



## Study of the $\omega\pi^0$ system

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(Joint CERN - IHEP experiment)

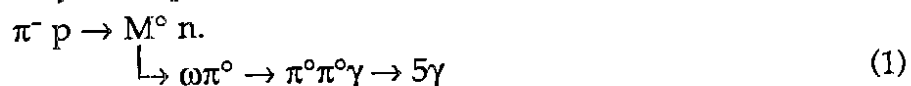
### Abstract

The mass spectra of the  $\omega\pi^0$  system, produced in  $\pi^-p$  charge-exchange reactions at 38 GeV/c and 100 GeV/c, show a new neutral state with a mass of  $2200 \pm 20$  MeV and a width of  $260 \pm 50$  MeV. Like  $\rho_3(1690)$ , this meson, X(2200), is produced through one-pion-exchange. X(2200)  $J^{PC}$  is likely  $1^{--}$ . It might be a  $2^3D_1$  radial-orbital excitation of  $\rho^0$ . Its production cross section falls with energy quadratically, in agreement with one-pion-exchange. The cross sections for  $b_1(1235)$  and  $\rho_3(1690)$  production have also been measured; their energy dependences have been determined.

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The study of  $\omega\pi^0$  systems produced in the charge-exchange reaction



offers the possibility to investigate isovector  $q\bar{q}$  mesons and to search for exotic states which do not contain strange quarks. This reaction has been studied previously at low energies [1,2].

The region of the mass spectrum above  $\rho_3^0(1690)$  which has not yet been thoroughly investigated is the main subject of this work. Measurements have been carried out with 38 GeV/c pions at IHEP using the multiphoton hodoscope spectrometer GAMS-2000 [3] and with 100 GeV/c pions at CERN with GAMS-4000 [4].

The results of the analysis of the 100 GeV/c data have been preliminary reported [5]. Higher statistics have been accumulated during four 38 GeV/c runs in 1983, 1984 and 1987. This has allowed in particular a more thorough study of the high mass spectrum produced in reaction (1).

The experimental setup, measurement conditions, calibration technique and geometrical reconstruction of events have been described earlier [3,4,6]. Some additional cuts have been introduced to select  $\omega\pi^0$  systems at 38 GeV/c. First, events with more than two out of ten possible  $2\gamma$  combinations having an effective mass in a rather wide range around the  $\pi^0$  mass ( $50 \div 250$  MeV) have been rejected. Such a cut suppresses the background from multiphoton systems which have lost one or more gammas. It also excludes events in which the selection of the  $\pi^0$  pair may be ambiguous. Second, events with two  $\pi^0\gamma$  compatible combinations, that have a mass close to that of  $\omega$ , have been excluded from further analysis. The  $\omega$  mass interval for this cut is also chosen to be wide,  $630 \div 930$  MeV. These two cuts eliminate the combinatorial background arising from misidentified  $\omega$  or  $\pi^0$ . Third, an energy threshold has been required for the free  $\gamma$  in the decay of  $\omega$  to  $\pi^0\gamma$  ( $E_\gamma > 2$  to 3 GeV). This cut suppresses the background of  $\pi^0\pi^0$  events associated with an additional noisy low energy gamma due to the high load of GAMS-2000 counters close to the beam [6].

Similar selection criteria had also been applied to the 100 GeV/c data. The main difference is that the suppression of the combinatorial background relies on a choice based on a fit of each event with constraints on the masses of the decay particles amongst several possible channels ( $\pi^0\pi^0\gamma$ ,  $\pi^0\eta\gamma$ ,  $\eta\eta\gamma$  etc.). The most probable channel is chosen. Only those  $\pi^0\pi^0\gamma$  events in which the mass of at least one  $\pi^0\gamma$  combination was found within the  $\omega$ -range ( $745 \div 825$  MeV) after the fit are kept as  $\omega\pi^0$ .

The efficiency of the  $\omega$  selection procedure is illustrated in fig. 1. The  $\omega$ -peak width is determined by the resolution of GAMS. This is equal to 50 MeV (FWHM) at 100 GeV/c and 100 MeV at 38 GeV/c. The peak-to-background ratio equals 5 on the average. It improves with increasing  $\omega\pi^0$  mass (up to 15 for masses larger than 2 GeV, fig. 1b).

The  $\omega\pi^0$  mass spectrum in reaction (1) at 100 GeV/c (fig. 2a) has been obtained after background subtraction under the  $\omega$ -peak in each 50 MeV  $\omega\pi^0$  mass bin and correction for efficiency. It shows, alongside the well-known  $b_1(1235)$  and  $\rho_3(1690)$  mesons, a peak in the 2.2 GeV region which is more pronounced at low  $|t|$  (fig. 2c). This  $3.5\sigma$  peak [5] is the first evidence of the presence of a vector meson that is approximately 300 MeV wide and has a mass of  $2170 \pm 30$  MeV.

The  $\omega\pi^0$  mass spectrum at 38 GeV/c (fig. 2b,d) has been obtained after a 4C-fit (the masses of the recoil neutron, of the two  $\pi^0$  and  $\omega$  being fixed) and correction for efficiency. It shows the same peak at 2.2 GeV with a much better statistical significance, larger than  $10\sigma$ , for the low momentum transfer events. The peak (hereafter named X(2200)) contains  $3500 \pm 200$   $\omega\pi^0$  events. The mass and width of X, evaluated by fitting the  $\omega\pi^0$  mass spectrum of fig. 2d with two Breit-Wigner resonances and a polynomial continuum, are :  $M = 2220 \pm 20$  MeV,  $\Gamma = 240 \pm 60$  MeV (corrected for GAMS resolution).

Taking the 100 GeV/c data into account, the parameters of X are

$$\begin{aligned} M &= 2200 \pm 20 \text{ MeV,} \\ \Gamma &= 260 \pm 50 \text{ MeV.} \end{aligned} \quad (2)$$

The distribution of 4-momentum transfer squared (fig.3a), the angular distribution in the Gottfried-Jackson frame of reaction (1) and that in the  $\omega$  helicity frame (fig.4) of the events in the X-peak have been studied.

The  $t$ -distributions of both  $\rho_3$  and X events have been plotted after fitting the  $\omega\pi^0$  mass spectrum at 38 GeV/c in each  $t$ -bin with the sum of a polynomial background and two Breit-Wigner resonances with the listed parameters of these two states. Both show a typical one-pion-exchange (OPE) exponential behaviour  $|t|e^{bt}/(m_\pi^2 - t)^2$ , with a slope

$$b = 7.2 \pm 0.8 \text{ (GeV/c)}^{-2} \quad (3)$$

for X, in good agreement with the slope for  $\rho_3$  which is observed simultaneously (fig. 3a). A similar  $t$ -distribution is observed at 100 GeV/c in the X mass corridor.

These typical OPE distributions should be compared with those of  $b_1$  (fig. 3b) obtained using the same bin-by-bin technique (at 38 GeV/c) or  $b_1$  mass interval (at 100 GeV/c). The  $b_1$   $t$ -distribution follows an exponential function  $e^{ct}$  with

$$c = 3.4 \pm 0.2 \text{ (GeV/c)}^{-2}. \quad (4)$$

The assumption that OPE is the dominating mechanism in the X(2200) production is also confirmed by the isotropic Treiman-Yang angular distribution and by the characteristic angular distributions in the  $\omega$  helicity frame (fig. 4). In the case of OPE without absorption, the following distributions are expected :  $d\sigma/d\cos\theta_\pi \sim 1 + \cos^2\theta_\pi$ ,  $d\sigma/d\phi_\pi \sim 1/2 + \cos^2\phi_\pi$  (here  $\theta_\pi$ ,  $\phi_\pi$  define the  $\pi^0$  direction in the  $\omega \rightarrow \pi^0\gamma$  decay frame). The experimental data follow nicely these formulas.

The conclusion is that the observed  $\omega\pi^0$  resonance, X(2220), is produced in reaction (1) through OPE exchange, which limits its quantum numbers to the following :  $J^{PC} = 1^{--}, 3^{--}, \dots, I^G = 1^+$ .

The angular distributions in the Gottfried-Jackson frame have been fitted with the sum of a constant background and the theoretical distributions expected for  $J = 1$  and for  $J = 3$  in OPE. In this model, which does not take into account a possible interference of X(2200) with the background and with the  $\rho_3$  "tail", the fit favours the hypothesis  $J^{PC} = 1^{--}$ . It is incompatible with spin 3 (fig. 4a).

Further evidence that the X(2200) spin is not 3 is shown on fig. 5 (38 GeV/c data). Here the  $\omega\pi^0$  spectrum that includes selected events with  $0.3 < |\cos \theta_{GJ}| < 0.6$ , an interval where the contribution of the  $3^{--}$  wave is small, shows a X peak that is considerably enhanced with respect to  $\rho_3$  and the continuum.

The peak also manifests itself within the  $0.2 < |\cos \theta_{GJ}| < 0.4$  range, situated in the vicinity of zeros of the  $5^{--}$  wave. The angular distribution (fig. 4a) does not exclude a small contribution of the  $5^{--}$  wave (at a level less than 1/4 of that of the  $1^{--}$  wave). Such a wave in the 2.2 GeV to 2.5 GeV mass region might be due to  $\rho_3^0(2350)$ . Contributions of this meson as well as of  $\rho_1^0(1700)$  to the  $\omega\pi^0$  mass spectrum will be evaluated in the partial wave analysis of the whole set of the GAMS data currently in progress.

The cross section of the X meson production in reaction (1) is equal to

$$\sigma(\pi^- p \rightarrow X n) \cdot BR(X \rightarrow \omega\pi^0) = \begin{cases} 120 \pm 25 \text{ nb} & \text{at 38 GeV/c,} \\ 20 \pm 6 \text{ nb} & \text{at 100 GeV/c.} \end{cases} \quad (5)$$

The cross section decreases with energy quadratically, in agreement with OPE.

The observed resonance is a serious candidate for a  $J^{PC} = 1^{--}, I^G = 1^+$  multiplet. It might be a radial and orbital  $2^3D_1$  excitation of  $\rho^0$ . In the Godfrey-Isgur model [7], such a meson would have a mass near 2150 MeV (cf. e.g. [8]). A similar value ( $M = 2200 \pm 60$  MeV) follows from the Veneziano formula [9] for the fourth  $\rho^0$  excitation [10]. The predictions of these two models together with experimental data are shown in fig.6. The possible existence of a  $1^{--}$  state around 2.15 GeV had also been suggested in an analysis of the  $\pi\pi$  system [11].

A preliminary analysis of the  $\rho_3$  production in reaction (1) has been performed. The  $\cos\theta_{GJ}$  distributions both for 38 GeV/c and 100 GeV/c are presented in fig. 7. They are well fitted by the spin-3 OPE curve plus a small constant term. The parameters of the  $\rho_3$  resonance determined from the 38 GeV/c and 100 GeV/c data,

$$M = 1690 \pm 20 \text{ MeV}, \Gamma = 230 \pm 30 \text{ MeV}, \quad (6)$$

are in agreement with the tabulated values, excluding the narrow width found in some older experiments [12].

The measured  $\rho_3$  production cross section

$$\sigma(\pi^- p \rightarrow \rho_3 n) \cdot \text{BR}(\rho_3 \rightarrow \omega\pi^0) = \begin{cases} 0.8 \pm 0.2 \mu\text{b} & \text{at } 38 \text{ GeV}/c, \\ 0.15 \pm 0.04 \mu\text{b} & \text{at } 100 \text{ GeV}/c, \end{cases} \quad (7)$$

together with the low-energy cross section [13], considering the different branching ratios [12], fits the OPE quadratic energy dependence well in the momentum range from 10 GeV/c to 100 GeV/c:  $\sigma \sim s^{-\alpha/2}$ ,  $\alpha = 2.0 \pm 0.3$ .

The high- $|t|$  region of reaction (1) is dominated by  $b_1(1235)$  production. The  $b_1$  parameters obtained from the 38 GeV/c and 100 GeV/c data,

$$M = 1235 \pm 15 \text{ MeV}, \Gamma = 160 \pm 30 \text{ MeV}, \quad (8)$$

agree with the tabulated values [12]. The  $b_1$  production cross section values,

$$\sigma(\pi^- p \rightarrow b_1 n) \cdot \text{BR}(b_1 \rightarrow \omega\pi^0) = \begin{cases} 1.0 \pm 0.2 \mu\text{b} & \text{at } 38 \text{ GeV}/c, \\ 0.28 \pm 0.07 \mu\text{b} & \text{at } 100 \text{ GeV}/c, \end{cases} \quad (9)$$

together with previous data at 8.45 GeV/c [2] show an energy decrease with a slope  $\alpha = 1.50 \pm 0.25$ .

The 100 GeV/c mass spectrum of the  $\omega\pi^0$  events measured with  $|t| > 0.2 \text{ (GeV}/c)^2$  suggests a bump near 1.4 GeV [5]. It is not confirmed by the 38 GeV/c data.

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## Figure captions

Fig. 1 Identification of  $\omega$  produced in reaction (1) events.

- a) Mass spectrum of  $\pi^0\gamma$  systems after identification of one  $\pi^0$  (100 GeV/c data, 38 GeV/c spectrum is similar);
- b) idem for 38 GeV/c in the  $2100 \text{ MeV} < M_{\pi^0\pi^0\gamma} < 2350 \text{ MeV}$  range. Arrows indicate the PDG value of the  $\omega$  mass. Curves show the results of the fit.

Fig. 2 Mass spectrum of the  $\omega\pi^0$  system.

- a,c) 100 GeV/c, all events and  $|t| < 0.04 \text{ (GeV/c)}^2$  events.
- b,d) 38 GeV/c, all events and  $|t| < 0.05 \text{ (GeV/c)}^2$  events. Here and further data are shown corrected for the mass and angular dependent efficiency. The curves are the sums of a polynomial background and three Breit-Wigner functions, respectively for  $b_1$  and  $\rho_3$  with their tabulated parameters and for X(2200) with the measured parameters.
- e) Id., after continuum and  $\rho_3$  subtraction.

Fig.3 a) The t-distributions (arbitrary units) of the  $\omega\pi^0$  events for X(2200) (histogram) and  $\rho_3$  (open circles) at 38 GeV/c and for  $\rho_3$  at 100 GeV/c (crosses). The curve shows the t-dependence for OPE with the exponential slope (3).

b) The same for  $b_1$ . The straight line has the slope (4).

Fig. 4 Angular distributions in the X-peak region.

a,b) Distributions in the Gottfried-Jackson frame at 38 GeV/c (histogram) and 100 GeV/c (crosses) fitted with the sum of a constant background and OPE functions for  $J^{PC} = 1^{--}$  (solid curve). The behaviour of the  $3^{--}$  wave is shown by a dashed curve (the fit gives zero contribution for this wave).

c,d) Distributions in the  $\omega$  helicity frame at 38 GeV/c fitted simultaneously with the distributions in fig. 4a,b (cf. text for curves definition).

Fig. 5 38 GeV/c mass spectrum of the  $|t| < 0.05 \text{ (GeV/c)}^2$   $\omega\pi^0$  events in the  $0.3 < |\cos\theta_{GJ}| < 0.6$  range where the  $J = 3$  contribution is small.

Fig. 6 Comparison of the experimental data with theoretical models. Dots with error bars are experimental points, open circles are Godfrey-Isgur predictions [7]. Straight lines show the Veneziano formula prediction [9] corresponding to the slope  $1.07 \pm 0.07 \text{ GeV}^2$ .

Fig. 7 Same as in fig. 4a but for the  $\rho_3$  region.

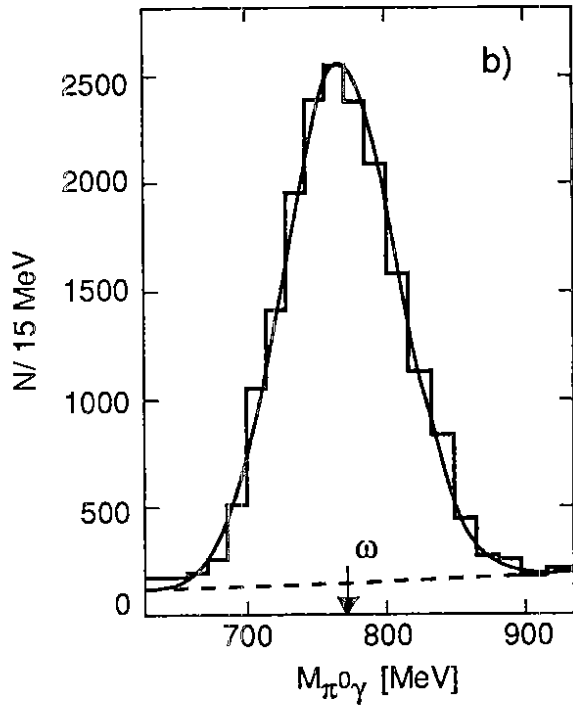
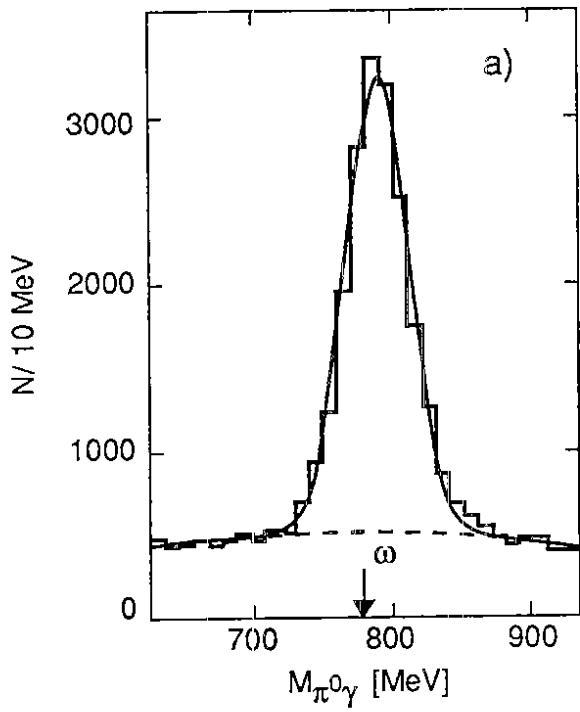


Fig. 1

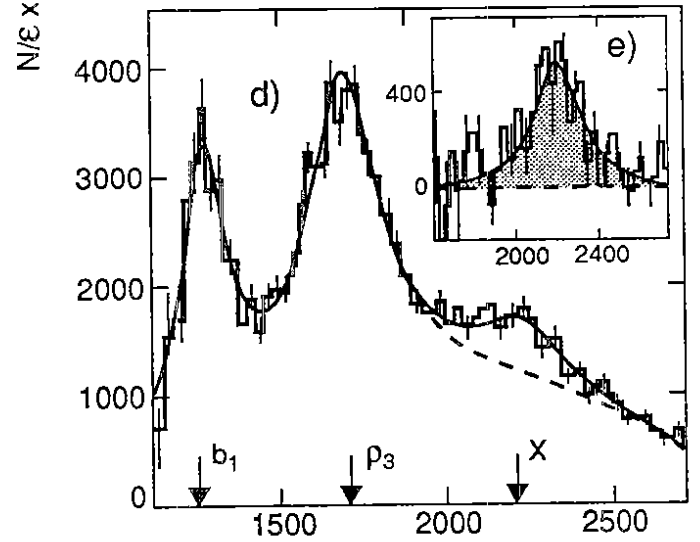
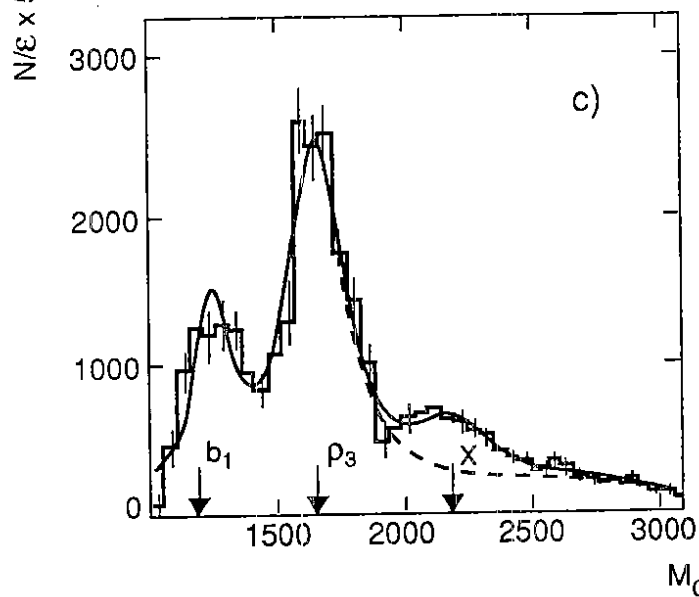
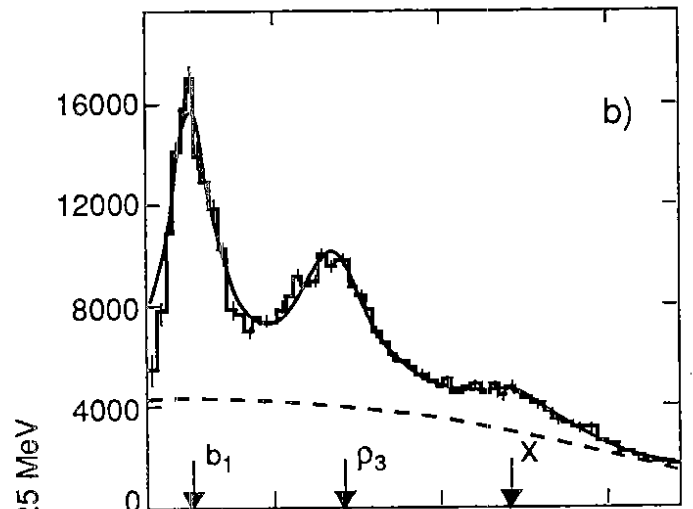
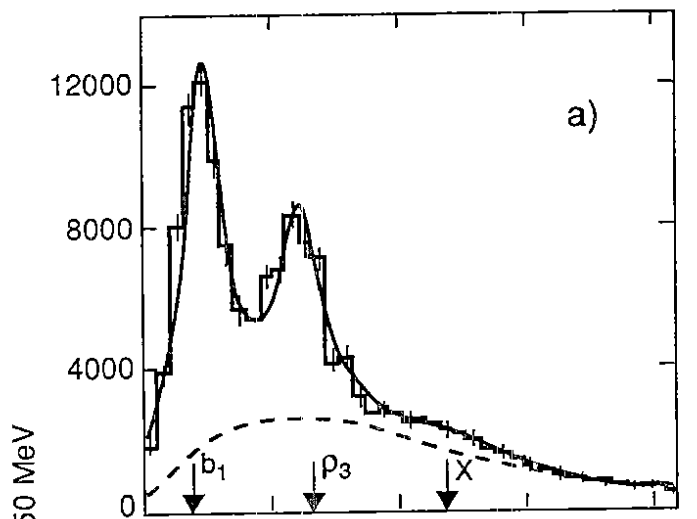


Fig. 2



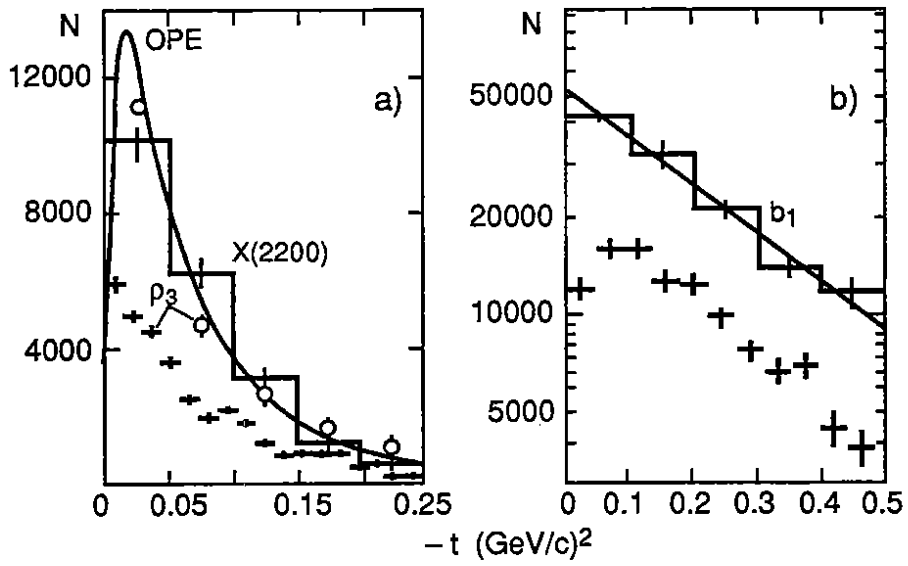


Fig. 3

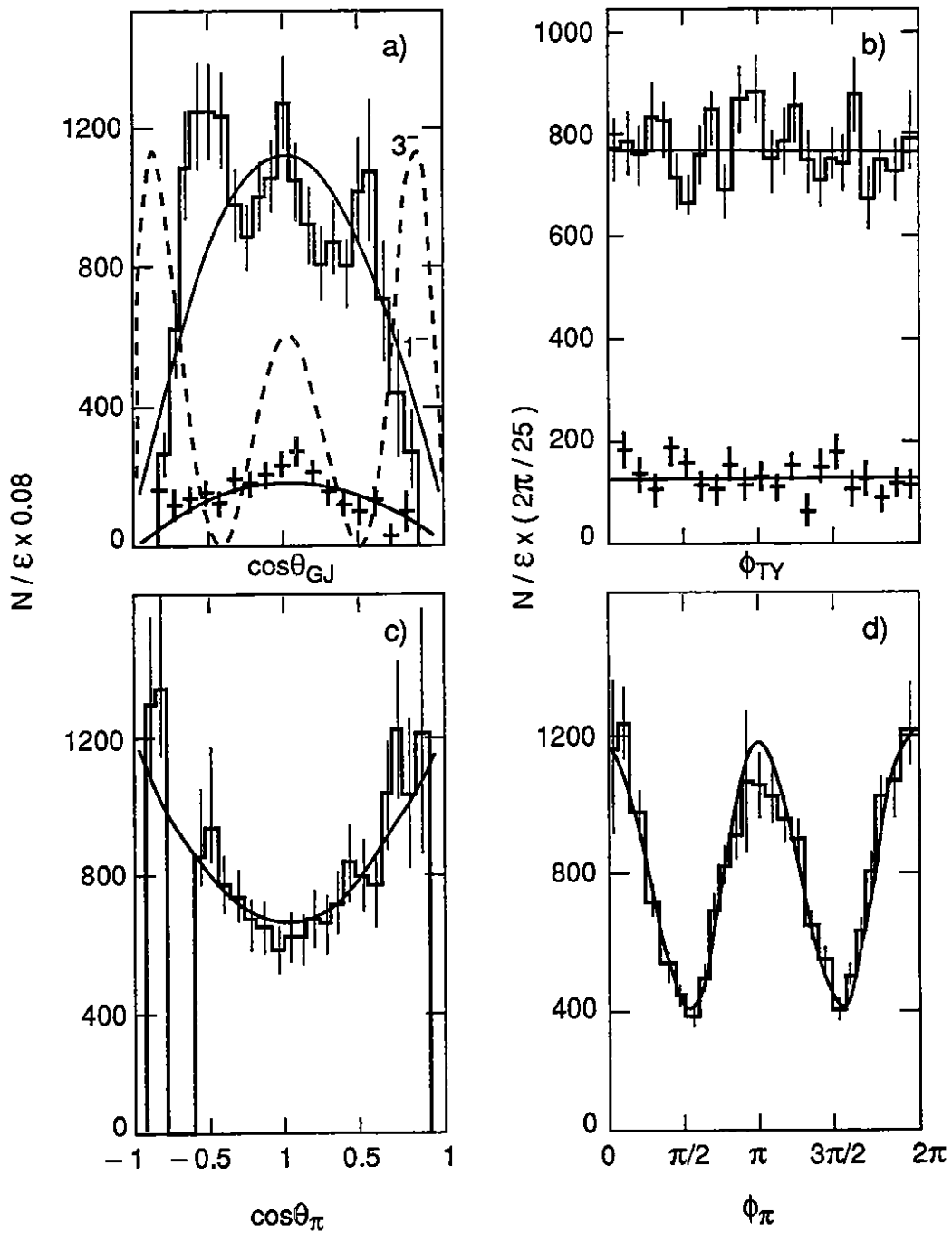


Fig. 4

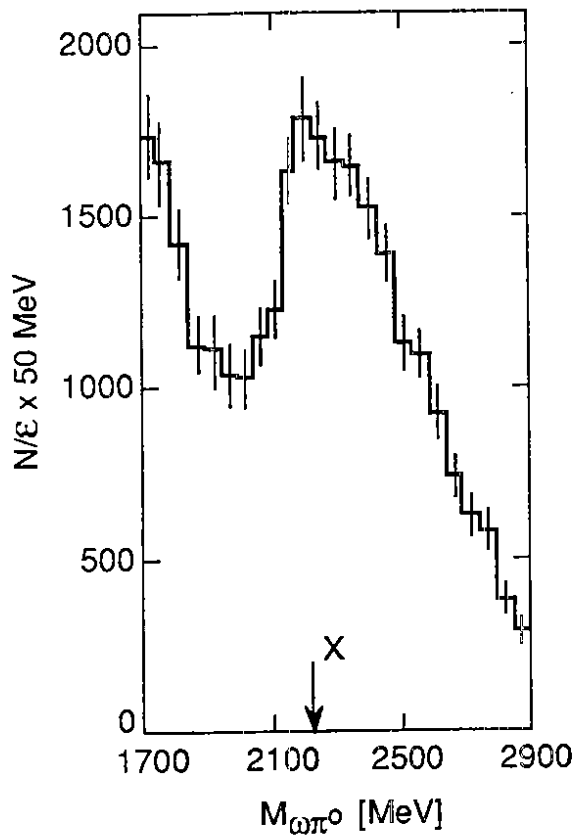


Fig. 5

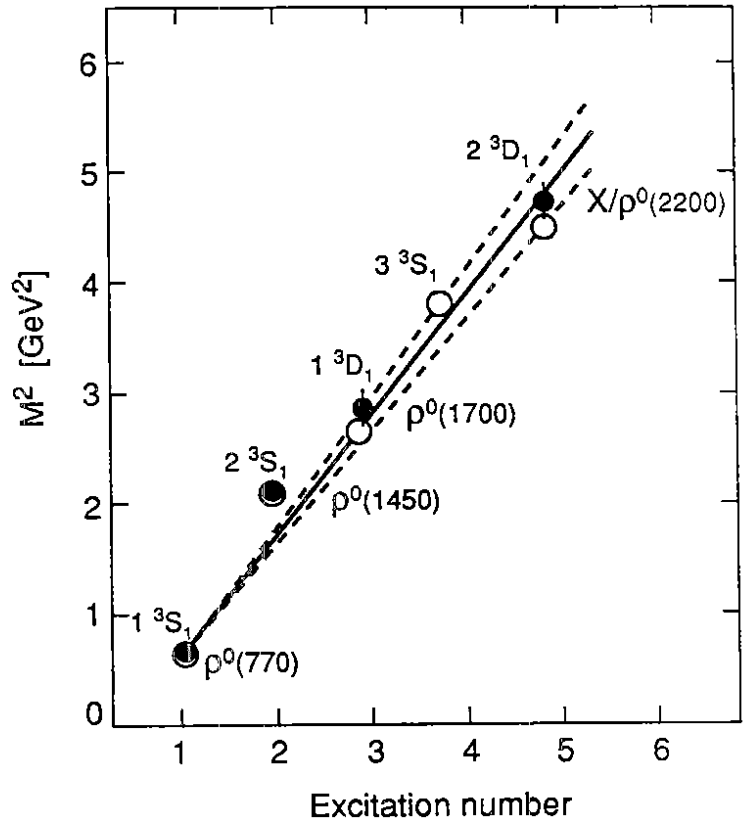


Fig. 6

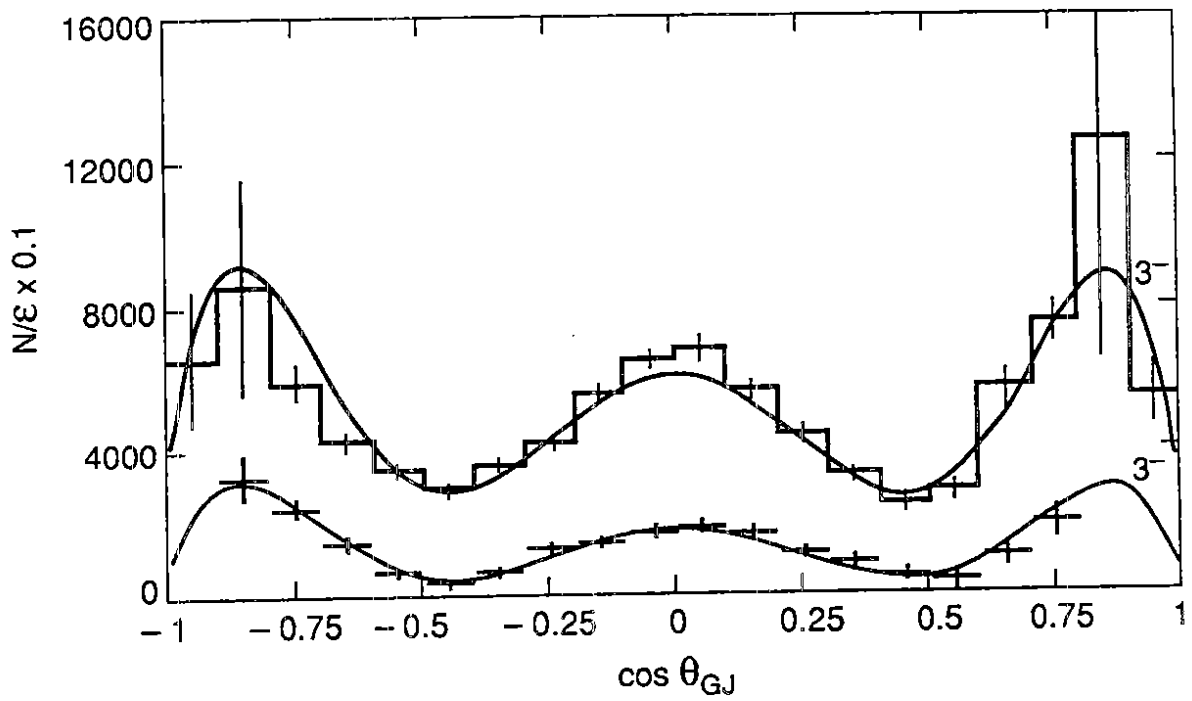


Fig. 7