a way to validate the modeling of large-R jets in data:
  - Any deviation from 1 is taken to be a scale uncertainty in the measurement.
- This method is used to determine uncertainties on pT, mass and substructure information of R=1.0 jets.

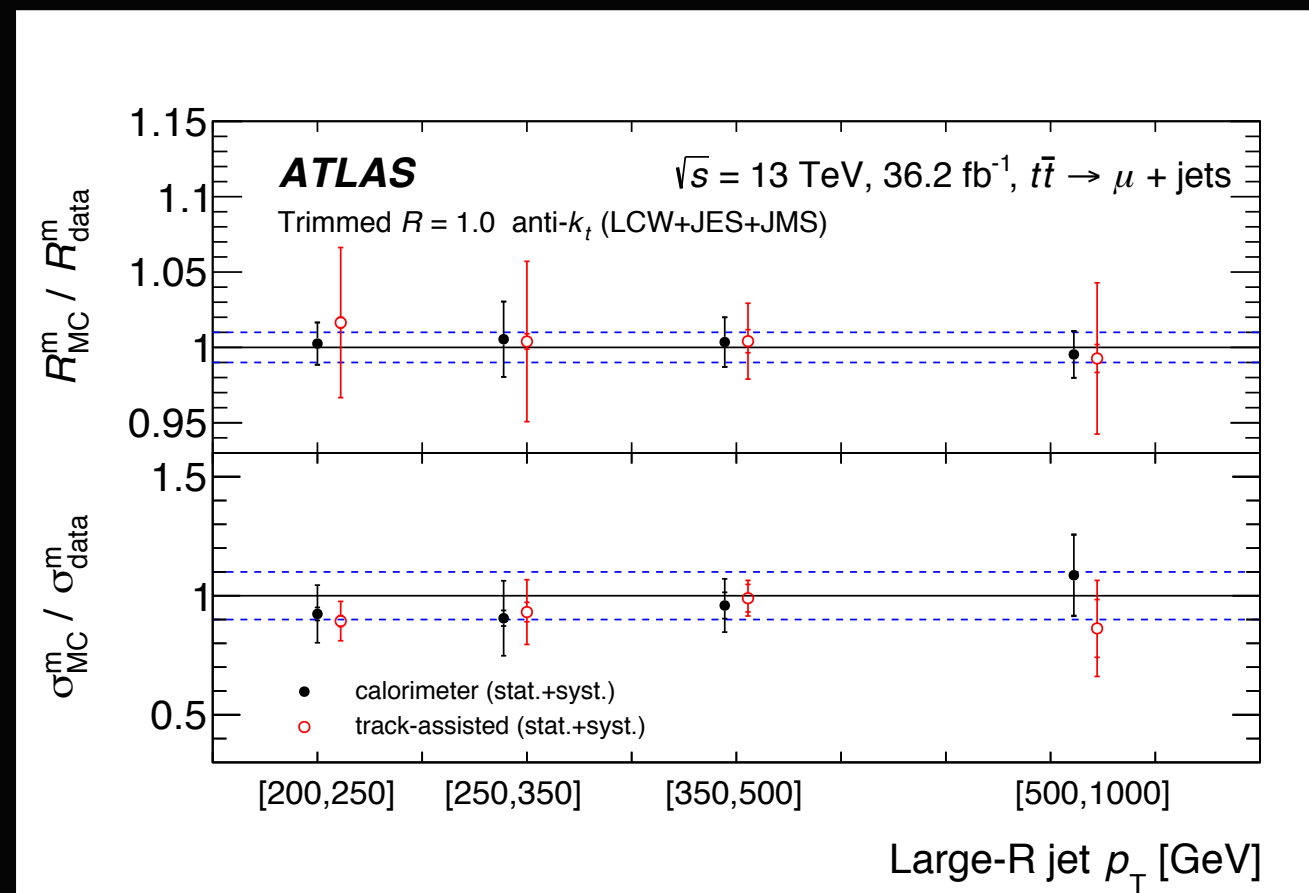
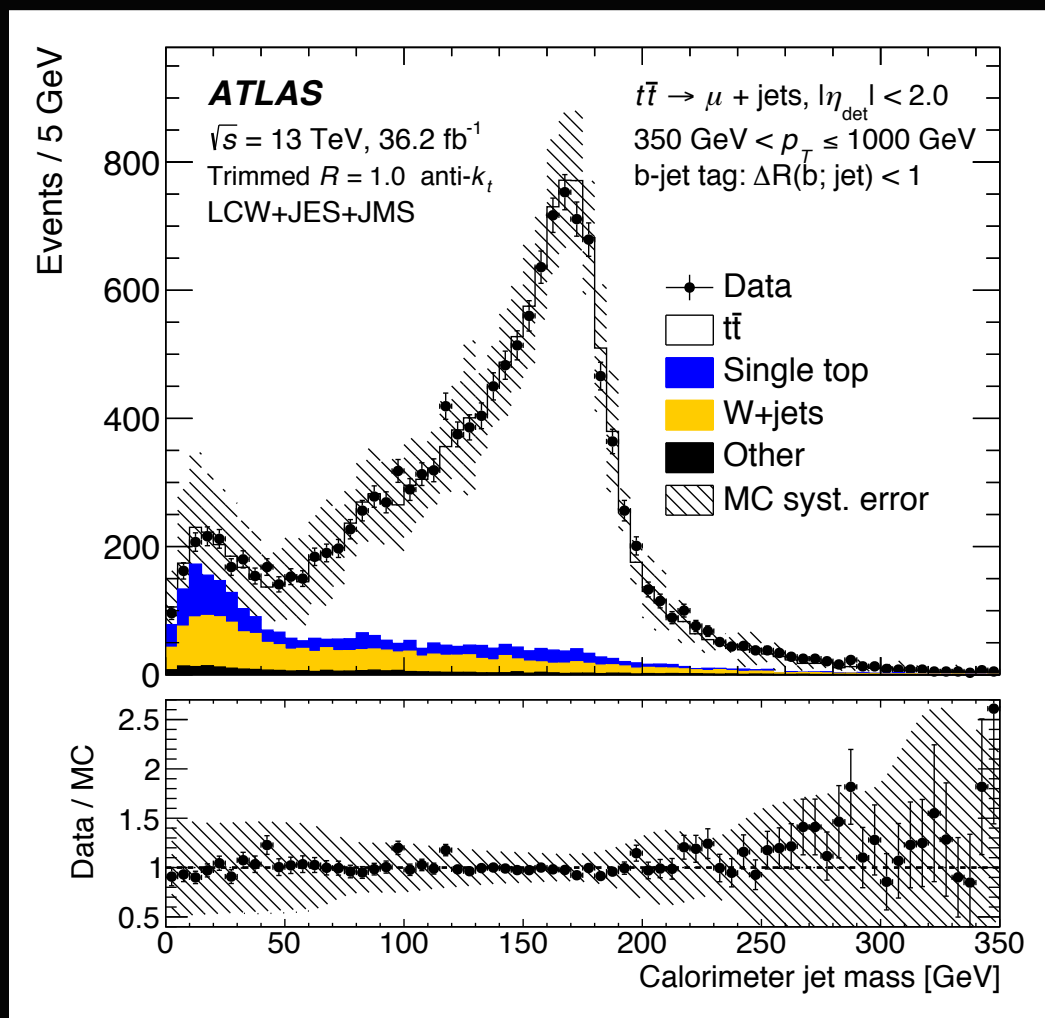
$$R_{\text{trk}} = \left\langle \frac{p_T^{\text{calo}}}{p_T^{\text{track}}} \right\rangle$$



Breakdown of uncertainties on pT and mass response for large-R jets with values of  $m/p_T \approx 0.2$  using the Rtrk method



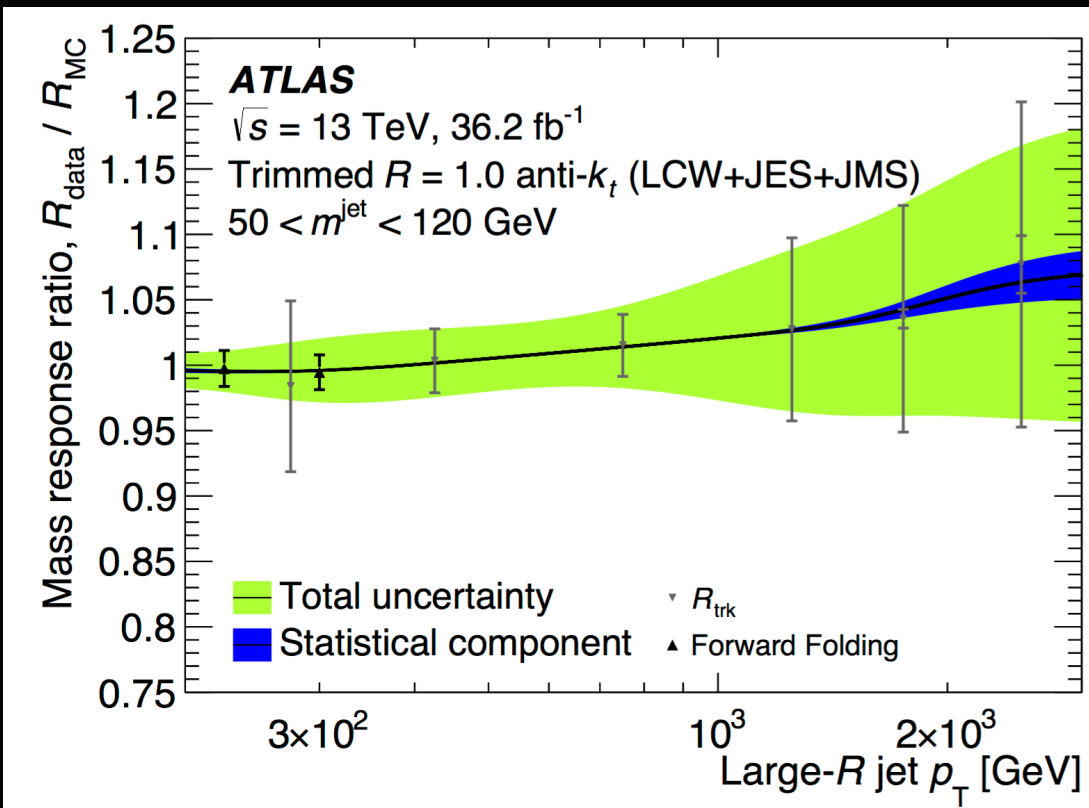
- Produces simulation predictions of the jet mass spectrum with variable response and resolution.
- Ratio of the mass response in data and simulations ( $s = R_{\text{mdata}}/R_{\text{mMC}}$ ) and of the mass resolution in data and simulations ( $r = \sigma_{\text{mdata}}/\sigma_{\text{mMC}}$ ) are extracted from the jet mass spectrum.
- Done by folding particle-level jets with a response function.



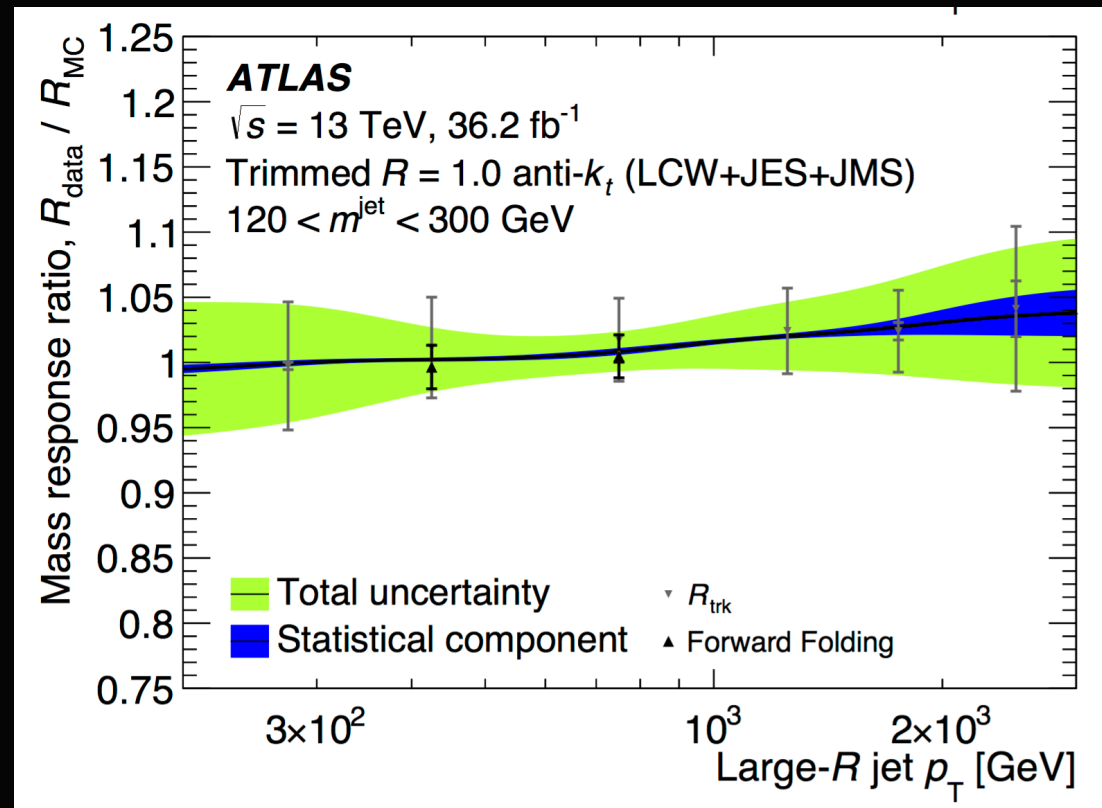


# R=1.0 jets: Combined (Rtrk and fwd folding) measurement of Jet mass scale

- Forward folding provides four measurements for  $p_T < 1$  TeV.
- Rtrk method extends the measurement to  $\sim 2$  TeV.
- Found to be consistent with one.



Fit in W mass range



Fit in top mass range



# Summary

- Jets can be used to study quark-gluon initiated jets or to reconstruct hadronic decays of massive particles like top, W/Z and Higgs.
- Jets can be reconstructed from various inputs
  - Using tracking information can improve pT or substructure resolution.
- Several in situ calibration methods are used to measure the response of the ATLAS detector to small-R and trimmed large-R jets.
- For R=0.4 jets, the uncertainty on the jet energy scale derived from data is  $< 1\%$  for  $0.1 < p_T < 1$  TeV in the central region
- For R=1.0 jets:
  - Uncertainty on the jet energy scale derived from data is 1-2% for pT from 150 GeV to 2 TeV.
  - Jet mass scale precision varies from 2% to 10% over the same pT range.
  - The in situ JES calibration, derived from light quark and gluon jets, is found to fully correct the energy and mass scales of high pT W bosons and top quarks to within the precision of the present measurement (1-3%).
- Effort continues to measure JES more precisely using the larger full Run 2 dataset.

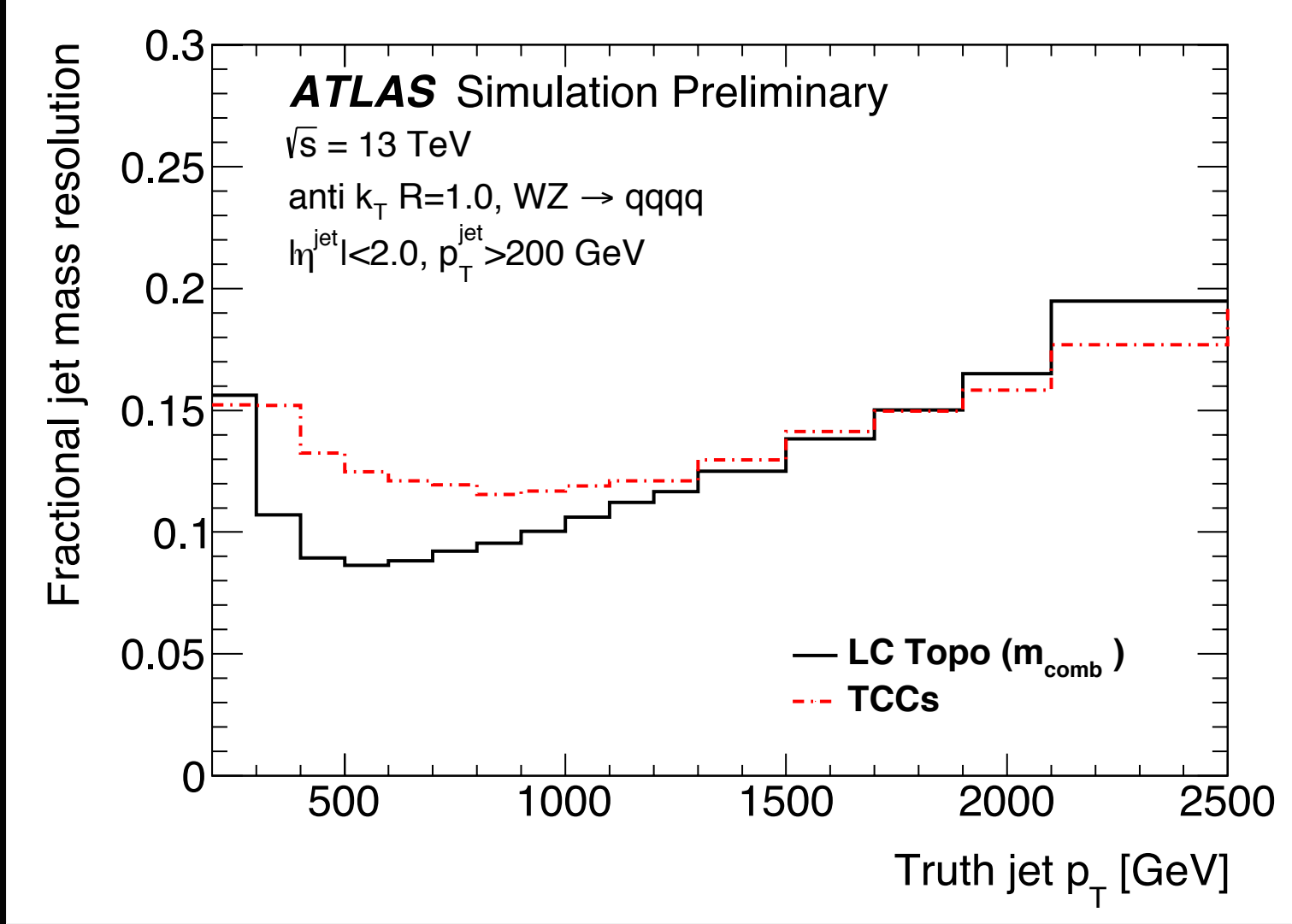


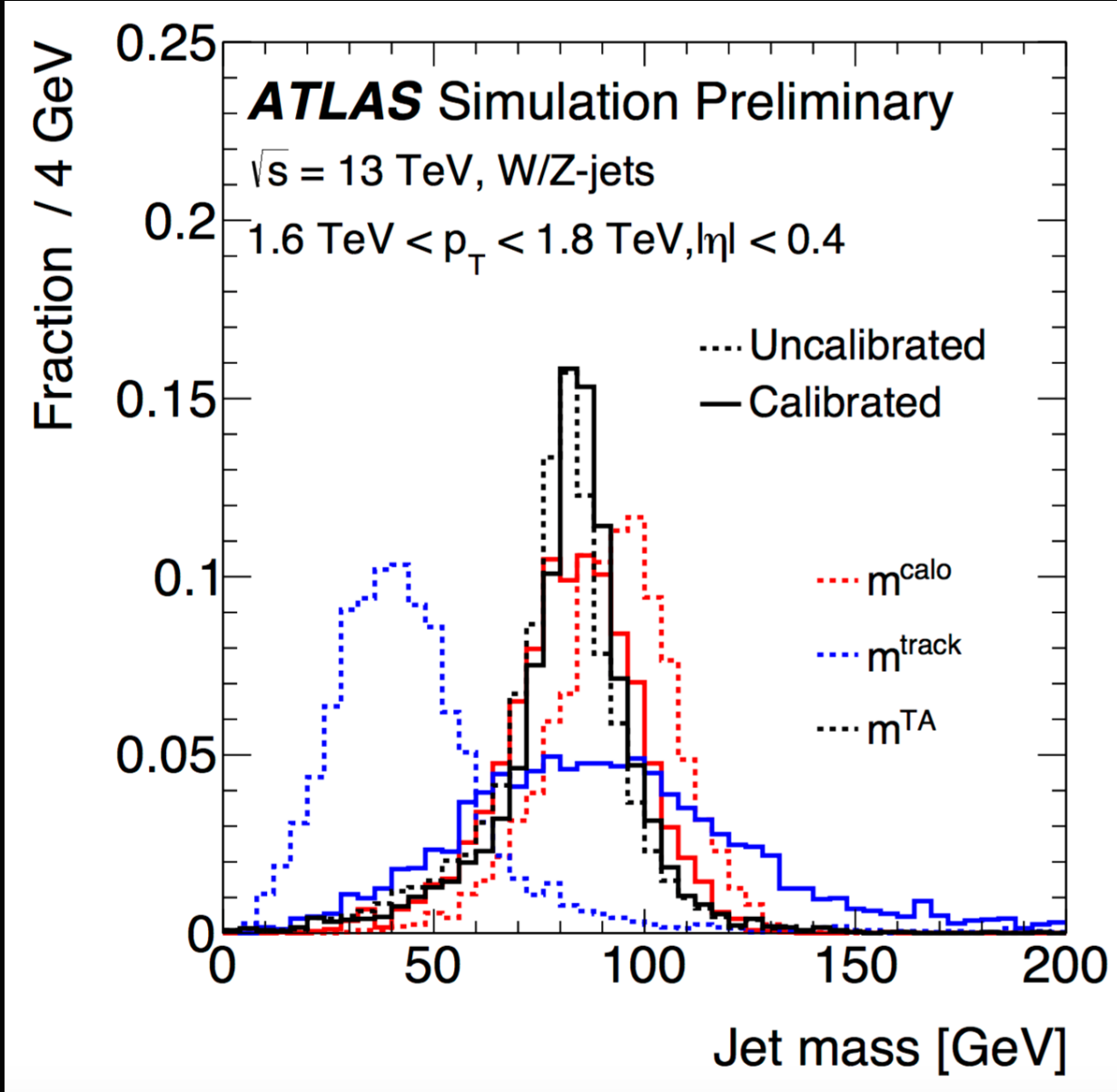
# flow subtraction procedure for different cases (1)

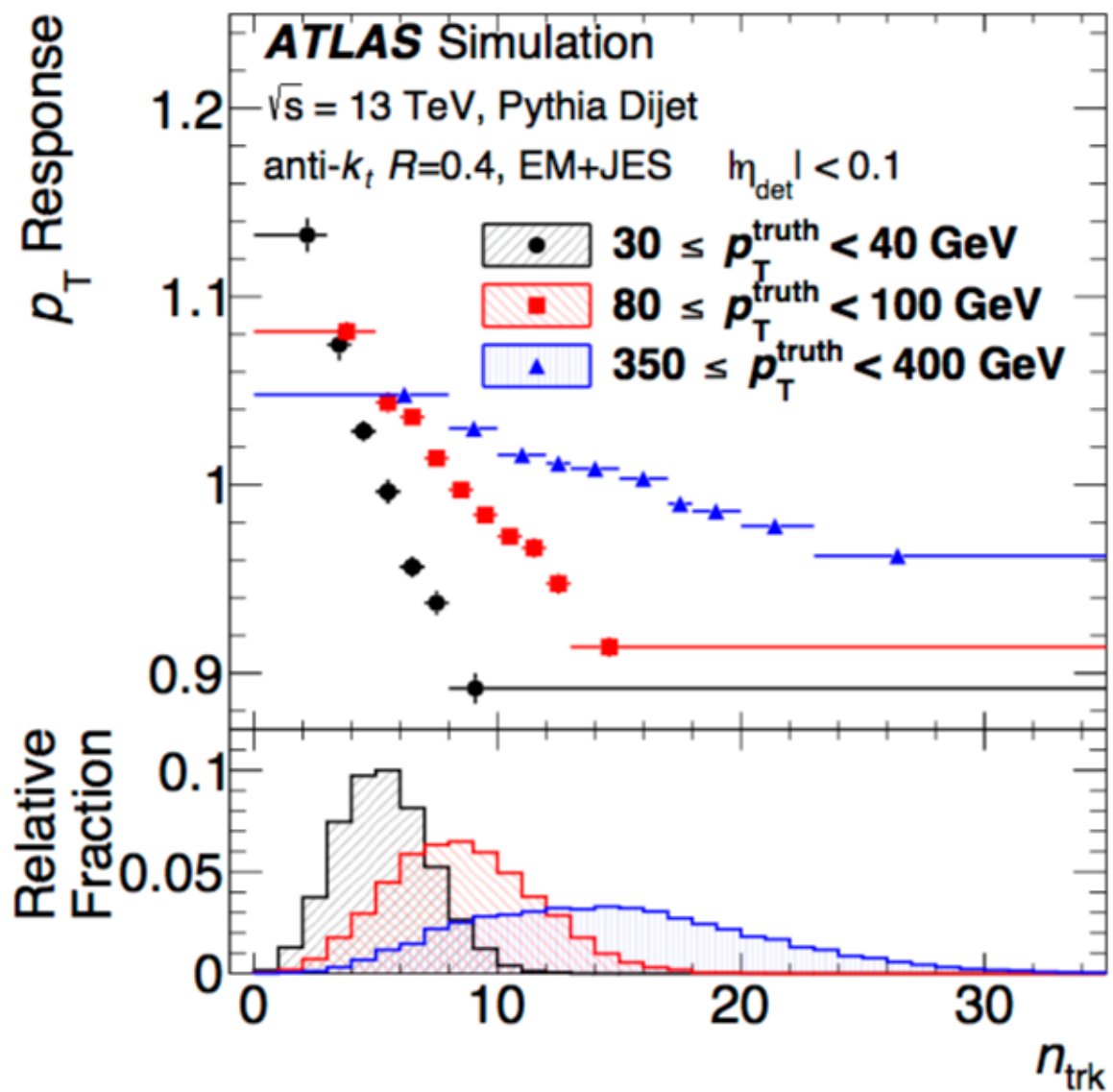
	Track/topo-cluster matching	Split shower recovery	Cell subtraction	Remnant removal
1 particle, 1 topo-cluster				
1 particle, 2 topo-clusters				

# flow subtraction procedure for different cases (2)

	Track/topo-cluster matching	Split shower recovery	Cell subtraction	Remnant removal
2 particles, 2 topo-clusters				
2 particles, 1 topo-cluster				







# Calibration procedure of calorimeter signals

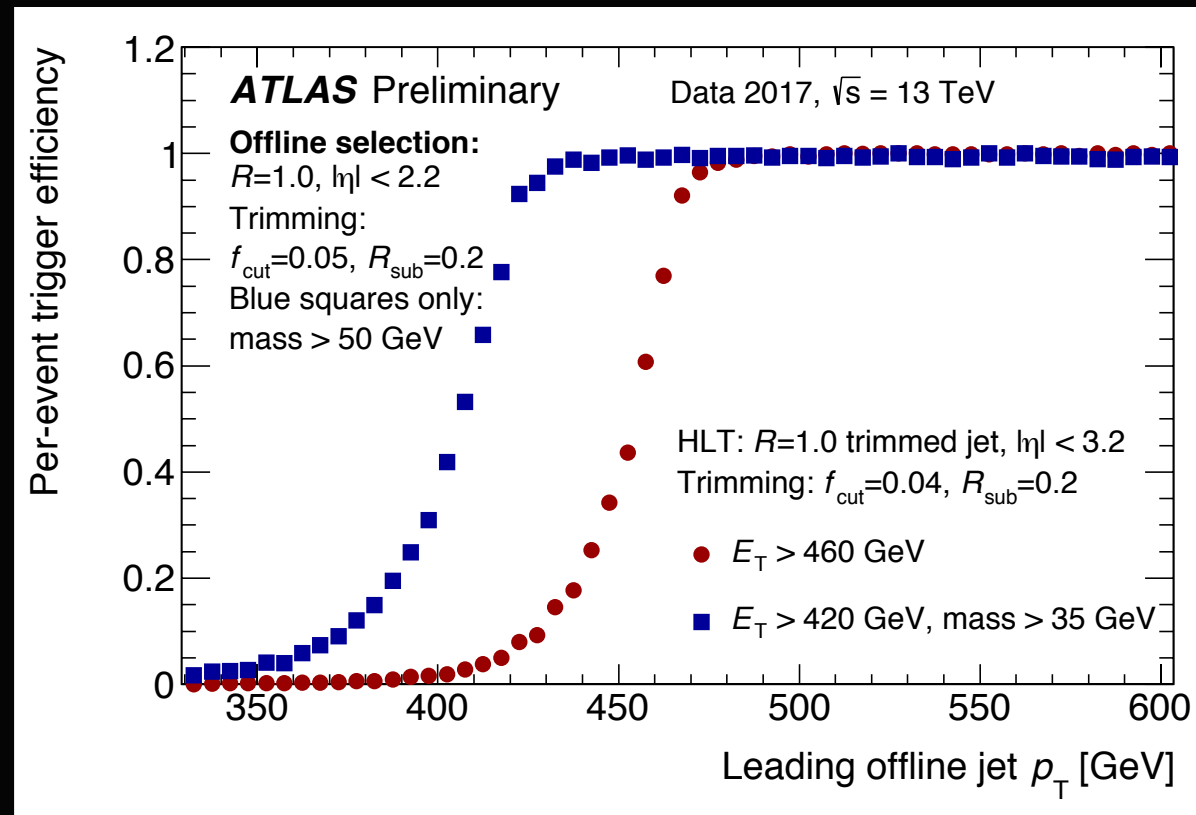
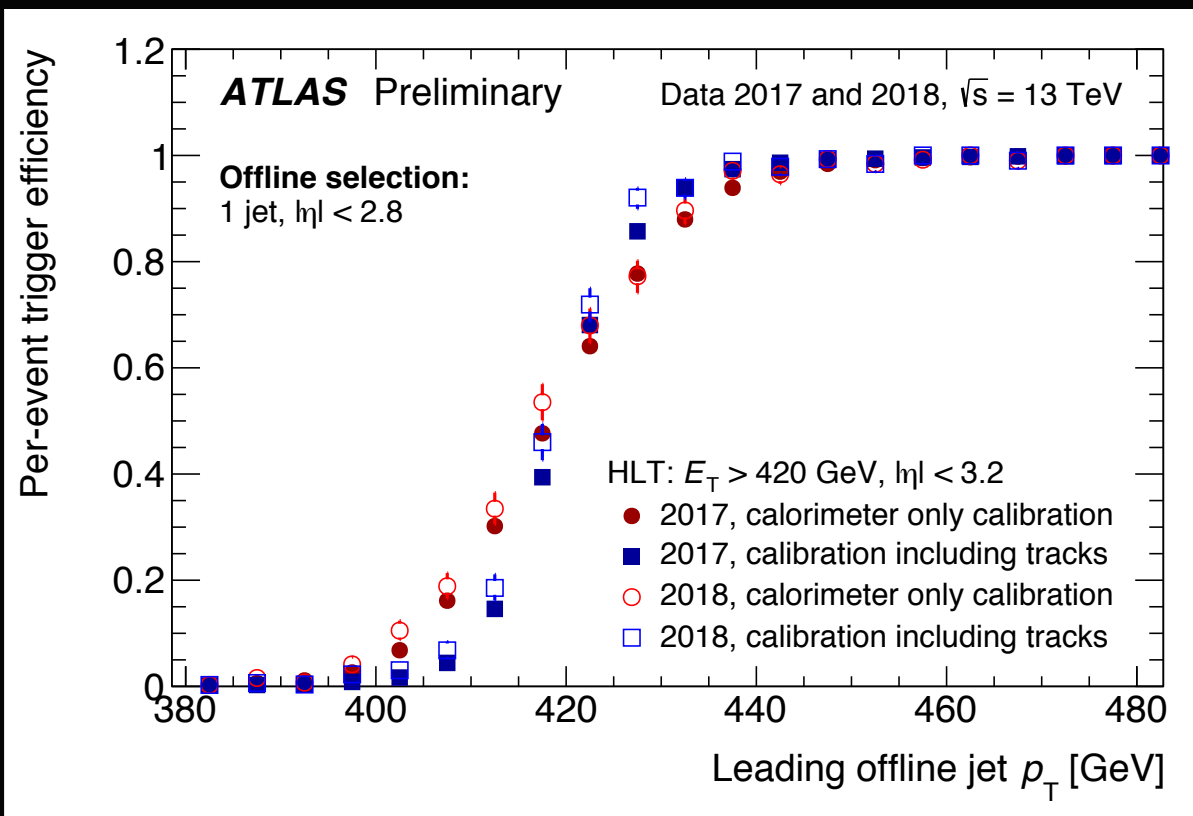
---



- Three types of calibration runs for extraction of electronic calibration constants: pedestal, ramp and delay
  - Pedestal: measurement of baseline level and noise properties of the readout electronics
  - ramp: measurement of readout gain
  - delay: measurement of pulse shape as a function of time.
- These special calibration runs are acquired between LHC fills, in absence of collisions

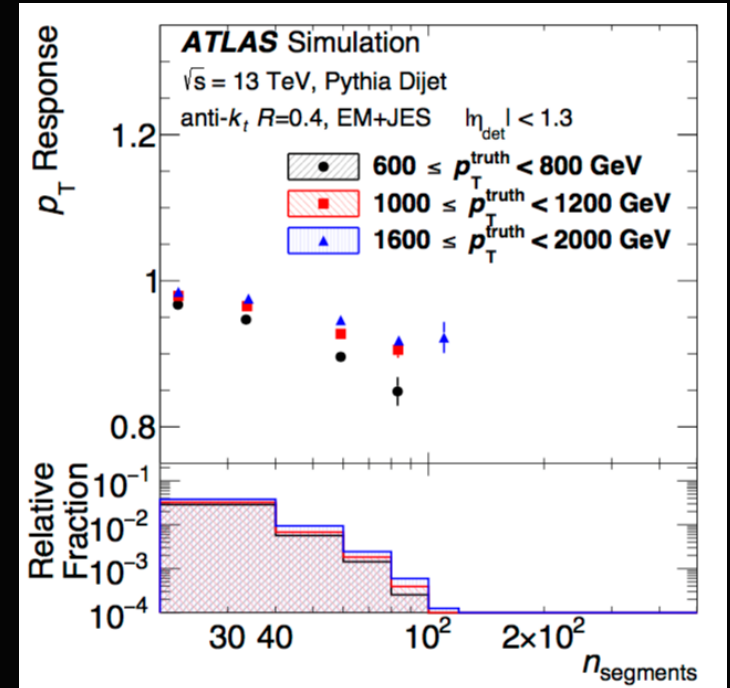
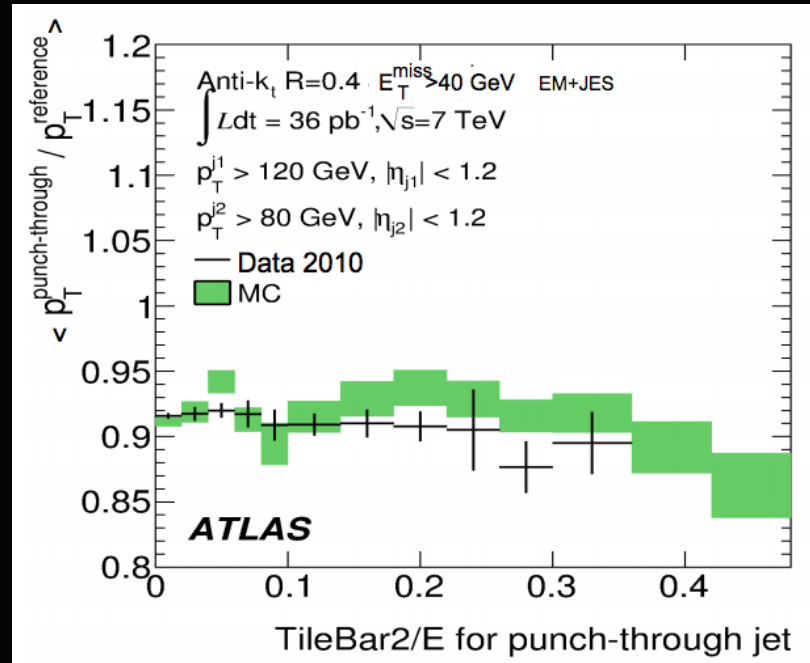
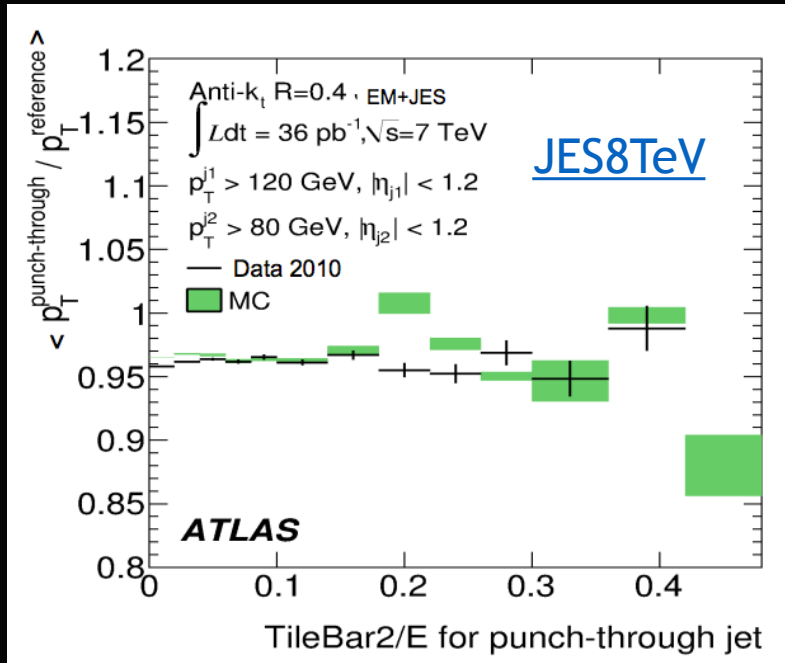


Efficiencies for single  $R=0.4$  and  $R=1.0$  jet triggers as a function of the leading offline trimmed jet  $p_T$



# Study of jet punch-through

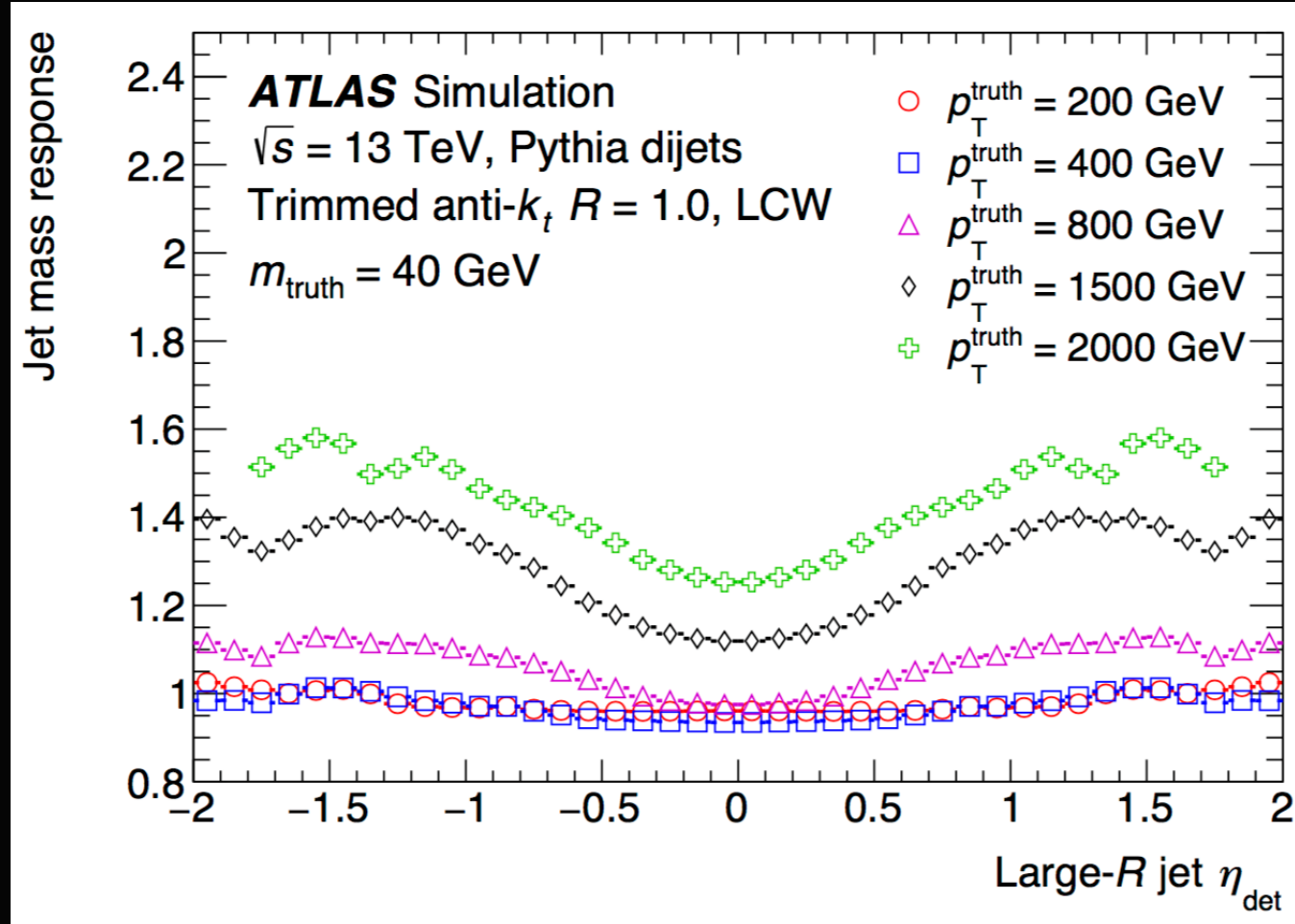
- nsegments: number of muon track segments ghost-associated with the jet ( $|\eta(\text{det})| < 2.7$ ): targeting jets that are not fully contained in the calorimeter (punch-through jets)
- Jets that deposit energy beyond the hadronic Tile calorimeter and in the muon system
  - Systematic reduction of measured jet energy
  - Can happen in any detector pseudo rapidity region
- Dijet  $p_T$  balance technique:
  - Asymmetry between transverse momentum of reference jet ( $p_{T\text{reference}}$ ) and punch-through jet as a function of energy deposition of the latter jet
  - Cannot know *a priori* which jet will be affected by punch-through effect
  - Use missing transverse energy ( $E_{T\text{miss}}$ ): energy lost beyond calorimeter creates a component of missing transverse energy in the direction of punch-through
    - Jet closest to  $E_{T\text{miss}}$  Phi direction selected as punch-through jet



# Simulation based mass response correction ( $m=40$ GeV)

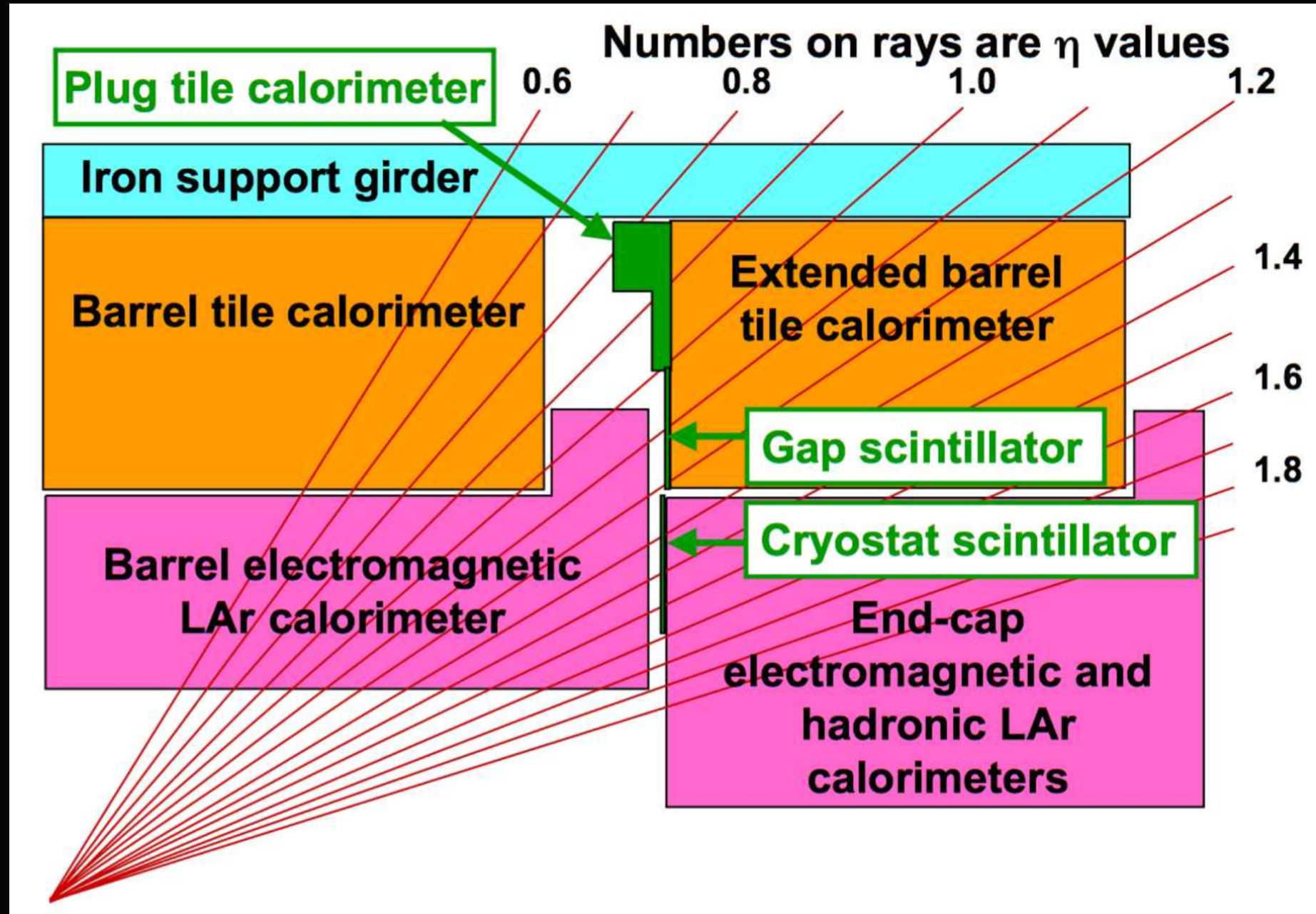


- Jet mass more sensitive to soft, wide angle radiation and topocluster merging and splitting and calorimeter geometry than  $p_T$ .
- Correction (cJMS) applied after Jet energy response correction, R jet energy kept fixed and  $p_T$  allowed to vary.



$m=40$  represents a typical value for quark or gluon jets

# Calo transition regions



# Summary of in-situ calibration factors applied in data



Full calibration applied to jets in data impacts the reconstructed jet energy, mass, eta, and pT:

$$E_{\text{reco}} = c_s \sqrt{E_0^2 + c_{\text{JMS}} m_0 (c_m^2 - 1)}, \quad m_{\text{reco}} = c_s c_{\text{JMS}} c_m m_0, \quad \eta_{\text{reco}} = \eta_0 + \Delta\eta,$$
$$p_{\text{T}}^{\text{reco}} = c_s \sqrt{(E_0^2 - c_{\text{JMS}}^2 m_0^2) \cosh(\eta + \Delta\eta)},$$

Pileup correction: 0.4 jets

$$p_{\text{T}}^{\text{corr}} = p_{\text{T}}^{\text{reco}} - \rho \times A - \alpha \times (N_{\text{PV}} - 1) - \beta \times \mu$$



# Uncertainty on R-1.0 jet fractional pT resolution

